

# DEVELOPMENT OF ROLLING TEXTURES IN AN AUSTENITIC STAINLESS STEEL

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Three dimensional texture analysis by means of orientation distribution functions (ODF) was used to examine the texture development during rolling at 473 K in an austenitic stainless steel. With the help of ODFs results, the different stages of texture development could be assigned to the existing theories of heterogeneous deformation mechanisms of low SFE face-centred cubic metals. The texture at very low degree of rolling consists of two limited orientation tubes with their fibre axes  $\langle 110 \rangle // ND$  and  $\langle 110 \rangle 60^\circ ND$  and agrees with the predictions made by Taylor model. With further deformation, twinning causes the reduction of  $\approx \{112\}\langle 111 \rangle$  component and leads to the formation of twin  $\{552\}\langle 115 \rangle$ . Abnormal slip on slip planes parallel to the twin boundaries rotates the twins into the  $\{332\}\langle 113 \rangle$  and  $\{111\}\langle 110 \rangle$  positions. The shear bands formation in the rotated twin-matrix lamellae changes their orientations near to  $\{011\}\langle 100 \rangle$  and  $\{011\}\langle 112 \rangle$  positions. Finally, normal slip again continues and sharpens the brass-type rolling texture.

**KEY WORDS** Rolling texture, copper-type, brass-type, slip, mechanical twinning, shear bands

## 1. INTRODUCTION

The study of the development of rolling textures in cubic metals in general and austenitic stainless steel in particular at ambient temperature, is of great scientific and technological importance, since these control further processing stages and thereby the resulting properties (i.e. Young's modulus or plastic anisotropy). It also gives more information on the active deformation mechanisms in polycrystalline materials.

There is a general consent in the literatures (Wassermann and Grewen, 1962; Lücke, 1981; Bacroix and Jonas, 1988; Hirsch and Lücke, 1988) that face-centred cubic (FCC) metals and alloys possess different rolling texture (i.e. copper-type and brass-type). The transition texture in between copper-type and brass-type is also perceptible. The possible causes for different rolling textures are mainly rolling degree, rolling temperature, starting texture, grain shape and finally the stacking fault energy (SFE). As the SFE increases the texture transition from brass to copper-type takes place and has been systematically surveyed as a function of solute element (Zn) in copper (Alam, Mengelberg and Lücke, 1967) and as a function of rolling temperature in austenitic stainless steel (Goodman

and Hu, 1964), copper (Hu and Goodman, 1963) and silver (Hu and Cline, 1961; Hu, Cline and Goodman, 1961). The FCC rolling textures may be described either in terms of ideal orientations (Hu, Sperry and Beck, 1952) or in terms of complete or limited fibre axis (Grewen and Wassermann, 1955; Wassermann, 1963). According to the concept of ideal orientations, the copper-type texture is characterised mainly by the orientations  $\{112\}\langle 111\rangle$ ,  $\{123\}\langle 634\rangle$  and  $\{011\}\langle 112\rangle$ , in doing so the indexing of  $\{123\}\langle 634\rangle$  component can adopt somewhat  $\{123\}\langle 412\rangle$  or  $\{124\}\langle 211\rangle$ , whereas the brass-type texture is dominated by  $\{011\}\langle 112\rangle$  and to a lesser extent by  $\{011\}\langle 100\rangle$ . The transition from copper to brass-type textures is described by the disappearance of  $\{112\}\langle 111\rangle$  component. On the contrary, the copper-type texture has been described by taking into consideration not only the ideal orientations  $\{112\}\langle 111\rangle$  and  $\{011\}\langle 112\rangle$  but also the spread ranges extended between them with the aid of two limited  $\langle 111\rangle$  fibres formed by rotations upto  $\pm 30^\circ$  about certain  $\{111\}$  poles of two ideal components. Out of the two fibre axes, one is taken under  $19^\circ$  to RD on the periphery of the  $\{111\}$  pole figure and the other under  $19^\circ$  from ND on the ND/RD radius of the pole figure.

The results obtained from the pole figure analysis were subsequently corroborated after the advent of orientation distribution functions (ODF) technique (Bunge, 1969) by many investigators (Bunge and Tobisch, 1968; Bunge and Haessner, 1968; Bunge, 1971; Davies, Kallend and Ruberg, 1973; Truszkowski *et al.*, 1973; Pospiech and Lücke, 1975; Virnich and Lücke, 1978; Hirsch and Lücke, 1988). According to these investigations, the rolling textures of FCC metals and alloys have been described in terms of two limited fibres (i.e. a fibre axis  $\langle 110\rangle//\text{ND}$  extending from  $\{011\}\langle 100\rangle$  to  $\{011\}\langle 112\rangle$  and the other fibre  $\langle 110\rangle 60^\circ\text{ND}$  running from  $\{112\}\langle 111\rangle$  through  $\{123\}\langle 634\rangle$  to  $\{011\}\langle 112\rangle$  with  $\langle 110\rangle$  axis inclined about  $60^\circ$  from ND towards RD). While the brass-type texture (low SFE materials) is characterised by  $\langle 110\rangle//\text{ND}$  fibre with  $\{011\}\langle 112\rangle$  as a major component, the copper-type (high and medium SFE materials) is described mainly by  $\langle 110\rangle 60^\circ\text{ND}$  fibre with  $\{112\}\langle 111\rangle$  or  $\{123\}\langle 634\rangle$  as major orientation.

Extensive metallographic studies of the deformation structure by optical and transmission electron microscopy (TEM) have been performed on FCC metals like copper and brass and these microscopical results lead to a number of new aspects in the interpretation of deformation structure and texture development (Duggan *et al.*, 1978; Hatherly, 1978; Malin and Hatherly, 1979; Wakefield and Hatherly, 1981). According to Malin and Hatherly (1979) the deformation sequence in high and medium SFE metals (i.e. copper) involves the formation of equiaxed cells, microbands, clustering of microbands and shear bands. In these metals copper-type rolling texture forms which can be determined theoretically from the calculations based on Taylor model (Dillamore, Butler and Green, 1968; Bunge, 1970; Dillamore and Katoh, 1971; Dillamore and Katoh, 1974), Taylor model adapted to both slip and twinning (Van Houtte and Aernoudt, 1976; Van Houtte, 1978a; Van Houtte, 1978b) and relaxed constraints model (Honneff and Mecking, 1981; Van Houtte, 1981). With decrease in SFE, an increasing tendency of mechanical twinning during deformation is often considered as the mechanism (Wassermann, 1963) which causes the change from the copper-type to the brass-type rolling texture. In metals of low SFE such as 70–30 brass (Fargette and Whitwham, 1976; Grewen, Noda and Sauer, 1977; Duggan *et al.*, 1978;

Hutchinson, Duggan and Hatherly, 1979) and stainless steel (Blicharski and Gorczyca, 1978), deformation begins by the glide of partial dislocations, followed by mechanical twinning, the rotation of the twin-parent lamellae into alignment with the rolling plane, the development of shear bands and the resumption of normal octahedral slip in the recovered crystallites of the shear bands. But the effects of the mechanical twinning on further deformation and on the formation of rolling textures in austenitic stainless steel, which become even more complicated by shear band formation, are not yet completely understood. Furthermore, the characterisation of the rolling textures developed in austenitic stainless steel above ambient temperature was based on the pole figure analysis (Goodman and Hu, 1964) which possesses very limited resolving power in revealing the differences which are more of quantitative than of qualitative nature.

The aim of the present investigation was to clarify the microscopical observations of orientations and their statistical appearance in the textures of austenitic stainless steel cold rolled at 473 K by means of the three-dimensional ODF analysis which describes the texture in a more complete, exact and explicit fashion. The results also outline the process of texture development during cold rolling.

## 2. EXPERIMENTAL PROCEDURE

A commercially produced hot band (HB) of an austenitic stainless steel of nominal composition (by wt%): 19.25% Cr, 8.40% Ni, 0.05% C, 1.24% Mn, 0.48% Si, 0.02% P, 0.02% S was used for the present investigation. As per supplier's information 2.95 mm thick HB strip was finished at about 1173 K and coiled at about 1023 K. In order to remove any oxides and scales from the surfaces, the hot band was pickled in a pickling solution (by volume %) of 10%  $\text{HNO}_3$  + 2% HF + 88% distilled water at a temperature 323 K for a period of 5 minutes and then rinsed with water and dried.

The hot band was cold rolled (CR) at a temperature 473 K without lubrication. The cold rolling operation was oriented parallel to the original hot rolling direction and carried out in such a way that the true strain per pass was about 0.05. In the interval between the passes specimens were kept in the furnace at 473 K for a period of 5 minutes in order to maintain the temperature. The two-high mill used had 55 mm dia rollers. The cumulative rolling reductions investigated were from 10 to 90% at intervals of 10%.

Optical metallography was performed on specimen of HB, which was mechanically polished and electrolytically etched in saturated oxalic acid solution at 10 volts for 2 minutes, at section ( $S = 0$ ) cut perpendicular to the transverse direction containing (R) and normal direction (N) RN. Here  $S = 0$  denotes the mid-thickness of the specimen.

X-ray texture measurements were performed for HB rolled at 0, 30, 50, 70 and 90% reduction in thickness on an automatic texture goniometer, using  $\text{MoK}_\alpha$  radiation. For each ground and etched specimen (20 mm  $\times$  14 mm), the texture was determined at the section ( $S = 0$ ) parallel to the rolling plane by measuring four incomplete pole figures of the plane {111}, {200}, {220} and {113}, using

Schulz back reflection technique (Schulz, 1949). The diffracted intensity was recorded continuously every  $5^\circ$  along concentric circles in the angular range from  $0$  to  $75^\circ$  in steps of  $5^\circ$ . The measured intensity was subjected to background, geometrical and defocussing corrections using a random specimen of pressed and sintered austenitic stainless steel powders. Each ODF was calculated from the data of four incomplete pole figures following the series expansion method of Bunge (1982) and using the pseudo-normalisation technique of Kern and Bergmann (1978). The series were extended to the degree  $l=22$ . The texture strength as expressed by the texture index  $J$  (Truszkowski *et al.*, 1978) which defines the mean square deviation of the ODF from random distribution, was also determined in each case. The inverse pole figure measurements were also made for the cold rolled specimens parallel to the rolling plane. For these, the measured integrated intensities of selected reflections were used for determining the pole densities of austenite grains having their plane normals parallel to normal direction (ND) of the sample (Singh, 1989) and the pole density of austenite grains was expressed as

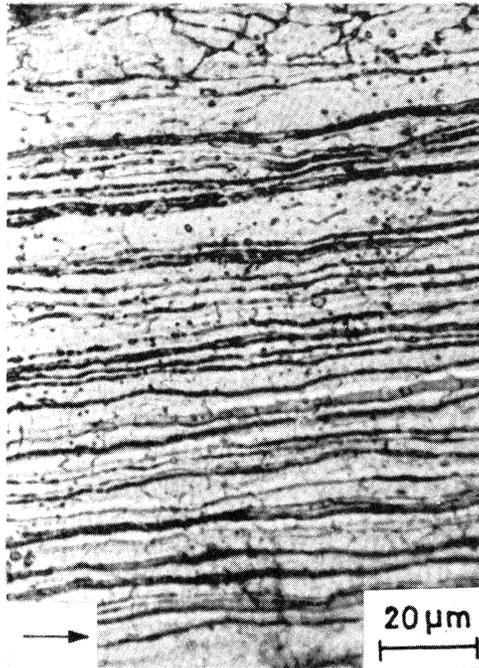
$$\frac{I_{hkl}/R_{hkl}}{\frac{1}{m} \sum I_{hkl}/R_{hkl}}$$

where  $I$ ,  $R$  and  $m$  represent the integrated intensity, theoretically calculated factor and number of reflections respectively.

### 3. EXPERIMENTAL RESULTS

The industrially produced hot band of austenitic stainless steel was fully austenitic. The optical micrograph of the HB in the RN section near the centre ( $S=0$ ) is shown in Figure 1 and indicates mostly the elongated grains. Some recrystallised grains are also present. The dark bands in the micrograph appear to be thin sheet like bands of heavily deformed austenite grains. The average grain size measured on the section ( $S=0$ ) parallel to the rolling plane was about  $7 \mu\text{m}$ .

Figure 2 shows the three dimensional orientation density distribution (i.e. ODF) of crystallites for the HB in constant  $\phi_2$  sections through Euler space. This ODF exhibits maxima of orientations which are given in Table 1 along with their Euler angles ( $\phi_1$ ,  $\phi$ ,  $\phi_2$ ). The main features of the texture components present in this ODF are rolling (copper-type) and recrystallisation textures. Singh, Ramaswamy and Suryanarayana (1991) have discussed the texture evolution in the present steel during hot rolling in detail and characterised the rolling texture components, which were retained, as  $\{011\}\langle 112\rangle$ ,  $\approx\{123\}\langle 634\rangle$  and  $\{112\}\langle 111\rangle$  orientations and also described the orientations  $\{001\}\langle 100\rangle$ , RD rotated cubes  $\{013\}\langle 100\rangle$ ,  $\{012\}\langle 100\rangle$  and Goss  $\langle 011\rangle\langle 100\rangle$  as the recrystallised components. The orientation elements of the retained rolling texture are distributed along an orientation tube which is inclined to constant  $\phi_2$  sections and extends from  $\{011\}\langle 112\rangle$  at  $\phi_1=35^\circ$ ,  $\phi=45^\circ$  and  $\phi_2=0^\circ$  to  $\{112\}\langle 111\rangle$  at  $\phi_1=39^\circ$ ,  $\phi=66^\circ$  and  $\phi_2=27^\circ$ , or symmetrically equivalent and more convenient, from  $\{112\}\langle 111\rangle$  at  $\phi_1=90^\circ$ ,  $\phi=35^\circ$ , and  $\phi_2=45^\circ$  to  $\{011\}\langle 112\rangle$  at  $\phi_1=35^\circ$ ,  $\phi=45^\circ$  and  $\phi_2=90^\circ$ . The skeleton line of this tube also includes the orientation  $\approx\{123\}\langle 634\rangle$



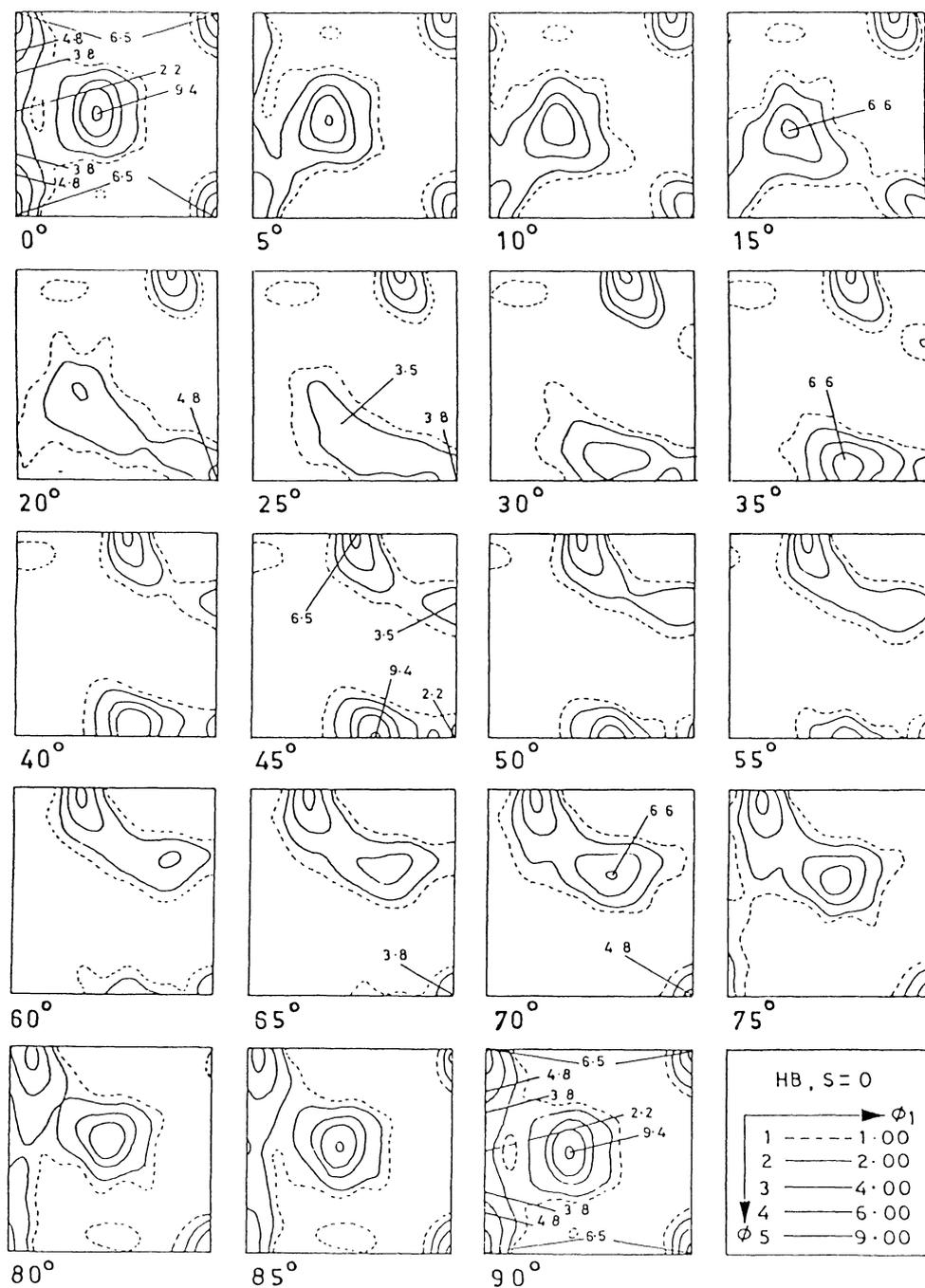
**Figure 1** Optical micrograph of RN section of the hot band (HB) at centre ( $S = 0$ ): arrow parallel to rolling direction.

at  $\phi_1 = 59^\circ$ ,  $\phi = 37^\circ$  and  $\phi_2 = 63^\circ$  and has its  $\langle 110 \rangle$  fibre axis inclined  $\sim 60^\circ$  from ND towards RD ( $\langle 110 \rangle 60^\circ \text{ND}$ ). On the other hand, the components of the recrystallised texture are scattered along  $\langle 100 \rangle // \text{RD}$  fibre which runs from cube  $\{001\}\langle 100 \rangle$  to Goss  $\{011\}\langle 100 \rangle$  positions for the constant  $\phi_2 = 0^\circ$  section (Figure 2). The major components of the rolling texture and the recrystallisation texture are  $\{011\}\langle 112 \rangle$  and  $\{001\}\langle 100 \rangle$  respectively.

Upon rolling the HB at 473 K, very limited amount of  $\alpha'$ -martensite (about 16.0 vol.% at 90% reduction) was produced at the centre, which was not sufficient for its texture measurements (Singh, 1989).

The ND inverse pole figure densities of the selected planes [i.e.  $(111)_\gamma$ ,  $(200)_\gamma$ ,  $(220)_\gamma$  and  $(311)_\gamma$ ] parallel to the rolling plane of the cold rolled specimens, with increasing rolling reductions (range: 10–90%) are shown in Figure 3. Although, the pole density of the plane  $(220)_\gamma$  is relatively higher than that of the other planes, it increases rapidly between the reductions 10 to 50%, and thereafter, remains more or less constant upto 80% reduction and then there is an increasing trend after 80% reduction. In contrast, the pole density of  $(111)_\gamma$  remains relatively constant upto 40% reduction, and thereafter, increases to a maximum value at 80% reduction and then, there is a decreasing trend after 80% reduction. The pole densities of other two planes  $(311)_\gamma$  and  $(200)_\gamma$  are continuously decreasing with increasing reduction.

The three dimensional orientation density distributions (i.e. ODFs) of the HB cold rolled by 30, 50, 70 and 90% reductions, are presented in constant  $\phi_2$

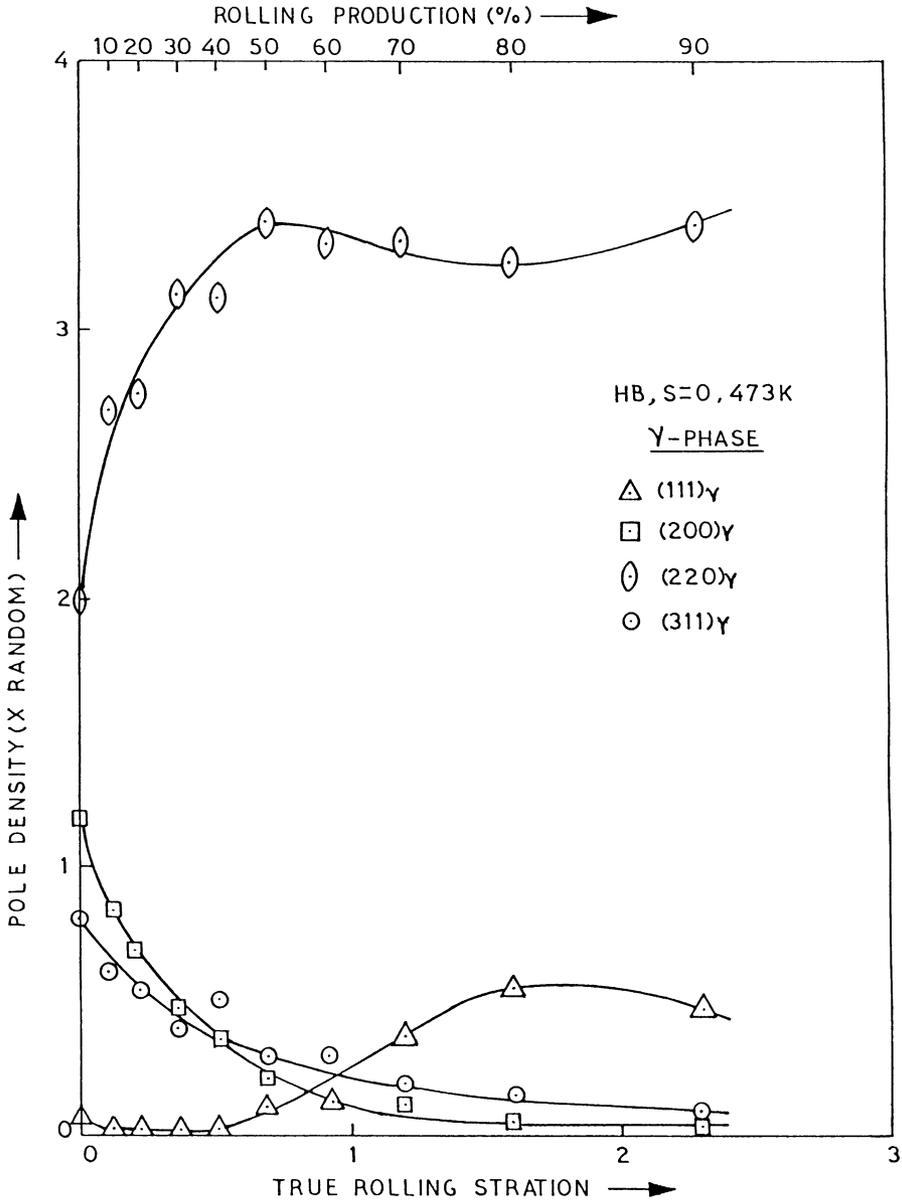


**Figure 2** ODF for central level ( $S=0$ ) of the hot band (HB) in constant  $\phi_2$  sections.  $J=3.7$ .

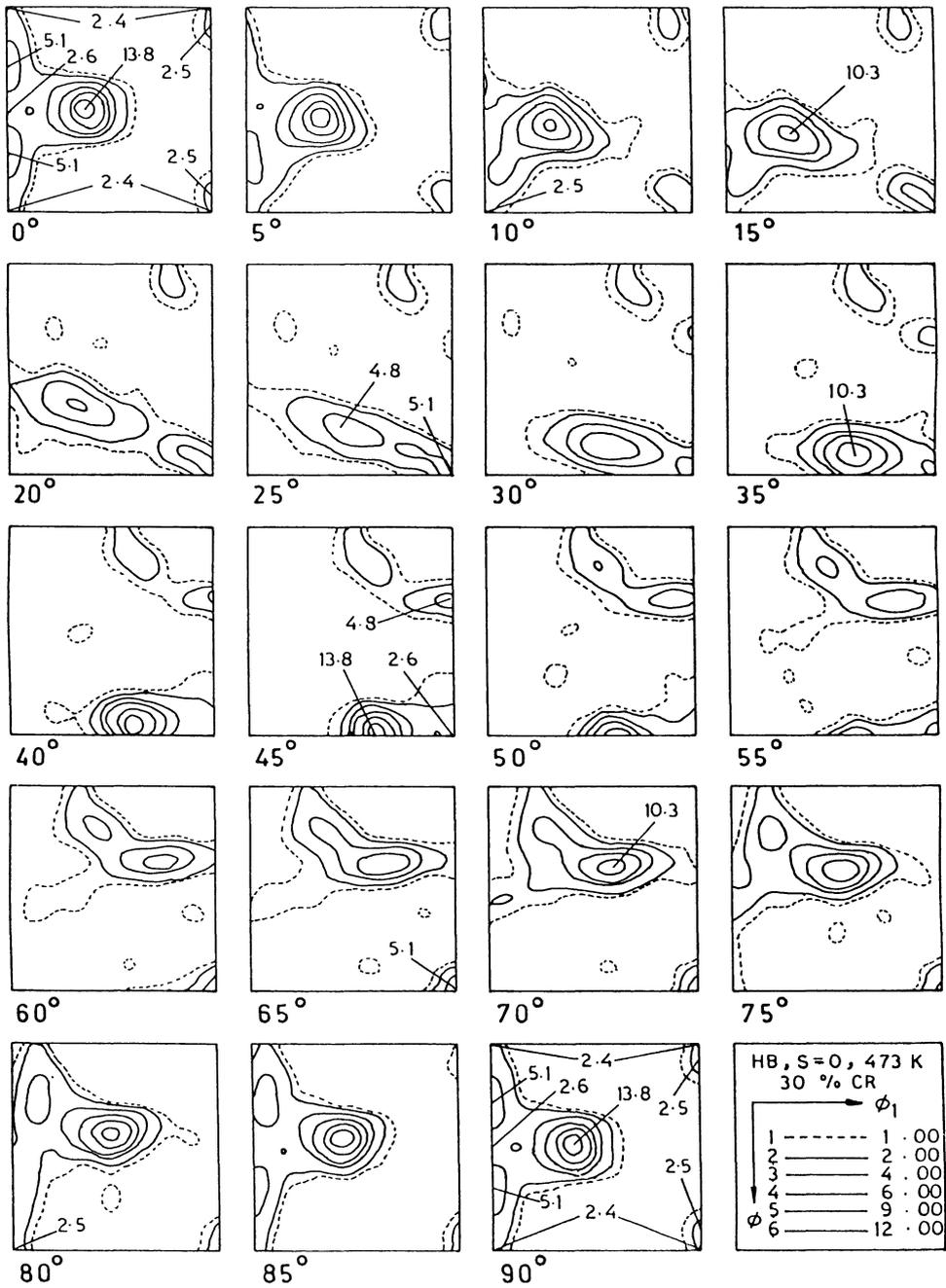
**Table 1** Orientation density of the texture components  $\{HKL\}\langle UVW \rangle$  at the centre of the hot band

<i>HKL</i>	<i>UVW</i>	<i>Euler Angles</i>			<i>Orientation density</i>
		$\phi_1$	$\phi$	$\phi_2$	
<i>Retained rolling texture</i>					
{011}	$\langle 112 \rangle$	35	45	0	9.4
{134}	$\langle 956 \rangle$	54	38	72	6.6
(equal to)					
{123}	$\langle 634 \rangle$	59	37	63	
{112}	$\langle 111 \rangle$	90	35	45	3.5
<i>Recrystallised texture</i>					
{001}	$\langle 100 \rangle$	0	0	0	6.5
{013}	$\langle 100 \rangle$	0	18	0	4.8
{012}	$\langle 100 \rangle$	0	27	0	3.8
{011}	$\langle 100 \rangle$	0	45	0	2.2

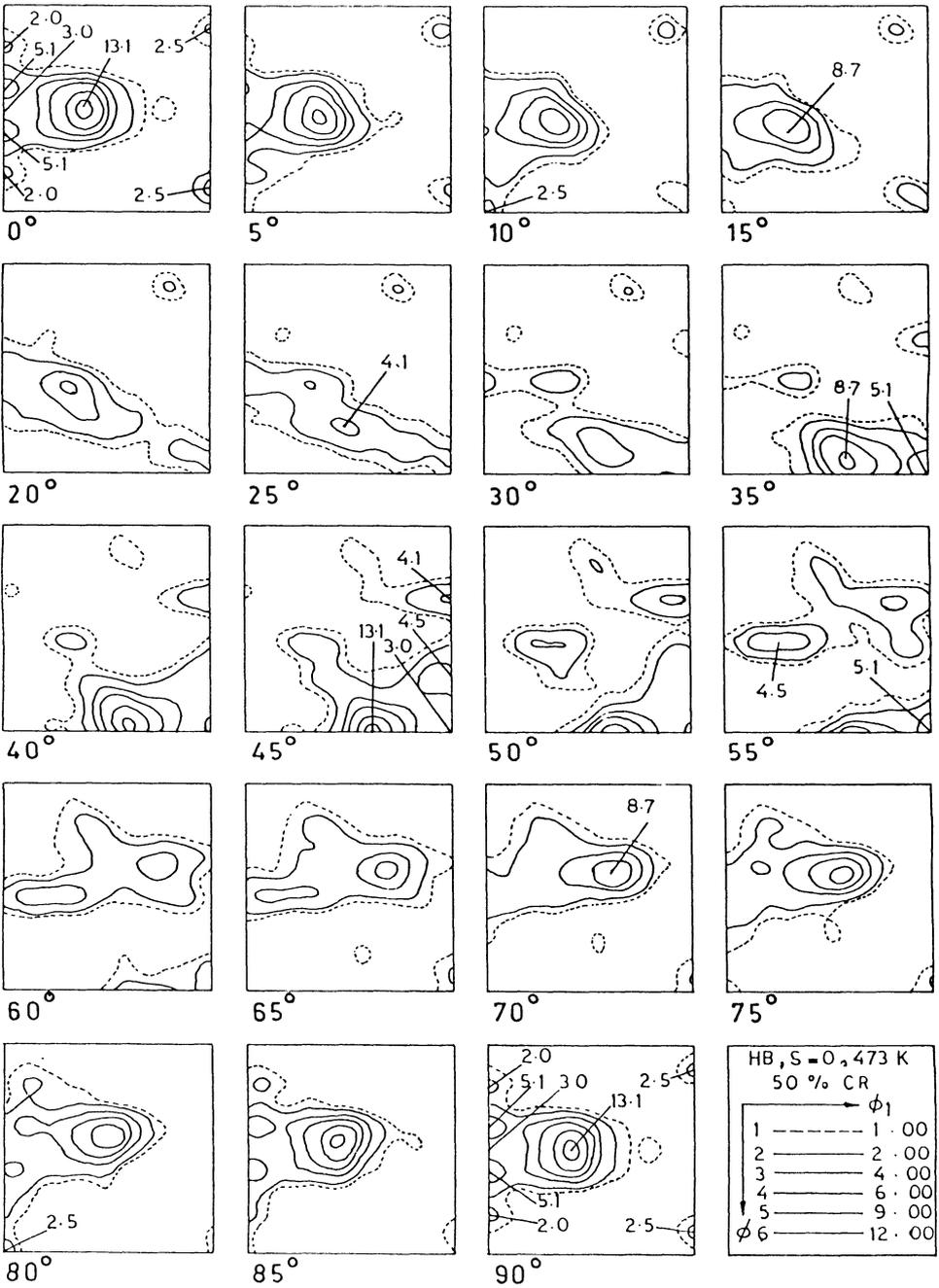
sections through Euler space in Figures 4(a), (b), (c) and (d) respectively. These ODF's exhibit maxima which are listed in Table 2 along with their Euler angles. While the exact orientations of the maxima are listed in the table, in the text approximate orientations are used (indicated by  $\approx$ ) if the indices are clearly simpler. At low degrees of rolling [Figures 4(a) and (b)] the texture components (i.e.  $\{011\}\langle 112 \rangle$ ,  $\approx\{123\}\langle 634 \rangle$ ,  $\approx\{112\}\langle 111 \rangle$  and  $\{011\}\langle 100 \rangle$ ) are distributed mainly along the two limited orientation tubes in the Euler space. One of the tubes runs parallel to  $\phi_1$  through  $\phi_2 = 0^\circ$ ,  $\phi = 45^\circ$  and includes the orientations  $\{011\}\langle 100 \rangle$  at  $\phi_1, \phi, \phi_2 = 0, 45$  and  $0^\circ$  and  $\{011\}\langle 112 \rangle$  at  $\phi_1, \phi, \phi_2 = 35, 45$  and  $0^\circ$ . This tube has its axis  $\langle 110 \rangle // ND$  and the maxima in the orientation density along this tube is observed at  $\{011\}\langle 112 \rangle$  component. This maximum in orientation density at  $\{011\}\langle 112 \rangle$  is partly due to the effect of initial starting texture of the HB (Figure 2) which shows the maximum in density at  $\{011\}\langle 112 \rangle$  component. The texture index  $J$  has increased from the value of 3.7 (HB) to a value of 5.6 after cold rolling the HB by 50% reduction in thickness. Figure 5 shows the distribution of orientation densities along  $\langle 110 \rangle // ND$  fibre axis of the above tube. The other tube, inclined to constant  $\phi_2$  sections, runs from  $\{011\}\langle 112 \rangle$  at  $\phi_1, \phi, \phi_2 = 35, 45$  and  $0^\circ$  to  $\approx\{112\}\langle 111 \rangle$  at  $\phi_1, \phi, \phi_2 = 39, 66$  and  $27^\circ$ , or symmetrically equivalent from  $\approx\{112\}\langle 111 \rangle$  at  $\phi_1, \phi, \phi_2 = 90, 35$  and  $45^\circ$  to  $\{011\}\langle 112 \rangle$  at  $\phi_1, \phi, \phi_2 = 35, 45, 90^\circ$ . The skeleton line of this tube also includes the orientation  $\approx\{123\}\langle 634 \rangle$  at  $\phi_1, \phi, \phi_2 = 59, 37$  and  $63^\circ$  and has its  $\langle 110 \rangle$  fibre axis inclined  $\sim 60^\circ$  from ND towards RD ( $\langle 110 \rangle 60^\circ ND$ ). Figure 6 depicts the orientation density distribution along this tube. It can be inferred from the Figures 5 and 6 that for very low degrees of rolling (e.g.  $\leq 30\%$ ), the densities for all the texture components along the two limited tubes (i.e.  $\langle 110 \rangle // ND$  and  $\langle 110 \rangle 60^\circ ND$  fibres) increase over their whole length. The increase in density is more at  $\{011\}\langle 112 \rangle$  component than that at  $\{011\}\langle 100 \rangle$ ,  $\approx\{112\}\langle 111 \rangle$  and  $\approx\{123\}\langle 634 \rangle$  components. At 50% rolling reduction the densities at all texture components along  $\langle 110 \rangle // ND$  fibre increase very slightly and at  $\{011\}\langle 112 \rangle$  component it remains more or less constant, whereas the densities of  $\approx\{112\}\langle 111 \rangle$  and  $\approx\{123\}\langle 634 \rangle$  components of  $\langle 110 \rangle 60^\circ ND$  fibre have started decreasing. At intermediate degrees of rolling (e.g.  $\leq 70\%$ ), the orientation



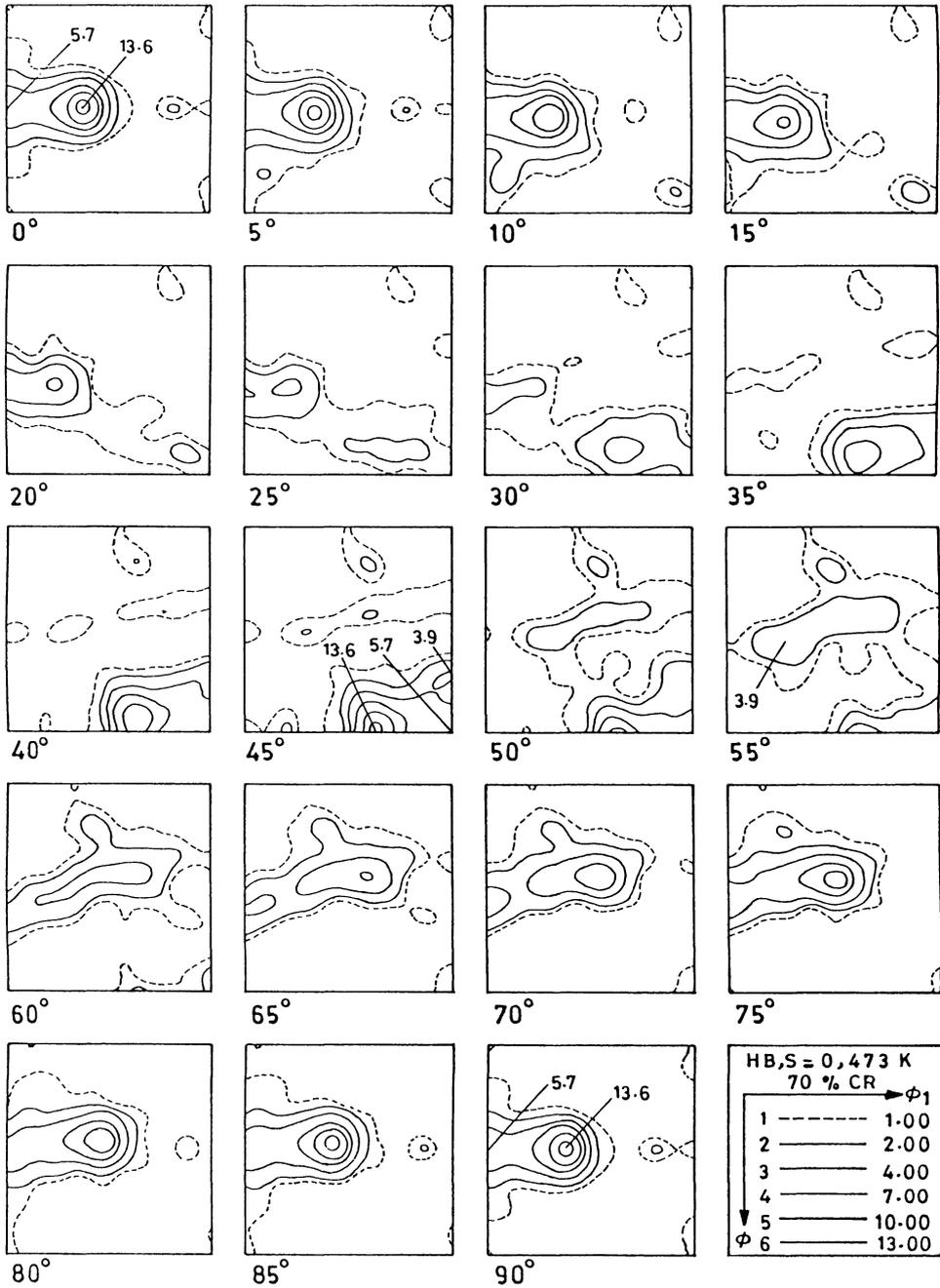
**Figure 3** ND inverse pole figure data showing densities of selected planes parallel to the sheet surface of the hot band (HB) at centre (S=0) with increasing rolling reduction at 473 K.



**Figure 4** ODFs for central level ( $S=0$ ) of the hot band (HB) cold rolled at 473 K in constant  $\phi_2$  sections: (a) 30% reduction,  $J=5.6$ , (b) 50% reduction,  $J=5.6$ , (c) 70% reduction,  $J=5.7$  and (d) 90% reduction,  $J=6.8$ .

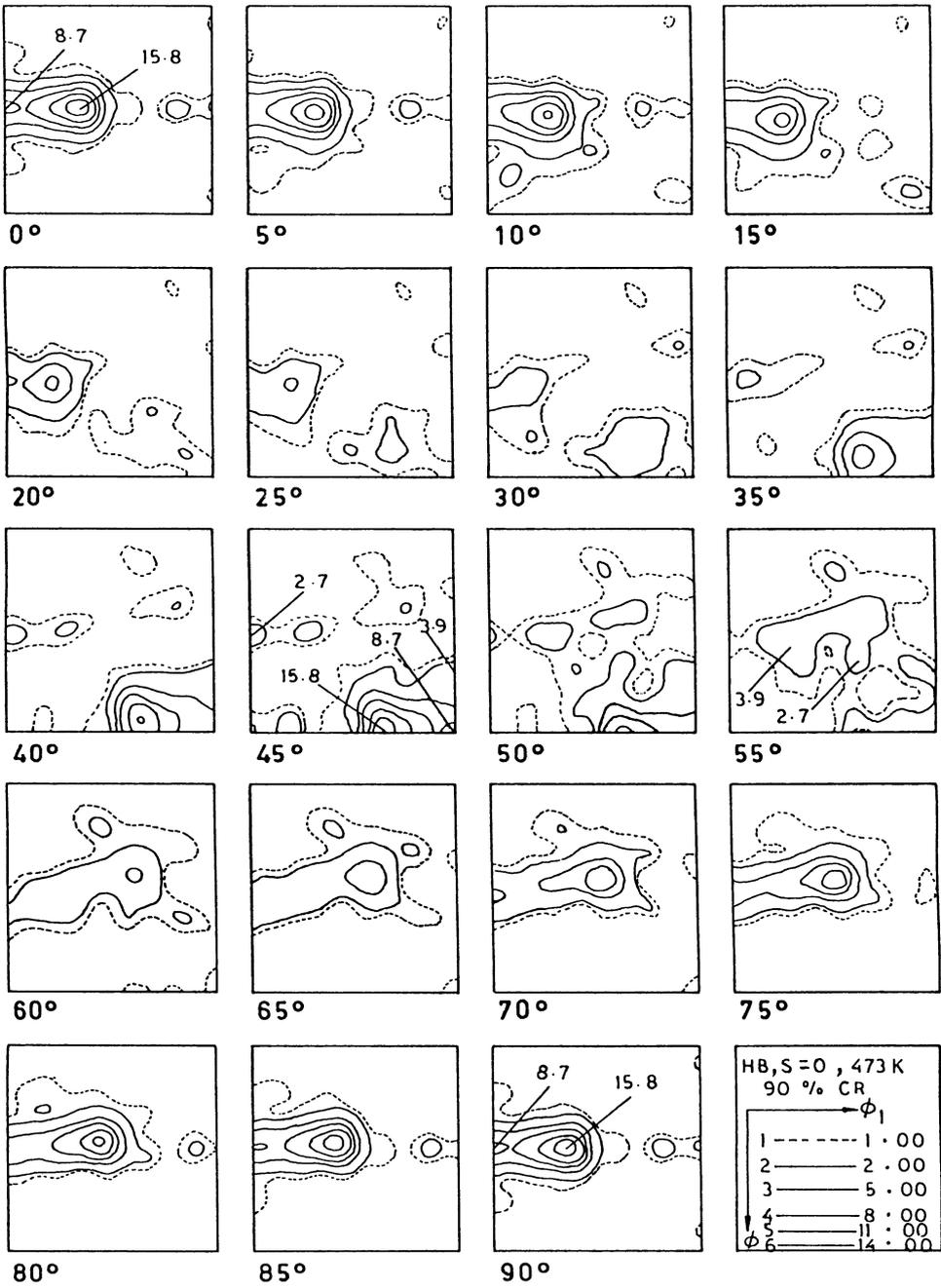


**Figure 4** ODFs for central level ( $S=0$ ) of the hot band (HB) cold rolled at 473 K in constant  $\phi_2$  sections: (a) 30% reduction,  $J=5.6$ , (b) 50% reduction,  $J=5.6$ , (c) 70% reduction,  $J=5.7$  and (d) 90% reduction,  $J=6.8$ .



(c)

**Figure 4** ODFs for central level ( $S=0$ ) of the hot band (HB) cold rolled at 473 K in constant  $\phi_2$  sections: (a) 30% reduction,  $J=5.6$ , (b) 50% reduction,  $J=5.6$ , (c) 70% reduction,  $J=5.7$  and (d) 90% reduction,  $J=6.8$ .



(d)

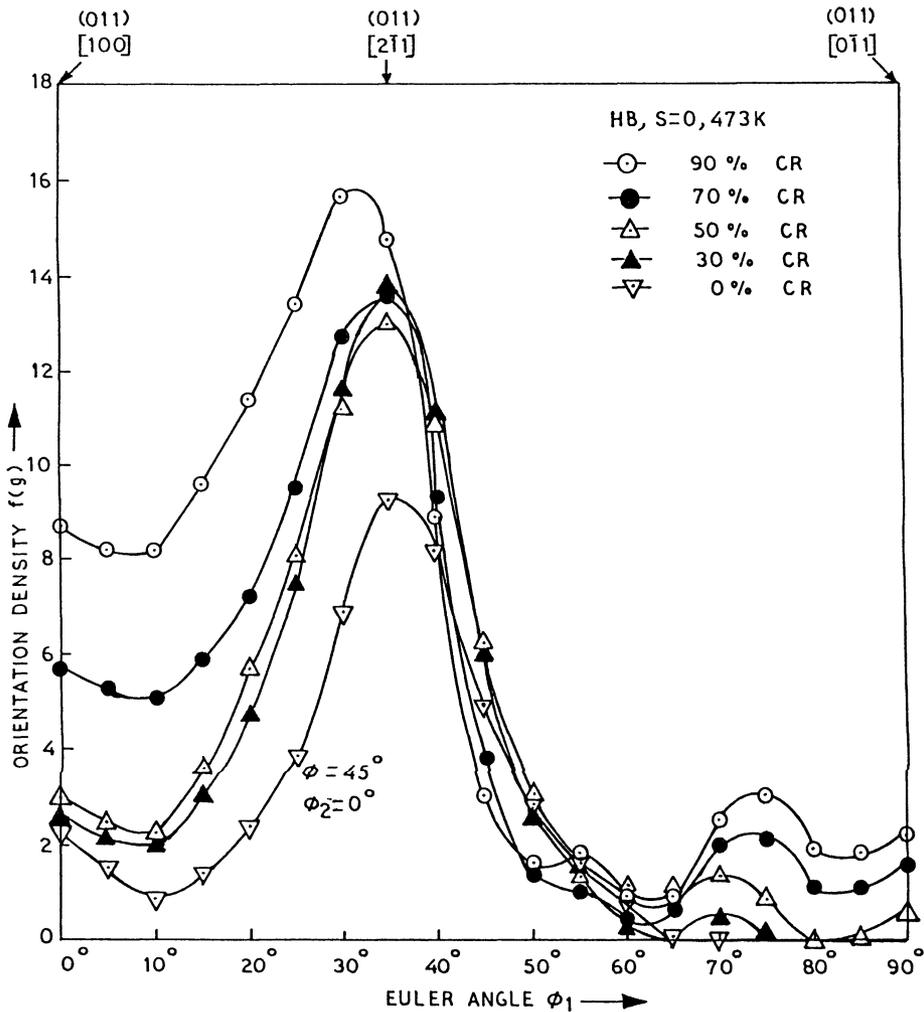
**Figure 4** ODFs for central level ( $S=0$ ) of the hot band (HB) cold rolled at 473 K in constant  $\phi_2$  sections: (a) 30% reduction,  $J=5.6$ , (b) 50% reduction,  $J=5.6$ , (c) 70% reduction,  $J=5.7$  and (d) 90% reduction,  $J=6.8$ .

**Table 2** Orientation density of texture components of austenite at the centre of the hot band cold rolled by 30, 50, 70 and 90% reduction at 473 K

HKL	UVW	Euler Angles			Orientation density			
		$\phi_1$	$\phi$	$\phi_2$	30% CR	50% CR	70% CR	90% CR
<i>Rolling texture components</i>								
{011}	$\langle 112 \rangle$	35	45	0	13.8	13.1	13.6	15.8
{134}	$\langle 734 \rangle$	49	38	72	10.3	8.7	—	—
	(equal to)							
{123}	$\langle 634 \rangle$	59	37	63				
{225}	$\langle 554 \rangle$	90	30	45	4.8	4.1	1.6	1.3
	(equal to)							
{112}	$\langle 111 \rangle$	90	35	45				
{011}	$\langle 100 \rangle$	0	45	0	2.6	3.0	5.7	8.7
{552}	$\langle 115 \rangle$	90	74	45	1.9	—	—	—
{332}	$\langle 113 \rangle$	90	65	45	1.2	5.0	4.1	4.6
{334}	$\langle 110 \rangle$	0	47	45	—	—	1.5	2.7
	(equal to)							
{111}	$\langle 110 \rangle$	0	55	45				
{8 8 11}	$\langle 4 4 11 \rangle$	90	63	45	—	4.5	3.9	3.9
	(equal to)							
{111}	$\langle 112 \rangle$	90	55	45				
<i>Recrystallised texture components</i>								
{001}	$\langle 100 \rangle$	0	0	0	2.4	—	—	—
{013}	$\langle 100 \rangle$	0	18	0	—	2.0	—	—
{025}	$\langle 100 \rangle$	0	22	0	—	—	1.4	1.4
{012}	$\langle 100 \rangle$	0	27	0	5.1	—	—	—
{023}	$\langle 100 \rangle$	0	34	0	—	5.1	—	—
{0 2 11}	$\langle 11 2 0 \rangle$	90	10	0	2.5	2.5	1.6	1.0

density at  $\{011\}\langle 100 \rangle$  component increases more than that at  $\{011\}\langle 112 \rangle$  component, while the densities at  $\approx\{112\}\langle 111 \rangle$  and  $\approx\{123\}\langle 634 \rangle$  components of  $\langle 110 \rangle 60^\circ \text{ND}$  fibre have decreased considerably. At high degrees of rolling (e.g. 90%) the orientation densities for all texture components of  $\langle 110 \rangle // \text{ND}$  fibre have increased considerably.

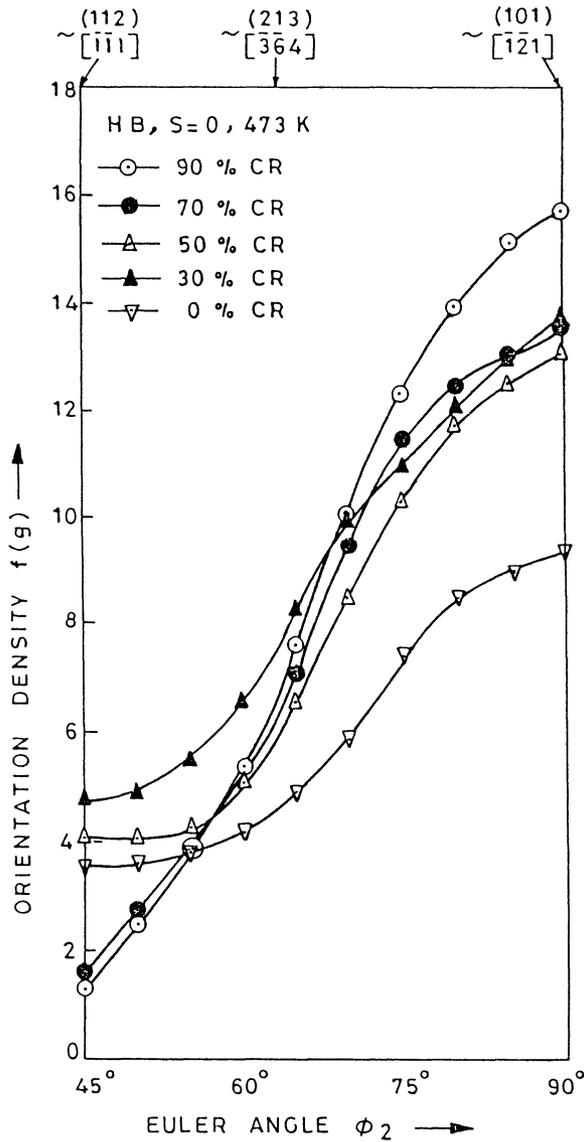
The reduction in density of  $\approx\{112\}\langle 111 \rangle$  component at  $\sim 50\%$  reduction starts the transition of rolling texture from copper-type to brass-type. This transition is known to be caused by mechanical twinning by which the component  $\approx\{112\}\langle 111 \rangle$  twins to a position  $\{552\}\langle 112 \rangle$  near  $\{011\}\langle 100 \rangle$  component (Wassermann, 1963). The distribution of orientation densities of the component  $\approx\{112\}\langle 111 \rangle$ , its twin  $\{552\}\langle 115 \rangle$ ,  $\{011\}\langle 100 \rangle$  and other orientations associated with further deformation by slip is shown in Figure 7 which mainly reveals the plot of orientation densities of components with a common axis  $\langle 110 \rangle // \text{TD}$  fibre axis against  $\phi$  at  $\phi_1 = 90^\circ$ ,  $\phi_2 = 45^\circ$ . The  $\langle 110 \rangle // \text{TD}$  fibre axis includes the component  $\approx\{112\}\langle 111 \rangle$  at  $\phi = \sim 30^\circ$ , its twin at  $\phi = 75^\circ$ , the component  $\{011\}\langle 100 \rangle$  at  $\phi = 90^\circ$  and the orientation  $\{111\}\langle 112 \rangle$  which is an intermediate component at  $\phi = 55^\circ$ . For lower degrees of rolling, the density of the component  $\approx\{112\}\langle 111 \rangle$  increases upto  $\sim 30\%$  reduction and then starts decreasing at about 50% reduction. Simultaneously with the reduction in the density of the component  $\approx\{112\}\langle 111 \rangle$  the orientation density of its twin position  $\{552\}\langle 115 \rangle$  starts increasing and also the twin orientation  $\{552\}\langle 115 \rangle$  has shifted its position



**Figure 5** Orientation densities along  $\langle 110 \rangle // ND$  fibre of the orientation tube at the centre level ( $S=0$ ) of the hot band (HB) cold rolled by 0, 30, 50, 70 and 90% reductions at 473 K.

towards the orientation  $\{332\}\langle 113 \rangle$  at  $\phi = 65^\circ$ , whereas there is very little increase in the density of the component  $\{011\}\langle 100 \rangle$ . On further deformation the density of the component  $\{332\}\langle 113 \rangle$  decreases (e.g. 70% reduction) and then increases (e.g. 90% reduction), while that of the component  $\{011\}\langle 100 \rangle$  progressively increases. In the same deformation range (e.g. 70–90%), an increase in the density of the component  $\{111\}\langle 110 \rangle$  has also been detected [Figures 4(c) and (d)].

Furthermore, it is clearly visible from the ODFs for HB cold rolled by 0, 30 and 50% reductions [Figures 2, 4(a) and 4(b)] that the initial scattering of the cube orientation  $\{001\}\langle 100 \rangle$  towards Goss orientation  $\{011\}\langle 100 \rangle$  in  $\phi$  direction at  $\phi_1 = \phi_2 = 0^\circ$ , has started shrinking towards the component  $\{011\}\langle 100 \rangle$  with

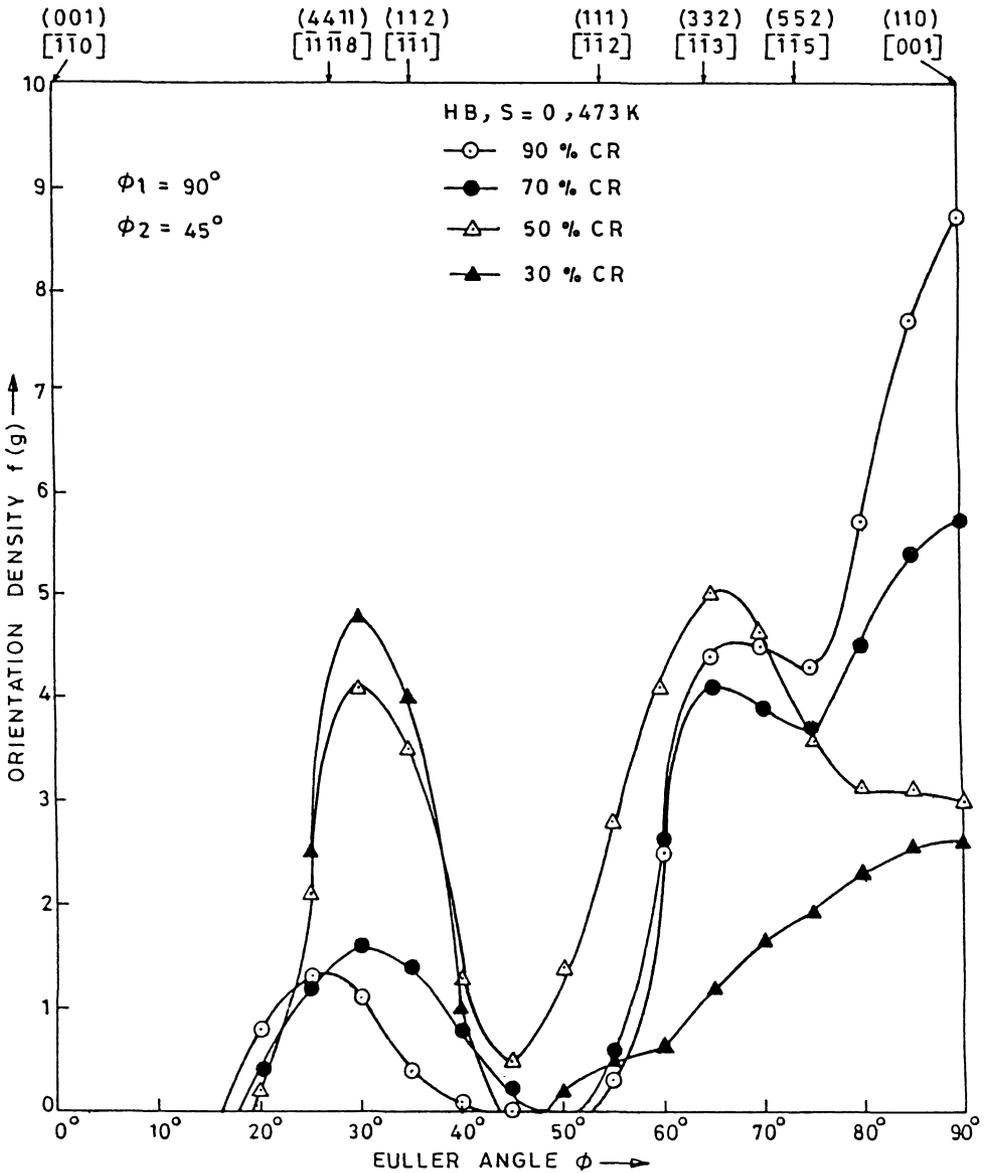


**Figure 6** Orientation densities along  $\langle 110 \rangle 60^\circ \text{ND}$  fibre of the orientation tube at the centre level ( $S=0$ ) of the hot band (HB) cold rolled by 0, 30, 50, 70 and 90% reductions at 473 K.

increasing deformation and also there is a decrease in the density of the cube orientation with deformation.

#### 4. DISCUSSION

The microstructure [Figure (1)] at the centre of the HB reveals the presence of deformed as well as recrystallised grains. Further, the interior layers of the HB



**Figure 7** Orientation densities along  $\langle 110 \rangle // \text{TD}$  fibre at the centre level ( $S=0$ ) of the hot band (HB) cold rolled by 30, 50, 70 and 90% reductions at 473 K.

retain considerably a deformation structure. The ODF analysis of the HB also indicates the presence of textural elements of both the retained rolling textures (i.e. copper-type) and recrystallisation texture (i.e. cube and RD rotated cubes). The  $\langle 110 \rangle 60^\circ \text{ND}$  fibre of the copper-type texture extends from  $\approx \{112\} \langle 111 \rangle$  through the region of  $\approx \{123\} \langle 634 \rangle$  to the orientation  $\{011\} \langle 112 \rangle$  with a maximum at  $\{011\} \langle 112 \rangle$  and minimum at  $\approx \{112\} \langle 111 \rangle$ . The recrystallised

texture components exhibits a scattering from cube position  $\{001\}\langle 100\rangle$  towards Goss position  $\{011\}\langle 100\rangle$  along with other elements of RD rotated cubes (i.e.  $\{013\}\langle 100\rangle$  and  $\{012\}\langle 100\rangle$ ). All the evidences (i.e. the presence of deformed grains and copper-type texture as well as some recrystallised grains and recrystallised texture) indicate that the HB is partially recrystallised by the process of dynamic recrystallisation during hot rolling (Singh, Ramaswamy and Suryanarayana, 1991).

The results of the ODFs analysis of the deformed austenite at the centre of the HB cold rolled at 473 K show that at low degrees of rolling ( $\leq 50\%$ ), the observed texture components are distributed along the two limited orientation tubes. The first tube with its fibre axis  $\langle 110\rangle//ND$  stretches along  $\phi_1$  at  $\phi_2 = 0^\circ$  and  $\phi = 45^\circ$  and includes the orientations  $\{011\}\langle 110\rangle$  at  $\phi_1 = \phi = \phi_2 = 0, 45$  and  $0^\circ$  and  $\{011\}\langle 112\rangle$  at  $\phi_1 = \phi = \phi_2 = 35, 45$  and  $0^\circ$ . The other tube inclined to constant  $\phi_2$  sections runs through the Euler space from  $\{225\}\langle 554\rangle$  ( $\approx \{112\}\langle 111\rangle$ ) at  $\phi_1, \phi, \phi_2 = 90, 30$  and  $45^\circ$  over  $\approx \{123\}\langle 634\rangle$  at  $\phi_1, \phi, \phi_2 = 59, 37$  and  $63^\circ$ , to  $\{011\}\langle 112\rangle$  at  $\phi_1, \phi, \phi_2 = 35, 45$  and  $90^\circ$  and has its fibre axis  $\langle 110\rangle$  inclined  $\sim 60^\circ$  from ND towards RD. The orientation element  $\{011\}\langle 112\rangle$  is at the intersection point of the above two limited tubes and has the highest density amongst all the texture components. The exceptional high density of  $\{011\}\langle 112\rangle$  component is partly due to the effect of starting texture of the HB (Figure 2) which shows the highest density of the orientation  $\{011\}\langle 112\rangle$  and partly due to the flow of orientations of the scattered zone between the cube components  $\{001\}\langle 100\rangle$  at  $\phi_1, \phi, \phi_2 = 0, 0$  and  $0^\circ$  and Goss component  $\{011\}\langle 100\rangle$  at  $\phi_1, \phi, \phi_2 = 0, 45$  and  $0^\circ$ , along the line parallel to  $\phi$  at  $\phi_1 = \phi_2 = 0^\circ$  from the cube to Goss component (i.e. metastable orientation) and then along the line parallel to  $\phi_1$  at  $\phi = 45^\circ$  and  $\phi_2 = 0^\circ$ , from Goss to the component  $\{011\}\langle 112\rangle$  (i.e. stable orientation) at  $\phi_1, \phi, \phi_2 = 35, 45$  and  $0^\circ$ , obtained as a result of rolling. The orientation  $\{225\}\langle 554\rangle$  at which the orientation tube with  $\langle 110\rangle 60^\circ ND$  fibre starts, differs by  $5^\circ$  from the orientation  $\{112\}\langle 111\rangle$  at  $\phi_1, \phi, \phi_2 = 90, 35$  and  $45^\circ$  and by  $3^\circ$  from the theoretically calculated orientation  $\{4411\}\langle 11118\rangle$  at  $\phi_1, \phi, \phi_2 = 90, 27$  and  $45^\circ$  by Dillamore, Butler and Green (1968) for Taylor type deformation. An increase in orientation densities for all texture components of the two limited tubes i.e.  $\langle 110\rangle//ND$  and  $\langle 110\rangle 60^\circ ND$  (Figures 5 and 6) occurs with increasing strain (e.g.  $\leq 30\%$  reduction). All these evidences suggest that the development of rolling textures at low degree of rolling ( $\leq 50\%$ ) agrees fairly well with the predictions of Taylor model (Dillamore, Butler and Green, 1968).

At 50% rolling reduction, the decrease in the otherwise stable  $\approx \{112\}\langle 111\rangle$  orientation marks the beginning of mechanical twinning (Wassermann, 1963) whereby the orientation  $\approx \{112\}\langle 111\rangle$  undergoes mechanical twinning during rolling to form  $\{552\}\langle 115\rangle$  at  $\phi_1, \phi, \phi_2 = 90, 75$  and  $45^\circ$ . Wassermann (1963) has indicated that the twinning is the most favourable deformation process which causes the beginning of transition of the rolling texture from copper to brass-type. But instead of rotating to the  $\{011\}\langle 100\rangle$  position and from there finally rotating to  $\{011\}\langle 112\rangle$  on further rolling, the twin orientation  $\{552\}\langle 115\rangle$  rotates in the opposite direction towards  $\{111\}\langle 112\rangle$  on favourable slip system. The change in the orientation densities of the components  $\approx \{112\}\langle 111\rangle$ , twin  $\{552\}\langle 115\rangle$ , Goss  $\{011\}\langle 100\rangle$  and other orientations obtained as a result of further deformation (Figure 7) illustrates clearly that the reverse rotation of the twin orientation  $\{552\}\langle 115\rangle$  towards  $\{111\}\langle 112\rangle$  at  $\phi_1, \phi, \phi_2 = 90, 55$  and  $45^\circ$  occurs. These

orientation elements are mainly distributed along a fibre axis  $\langle 110 \rangle // \text{TD}$  which includes the orientation  $\approx \{112\}\langle 111 \rangle$  at  $\phi = \sim 30^\circ$ , the twin  $\{552\}\langle 115 \rangle$  at  $\phi = 75^\circ$ , the Goss component  $\{011\}\langle 100 \rangle$  at  $\phi = 90^\circ$ , the orientation  $\{111\}\langle 112 \rangle$  at  $\phi = 55^\circ$  and the orientation  $\{332\}\langle 113 \rangle$  at  $\phi = 65^\circ$ .

It has been reported by TEM studies of the microstructures in rolled polycrystalline 70–30 brass (Duggan *et al.*, 1978) and silver (Pospiech *et al.*, 1975) that twinning becomes an important deformation mechanism at certain strain level but instead of normal rotation of the twin component  $\{552\}\langle 115 \rangle$  towards  $\{011\}\langle 100 \rangle$ , an abnormal rotation of the twin-matrix lamellae towards  $\{111\}\langle 112 \rangle$  position occurs. The orientation  $\{111\}\langle 112 \rangle$  is a position of the twin planes parallel to the rolling plane. The results of ODF analysis (Figure 7) where the twin orientation  $\{552\}\langle 115 \rangle$  at intermediate degrees of rolling ( $\geq 50\%$ ), shifts to lower  $\phi$  angles instead of rotating towards  $\{011\}\langle 100 \rangle$  position, prove that a rotation towards  $\{111\}\langle 112 \rangle$  position takes place, whereas the orientation density of the orientation  $\{011\}\langle 100 \rangle$  remains more or less constant (i.e. rather increases very slightly).

This abnormal rotation towards  $\{111\}\langle 112 \rangle$  is mainly due to the preferred slip on slip planes parallel to the coherent twin boundaries of the twin lamellae. In metals of low SFE, the normal slip on the slip planes of highest critical resolved shear stress, is severely limited mainly because of the very small width of the twin lamellae (Peissker, 1965). Further, the shear stress for the abnormal slip starts decreasing and finally becomes zero as the twin-parent lamellae become oriented with the rolling plane, and ultimately the abnormal rotation stops and, as detected in Figure 7, the statistical maxima of the rotated orientations is observed near  $\{332\}\langle 113 \rangle$ , which is  $10^\circ$  away from  $\{111\}\langle 112 \rangle$  position. The rotation of the matrix orientation  $\approx \{112\}\langle 111 \rangle$  towards  $\{111\}\langle 112 \rangle$  could not be observed in the ODF results (Figure 7) mainly because the obtained ODF is the reduced ODF. However, Hirsch, Virnich and Lücke (1981) have reported the rotation of matrix orientation towards  $\{111\}\langle 112 \rangle$  position in 70–30 brass with the help of complete ODF.

Besides the component  $\{332\}\langle 113 \rangle$  at intermediate degree of rolling, an orientation near  $\{111\}\langle 110 \rangle$ , also with  $\{111\}$  plane nearly parallel to the rolling plane appears. Since  $\{111\}\langle 110 \rangle$  develops simultaneously with the  $\{332\}\langle 113 \rangle$ , it can be assumed that this position is formed by the preferential slip on one of the two active slip systems parallel to the twin boundary that rotates the twin component  $\{552\}\langle 115 \rangle$  and other possible twin orientation of  $\approx \{123\}\langle 634 \rangle$  component. The twin of  $\approx \{123\}\langle 634 \rangle$  component is symmetrically equivalent to  $\approx \{123\}\langle 634 \rangle$  position. This causes the rotation of the preferred slip direction parallel to the rolling plane.

Since the rotated twin-matrix lamellae finally no longer contribute to deformation process by slip, on further rolling shear bands form in the highly unstable structure of aligned twin-matrix lamellae. Formation of shear bands has been observed by many investigators (Duggan *et al.*, 1978; Hatherly, 1978; Blicharski and Gorczyca, 1978; Wakefield and Hatherly, 1981) in their TEM studies. According to these investigators shear bands usually form at  $35^\circ$  to the rolling plane and the formation of these bands is mainly responsible for the decrease of  $\{111\}$  orientations at high degree of rolling (i.e.  $\geq 70\%$ ). Further, such a rotation by  $\pm 35^\circ$  around TD would lead  $\{111\}\langle 112 \rangle$  (i.e. observed orientation  $\{332\}\langle 113 \rangle$ ) to  $\{011\}\langle 100 \rangle$  position at  $\phi_1, \phi, \phi_2 = 90, 90$  and  $45^\circ$ , and to  $\approx \{112\}\langle 111 \rangle$  position at  $\phi_1, \phi, \phi_2 = 90, 20$  and  $45^\circ$ , and also the orientation

$\{111\}\langle 110\rangle$  to a position in between  $\approx\{123\}\langle 634\rangle$  and  $\{011\}\langle 112\rangle$  positions. The observations, an increase in the density of  $\{011\}\langle 100\rangle$  component and the shifting of  $\approx\{112\}\langle 111\rangle$  component towards lower  $\phi$  values after intermediate degree (i.e. 70% reduction) and high degree (i.e. 90% reduction) of rolling (Figure 7) and also an increase in the orientation density of  $\{011\}\langle 112\rangle$  component (Figure 6) after 90% reduction, establish the formation of shear bands in the above deformation range (i.e. intermediate and high degrees of rolling). Further, the cyclic variation in the orientation density of  $\{332\}\langle 113\rangle$  component after 50% reduction indicates a competition between mechanical twinning and the associated abnormal slip on one hand and shear bands formation in the rotated twin-matrix lamellal on the other. Twinning and abnormal rotation cause an increase whereas the shear bands formation leads to a decrease in the density of  $\{332\}\langle 113\rangle$  component. At 90% reduction, the observation of highest density at  $\{011\}\langle 112\rangle$  component also indicates that normal slip is re-occurring in the crystallites of the shear bands and sharpens the brass-type rolling texture (i.e.  $\langle 110\rangle//\text{ND}$  orientation tube).

The variation in the ND inverse pole figure densities of selected planes parallel to rolling plane of the cold rolled HB, with increasing reduction (Figure 3) can be correlated with ODF results. The  $\{220\}_\gamma$  reflection may be considered as representative of the orientations  $\{011\}\langle 100\rangle$  and  $\{011\}\langle 112\rangle$ . The twin rotated orientations  $\{332\}\langle 113\rangle$  and  $\{111\}\langle 110\rangle$  approximate the reflection  $\{111\}_\gamma$ , whereas the reflection  $\{311\}_\gamma$  represents approximately  $\approx\{112\}\langle 111\rangle$  component. At low degrees of rolling (<50% reduction), the deformation is occurring primarily by normal octahedral slip process which causes the flow of recrystallised texture components of the HB from cube position  $\{001\}\langle 100\rangle$  through Goss orientation  $\{011\}\langle 100\rangle$  to  $\{011\}\langle 112\rangle$  position and simultaneously there is an increase in the density of  $\approx\{112\}\langle 111\rangle$  and  $\approx\{123\}\langle 634\rangle$  components. The rapid increase in the density of  $\{220\}_\gamma$  reflection indicates the flow of orientations due to normal slip. The reduction in the density of  $\approx\{112\}\langle 111\rangle$  component at 50% reduction establishes the occurrence of mechanical twinning and the twin  $\{552\}\langle 115\rangle$ -matrix lamellae get rotated towards  $\{111\}\langle 112\rangle$  at intermediate degrees of rolling (i.e.  $\geq 50\%$  reduction). This abnormal rotation of twin-matrix lamellae finally stops near  $\{332\}\langle 113\rangle$  orientation which is  $10^\circ$  apart from  $\{111\}\langle 112\rangle$  and simultaneously there is the development of  $\{111\}\langle 110\rangle$  orientation. The increase in the pole density of  $\{111\}_\gamma$  reflection to a maximum value at about 80% reduction represents this stage (i.e. the orientations  $\{332\}\langle 113\rangle$  and  $\{111\}\langle 110\rangle$ ). The formation of the shear bands in the rotated twin-matrix lamellae causes the destruction of the orientations  $\{332\}\langle 113\rangle$  and  $\{111\}\langle 110\rangle$  and hence decreases the pole density of  $\{111\}_\gamma$  reflection and finally rotate these orientations to  $\{011\}\langle 100\rangle$  and near  $\{011\}\langle 112\rangle$  positions. In addition, again normal slip occurs in the grains of shear bands which sharpens the texture components  $\{011\}\langle 112\rangle$  and  $\{011\}\langle 100\rangle$  of the brass-type rolling texture (i.e.  $\langle 110\rangle//\text{ND}$  orientation tube) as reflected by the increasing tendency of the pole density of  $\{220\}_\gamma$  reflection after 80% rolling reduction.

## 5. CONCLUSIONS

- (i) At very low degree of rolling (i.e. 30% reduction), the texture consists of two limited fibres i.e.  $\langle 110\rangle//\text{ND}$  and  $\langle 110\rangle 60^\circ\text{ND}$ . The  $\langle 110\rangle//\text{ND}$  fibre runs

from  $\{011\}\langle 100 \rangle$  to  $\{011\}\langle 112 \rangle$  and the  $\langle 110 \rangle 60^\circ \text{ND}$  fibre stretches from  $\approx \{112\}\langle 111 \rangle$  through  $\approx \{123\}\langle 634 \rangle$  to  $\{011\}\langle 112 \rangle$ . The development of these two limited fibres agrees to the predictions of Taylor model based on normal octahedral slip.

(ii) At low degree of rolling (i.e. 50% reduction) mechanical twinning leads to the reduction of  $\approx \{112\}\langle 111 \rangle$  component and the formation of twin component  $\{552\}\langle 115 \rangle$  in form of the fine twin-matrix lamellae.

(iii) At the intermediate degrees of rolling (i.e.  $\geq 50\%$  reduction) slip is preferred on planes parallel to the twin boundaries of the fine twin lamellae, leading to an abnormal rotation of twin-matrix lamellae into  $\{332\}\langle 113 \rangle$  and  $\{111\}\langle 110 \rangle$  and finally this abnormal rotation stops when the slip plane becomes parallel to the rolling plane.

(iv) At intermediate (i.e. 70% reduction) and high (i.e. 90% reduction) degrees of rolling, the formation of shear bands in the rotated twin-matrix regions changes the orientations  $\{332\}\langle 113 \rangle$  and  $\{111\}\langle 110 \rangle$  near to  $\{011\}\langle 100 \rangle$  and  $\{011\}\langle 112 \rangle$  positions respectively.

(v) At high degree of rolling (i.e. 90% reduction) the re-establishment of homogeneous slip in the fine grains of shear bands causes sharpening of the brass-type rolling texture by the rotation of  $\{011\}\langle 100 \rangle$  component to  $\{011\}\langle 112 \rangle$  component.

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