

THE EVOLUTION OF TEXTURE DURING ANNEALING OF NANOCRYSTALLINE ELECTROLESS NICKEL ALLOY DEPOSITS

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The electroless Ni–Co–P alloy films containing up to 13.4% Co and 7.6% P were deposited on a 5086 aluminum alloy sheet using baths consisting of NiSO₄, CoSO₄, (NH₄)₂SO₄, NaH₂PO₂, sodium citrate and thiourea. The deposits were solid solutions having a grain size of 6–7 nm and strong $\langle 111 \rangle$ texture. When the deposits were annealed for 2 h, Ni₃P₂ was precipitated at 325°C and transformed into the stable Ni₃P phase at higher temperatures. During annealing the texture of deposit changed from $\langle 111 \rangle$ to $\langle 100 \rangle$. The texture change was attributed to preferential growth of $\langle 100 \rangle$ grains at the expense of $\langle 111 \rangle$ grains to decrease the thermal strain energy of the deposits. On the other hand, the electroless Ni–Cu–P alloy films containing Cu and P deposited on 5086 aluminum alloy sheet had the microstructures and textures which are similar to these of the Ni–Co–P alloy films. However, their initial $\langle 111 \rangle$ texture remained relatively stable even after annealing. The differences in annealing textures of Ni–Co–P and Ni–Cu–P alloy films have been discussed.

Keywords: Electroless nickel alloy; Deposition texture; Annealing texture; Nanocrystalline

INTRODUCTION

The Ni–Co–P alloy deposits can be obtained from alkaline solutions over entire Co compositions (Pearlstein, 1974). The deposits with higher Co concentrations are applied to magnetic recording media and vertical

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magnetic recording media by adding Mn and Re (Osaka *et al.*, 1983a,b; Koiwa and Osaka, 1987). The increase in the Co concentration causes the decrease in the P content and the increase in grain size (Chow *et al.*, 1972). Little study has been made of texture changes during annealing of the Ni-Co-P alloy deposits (Lee and Hur, 1999). Microstructures and crystallization of the electroless Ni-P deposits and magnetic properties and microstructures of electroless Ni-Cu-P deposits were studied by the present authors (Park and Lee, 1988; Hur *et al.*, 1990; 1991).

In this paper, the evolution of texture and phase transformation behavior during annealing of the Ni-Co-P and Ni-Cu-P deposits have been discussed.

EXPERIMENTAL PROCEDURE

The electroless Ni-Co-P alloy films were deposited on a 5086 aluminum alloy sheet which had been preplated. The preplating process included cleaning, and nitric acid and double zincating treatments. The zincating treatments are performed to remove oxide film formed after alkali cleaning and acid etching, to prevent reoxidation and to form catalytic surface, which is necessary for electroless nickel plating. The bath used to obtain the electroless deposits consisted of 0.07–0.1 mol/l Ni^{2+} (NiSO_4), 0–0.03 mol/l Co^{2+} (CoSO_4), 0.5 mol/l $(\text{NH}_4)_2\text{SO}_4$, 0.2 mol/l NaH_2PO_2 , 0.2 mol/l sodium citrate and 1 ppm thiourea. Ammonia solution was used to adjust the pH of the bath to 9.0. The temperature of the bath was maintained at 90°C with a thermoregulator and agitated by magnetic stirring.

The thickness of deposits ranged from 30 to 40 μm . The P and Co contents in the deposits were analyzed by the inductively coupled plasma (ICP) emission spectrometry. Adherence between the electroless deposit and the aluminum substrate depends on the zinc layer thickness. If the thickness is too thick, peeling occurs between the deposit and the substrate. The appropriate thickness is known to be about 100 nm. Changes in structures of as-deposited and heat-treated films were measured by the usual continuous scanning X-ray diffraction method. The heat treatment of the deposits used in X-ray diffraction was carried out isothermally in a salt bath for 2 h at temperatures from 250°C to 450°C with 50°C intervals. The diffusion depth due to 100 nm thick zinc

layer was estimated to be less than 1 μm when the deposits are heated at 4500°C for 2 h, and the thickness of deposits ranged from 30 to 40 μm . Therefore, the X-ray diffraction results were not influenced by the zinc layer. The structure analysis and heat treatments of the films were conducted without being separated from the substrate.

RESULTS AND DISCUSSION

Deposition Rate and Composition

Figure 1 shows the deposition rate of the electroless In–Co–P deposit as a function of $[\text{Co}^{2+}]/[\text{Ni}^{2+}]$ in deposition solution. The deposition rate remained constant until a cobalt ion to nickel ion ratio of 0.1, beyond which the deposition rate decreased slowly with increasing ratio. Figure 2 shows the cobalt and phosphorous contents of the deposits as a function of $[\text{Co}^{2+}]/[\text{Ni}^{2+}]$ in the deposition solution. The cobalt content in the deposit increased with increasing ratio of $[\text{Co}^{2+}]/[\text{Ni}^{2+}]$. However, the increasing rate of cobalt content in the deposit was not so large as the increasing rate of cobalt ion in the solution.

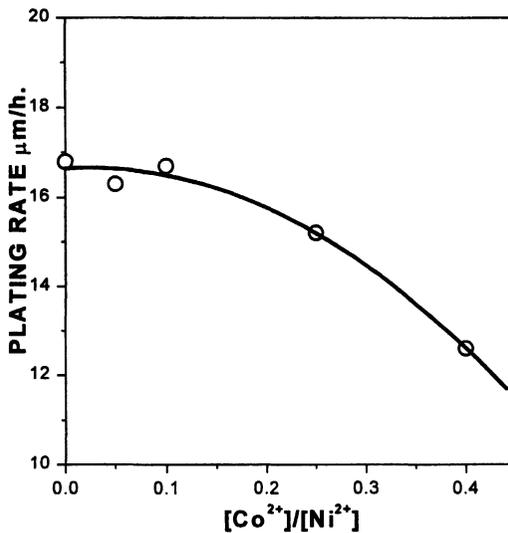


FIGURE 1 Plating rate of electroless Ni–Co–P deposit as a function of $[\text{Co}^{2+}]/[\text{Ni}^{2+}]$ in solution.

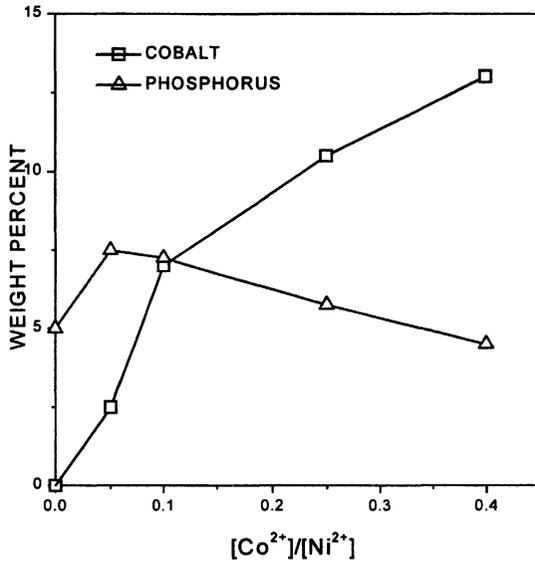


FIGURE 2 Cobalt and phosphorous content in electroless Ni-Co-P deposit as a function of $[\text{Co}^{2+}]/[\text{Ni}^{2+}]$ in solution.

The phosphorous content in the deposit increased initially and decreased slowly with increasing $[\text{Co}^{2+}]/[\text{Ni}^{2+}]$ ratio.

Microstructures and Textures of Deposits

Transmission electron micrographs (dark field image) and diffraction patterns of various electroless Ni-Co-P deposits are shown in Fig. 3. The micrographs indicate that fine grains are imbedded in amorphous matrix. The amorphous phase is supposed to have a higher phosphorous content than crystalline grains from the well known fact that the electroless Ni-P deposits tend to become amorphous with increasing phosphorous content. X-ray diffraction patterns of electroless Ni-Co-P alloy deposits in Fig. 4(a) shows strong (111) and very weak (200) reflection peaks, which indicate that the deposits have the (111) texture. Their grain size is estimated to be 6–7 nm by the Scherrer equation. This size is approximately the same as that of white spots as in Fig. 3. It can be seen that the (200) peak intensity relative to the (111) peak increases with increasing cobalt content in the deposits. The Ni-5.1P deposit also shows the (200) peak.

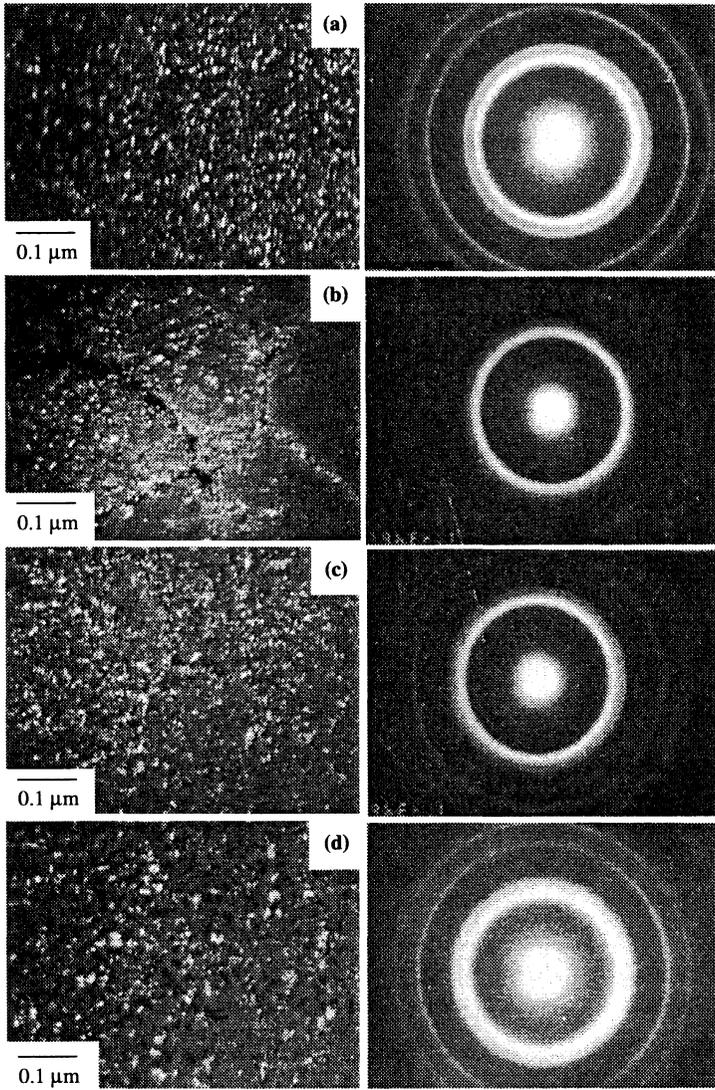


FIGURE 3 Dark field micrographs and SAD patterns of electroless Ni-Co-P deposits. (a) Ni-2.4Co-7.6P, (b) Ni-7.0Co-7.2P, (c) Ni-10.5Co-5.8P, (d) Ni-13.4Co-4.5P.

Annealing Effects

Figure 4 shows X-ray diffraction patterns of electroless Ni–Co–P deposits annealed at various temperatures for 2 h. The (200) peak intensity increases more rapidly than the (111) intensity with increasing annealing temperature and cobalt content in the deposits. At an annealing temperature of 325°C, Ni₅P₂ forms especially in the higher phosphorous deposits, but it disappears and stable Ni₃P appeared at higher temperatures. A similar phenomenon was observed in electroless Ni–Cu–P alloys (Hur *et al.*, 1991). The appearance of Ni₅P₂ from the annealed Ni–Cu–P deposit was attributed to a higher phosphorous segregation in grain boundary region. When annealed, the higher phosphorous region was supposed to be further increased by extraction of P dissolved in nickel grains which in turn gave rise to precipitation of the Ni₅P₂ phase. The same explanation may be applicable to the Ni–Co–P deposits.

The evolution of $\langle 100 \rangle$ in a fcc metal film with $\langle 111 \rangle$ orientation on a substrate during annealing can originate from a few sources. A lattice matching at the film–substrate interface may lead to the $\langle 100 \rangle$ orientation. But there is little possibility of lattice matching between the substrate and the nickel alloy deposit.

When $\langle 111 \rangle$ fcc copper electrodeposits (Lee *et al.*, 1995; Lee, 1996a,b) and silver electrodeposits (Nam and Lee, 1999) were recrystallized, the $\langle 100 \rangle$ orientation developed (Lee *et al.*, 1995; Lee, 1996a,b). However, this is possible when the dislocation density in the deposit is high enough to cause recrystallization. Since the present deposits have a grain size of several nanometer, they are likely to have little dislocation. Therefore, the explanation cannot be applied to this system.

The texture evolution during annealing may be due to grain growth. A duplex texture of $\langle 111 \rangle$ and $\langle 100 \rangle$ in fcc metals tends to develop $\langle 111 \rangle$ during grain growth. The grain boundaries in a material with a duplex fiber texture of $\langle 111 \rangle$ and $\langle 100 \rangle$ may be described approximately as tilt boundaries of $\langle 111 \rangle$ and $\langle 100 \rangle$ grains and boundaries between $\langle 111 \rangle$ and $\langle 100 \rangle$ grains as shown in Fig. 5. The $\langle 111 \rangle$ – $\langle 100 \rangle$ boundaries will not determine which grains grow favorably. If the tilt boundary energies of the two differently textured grains are different, the grains with lower boundary energy will grow at the expense of the grains with the higher boundary energies when the mobilities of the grain boundaries

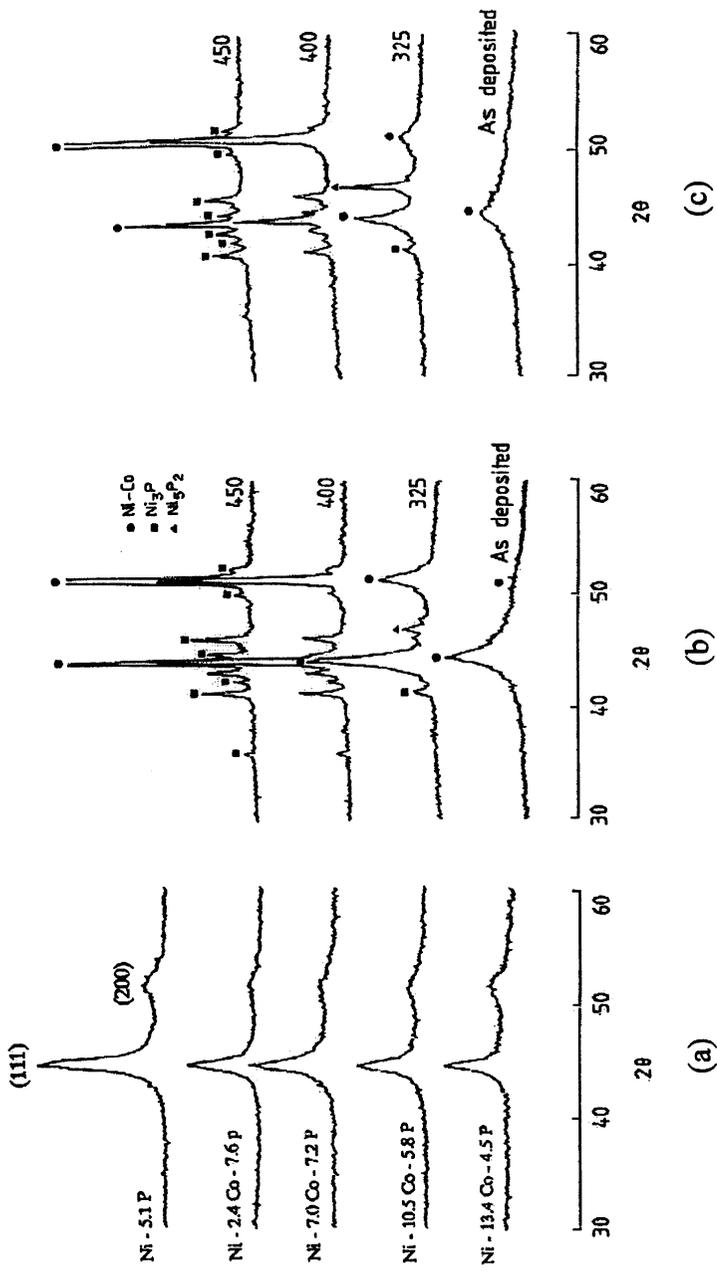


FIGURE 4(a)-(c)

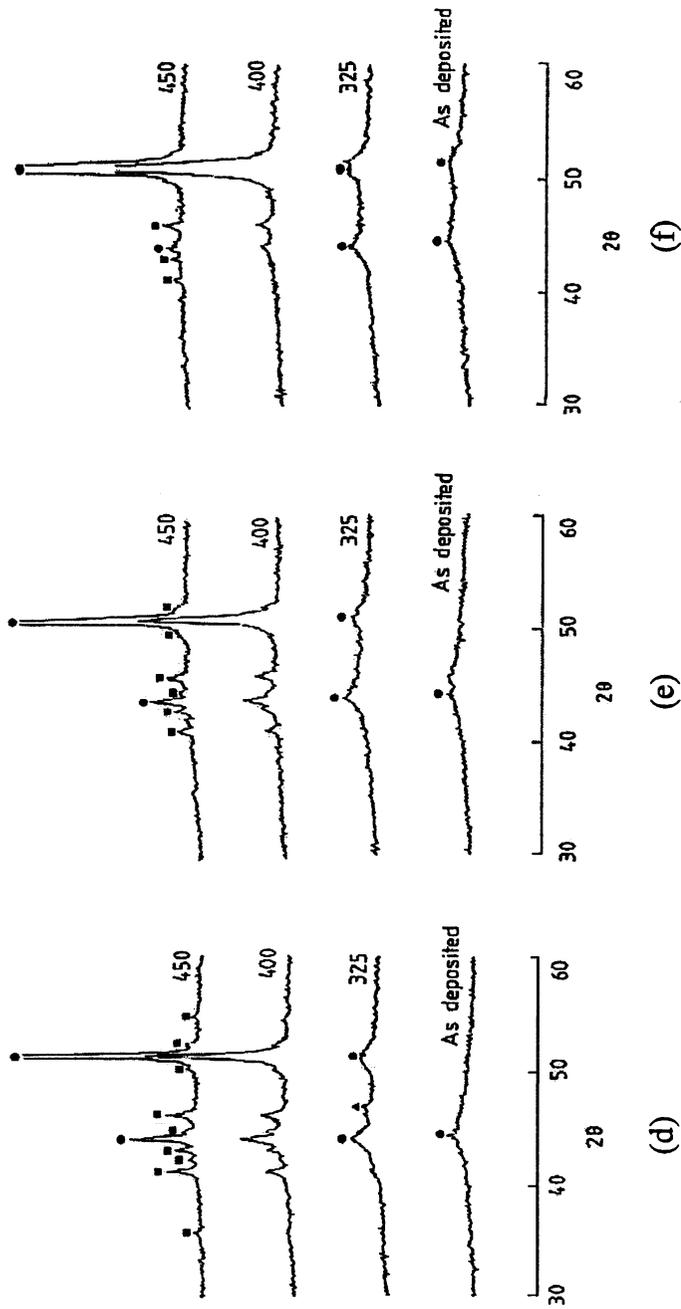


FIGURE 4(d)-(f)

FIGURE 4 X-ray diffraction patterns of electroless Ni-Co-P deposits. (a) as-deposited, (b) Ni-5.1P, (c) Ni-2.4Co-7.6P, (d) Ni-7.0Co-7.2P, (e) Ni-10.5Co-5.8P and (f) Ni-13.4Co-4.5P are annealed at various temperatures for 2 h. Peak intensities in (b) to (f) are attenuated.

are the same. For copper and gold, the calculated average boundary energy of $\langle 111 \rangle$ grains is slightly higher than $\langle 100 \rangle$ grains as shown in Fig. 6 (Wolf, 1990). For aluminum, the mobility of tilt grain boundaries of $\langle 111 \rangle$ grains is higher than that of $\langle 100 \rangle$ grains at high temperatures as shown in Fig. 7 (Aleshin *et al.*, 1978). In this case, neighboring $\langle 111 \rangle$ grains will grow faster than neighboring $\langle 100 \rangle$ grains, resulting in the $\langle 111 \rangle$ grain size being larger than the $\langle 100 \rangle$ grain size. The $\langle 111 \rangle$ grains will grow at the expense of the $\langle 100 \rangle$ grains because of the size advantage, and develop the $\langle 111 \rangle$ texture during grain growth. Indeed, nanocrystalline Fe–Ni alloy electrodeposits developed the $\langle 111 \rangle$ texture when annealed (Czerwinski *et al.*, 1997). The recrystallization textures of silver (Shin *et al.*, 2000) and copper (Backofen, 1952) wires are $\langle 100 \rangle$, whereas their growth textures are $\langle 111 \rangle$ (Shin *et al.*, 2000; Grant *et al.*, 1984). Therefore, the $\langle 100 \rangle$ annealing texture in this study cannot be caused by grain growth.

The texture change from $\langle 111 \rangle$ to $\langle 100 \rangle$ may be due to the fact that the $\langle 111 \rangle$ grains have the highest strain energy while the $\langle 100 \rangle$ grains the

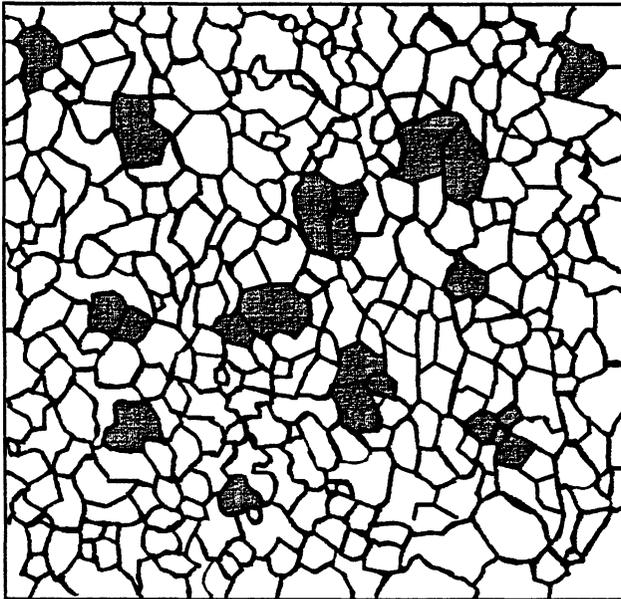


FIGURE 5 Schematic microstructure of section normal to growth direction of deposit which consists of $\langle 111 \rangle$ grains and $\langle 100 \rangle$ grains (dark).

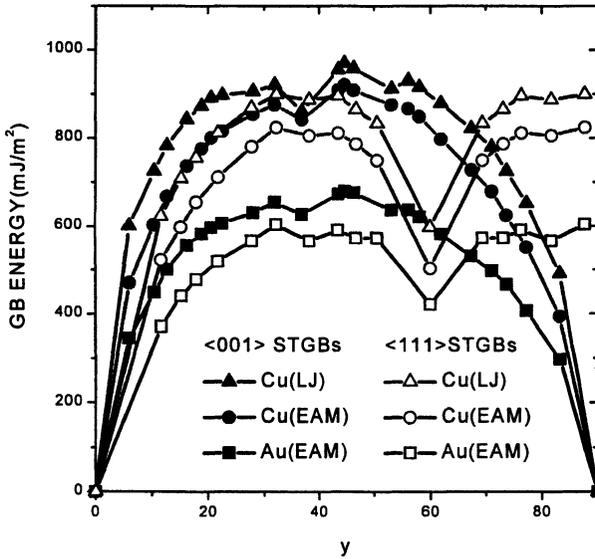


FIGURE 6 Calculated energy (in mJ/m^2) for symmetrical tilt grain boundaries of fcc metals with interface planes perpendicular to $\langle 111 \rangle$ and $\langle 100 \rangle$ as function of tilt angle (Wolf, 1990). (LJ and (EAM) indicate energies have been calculated using LJ and EAM potentials.)

lowest strain energy (Carel *et al.*, 1996). The deposit and substrate, which have different thermal expansion coefficients, undergo different volume changes due to a temperature change. In addition, the formation of nickel phosphide could lead to volume changes. These changes can produce strains and in turn stresses in the deposits during annealing. The experimental results (Fig. 4) indicate that the texture change is more influenced by cobalt concentration than by phosphorous. This implies that the volume changes due to phosphide formation are not a dominant factor leading to the $\langle 100 \rangle$ texture. Therefore this factor is not considered here.

The strains and stresses due to the different thermal expansion coefficients are calculated. The thermal strain ϵ_{th} due to the different thermal expansions is calculated by

$$\epsilon_{\text{th}} = \int_{T_0}^T (\alpha_s - \alpha_f) dT \quad (1)$$

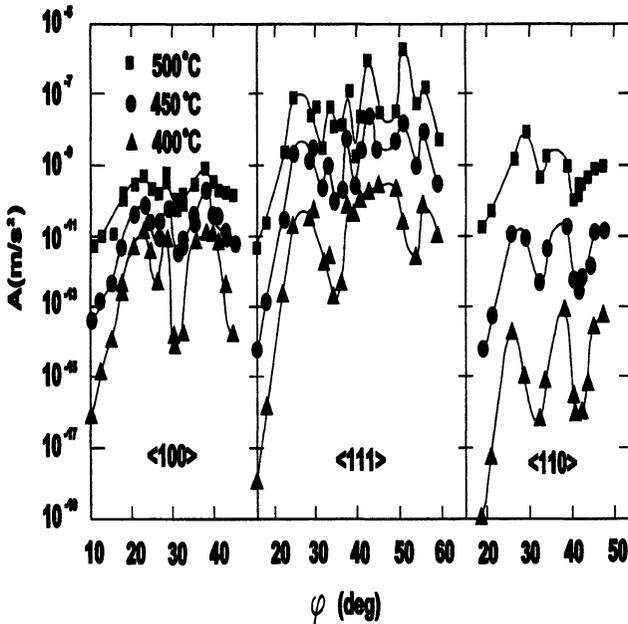


FIGURE 7 Misorientation dependence of the tilt grain boundary mobility in aluminum of misorientation axes (100), (110) and (111) (Aleshin *et al.*, 1978).

where T and α are temperature and thermal expansion coefficient, respectively. Subscripts 0, s and f indicate initial state, substrate and film, respectively. Setting $T = 250\text{--}450^\circ\text{C}$, $T_0 = 90^\circ\text{C}$, $\alpha_s = \alpha_{\text{Al}} = 25.8 \times 10^{-6} \text{K}^{-1}$ (Metals Handbook, 1979), $\alpha_f = \alpha_{\text{Ni}} = 13.3 \times 10^{-6} \text{K}^{-1}$ (Metals Handbook, 1979), we obtain

$$\varepsilon_{\text{th}} = (2 \text{ to } 4.5) \times 10^{-3} \quad (2)$$

To be more rigorous, the thermal expansion coefficients of fcc Ni–Co solid solutions should be used. The thermal expansion coefficient of cobalt (hcp) is $13.8 \times 10^{-6} \text{K}^{-1}$ (Metals Handbook, 1979) which is very close to that of nickel. Therefore, the thermal expansion coefficient of nickel can be a good approximation of the alloys in the absence of experimental values.

A polycrystalline deposit which has a strong fiber texture is expected to have a normal anisotropy, even though each grain may not have the

normal anisotropy. Therefore, the stress state of deposit can be approximated by an equibiaxial plane stress state, which may be expressed as $\sigma_{11} = \sigma_{22} = \sigma$ and $\sigma_{12} = 0$. It follows from the generalized Hooke's law for anisotropic materials that

$$\varepsilon_{11} = S'_{1111}\sigma'_{22} = (S'_{1111} + S'_{1122})\sigma \quad (3)$$

where the prime means the quantity along an arbitrary axis in plane and S_{ijkl} is the compliances. Using the transformations law, the following equations are obtained for cubic crystals

$$\sigma = \frac{\varepsilon'_{11}}{S_{11} + S_{12}} \quad \text{for } \{100\} \text{ planes} \quad (4)$$

$$\sigma = \frac{\varepsilon'_{11}}{S_{11} + S_{12} + (1/6)(S_{44} - 2S_{11} + 2S_{12})} \quad \text{for } \{111\} \text{ planes} \quad (5)$$

$$\sigma = \frac{E}{1 - \nu_{11}} \varepsilon'_{11} \quad \text{for isotropic materials} \quad (6)$$

where E and ν are Young's modulus and Poisson's ratio, respectively. The above stresses are independent of x'_1 direction, which in turn means that ε'_{11} is also independent of x'_1 direction. That is, the above equations hold for the equibiaxial strain state. It follows that

$$\sigma_{\text{th}}(100) = \frac{\varepsilon_{\text{th}}}{S_{11} + S_{12}} \quad (7)$$

$$\sigma_{\text{th}}(111) = \frac{\varepsilon_{\text{th}}}{S_{11} + S_{12} + (1/6)(S_{44} - 2S_{11} + 2S_{12})} \quad (8)$$

$$\sigma_{\text{th}}(\text{iso}) = \frac{E}{1 - \nu} \varepsilon_{\text{th}} \quad (9)$$

where $\sigma_{\text{th}}(100)$, $\sigma_{\text{th}}(111)$ and $\sigma_{\text{th}}(\text{iso})$ indicate thermal stresses of the film having (100), (111) and random orientations, respectively.

Setting $S_{11} = 0.00734$, $S_{44} = 0.00802$, $S_{12} = -0.00274 \text{ GPa}^{-1}$ (Reid, 1973), $E = 201$, and $\nu = 0.31$, $\sigma_{\text{th}}(100) = 217 \times (2-4.5) \text{ MPa}$, $\sigma_{\text{th}}(111) = 388 \times (2-4.5) \text{ MPa}$ and $\sigma_{\text{th}}(\text{iso}) = 291 \times (2-4.5) \text{ MPa}$. $\sigma_{\text{th}}(100)$ is lower than $\sigma_{\text{th}}(111)$. Therefore, the $\langle 100 \rangle$ grains have lower strain energy than the $\langle 111 \rangle$ grains. These values are based on room temperature data. The stresses at experimental temperatures of 250–450°C would be lower, but still substantial. In conclusion, the evolution of the $\langle 100 \rangle$ orientation during annealing is suggested to be due to favorable growth of $\langle 100 \rangle$ grains at the expense of $\langle 111 \rangle$ grains under thermal stresses.

Electroless Ni–Cu–P deposits showed similar grain size and textures at the initial state, but the $\langle 111 \rangle$ texture remained relatively stable even after annealing within the experimental range as shown in Fig. 8 (Hur *et al.*, 1991). The thermal expansion coefficient of copper (19.6×10^{-6} at 700 K (Metals Handbook, 1979)) is higher than that of cobalt.

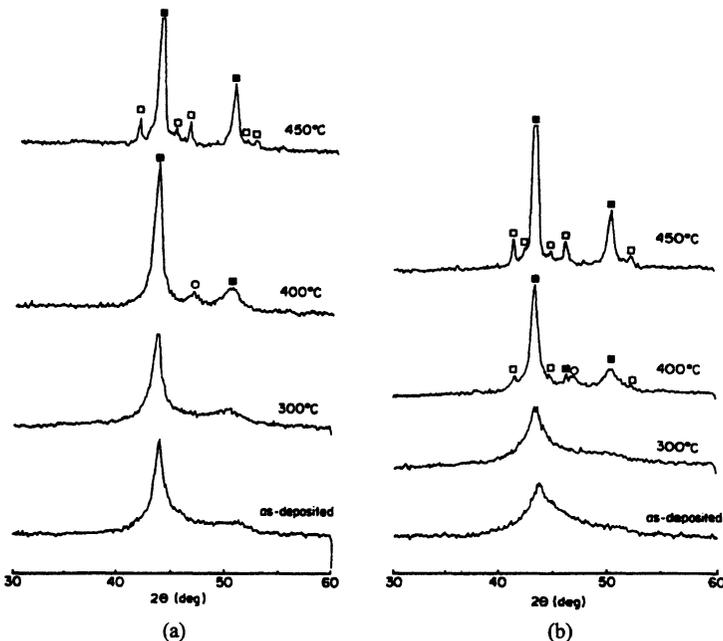


FIGURE 8 X-ray diffraction patterns of electroless (a) Ni–44.5Cu–7.5P and (b) Ni–50.3Cu–5.7P deposits. Annealing was performed at 300°C, 400°C and 450°C for 2 h.

Therefore, the less thermal stress would be developed, resulting in the slower evolution of the $\langle 100 \rangle$ orientation.

CONCLUSIONS

- (1) The plating rate of electroless Ni–Co–P alloys decreased with increasing ratio of cobalt ion concentration to nickel ion concentration in bath.
- (2) The phosphorous content in the deposit slightly decreased with increasing cobalt content.
- (3) When the deposit was annealed for 2 h, Ni₅P₂ was precipitated at 325°C and replaced by stable Ni₃P at higher temperatures.
- (4) The as-deposited Ni–Co–P films had strong $\langle 111 \rangle$ texture, which changed to the $\langle 100 \rangle$ texture on annealing. The rate of texture change increased with increasing cobalt concentration. The texture change was due to preferential growth of $\langle 100 \rangle$ grains at the expense of $\langle 111 \rangle$ grains to reduce the thermal strain energy of the deposits.
- (5) The Ni–Cu–P films had the texture and microstructure similar to the Ni–Co–P films. However, the texture remained relatively stable possibly due to the less thermal stresses.

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