A New Method for Inextensible Flows of Timelike Curves in Minkowski Space-Time $\mathbb{E}^4_1$

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We construct a new method for inextensible flows of timelike curves in Minkowski space-time $\mathbb{E}^4_1$. Using the Frenet frame of the given curve, we present partial differential equations. We give some characterizations for curvatures of a curve in Minkowski space-time.

1. Introduction

Numerous processing operations of complex fluids involve free surface deformations; examples include spraying and atomization of fertilizers and pesticides, fiber-spinning operations, paint application, roll-coating of adhesives, and food processing operations such as container- and bottle-filling. Systematically understanding such flows can be extremely difficult because of the large number of different forces that may be involved, including capillarity, viscosity, inertia, gravity, and the additional stresses resulting from the extensional deformation of the microstructure within the fluid. Consequently, many free-surface phenomena are described by heuristic and poorly quantified words such as “spinnability,” “tackiness,” and “stringiness.” Additional specialized terms used in other industries include “pituity” in lubricious aqueous coatings, “body” and “length” in the printing ink business, “ropiness” in yogurts, and “long/short textures” in starch processing [1].

The flow of a curve or surface is said to be inextensible if, in the former case, the arc length is preserved, and, in the latter case, if the intrinsic curvature is preserved [2–7]. Physically, inextensible curve and surface flows are characterized by the absence of any strain energy induced from the motion. Kwon investigated inextensible flows of curves and developable surfaces in $\mathbb{R}^3$. Necessary and sufficient conditions for an inextensible curve flow first are expressed as a partial differential equation involving the curvature and torsion. Then, they derived the corresponding equations for the inextensible flow of a developable surface and showed that it suffices to describe its evolution in terms of two inextensible curve flows [8]. Additionally, there are many works related with inextensible flows [1, 8–15].

In the past two decades, for the need to explain certain physical phenomena and to solve practical problems, geometers and geometric analysis have begun to deal with curves and surfaces which are subject to various forces and which flow or evolve with time in response to those forces so that the metrics are changing. Now, various geometric flows have become one of the central topics in geometric analysis. Many authors have studied geometric flow problems [1, 12, 16, 17].

This study is organized as follows: firstly, we study inextensible flows of timelike curves in Minkowski space-time. Secondly, using the Frenet frame of the given curve, we present partial differential equations. Finally, we give some characterizations for curvatures of a curve in Minkowski space-time.

2. Preliminaries

A “particle” in special relativity means a curve $\alpha$ with a timelike unitary tangent vector [18, 19].

Since $g$ is an indefinite metric, recall that a vector $v$ can have one of the following three casual characterizations:

(i) it can be space-like if $g(v,v) > 0$ or $v = 0$;
(ii) it can be timelike if \( g(v, v) < 0 \);

(iii) it can be null (light-like) if \( g(v, v) = 0 \) and \( v \neq 0 \).

Similarly, an arbitrary curve \( \alpha \) can be locally space-like, timelike, or null (light-like), if all of its velocity vectors \( \alpha' \) are, respectively, space-like, timelike, or null. Also, recall that the norm of a vector \( v \) is given by

\[
\|v\| = \sqrt{g(v, v)}.
\]

(1)

Therefore, \( v \) is a unit vector if \( g(v, v) = \pm 1 \). Next, vectors \( v, w \) are said to be orthogonal if \( g(v, w) = 0 \) (see [20]). The velocity of the curve \( \alpha \) is given by \( \|\alpha'(s)\| \).

Denote by \( \{T, N, B_1, B_2\} \) the moving Frenet frame along the curve \( \alpha(s) \) in the space-time. Then \( T, N, B_1, B_2 \) are, respectively, the tangent, the principal normal, the binormal, and the trinormal vector fields. A space-like or timelike curve \( \alpha(s) \) is said to be parameterized by arc length function \( s \), if

\[
g\left(\alpha'(s), \alpha'(s)\right) = \pm 1.
\]

(2)

Let \( \alpha(s) \) be a timelike curve in the space-time, parameterized by arc length function \( s \).

The timelike curve \( \alpha \) is called timelike Frenet curve if there exist three smooth functions \( k_1, k_2, k_3 \) on \( \alpha \) and smooth nonnull frame fields \( \{T, N, B_1, B_2\} \) along the curve \( \alpha \). Also, the functions \( k_1, k_2, \) and \( k_3 \) are called the first, the second, and the third curvature function on \( \alpha \), respectively. Then, for the unit speed timelike curve \( \alpha \) with nonnull frame vectors [16], the following Frenet equations are given,

\[
\begin{align*}
T' &= k_1 N, \\
N' &= k_1 T + k_3 B_1, \\
B'_1 &= -k_3 N + k_2 B_2, \\
B'_2 &= -k_2 B_1.
\end{align*}
\]

(3)

Here, due to characters of Frenet vectors of the timelike curve, \( T, N, B_1, \) and \( B_2 \) are mutually orthogonal vector fields satisfying equations

\[
\begin{align*}
g(T, T) &= -1, \\
g(N, N) &= g(B_1, B_1) = g(B_2, B_2) = 1.
\end{align*}
\]

(4)

3. A New Method for Inextensible Flows of Timelike Curves in \( \mathbb{E}^4 \)

Physically, inextensible curve and surface flows give rise to motions in which no strain energy is induced. The swinging motion of a cord of fixed length or, for example, of a piece of paper carried by the wind, can be described by inextensible curve and surface flows. Such motions arise quite naturally in a wide range of physical applications [8, 11, 12].

Let \( \alpha(u, t) \) be a one parameter family of smooth timelike curves in \( \mathbb{E}^4 \).

Any flow of \( \alpha \) can be represented as

\[
\frac{\partial \alpha}{\partial t} = \mathcal{U}_1 T + \mathcal{U}_2 N + \mathcal{U}_3 B_1 + \mathcal{U}_4 B_2,
\]

(5)

where \( \mathcal{U}_1, \mathcal{U}_2, \mathcal{U}_3, \mathcal{U}_4 \) are smooth functions.

Let the arc length variation be

\[
s(u, t) = \int_0^u v du.
\]

(6)

In the \( \mathbb{E}^4 \) the requirement that the curve be not subject to any elongation or compression can be expressed by the condition

\[
\frac{\partial}{\partial t} s(u, t) = \int_0^u \frac{\partial v}{\partial u} du = 0.
\]

(7)

**Definition 1.** The flow \( \frac{\partial \alpha}{\partial t} \) in \( \mathbb{E}^4 \) is said to be inextensible if

\[
\frac{\partial}{\partial t} \left\| \frac{\partial \alpha}{\partial u} \right\| = 0.
\]

(8)

**Theorem 2.** Let \( \frac{\partial \alpha}{\partial t} \) be a smooth flow of \( \gamma \). The flow is inextensible if and only if

\[
\frac{\partial \mathcal{U}_1}{\partial u} = -\mathcal{U}_2 v k_1.
\]

(9)

**Proof.** Assume that \( \frac{\partial \alpha}{\partial t} \) is inextensible. Then,

\[
\frac{\partial}{\partial t} s(u, t) = \int_0^u \frac{\partial v}{\partial u} du = \int_0^u \left( \frac{\partial \mathcal{U}_1}{\partial u} + \mathcal{U}_2 v k_1 \right) du = 0.
\]

(10)

Substituting (8) in (10) completes the proof of the theorem.

We now restrict ourselves to arc length parameterized curves. That is, \( v = 1 \) and the local coordinate \( u \) corresponds to the curve arc length \( s \). We require the following lemma.

**Lemma 3.** If the flow is inextensible, then

\[
\begin{align*}
\frac{\partial T}{\partial t} &= \left( \mathcal{U}_1 k_1 + \frac{\partial \mathcal{U}_2}{\partial s} - \mathcal{U}_3 k_2 \right) N \\
&\quad + \left( \mathcal{U}_2 k_2 + \frac{\partial \mathcal{U}_3}{\partial s} - \mathcal{U}_4 k_3 \right) B_1 \\
&\quad + \left( \mathcal{U}_3 k_3 + \frac{\partial \mathcal{U}_4}{\partial s} \right) B_2,
\end{align*}
\]

(11)

where \( \mathcal{U}_1, \mathcal{U}_2, \mathcal{U}_3, \mathcal{U}_4 \) are smooth functions of time and arc length.

**Proof.** Using definition of \( \alpha \), we have

\[
\frac{\partial}{\partial t} T = \left( \frac{\partial \mathcal{U}_1}{\partial s} + \mathcal{U}_2 k_1 \right) T + \left( \mathcal{U}_1 k_1 + \frac{\partial \mathcal{U}_2}{\partial s} - \mathcal{U}_3 k_2 \right) N \\
&\quad + \left( \mathcal{U}_2 k_2 + \frac{\partial \mathcal{U}_3}{\partial s} - \mathcal{U}_4 k_3 \right) B_1 \\
&\quad + \left( \mathcal{U}_3 k_3 + \frac{\partial \mathcal{U}_4}{\partial s} \right) B_2.
\]

(12)

Substituting (9) in (12), we obtain (11). This completes the proof.
Now we give the characterization of evolution of first curvature as below.

**Theorem 4.** Let $\partial \alpha / \partial t$ be inextensible flow of timelike $\alpha$ in $\mathbb{E}_1^4$. Then, the evolution of $k_1$ is given by

$$
\frac{\partial k_1}{\partial t} = \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) - k_2 \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right],
$$

where $\mathcal{A}_1$, $\mathcal{A}_2$, $\mathcal{A}_3$, $\mathcal{A}_4$ are smooth functions of time and arc length.

**Proof.** Assume that $\partial \alpha / \partial t$ is inextensible in $\mathbb{E}_1^4$. Thus it is easy to obtain that

$$
\frac{\partial}{\partial s} \frac{\partial}{\partial t} T = k_1 \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) T
$$

$$
+ \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) - k_2 \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] N
$$

$$
+ \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) - k_3 \left( \mathcal{A}_3 k_3 + \mathcal{A}_4 \frac{\partial k_4}{\partial s} \right) \right] B_1
$$

$$
+ \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_3 k_3 + \mathcal{A}_4 \frac{\partial k_4}{\partial s} \right) + k_2 \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) \right] B_2.
$$

By the Frenet equations we have

$$
\frac{\partial k_1}{\partial t} = \frac{\partial}{\partial t} \left( \frac{\partial}{\partial s} T, N \right).
$$

(15)

Also,

$$
\frac{\partial k}{\partial t} = \left( \frac{\partial}{\partial t} \frac{\partial}{\partial s} T, N \right) + \left( \frac{\partial}{\partial s} T, \frac{\partial}{\partial t} N \right).
$$

(16)

By the definition of flow, we have

$$
\left( N, \frac{\partial}{\partial t} N \right) = 0.
$$

(17)

Combining these we have

$$
\frac{\partial k_1}{\partial t} = \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) - k_2 \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right].
$$

(18)

Thus, we obtain the theorem. This completes the proof. □

By this theorem we immediately have the following.

**Theorem 5.** Consider

$$
\frac{\partial N}{\partial t} = \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) T
$$

$$
+ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_4 k_4 + \mathcal{A}_3 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) - k_3 \left( \mathcal{A}_4 k_4 + \mathcal{A}_3 \frac{\partial k_4}{\partial s} \right) \right] \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right),
$$

(21)

where $\mathcal{A}_1$, $\mathcal{A}_2$, $\mathcal{A}_3$, $\mathcal{A}_4$ are smooth functions of time and arc length.

**Proof.** Using Frenet equations, we have

$$
\frac{\partial}{\partial t} \frac{\partial}{\partial s} T = \frac{\partial k_1}{\partial t} N + k_1 \frac{\partial N}{\partial t}.
$$

(20)

This implies

$$
\frac{\partial N}{\partial t} = \left( \mathcal{A}_1 k_1 + \mathcal{A}_2 \frac{\partial k_1}{\partial s} - \mathcal{A}_3 k_2 \right) T
$$

$$
+ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_4 k_4 + \mathcal{A}_3 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) - k_3 \left( \mathcal{A}_4 k_4 + \mathcal{A}_3 \frac{\partial k_4}{\partial s} \right) \right] \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right)
$$

$$
\left( \mathcal{A}_3 k_3 + \mathcal{A}_4 \frac{\partial k_4}{\partial s} \right) \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right),
$$

(21)

which completes the proof. □

**Theorem 6.** Let $\partial \alpha / \partial t$ be inextensible flow of $\alpha$ in $\mathbb{E}_1^4$. Then,

$$
\frac{\partial B_1}{\partial t} = \frac{1}{k_2} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_4 k_4 + \mathcal{A}_3 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) - \frac{\partial k_1}{\partial t} \right] T
$$

$$
+ \frac{1}{k_2} \left[ k_1 \left( \mathcal{A}_4 k_4 + \mathcal{A}_3 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right)
$$

$$
\left( \mathcal{A}_3 k_3 + \mathcal{A}_4 \frac{\partial k_4}{\partial s} \right) \left( \mathcal{A}_2 k_2 + \mathcal{A}_4 \frac{\partial k_3}{\partial s} - \mathcal{A}_4 k_3 \right),
$$

which completes the proof. □
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[-k_3 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} \right) \]
\[+ k_2 \left( \bar{A}_4k_1 + \frac{\partial \bar{A}_2}{\partial s} - \bar{A}_3k_2 \right) \]
\[\left. - k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right] \right] \mathbf{N} \]
\[+ \frac{1}{k_2} \left[ k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[-k_3 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} \right) \]
\[+ k_2 \left( \bar{A}_4k_1 + \frac{\partial \bar{A}_2}{\partial s} - \bar{A}_3k_2 \right) \]
\[\left. + \frac{\partial}{\partial s} \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_3k_3 + \frac{\partial \bar{A}_4}{\partial s} \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \mathbf{B}_2, \quad (22) \]

where \( \bar{A}_1, \bar{A}_2, \bar{A}_3, \bar{A}_4 \) are smooth functions of time and arc length.

**Proof.** Assume that \( \partial \alpha / \partial t \) is inextensible flow of \( \alpha \) in \( \mathbb{E}^4 \).

Consider
\[
\frac{\partial}{\partial t} \frac{\partial \mathbf{N}}{\partial s} = \frac{\partial}{\partial s} \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \mathbf{T} \]
\[+ \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \mathbf{N} \]
\[+ \frac{\partial}{\partial s} \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_3k_3 + \frac{\partial \bar{A}_4}{\partial s} \right) \right] \right. \]
\[-k_3 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} \right) \]
\[+ k_2 \left( \bar{A}_4k_1 + \frac{\partial \bar{A}_2}{\partial s} - \bar{A}_3k_2 \right) \]
\[\left. \right] \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \mathbf{B}_2, \quad (23) \]

Thus we compute
\[
\frac{\partial}{\partial t} \frac{\partial \mathbf{N}}{\partial s} = \frac{\partial}{\partial s} \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \mathbf{T} + k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \mathbf{B}_1 \]
\[+ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \mathbf{B}_1 + k_2 \frac{\partial}{\partial t} \mathbf{B}_1. \quad (24) \]

Then we can easily see that
\[
k_2 \frac{\partial \mathbf{B}_1}{\partial t} \]
\[= \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right. \]
\[-k_3 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \]
\[+ k_2 \left( \bar{A}_4k_1 + \frac{\partial \bar{A}_2}{\partial s} - \bar{A}_3k_2 \right) \]
\[\left. \right] \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \left[ k_1 \left( \bar{A}_1k_1 + \frac{\partial \bar{A}_4}{\partial s} - \bar{A}_3k_2 \right) \right. \]
\[-k_2 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \right] \right. \]
\[+ k_3 \left( \bar{A}_2k_2 + \frac{\partial \bar{A}_3}{\partial s} - \bar{A}_4k_3 \right) \]
\[\left. \right] \mathbf{B}_2, \quad (25) \]
\[-k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] + k_3 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] \] \\
\[-k_1 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \] \\
+ \left[ k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] - k_3 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] \right] \] \\
- k_1 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \] \\
\right] B_2.

(25)

From definition of flow, we have

\[ \left\langle B_1, \frac{\partial B_1}{\partial t} \right\rangle = 0. \]  

(26)

Thus, we obtain the theorem. The proof of theorem is completed.

Now we give the characterization of evolution of second curvature as below.

**Theorem 7.** The evolution of $k_2$ is given by

\[ \frac{\partial k_2}{\partial t} = \frac{\partial}{\partial s} \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] - k_3 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] \] \\
+ \left[ k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] - k_3 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \right] \right] \] \\
- k_1 \left( \mathcal{A}_2 k_2 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 k_3 \right) \] \\
\right\rangle N

(27)
where $\mathcal{A}_1$, $\mathcal{A}_2$, $\mathcal{A}_3$, $\mathcal{A}_4$ are smooth functions of time and arc length.

**Proof.** It is obvious from Theorem 6. This completes the proof. \[ \square \]
\[-k_3 \left[ \frac{1}{k_2} k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( A_{1k_1} + \frac{\partial A_{1k_2}}{\partial s} - A_{1k_3} \right) \right] - k_1 \left( A_{1k_3} + \frac{\partial A_{1k_4}}{\partial s} \right) \right] + k_2 \left( A_{1k_1} + \frac{\partial A_{1k_2}}{\partial s} - A_{1k_3} \right) \right] \right] \left[ \frac{1}{k_2} k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( A_{2k_1} + \frac{\partial A_{2k_2}}{\partial s} - A_{2k_3} \right) \right] - k_1 \left( A_{2k_3} + \frac{\partial A_{2k_4}}{\partial s} \right) \right] + k_2 \left( A_{2k_1} + \frac{\partial A_{2k_2}}{\partial s} - A_{2k_3} \right) \right] \right] B_1, \]

where $A_1$, $A_2$, $A_3$, $A_4$ are smooth functions of time and arc length.

**Proof.** Differentiating (22) with respect to $s$,

\[
\frac{\partial}{\partial s} \frac{\partial B_2}{\partial t} = \left[ \frac{\partial}{\partial s} \left[ \frac{1}{k_2} k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( A_{1k_1} + \frac{\partial A_{1k_2}}{\partial s} - A_{1k_3} \right) \right] - k_1 \left( A_{1k_3} + \frac{\partial A_{1k_4}}{\partial s} \right) \right] + k_2 \left( A_{1k_1} + \frac{\partial A_{1k_2}}{\partial s} - A_{1k_3} \right) \right] \right] \frac{\partial B_2}{\partial t} \]

\[
+ \left[ \frac{1}{k_2} \left[ k_3 \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( A_{2k_1} + \frac{\partial A_{2k_2}}{\partial s} - A_{2k_3} \right) \right] - k_1 \left( A_{2k_3} + \frac{\partial A_{2k_4}}{\partial s} \right) \right] + k_2 \left( A_{2k_1} + \frac{\partial A_{2k_2}}{\partial s} - A_{2k_3} \right) \right] \right] \frac{\partial B_2}{\partial t} \]

\[
+ \frac{\partial}{\partial s} \left[ \frac{1}{k_2} \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( A_{3k_1} + \frac{\partial A_{3k_2}}{\partial s} - A_{3k_3} \right) \right] - k_1 \left( A_{3k_3} + \frac{\partial A_{3k_4}}{\partial s} \right) \right] + k_2 \left( A_{3k_1} + \frac{\partial A_{3k_2}}{\partial s} - A_{3k_3} \right) \right] \right] \frac{\partial B_2}{\partial t} \]

\[
+ \frac{\partial}{\partial s} \left[ \frac{1}{k_2} \left[ \frac{1}{k_1} \left[ \frac{\partial}{\partial s} \left( A_{4k_1} + \frac{\partial A_{4k_2}}{\partial s} - A_{4k_3} \right) \right] - k_1 \left( A_{4k_3} + \frac{\partial A_{4k_4}}{\partial s} \right) \right] + k_2 \left( A_{4k_1} + \frac{\partial A_{4k_2}}{\partial s} - A_{4k_3} \right) \right] \right] \frac{\partial B_2}{\partial t} \]

where $A_1$, $A_2$, $A_3$, $A_4$ are smooth functions of time and arc length.

**Proof.** Differentiating (22) with respect to $s$,
Thus we easily obtain that
\[
\frac{\partial}{\partial t} \frac{\partial B}{\partial s} = k_2 \frac{\partial B_2}{\partial t} - k_2 \left( \mathcal{A}_1 \mathcal{K}_1 + \frac{\partial \mathcal{A}_1}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) T - \frac{\partial k_2}{\partial t} N
\]
\[
= -k_2 \frac{\partial}{\partial s} \left( \mathcal{A}_2 \mathcal{K}_2 + \frac{\partial \mathcal{A}_2}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) + k_2 \left( \mathcal{A}_1 \mathcal{K}_1 + \frac{\partial \mathcal{A}_1}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) B_1
\]
\[
+ k_2 \left( \mathcal{A}_2 \mathcal{K}_2 + \frac{\partial \mathcal{A}_2}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) B_2.
\]

(30)

Hence, the proof is complete.

Now we give the characterization of evolution of third curvature as below.

**Theorem 9.** If the flow is inextensible, then
\[
\frac{\partial k_3}{\partial t} = \frac{k_2}{k_1} \left( \frac{\partial}{\partial s} \left( \mathcal{A}_1 \mathcal{K}_1 + \frac{\partial \mathcal{A}_1}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) + k_3 \left( \mathcal{A}_2 \mathcal{K}_2 + \frac{\partial \mathcal{A}_2}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) \right)
\]
\[
+ \frac{\partial}{\partial s} \left[ \frac{1}{k_2} \left( \mathcal{A}_3 \mathcal{K}_3 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) \right]
\]
\[
- k_3 \left( \mathcal{A}_4 \mathcal{K}_4 + \frac{\partial \mathcal{A}_4}{\partial s} \right)
\]
\[
+ k_2 \left( \mathcal{A}_1 \mathcal{K}_1 + \frac{\partial \mathcal{A}_1}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right)
\]
\[
+ \frac{\partial}{\partial s} \left[ \frac{1}{k_1} \left( \mathcal{A}_3 \mathcal{K}_3 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) \right]
\]
\[
+ k_3 \left( \mathcal{A}_2 \mathcal{K}_2 + \frac{\partial \mathcal{A}_2}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right)
\]
\[
- k_1 \left( \mathcal{A}_3 \mathcal{K}_3 + \frac{\partial \mathcal{A}_3}{\partial s} - \mathcal{A}_4 \mathcal{K}_4 \right) \right)
\]
\[
(31)
\]

where \( \mathcal{A}_1, \mathcal{A}_2, \mathcal{A}_3, \mathcal{A}_4 \) are smooth functions of time and arc length.

**Proof.** It is obvious from Theorem 8. This completes the proof. □

**Conflict of Interests**

The author declares that there is no conflict of interests regarding the publication of this paper.

**References**


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