Research Article

Existence and Multiplicity of Positive Solutions for Dirichlet Problems in Unbounded Domains

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We consider the elliptic problem $-\Delta u + u = b(x)|u|^{p-2}u + h(x)$ in Ω , $u \in H_0^1(\Omega)$, where $2 < p < (2N/(N-2))$ ($N \geq 3$), $2 < p < \infty$ ($N = 2$), Ω is a smooth unbounded domain in $\mathbb{R}^N$, $b(x) \in C(\Omega)$, and $h(x) \in H^{-1}(\Omega)$. We use the shape of domain Ω to prove that the above elliptic problem has a ground-state solution if the coefficient $b(x)$ satisfies $b(x) \rightarrow b^\infty > 0$ as $|x| \rightarrow \infty$ and $b(x) \geq c$ for some suitable constants $c \in (0, b^\infty)$, and $h(x) \equiv 0$. Furthermore, we prove that the above elliptic problem has multiple positive solutions if the coefficient $b(x)$ also satisfies the above conditions, $h(x) \geq 0$ and $0 < \|h\|_{H^{-1}} < (p-2)(1/(p-1))^{(p-1)/(p-2)} [b_{\sup} S^p(\Omega)]^{1/(2-p)}$, where $S(\Omega)$ is the best Sobolev constant of subcritical operator in $H_0^1(\Omega)$ and $b_{\sup} = \sup_{x \in \Omega} b(x)$.

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1. Introduction

In this paper, we are concerned with the existence and multiplicity of positive solutions of the following elliptic problems:

$$\begin{aligned} -\Delta u + u &= b(x)|u|^{p-2}u + h(x) \quad \text{in } \Omega, \\ u &\in H_0^1(\Omega), \end{aligned} \tag{1.1}$$

where $2 < p < (2N/(N-2))$ ($N \geq 3$), $2 < p < \infty$ ($N = 2$), and Ω is a smooth unbounded domain in $\mathbb{R}^N$. We assume that $b(x) \in C(\Omega) \cap L^\infty(\Omega)$ satisfies

$$b(x) > 0, \quad \forall x \in \Omega, \tag{1.2}$$

and $h(x)$ satisfies

$$h(x) \in H^{-1}(\Omega), \quad h(x) \geq 0. \quad (1.3)$$

Associated with (1.1), we consider the energy functional J_h^b in the Sobolev space $H_0^1(\Omega)$:

$$J_h^b(u) = \frac{1}{2} \|u\|_{H^1}^2 - \frac{1}{p} \int_{\Omega} b(x) |u|^p - \int_{\Omega} h(x) u, \quad (1.4)$$

where $\|u\|_{H^1} = (\int_{\Omega} |\nabla u|^2 + u^2)^{1/2}$. By Rabinowitz [1, Proposition B.10], $J_h^b \in C^1(H_0^1(\Omega), \mathbb{R})$. It is well known that the solutions of (1.1) are the critical points of the energy functional J_h^b in $H_0^1(\Omega)$.

Under the assumption (1.3) and $h(x) \not\equiv 0$, (1.1) can be regarded as a perturbation problem of the following homogeneous elliptic equation:

$$\begin{aligned} -\Delta u + u &= b(x) |u|^{p-2} u \quad \text{in } \Omega, \\ u &\in H_0^1(\Omega). \end{aligned} \quad (1.5)$$

A typical approach for solving a problem of this kind is to use the minimax method:

$$\alpha_{\Gamma}^b(\Omega) = \inf_{\gamma \in \Gamma(\Omega)} \max_{t \in [0,1]} J_0^b(\gamma(t)), \quad (1.6)$$

where

$$\Gamma(\Omega) = \{\gamma \in C([0,1], H_0^1(\Omega)) \mid \gamma(0) = 0, \gamma(1) = e\}, \quad (1.7)$$

$J_0^b(e) = 0$, and $e \neq 0$. By the mountain pass lemma due to Ambrosetti and Rabinowitz [2], we called the nonzero critical point $u \in H_0^1(\Omega)$ of J_0^b is as ground-state solution of (1.5) in Ω if $J_0^b(u) = \alpha_{\Gamma}^b(\Omega)$. We note that the ground-state solutions of (1.5) in Ω can also be obtained by the Nehari minimization problem

$$\alpha_0^b(\Omega) = \inf_{v \in \mathbf{M}_0^b(\Omega)} J_0^b(v), \quad (1.8)$$

where $\mathbf{M}_0^b(\Omega) = \{u \in H_0^1(\Omega) \setminus \{0\} \mid \|u\|_{H^1}^2 = \int_{\Omega} b(x) |u|^p\}$. Note that $\mathbf{M}_0^b(\Omega)$ contains every nonzero solution of (1.5) in Ω , $\alpha_{\Gamma}^b(\Omega) = \alpha^b(\Omega) > 0$ (see Willem [3] and Wang and Wu [4]), and if $b(x) \equiv b^{\infty} > 0$ is a constant, then J_0^b and $\alpha_0^b(\Omega)$ are replaced by J_0^{∞} and $\alpha_0^{\infty}(\Omega)$, respectively.

That the existence of ground-state solutions of (1.5) is affected by the shape of the domain Ω and $b(x)$ that satisfies some suitable conditions has been the focus of a great deal of research in recent years. By the Rellich compactness theorem and the minimax method, it is easy to obtain a ground-state solution for (1.5) in bounded domains. When Ω is an unbounded domain and $b(x) \equiv b^{\infty}$, the existence of ground-state solutions has been established by several authors under various conditions. We mention, in particular, results by Berestycki and Lions [5], Lien et al. [6], Chen and Wang [7], and Del Pino and Felmer [8, 9]. In [5], $\Omega = \mathbb{R}^N$. Actually, Kwong [10] proved that the positive solution of (1.5) in $\mathbb{R}^N$ is unique. In [6], Ω is a periodic domain. In [7, 6], the domain Ω is required

to satisfy that

(Ω1) $\Omega = \Omega_1 \cup \Omega_2$, where Ω_1, Ω_2 are domains in $\mathbb{R}^N$ and $\Omega_1 \cap \Omega_2$ is bounded;

(Ω2) $\alpha_0^\infty(\Omega) < \min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2)\}$.

In [8, 9], for $1 \leq l \leq N-1$, $\mathbb{R}^N = \mathbb{R}^l \times \mathbb{R}^{N-l}$. For a point $x \in \mathbb{R}^N$, we have $x = (y, z)$, where $y \in \mathbb{R}^l$ and $z \in \mathbb{R}^{N-l}$. Let $y \in \mathbb{R}^l$, we denote by $\Omega^y \subset \mathbb{R}^{N-l}$ the projection of Ω onto $\mathbb{R}^{N-l}$, that is,

$$\Omega^y = \{z \in \mathbb{R}^{N-l} \mid (y, z) \in \Omega\}. \quad (1.9)$$

The domain Ω is required to satisfy that

(Ω3) Ω is a smooth subset of $\mathbb{R}^N$ and the projections Ω^y are bounded uniformly in $y \in \mathbb{R}^l$;

(Ω4) there exists a nonempty closed set $F \subset \mathbb{R}^{N-l}$ such that $F \subset \Omega^y$ for all $y \in \mathbb{R}^l$;

(Ω5) for each $\delta > 0$, there exists $K > 0$ such that

$$\Omega^y \subset \{z \in \mathbb{R}^{N-l} \mid \text{dist}(z, F) < \delta\} \quad (1.10)$$

for all $|y| \geq K$.

Moreover, when $\Omega = \mathbb{R}^N \setminus \omega$ is an exterior domain, where ω is a bounded domain. It is well known that (1.5) in $\mathbb{R}^N \setminus \omega$ does not admit any ground-state solution (see Benci and Cerami [12]). However, Bahri and Lions [11] and Benci and Cerami [12] asserted that (1.5) in $\mathbb{R}^N \setminus \omega$ has a higher-energy positive solution. As Ω is an Esteban-Lions domain, (1.5) in Ω does not admit any nontrivial solution (see Esteban and Lions [13]), where the definition of Esteban-Lions domain is as follows: for a proper unbounded domain Ω in $\mathbb{R}^N$, there exists $\chi \in \mathbb{R}^N$, $\|\chi\| = 1$ such that $n(x) \cdot \chi \geq 0$ and $n(x) \cdot \chi \not\equiv 0$ on $\partial\Omega$, where $n(x)$ is the unit outward normal vector to $\partial\Omega$ at the point x .

When $b(x) \not\equiv b^\infty$, which satisfies the condition (1.2), the existence of ground-state solutions of (1.5) has been established by the condition $b(x) \geq b^\infty$ and the existence of ground-state solutions of limit equation

$$\begin{aligned} -\Delta u + u &= b^\infty |u|^{p-2} u \quad \text{in } \Omega, \\ u &\in H_0^1(\Omega). \end{aligned} \quad (1.11)$$

On the other hand, for $\Omega = \mathbb{R}^N$ and $b(x) \leq b^\infty$ on $\mathbb{R}^N$ with a strict inequality on a set of positive measures, (1.5) in $\mathbb{R}^N$ does not admit any ground-state solution. However, Bahri and Lions [11], Cao [14], and Bahri and Li [15] asserted that (1.5) in $\mathbb{R}^N$ has a higher-energy positive solution under the coefficient $b(x)$ which satisfies conditions $b(x) \geq (1/2)^{(p-2)/2} b^\infty$ and $b(x) \rightarrow b^\infty$ as $|x| \rightarrow \infty$ such that the functional J_0^b in $H_0^1(\Omega)$ satisfies the Palais-Smale condition for energy level β with

$$\alpha_0^\infty(\mathbb{R}^N) < \beta < \alpha_0^\infty(\mathbb{R}^N) + \alpha_0^b(\mathbb{R}^N). \quad (1.12)$$

The first result of our paper is relaxing the condition $b(x) \geq b^\infty$ to show the existence of ground-state solution of (1.5) by the shape of domain Ω . First, we consider the following assumptions:

- ($\Omega 1'$) given $k \geq 0$ and $1 \leq m \leq k$, the domain $\Omega = \bigcup_{i=1}^k \Omega_i$, where $\Omega_i \cap \Omega_j$ is bounded for all $i \neq j$ and Ω_j is unbounded domain for all $j = 1, 2, \dots, m$;
- ($\Omega 2'$) the functional J_0^∞ in $H_0^1(\Omega)$ satisfies the Palais-Smale condition for energy level $\alpha_0^\infty(\Omega)$;
- (b1) $b(x) \geq (\alpha_0^\infty(\Omega)/\min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\})^{(p-2)/2} b^\infty$ and $b(x) \rightarrow b^\infty$ as $|x| \rightarrow \infty$.

Then we have the following result.

THEOREM 1.1. *If the domain Ω satisfies the conditions ($\Omega 1'$)-($\Omega 2'$) and $b(x)$ satisfies the condition (b1), then (1.5) in Ω has a ground-state solution.*

Remark 1.2. If the domain Ω satisfies the conditions ($\Omega 1$)-($\Omega 2$), then the functional J_0^∞ in $H_0^1(\Omega)$ satisfies the Palais-Smale condition for energy level $\alpha_0^\infty(\Omega)$, and we have

$$0 < \alpha_0^\infty(\Omega) < \min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\} \quad (1.13)$$

(see Lien et al. [6] and Chen and Wang [7]). Thus,

$$0 < \left(\frac{\alpha_0^\infty(\Omega)}{\min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}} \right)^{(p-2)/2} < 1. \quad (1.14)$$

It is known that the general unbounded domains in $\mathbb{R}^N$ can be classified into three kinds. If Ω is an unbounded domain in $\mathbb{R}^N$, then it satisfies one of the following conditions:

- (1) J_0^∞ in $H_0^1(\Omega)$ satisfies the Palais-Smale condition for energy level $\alpha_0^\infty(\Omega)$. In particular, (1.11) in Ω has a ground-state solution u_0 such that $J_0^\infty(u_0) = \alpha_0^\infty(\Omega)$;
- (2) J_0^∞ in $H_0^1(\Omega)$ does not satisfy the Palais-Smale condition for energy level $\alpha_0^\infty(\Omega)$, but (1.11) in Ω has a ground-state solution u_0 such that $J_0^\infty(u_0) = \alpha_0^\infty(\Omega)$;
- (3) equation (1.11) in Ω does not admit any ground-state solution.

In this motivation, consider a general unbounded domain Ω and its exterior domain $\Omega^c(r) = \Omega \setminus \overline{B^N(0; r)}$, and the following assumptions:

- ($\Omega 3'$) equation (1.11) in Ω has a ground state solution u_0 such that $J_0^\infty(u_0) = \alpha_0^\infty(\Omega)$.
- (b2) $b(x) \geq (\alpha_0^\infty(\Omega)/\lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)))^{(p-2)/2} b^\infty$ and $b(x) \rightarrow b^\infty$ as $|x| \rightarrow \infty$.

Then we have the following result.

THEOREM 1.3. *If the unbounded domain Ω satisfies the condition ($\Omega 3'$) and $b(x)$ satisfies the condition (b2), then (1.5) in Ω has a ground-state solution.*

Remark 1.4. (1) If the domain Ω satisfies the conditions ($\Omega 3$)-($\Omega 5$), J_0^∞ in $H_0^1(\Omega)$ satisfies the Palais-Smale condition for energy level $\alpha_0^\infty(\Omega)$. Then $\alpha_0^\infty(\Omega) < \alpha_0^\infty(\Omega^c(r))$ for all $r > 0$ (see Del Pino and Felmer [8, 9] or Wu [16]). Since $\alpha_0^\infty(\Omega^c(r))$ is nondecreasing as r is

increasing, we have

$$0 \leq \left(\frac{\alpha_0^\infty(\Omega)}{\lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r))} \right)^{(p-2)/2} < 1. \quad (1.15)$$

(2) If Ω is a periodic domain, then J_0^∞ in $H_0^1(\Omega)$ does not satisfy the Palais-Smale condition for energy level $\alpha_0^\infty(\Omega)$, but (1.11) in Ω has a ground-state solution u_0 such that $J_0^\infty(u_0) = \alpha_0^\infty(\Omega)$. Then $\alpha_0^\infty(\Omega) = \alpha_0^\infty(\Omega^c(r))$ for all $r > 0$ (see Lien et al. [6]). Thus,

$$\left(\frac{\alpha_0^\infty(\Omega)}{\lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r))} \right)^{(p-2)/2} \equiv 1. \quad (1.16)$$

Remark 1.5. If the domain $\Omega = \mathbb{R}^N$, coefficient $b(x)$ satisfies the condition (1.2) and $b(x) \leq b^\infty$ with a strict inequality on a set of positive measures, then (1.5) in $\mathbb{R}^N$ does not admit any ground-state solution and $\alpha_0^\infty(\mathbb{R}^N) = \alpha_0^b(\mathbb{R}^N)$. However, if the domain Ω satisfies the conditions $(\Omega 1)$ – $(\Omega 2)$ (or $(\Omega 3)$ – $(\Omega 5)$), $b(x)$ satisfies the condition $(b 1)$ (or $(b 2)$) and $b(x) \leq b^\infty$ with a strict inequality on a set of positive measure, then from Theorem 1.1 (or Theorem 1.3), we can conclude that (1.5) has a ground-state solution. Moreover, $\alpha_0^\infty(\Omega) < \alpha_0^b(\Omega)$.

Finally, we consider (1.1). For $\Omega = \mathbb{R}^N$, several authors have shown the existence of at least two positive solutions of (1.1) in $\mathbb{R}^N$ under some suitable conditions. In [17] by Zhu for $b(x) = b^\infty$, $h(x)$ is exponential decay and $\|h\|_{L^2}$ is sufficiently small. By Cao and Zhou in [18] and Jeanjean [19], for $b(x) \geq b^\infty$ and $\|h\|_{H^{-1}}$ sufficiently small. By Adachi and Tanaka in [20], for $b(x) \geq b^\infty - Ce^{-\lambda|x|}$ for some $C, \lambda > 0$ and $\|h\|_{H^{-1}}$ sufficiently small. Moreover, Adachi and Tanaka [21] used that (1.5) in $\mathbb{R}^N$ does not admit any ground-state solution for the condition $b(x) \leq b^\infty$ with a strict inequality on a set of positive measures, to show that (1.1) in $\mathbb{R}^N$ has at least four positive solutions for $\|h\|_{H^{-1}}$ sufficiently small. The second aim of our paper is also relaxing the condition $b(x) \geq b^\infty$ to show the existence of at least two positive solutions of (1.1) in Ω . Denote

$$b_{\sup} = \sup_{x \in \Omega} b(x) \quad (1.17)$$

and $S(\Omega) = [(2p/(p-2))\alpha_0^\infty(\Omega)]^{(2-p)/2p}$ is the best Sobolev constant of subcritical operator in $H_0^1(\Omega)$ (see Lin et al. [22] or Willem [3]). Then we have the following results.

THEOREM 1.6. *Suppose that the domain Ω satisfies the conditions $(\Omega 1')$ – $(\Omega 2')$ and $b(x)$ satisfies the condition $(b 1)$. If $h \geq 0$ and*

$$0 < \|h\|_{H^{-1}} < (p-2) \left(\frac{1}{p-1} \right)^{(p-1)/(p-2)} [b_{\sup} S^p(\Omega)]^{1/(2-p)}, \quad (1.18)$$

then (1.1) in Ω has at least two positive solutions.

THEOREM 1.7. *Suppose that the domain Ω satisfies the condition $(\Omega 3')$ and $b(x)$ satisfies the condition $(b2)$. If $h \geq 0$ and*

$$0 < \|h\|_{H^{-1}} < (p-2) \left(\frac{1}{p-1} \right)^{(p-1)/(p-2)} [b_{\sup} S^p(\Omega)]^{1/(2-p)}, \quad (1.19)$$

then (1.1) in Ω has at least two positive solutions.

This paper is organized as follows. In Section 2, we describe various preliminaries. In Section 3, we use the shape of the domain Ω to prove that (1.5) in Ω has a ground-state solution. In Section 4, we modify the proof of Adachi and Tanaka [21], Tarantello [23], Cao and Zhu [18], and Zhu [17] to prove that (1.1) in Ω has at least two positive solutions.

2. Preliminary

We define the Palais-Smale (PS) sequences, (PS) values, and (PS) conditions in $H_0^1(\Omega)$ for J_h^b as follows.

Definition 2.1. (i) For $\beta \in \mathbb{R}$, a sequence $\{u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b if $J_h^b(u_n) = \beta + o(1)$ and $(J_h^b)'(u_n) = o(1)$ strongly in $H^{-1}(\Omega)$ as $n \rightarrow \infty$;
(ii) $\beta \in \mathbb{R}$ is a (PS) value in $H_0^1(\Omega)$ for J_h^b if there is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b ;
(iii) J_h^b satisfies the $(PS)_\beta$ -condition in $H_0^1(\Omega)$ if every $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b contains a convergent subsequence;
(iv) J_h^b satisfies the (PS) condition in $H_0^1(\Omega)$ if for every $\beta \in \mathbb{R}$, J_h^b satisfies the $(PS)_\beta$ -condition in $H_0^1(\Omega)$.

We need the following lemmas.

LEMMA 2.2. *Let $u_n \rightharpoonup u$ weakly in $H_0^1(\Omega)$. Then there exists a subsequence $\{u_n\}$ such that*

- (i) $\{u_n\}$ is bounded in $H_0^1(\Omega)$ and $\|u\|_{H^1} \leq \liminf_{n \rightarrow \infty} \|u_n\|_{H^1}$;
- (ii) $u_n \rightharpoonup u$, $\nabla u_n \rightharpoonup \nabla u$ weakly in $L^2(\Omega)$, and $u_n \rightarrow u$ a.e. in Ω ;
- (iii) $\|u_n - u\|_{H^1}^2 = \|u_n\|_{H^1}^2 - \|u\|_{H^1}^2 + o(1)$.

The proof is clear by the routine arguments, and hence is omitted here.

LEMMA 2.3 (Brézis-Lieb lemma). *Suppose that $u_n \rightarrow u$ a.e. in Ω and there exists $c > 0$ such that $\|u_n\|_{L^p} \leq c$ for $n = 1, 2, \dots$. Then*

- (i) $\|u_n - u\|_{L^p}^p = \|u_n\|_{L^p}^p - \|u\|_{L^p}^p + o(1)$;
- (ii) $|u_n - u|^{p-2}(u_n - u) - |u_n|^{p-2}u_n + |u|^{p-2}u = o(1)$ in $L^{p/(p-1)}(\Omega)$.

For the proof, see Brézis and Lieb [24].

LEMMA 2.4. *Let $u_n \rightharpoonup u$ weakly in $H_0^1(\Omega)$ and*

$$(J_h^b)'(u_n) = -\Delta u_n + u_n - b(x)|u_n|^{p-2}u_n + h(x) = o(1) \quad \text{in } H^{-1}(\Omega). \quad (2.1)$$

Then

- (i) $|u_n - u|^{p-2}(u_n - u) - |u_n|^{p-2}u_n + |u|^{p-2}u = o(1)$ in $H^{-1}(\Omega)$;
- (ii) $(J_0^\infty)'(w_n) = -\Delta w_n + w_n - b^\infty |w_n|^{p-2}w_n = o(1)$ in $H^{-1}(\Omega)$, where $w_n = u_n - u$;
- (iii) if $\{u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b then $\{w_n\}$ is a $(PS)_{(\beta - J_h^b(u))}$ -sequence in $H_0^1(\Omega)$ for J_0^∞ .

Proof. For (i), (ii), see Bahri and Lions [11]. (iii) Since $u_n \rightharpoonup u$ weakly in $H_0^1(\Omega)$ and $\{u_n\}$ is a $(PS)_\beta$ -sequence for J_h^b in $H_0^1(\Omega)$, by Lemmas 2.2, 2.3, and the Sobolev embedding theorem, there exists a subsequence $\{u_n\}$ such that $w_n \rightharpoonup 0$ in $H_0^1(\Omega)$,

$$\begin{aligned} \|w_n\|_{H^1}^2 &= \|u_n\|_{H^1}^2 - \|u\|_{H^1}^2 + o(1), \\ \|w_n\|_{L^p}^p &= \|u_n\|_{L^p}^p - \|u\|_{L^p}^p + o(1). \end{aligned} \quad (2.2)$$

Thus,

$$J_0^\infty(w_n) = J_h^b(w_n) + o(1) = J_h^b(u_n) - J_h^b(u) + o(1) = \beta - J_h^b(u) + o(1). \quad (2.3)$$

Therefore, by part (ii), $\{p_n\}$ is a $(PS)_{(\beta - J_h^b(u))}$ -sequence in $H_0^1(\Omega)$ for J_0^∞ . $\square$

We need the following useful results.

LEMMA 2.5. *Let $\{u_n\}$ be a sequence in $H_0^1(\Omega)$. Then $\{u_n\}$ is a $(PS)_{\alpha_0^b(\Omega)}$ -sequence for J_0^b if and only if $J_0^b(u_n) = \alpha_0^b(\Omega) + o(1)$ and $\int_\Omega |\nabla u_n|^2 + u_n^2 = \int_\Omega b(x)|u_n|^p + o(1)$. In particular, every minimizing sequence $\{u_n\}$ in $\mathbf{M}_0^b(\Omega)$ of $\alpha_0^b(\Omega)$ is a $(PS)_{\alpha_0^b(\Omega)}$ -sequence in $H_0^1(\Omega)$ for J_0^b .*

The proof is almost the same as that by Wang and Wu in [4, Lemma 7], and is omitted here.

We introduce the Nehari minimization problem for (1.1) as

$$\alpha_h^b(\Omega) = \inf_{u \in \mathbf{M}_h^b(\Omega)} J_h^b(u), \quad (2.4)$$

where $\mathbf{M}_h^b(\Omega) = \{u \in H_0^1(\Omega) \setminus \{0\} \mid \langle (J_h^b)'(u), u \rangle = 0\}$. Define

$$\psi(u) = \langle (J_h^b)'(u), u \rangle = \|u\|_{H^1}^2 - \int_\Omega b(x)|u|^p - \int_\Omega h(x)u. \quad (2.5)$$

Then we have the following result.

LEMMA 2.6. *If $\|h\|_{H^{-1}} < (p-2)(1/(p-1))^{(p-1)/(p-2)}[b_{\sup} S^p(\Omega)]^{1/(2-p)}$, then for each $u \in \mathbf{M}_h^b(\Omega)$,*

$$\langle \psi'(u), u \rangle = \|u\|_{H^1}^2 - (p-1) \int_\Omega b(x)|u|^p \neq 0. \quad (2.6)$$

Proof. For $u \in \mathbf{M}_h^b(\Omega)$, we have

$$\|u\|_{H^1}^2 - \int_\Omega b(x)|u|^p - \int_\Omega h(x)u = 0. \quad (2.7)$$

Then

$$\begin{aligned}\langle \psi'(u), u \rangle &= 2\|u\|_{H^1}^2 - p \int_{\Omega} b(x)|u|^p - \int_{\Omega} h(x)u \\ &= \|u\|_{H^1}^2 - (p-1) \int_{\Omega} b(x)|u|^p.\end{aligned}\quad (2.8)$$

We claim that if $\|h\|_{H^{-1}} < (p-2)(1/(p-1))^{(p-1)/(p-2)}[b_{\sup}S^p(\Omega)]^{1/(2-p)}$, then $\langle \psi'(u), u \rangle \neq 0$ for all $u \in \mathbf{M}_h^b(\Omega)$. Let $I : \mathbf{M}_h^b(\Omega) \rightarrow \mathbb{R}$ be given by

$$I(u) = K(p) \left(\frac{\|u\|_{H^1}^{2p-2}}{\int_{\Omega} b(x)|u|^p} \right)^{1/(p-2)} - \int_{\Omega} h(x)u, \quad (2.9)$$

where $K(p) = (p-2)(1/(p-1))^{(p-1)/(p-2)}$. Then we have for $u \in \mathbf{M}_h^b(\Omega)$,

$$\begin{aligned}I(u) &= K(p) \left(\frac{\|u\|_{H^1}^{2p-2}}{\int_{\Omega} b(x)|u|^p} \right)^{1/(p-2)} - \int_{\Omega} h(x)u \\ &\geq K(p) \left(\frac{\|u\|_{H^1}^{2p-2}}{\int_{\Omega} b(x)|u|^p} \right)^{1/(p-2)} - \|h\|_{H^{-1}} \|u\|_{H^1} \\ &= \|u\|_{H^1} \left(K(p) \left(\frac{\|u\|_{H^1}^p}{\int_{\Omega} b(x)|u|^p} \right)^{1/(p-2)} - \|h\|_{H^{-1}} \right)\end{aligned}\quad (2.10)$$

since

$$\left(\frac{\|u\|_{H^1}^p}{\int_{\Omega} b(x)|u|^p} \right)^{1/(p-2)} \geq [b_{\sup}S^p(\Omega)]^{1/(2-p)} \quad \forall u \in H_0^1(\Omega) \setminus \{0\}. \quad (2.11)$$

Thus, for $\|h\|_{H^{-1}} < K(p)[b_{\sup}S^p(\Omega)]^{1/(2-p)}$, we have

$$I(u) > 0 \quad \forall u \in \mathbf{M}_h^b(\Omega). \quad (2.12)$$

Assume that there is a $w \in \mathbf{M}_h^b(\Omega)$ such that $\langle \psi'(w), w \rangle = 0$, then we have

$$\begin{aligned}\|w\|_{H^1}^2 &= (p-1) \int_{\Omega} b(x)|w|^p, \\ \int_{\Omega} h(x)w &= \|w\|_{H^1}^2 - \int_{\Omega} b(x)|w|^p = (p-2) \int_{\Omega} b(x)|w|^p.\end{aligned}\quad (2.13)$$

From (2.12) and (2.13),

$$\begin{aligned}0 < I(w) &= K(p) \left(\frac{\|w\|_{H^1}^{2p-2}}{\int_{\Omega} b(x)|w|^p} \right)^{1/(p-2)} - \int_{\Omega} h(x)w \\ &= \left(\frac{1}{p-1} \right)^{(p-1)/(p-2)} (p-2) \left(\frac{(p-1)^{p-1} [\int_{\Omega} b(x)|w|^p]^{p-1}}{\int_{\Omega} b(x)|w|^p} \right)^{1/(p-2)} - (p-2) \int_{\Omega} h(x)w = 0,\end{aligned}\quad (2.14)$$

which is a contradiction. Thus, we can conclude that for

$$\|h\|_{H^{-1}} < (p-2) \left(\frac{1}{p-1} \right)^{(p-1)/(p-2)} [b_{\sup} S^p(\Omega)]^{1/(2-p)}, \quad (2.15)$$

we have $\langle \psi'(u), u \rangle \neq 0$ for all $u \in \mathbf{M}_h^b(\Omega)$. $\square$

By Lemma 2.6, we write $\mathbf{M}_h^b(\Omega) = \mathbf{M}_h^{b+}(\Omega) \cup \mathbf{M}_h^{b-}(\Omega)$, where

$$\begin{aligned} \mathbf{M}_h^{b+}(\Omega) &= \left\{ u \in \mathbf{M}_h^b(\Omega) \mid \|u\|_{H^1}^2 - (p-1) \int_{\Omega} b(x)|u|^p > 0 \right\}, \\ \mathbf{M}_h^{b-}(\Omega) &= \left\{ u \in \mathbf{M}_h^b(\Omega) \mid \|u\|_{H^1}^2 - (p-1) \int_{\Omega} b(x)|u|^p < 0 \right\}, \end{aligned} \quad (2.16)$$

and define

$$\alpha_h^{b+}(\Omega) = \inf_{u \in \mathbf{M}_h^{b+}(\Omega)} J_h^b(u), \quad \alpha_h^{b-}(\Omega) = \inf_{u \in \mathbf{M}_h^{b-}(\Omega)} J_h^b(u). \quad (2.17)$$

For each $u \in H_0^1(\Omega) \setminus \{0\}$, we write

$$t_{\max} = \left(\frac{\|u\|_{H^1}^2}{(p-1) \int_{\Omega} b(x)|u|^p} \right)^{1/(p-2)} > 0. \quad (2.18)$$

Similar as the proof of some results by Tarantello in [23], we have the following two lemmas.

LEMMA 2.7. For each $u \in H_0^1(\Omega) \setminus \{0\}$,

- (i) there is a unique $t^- = t^-(u) > t_{\max} > 0$ such that $t^-u \in \mathbf{M}_h^{b-}(\Omega)$ and $J_h^b(t^-u) = \max_{t \geq t_{\max}} J_h^b(tu)$;
- (ii) $t^-(u)$ is a continuous function for nonzero u ;
- (iii) $\mathbf{M}_h^{b-}(\Omega) = \{u \in H_0^1(\Omega) \setminus \{0\} \mid (1/\|u\|_{H^1})t^-(u/\|u\|_{H^1}) = 1\}$;
- (iv) if $\int_{\Omega} hu > 0$, then there is a unique $0 < t^+ = t^+(u) < t_{\max}$ such that $t^+u \in \mathbf{M}_h^{b+}(\Omega)$ and $J_h^b(t^+u) = \min_{0 \leq t \leq t^-} J_h^b(tu)$.

LEMMA 2.8. (i) For each $u \in \mathbf{M}_h^{b+}(\Omega)$, $\int_{\Omega} h(x)u > 0$ and $J_h^b(u) < 0$. In particular, $\alpha_h(\Omega) \leq \alpha_h^{b+}(\Omega) < 0$;

(ii) J_h^b is coercive and bounded below on $\mathbf{M}_h^b(\Omega)$.

Proof. (i) For each $u \in \mathbf{M}_h^{b+}(\Omega)$, $\|u\|_{H^1}^2 - (p-1) \int_{\Omega} b(x)|u|^p > 0$ and

$$\|u\|_{H^1}^2 = \int_{\Omega} b(x)|u|^p + \int_{\Omega} h(x)u. \quad (2.19)$$

Thus,

$$\int_{\Omega} h(x)u = \|u\|_{H^1}^2 - \int_{\Omega} b(x)|u|^p > (p-2) \int_{\Omega} b(x)|u|^p > 0, \quad (2.20)$$

and hence

$$\begin{aligned}
 J_h^b(u) &= \left(\frac{1}{2} - \frac{1}{p}\right) \int_{\Omega} b(x)|u|^p - \frac{1}{2} \int_{\Omega} h(x)u \\
 &< \frac{p-2}{2p} \int_{\Omega} b(x)|u|^p - \frac{p-2}{2} \int_{\Omega} b(x)|u|^p \\
 &= -\frac{(p-1)(p-2)}{2p} \int_{\Omega} b(x)|u|^p < 0.
 \end{aligned} \tag{2.21}$$

(ii) Is similar to the proof of Theorem 1 by Tarantello in [23]. $\square$

3. Homogeneous problems

First, we present several (PS) conditions in $H_0^1(\Omega)$ for J_0^b which are used to prove our main results. As a consequence of Lemma 2.8(ii), for each $(PS)_{\beta}$ -sequence $\{u_n\}$ in $H_0^1(\Omega)$ for J_0^b , there exist a subsequence $\{u_n\}$ and u_0 in $H_0^1(\Omega)$ such that $u_n \rightharpoonup u_0$ weakly in $H_0^1(\Omega)$. Then u_0 is a solution of (1.5) in Ω . Moreover, we have the following lemma.

Let Ω be any unbounded domain and $\xi \in C^\infty([0, \infty))$ such that $0 \leq \xi \leq 1$ and

$$\xi(t) = \begin{cases} 0 & \text{for } t \in [0, 1] \\ 1 & \text{for } t \in [2, \infty). \end{cases} \tag{3.1}$$

Let

$$\xi_n(z) = \xi\left(\frac{2|z|}{n}\right). \tag{3.2}$$

Then we have the following result.

LEMMA 3.1. *Let $\{u_n\}$ be a $(PS)_{\beta}$ -sequence in $H_0^1(\Omega)$ for J_0^b satisfying $u_n \rightharpoonup 0$ weakly in $H_0^1(\Omega)$ and let $v_n = \xi_n u_n$. Then there exists a subsequence $\{u_n\}$ such that*

- (i) $\|u_n - v_n\|_{H^1} = o(1)$ as $n \rightarrow \infty$;
- (ii) $\int_{\Omega} b(x)|u_n|^p = \int_{\Omega} b(x)|v_n|^p + o(1) = \int_{\Omega} b^\infty |v_n|^p + o(1)$;
- (iii) $\int_{\Omega} |\nabla v_n|^2 + v_n^2 = \int_{\Omega} b^\infty |v_n|^p + o(1)$;
- (iv) $\{v_n\}$ is a $(PS)_{\beta}$ -sequence in $H_0^1(\Omega)$ for J_0^∞ .

Proof. By the fact that

$$\|u_n - v_n\|_{H^1}^2 = \|u_n\|_{H^1}^2 + \|v_n\|_{H^1}^2 - 2\langle u_n, v_n \rangle_{H^1}, \tag{3.3}$$

thus it suffices to show that $\langle u_n, v_n \rangle_{H^1} = \|u_n\|_{H^1}^2 + o(1) = \|v_n\|_{H^1}^2 + o(1)$. Since

$$\langle u_n, v_n \rangle_{H^1} = \int_{\Omega} \nabla u_n \nabla v_n + u_n v_n = \int_{\Omega} \xi_n [|\nabla u_n|^2 + u_n^2] + \int_{\Omega} u_n \nabla u_n \nabla \xi_n, \tag{3.4}$$

$|\nabla \xi_n| \leq c/n$ and $\{u_n\}$ is a $(PS)_{\beta}$ -sequence in $H_0^1(\Omega)$ for J_0^b , it follows that

$$\int_{\Omega} \xi_n^q u_n \nabla u_n \nabla \xi_n = o(1) \quad \text{for } q > 0. \tag{3.5}$$

Hence,

$$\langle u_n, v_n \rangle_{H^1} = \int_{\Omega} \xi_n [|\nabla u_n|^2 + u_n^2] + o(1). \quad (3.6)$$

Similarly, we have

$$\|v_n\|_{H^1}^2 = \int_{\Omega} \xi_n^2 [|\nabla u_n|^2 + u_n^2] + o(1). \quad (3.7)$$

Given $r \geq 1$, since $\{\xi_n^r u_n\}$ is bounded in $H_0^1(\Omega)$, we have

$$\begin{aligned} o(1) &= \langle (J_0^b)'(u_n), \xi_n^r u_n \rangle \\ &= \int_{\Omega} (\xi_n^r |\nabla u_n|^2 + r \xi_n^{r-1} u_n \nabla \xi_n \nabla u_n + \xi_n^r u_n^2) - \int_{\Omega} b(x) \xi_n^r |u_n|^p. \end{aligned} \quad (3.8)$$

From (3.5), we can conclude that

$$\int_{\Omega} \xi_n^r (|\nabla u_n|^2 + u_n^2) = \int_{\Omega} b(x) \xi_n^r |u_n|^p + o(1). \quad (3.9)$$

Since $u_n \rightharpoonup 0$ weakly in $H_0^1(\Omega)$ and $b(x) \rightarrow b^\infty$ as $|x| \rightarrow \infty$, there exists a subsequence $\{u_n\}$ such that $u_n \rightarrow 0$ strongly in $L_{\text{loc}}^p(\Omega)$, or there exists a subsequence $\{u_n\}$ such that

$$\int_{Q(n)} b(x) |u_n|^p = o(1), \quad (3.10)$$

where $Q(n) = \Omega \cap B^N(0; n)$. Clearly,

$$\int_{\Omega} b(x) |u_n|^p = \int_{\Omega} b(x) \xi_n^r |u_n|^p + o(1) = \int_{\Omega} b^\infty \xi_n^r |u_n|^p + o(1). \quad (3.11)$$

By (3.6), (3.7), (3.9), and (3.11),

$$\begin{aligned} \langle u_n, v_n \rangle_{H^1} &= \|u_n\|_{H^1}^2 + o(1) = \|v_n\|_{H^1}^2 + o(1), \\ \int_{\Omega} b(x) |u_n|^p &= \int_{\Omega} b(x) |v_n|^p + o(1) = \int_{\Omega} b^\infty |v_n|^p + o(1). \end{aligned} \quad (3.12)$$

Therefore, $\|u_n - v_n\|_{H^1} = o(1)$ as $n \rightarrow \infty$. The results of (iii) and (iv), from (i), (ii) and Lemmas 2.4, 2.5. $\square$

We need the following compactness results.

PROPOSITION 3.2. *Suppose that the domain Ω satisfies the conditions $(\Omega 1')$ – $(\Omega 2')$. If $\{u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_0^b with*

$$\alpha_0^b(\Omega) \leq \beta < \min \{\alpha_0^\infty(\Omega) + \alpha_0^b(\Omega), \alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}, \quad (3.13)$$

then there exist a subsequence $\{u_n\}$ and $u_0 \neq 0$ such that $u_n \rightarrow u_0$ strongly in $H_0^1(\Omega)$ and $J_0^b(u_0) = \beta$.

Proof. Let $\{u_n\}$ be a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_0^b with

$$\alpha_0^b(\Omega) \leq \beta < \min \{ \alpha_0^\infty(\Omega) + \alpha_0^b(\Omega), \alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m) \}. \quad (3.14)$$

Since $\{u_n\}$ is bounded, there exist a subsequence $\{u_n\}$ and u_0 in $H_0^1(\Omega)$ such that $u_n \rightharpoonup u_0$ weakly in $H_0^1(\Omega)$ and $u_n \rightarrow u_0$ a.e in Ω . Moreover, u_0 is a solution of (1.5) in Ω . If $u_0 \equiv 0$, by Lemma 3.1 there exists a subsequence $\{u_n\}$ such that $\{\xi_n u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_0^∞ , where ξ_n is as in (3.2). Let $v_n = \xi_n u_n$, and we obtain

$$J_0^\infty(v_n) = \beta + o(1), \quad (J_0^\infty)'(v_n) = o(1) \quad \text{in } H^{-1}(\Omega). \quad (3.15)$$

Since $\Omega_i \cap \Omega_j$ is bounded for $i \neq j$ and Ω_l is also bounded for $m+1 \leq l \leq k$, there exists $n_0 \in \mathbb{N}$ such that $v_n = 0$ in $\overline{\Omega(n_0)}$ for $n > 2n_0$ and $\Omega_l \subset \Omega(n_0)$ for all $l \in \{m+1, m+2, \dots, k\}$, where $\Omega(n) = \Omega \cap B^N(0; n)$. Moreover, $v_n = v_n^1 + v_n^2 + \dots + v_n^m$ and for $i = 1, 2, \dots, m$,

$$v_n^i(z) = \begin{cases} v_n(z) & \text{for } z \in \Omega_i, \\ 0, & \text{for } z \notin \Omega_i. \end{cases} \quad (3.16)$$

Then $v_n^i \in H_0^1(\Omega_i)$ and

$$\int_{\Omega_i} (|\nabla v_n^i|^2 + (v_n^i)^2) = \int_{\Omega_i} b^\infty |v_n^i|^p + o(1). \quad (3.17)$$

By (3.15), we obtain

$$\begin{aligned} (J_0^\infty)'(v_n^i) &= o(1) \quad \text{strongly in } H^{-1}(\Omega_i) \text{ for } i = 1, 2, \dots, m, \\ \beta &= J_0^\infty(v_n) + o(1) = \sum_{i=1}^m J_0^\infty(v_n^i) + o(1). \end{aligned} \quad (3.18)$$

Assume that

$$J_0^\infty(v_n^i) = c_i + o(1) \quad \text{for } i = 1, 2, \dots, m, \quad (3.19)$$

then $c_1 + c_2 + \dots + c_m = \beta$, since all of c_i are (PS) -values in $H_0^1(\Omega_i)$ for J_0^∞ and nonnegative. Thus, there exists $i_0 \in \{1, 2, \dots, m\}$ such that c_{i_0} are positive (PS) -values in $H_0^1(\Omega_{i_0})$ for J_0^∞ and

$$\alpha_0^\infty(\Omega_{i_0}) \leq c_{i_0} \leq \beta, \quad (3.20)$$

which contradicts (3.14). Consequently, $u_0 \neq 0$ and $\beta \geq J_0^b(u_0) \geq \alpha_0^b(\Omega)$. Let $p_n = u_n - u_0$. By Lemma 2.4, $\{p_n\}$ is a $(PS)_{(\beta - J_0^b(u_0))}$ -sequence in $H_0^1(\Omega)$ for J_0^∞ . Since $\beta < \alpha_0^\infty(\Omega) + \alpha_0^b(\Omega)$, $J_0^b(u_0) \geq \alpha_0^b(\Omega)$ and $\alpha_0^b(\Omega)$ is a smallest positive (PS) -value in $H_0^1(\Omega)$ for J_0^b . Thus, $\beta - J_0^b(u_0) = 0$. This implies that $u_n \rightarrow u_0$ strongly in $H_0^1(\Omega)$ and $J_0^b(u_0) = \beta$. $\square$

PROPOSITION 3.3. Suppose that the unbounded domain Ω satisfies the condition $(\Omega 3')$. If $\{u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_0^b with

$$\alpha_0^b(\Omega) \leq \beta < \min \left\{ \alpha_0^\infty(\Omega) + \alpha_0^b(\Omega), \lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)) \right\}, \quad (3.21)$$

then there exist a subsequence $\{u_n\}$ and $u_0 \neq 0$ such that $u_n \rightarrow u_0$ strongly in $H_0^1(\Omega)$ and $J_0^b(u_0) = \beta$.

Proof. Let $\{u_n\}$ be a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_0^b with

$$\alpha_0^b(\Omega) \leq \beta < \min \left\{ \alpha_0^\infty(\Omega) + \alpha_0^b(\Omega), \lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)) \right\}. \quad (3.22)$$

Since $\{u_n\}$ is bounded, there exist a subsequence $\{u_n\}$ and u_0 in $H_0^1(\Omega)$ such that $u_n \rightharpoonup u_0$ weakly in $H_0^1(\Omega)$ and $u_n \rightarrow u_0$ a.e in Ω . Moreover, u_0 is a solution of (1.5) in Ω . If $u_0 \equiv 0$, by Lemma 3.1 there exists a subsequence $\{u_n\}$ such that $\{\xi_n u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_0^∞ , where ξ_n is as in (3.2). Let $v_n = \xi_n u_n$, we obtain $v_n \in H_0^1(\Omega^c(n))$ for each n ,

$$J_0^\infty(v_n) = \beta + o(1), \quad (J_0^\infty)'(v_n) = o(1) \quad \text{in } H^{-1}(\Omega). \quad (3.23)$$

Moreover, there is an $s_n > 0$ such that $s_n v_n \in \mathbf{M}^\infty(\Omega^c(n))$ and $s_n = 1 + o(1)$. Then

$$J_0^\infty(s_n v_n) \geq \alpha_0^\infty(\Omega^c(n)). \quad (3.24)$$

By (3.23), (3.24), we obtain

$$\beta \geq \lim_{n \rightarrow \infty} \alpha_0^\infty(\Omega^c(n)), \quad (3.25)$$

which contradicts (3.22). Consequently, $u_0 \neq 0$ and $\beta \geq J_0^b(u_0) \geq \alpha_0^b(\Omega)$. Let $p_n = u_n - u_0$. By Lemma 2.4, $\{p_n\}$ is a $(PS)_{(\beta - J_0^b(u_0))}$ -sequence in $H_0^1(\Omega)$ for J_0^∞ . Since $\beta < \alpha_0^\infty(\Omega) + \alpha_0^b(\Omega)$, $J_0^b(u_0) \geq \alpha_0^b(\Omega)$ and $\alpha_0^b(\Omega)$ is smallest positive (PS) -value in $H_0^1(\Omega)$ for J_0^b . Thus, $\beta - J_0^b(u_0) = 0$. This implies that $u_n \rightarrow u_0$ strongly in $H_0^1(\Omega)$ and $J_0^b(u_0) = \beta$. $\square$

Now, we begin to show the proof of Theorem 1.1: since the domain Ω satisfies the conditions $(\Omega 1')$ – $(\Omega 2')$, we have (1.11), and there exists a ground-state solution u_0 such that $J_0^\infty(u_0) = \alpha_0^\infty(\Omega)$. Let $s_0 > 0$ with $s_0 u_0 \in \mathbf{M}_0^b(\Omega)$. Then

$$s_0^2 \int_{\Omega} (|\nabla u_0|^2 + u_0^2) = s_0^p \int_{\Omega} b(x) |u_0|^p. \quad (3.26)$$

Since $b(x) \geq b^\infty(\alpha_0^\infty(\Omega)/\min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\})^{(p-2)/2}$ and $b(x) \rightarrow b^\infty$ as $|x| \rightarrow \infty$, we apply (3.26) to obtain

$$s_0 < \left(\frac{\min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}}{\alpha_0^\infty(\Omega)} \right)^{1/2}. \quad (3.27)$$

Thus,

$$\begin{aligned}
 \alpha_0^b(\Omega) &\leq J_0^b(s_0 u_0) = \left(\frac{1}{2} - \frac{1}{p}\right) s_0^2 \int_{\Omega} (|\nabla u_0|^2 + u_0^2) \\
 &< \frac{\min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}}{\alpha_0^\infty(\Omega)} \left(\frac{1}{2} - \frac{1}{p}\right) \int_{\Omega} (|\nabla u_0|^2 + u_0^2) \\
 &= \min\{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}.
 \end{aligned} \tag{3.28}$$

By Proposition 3.2, (1.5) has a ground-state solution.

Now, we begin to show the proof of Theorem 1.3: since the domain Ω satisfies the condition $(\Omega 3')$, we have (1.11) in Ω , and there exists a ground-state solution u_0 such that $J_0^\infty(u_0) = \alpha_0^\infty(\Omega)$. Let $s_0 > 0$ with $s_0 u_0 \in \mathbf{M}_0^b(\Omega)$. Then

$$s_0^2 \int_{\Omega} (|\nabla u_0|^2 + u_0^2) = s_0^p \int_{\Omega} b(x) |u_0|^p. \tag{3.29}$$

Since $b(x) \geq b^\infty(\alpha_0^\infty(\Omega)/\lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)))^{(p-2)/2}$ and $b(x) \rightarrow b^\infty$ as $|x| \rightarrow \infty$, we apply (3.29) to obtain

$$s_0 < \left(\frac{\lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r))}{\alpha_0^\infty(\Omega)} \right)^{1/2}. \tag{3.30}$$

Thus,

$$\begin{aligned}
 \alpha_0^b(\Omega) &\leq J_0^b(s_0 u_0) = \left(\frac{1}{2} - \frac{1}{p}\right) s_0^2 \int_{\Omega} (|\nabla u_0|^2 + u_0^2) \\
 &< \frac{\lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r))}{\alpha_0^\infty(\Omega)} \left(\frac{1}{2} - \frac{1}{p}\right) \int_{\Omega} (|\nabla u_0|^2 + u_0^2) \\
 &= \lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)).
 \end{aligned} \tag{3.31}$$

By Proposition 3.3, (1.5) has a ground-state solution.

4. Nonhomogeneous problems

4.1. Existence of a local minimum. First, we establish the existence of a local minimum. Similar as the proof of Lemma 1.4 by Adachi and Tanaka in [21], we have the following lemma.

LEMMA 4.1. If $\|h\|_{H^{-1}} < (p-2)(1/(p-1))^{(p-1)/(p-2)} [b_{\sup} S^p(\Omega)]^{1/(2-p)}$, then

- (i) $\mathbf{M}_h^{b+}(\Omega) \subset B(0; r_0)$;
- (ii) $J_h^b(u)$ is strictly convex in $B(0; r_0)$,

where $B(0; r_0) = \{u \in H^1(\Omega) \mid \|u\|_{H^1} < r_0\}$ and $r_0 = [(p-1)b_{\sup} S^p(\Omega)]^{1/(2-p)}$.

Furthermore, we have the following theorem.

THEOREM 4.2. If r_0 is as in Lemma 4.1, then the functional J_h^b has a unique critical point $u_{\min}$ in $B(0; r_0)$ and it satisfies

- (i) $u_{\min} \in \mathbf{M}_h^{b+}(\Omega)$ and $J_h^b(u_{\min}) = \alpha_h^{b+}(\Omega) = \alpha_h^b(\Omega)$;
- (ii) $u_{\min}$ is a positive solution of (1.1).

Proof. Similar as the proof of Theorem 2.1 by Cao and Zhu in [18], there is a $u_{\min} \in \mathbf{M}_h^{b+}(\Omega)$ which is a critical point for J_h^b such that $J_h^b(u_{\min}) = \alpha_h^{b+} = \alpha_h^b$, since $\mathbf{M}_h^{b+}(\Omega) \subset B(0; r_0)$ and $J_h^b(u)$ is strictly convex in $B(0; r_0)$, so that $u_{\min}$ is a unique critical point of J_h^b in $B(0; r_0)$. Since $u_{\min}$ is a unique critical point of J_h^b in $B(0; r_0)$, we have that $u_{\min}$ is a nonnegative solution of (1.1). By the maximum principle, $u_{\min}$ is positive. $\square$

4.2. Multiple positive solutions. Throughout this section, we let $u_{\min}$ be the local minimum for J_h^b in $H_0^1(\Omega)$ in Theorem 4.2 and

$$\|h\|_{H^{-1}} < (p-2) \left(\frac{1}{p-1} \right)^{(p-1)/(p-2)} [b_{\sup} S^p(\Omega)]^{1/(2-p)}. \quad (4.1)$$

Then we have the following restricted (PS) conditions.

PROPOSITION 4.3. Suppose that the domain Ω satisfies the conditions $(\Omega 1')$ – $(\Omega 2')$. If $\{u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b with

$$\beta < \alpha_h^b(\Omega) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}, \quad (4.2)$$

then there exist a subsequence $\{u_n\}$ and u in $H_0^1(\Omega)$ such that $u_n \rightarrow u$ strongly in $H_0^1(\Omega)$ and $J_h^b(u) = \beta$.

Proof. Let $\{u_n\}$ be a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b . By Lemma 2.8(ii), $\{u_n\}$ is bounded. Then there exist a subsequence $\{u_n\}$ and a nonzero solution u of (1.1) such that $u_n \rightharpoonup u$ weakly in $H_0^1(\Omega)$. Suppose that $u_n \not\rightarrow u$ strongly in $H_0^1(\Omega)$. Let $w_n = u_n - u$ for $n = 1, 2, \dots$. Then, by Lemma 2.4, $\{w_n\}$ is a $(PS)_{\beta - J_h^b(u)}$ -sequence in $H_0^1(\Omega)$ for J_0^∞ , since $w_n \rightharpoonup 0$ and $w_n \not\rightarrow 0$ strongly in $H_0^1(\Omega)$. Similar as the proof of Proposition 3.2,

$$\beta - J_h^b(u) \geq \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}, \quad (4.3)$$

which is a contradiction. Thus $u_n \rightarrow u$ strongly in $H_0^1(\Omega)$. $\square$

PROPOSITION 4.4. Suppose that the domain Ω satisfies the condition $(\Omega 3')$. If $\{u_n\}$ is a $(PS)_\beta$ -sequence in $H_0^1(\Omega)$ for J_h^b with

$$\beta < \alpha_h(\Omega) + \lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)), \quad (4.4)$$

then there exist a subsequence $\{u_n\}$ and u in $H_0^1(\Omega)$ such that $u_n \rightarrow u$ strongly in $H_0^1(\Omega)$ and $J_h^b(u) = \beta$.

The proof is similar to the proof of Proposition 4.3.

LEMMA 4.5. Suppose that the domain Ω satisfies the conditions $(\Omega 1')$ – $(\Omega 2')$ and the coefficient $b(x)$ satisfies the condition $(b1)$. Let $\bar{u}$ be a positive solution of (1.11) in Ω such that $J_0^\infty(\bar{u}) = \alpha_0^\infty(\Omega)$ and let $u_{\min}$ be a local minimum in Theorem 4.2. Then

$$\sup_{t \geq 0} J_h^b(u_{\min} + t\bar{u}) < J_h^b(u_{\min}) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}. \quad (4.5)$$

Proof. Since $u_{\min}$ is a positive solution of (1.1). Let $f(s) = s^{p-1}$ for $s \geq 0$ and $F_b(u) = \int_\Omega b(x) \int_0^u f(s) ds dx = (1/p) \int_\Omega b(x) u^p$, then

$$\begin{aligned} J_h^b(u_{\min} + t\bar{u}) &= J_h^b(u_{\min}) + J_0^b(t\bar{u}) + t \left(\int_\Omega b(x) u_0^{p-1} \bar{u} + h(x) \bar{u} \right) - \int_\Omega h(x) t\bar{u} \\ &\quad + \frac{1}{p} \left[\int_\Omega b(x) u_0^p + \int_\Omega b(x) |t\bar{u}|^p - \int_\Omega b(x) |u_0 + t\bar{u}|^p \right] \\ &= J_h^b(u_{\min}) + J_0^b(t\bar{u}) - \int_\Omega b(x) \left\{ \int_0^{t\bar{u}} [f(u_0 + s) - f(s) - f(u_0)] ds \right\}. \end{aligned} \quad (4.6)$$

For $v > 0$ and $w > 0$, we have

$$\begin{aligned} f(v+w) &= (v+w)^{p-1} \\ &= (v+w)^{p-2}v + (v+w)^{p-2}w \\ &> v^{p-1} + w^{p-1} = f(v) + f(w). \end{aligned} \quad (4.7)$$

Thus, $J_h^b(u_{\min} + t\bar{u}) \leq J_h^b(u_{\min}) + J_0^b(t\bar{u})$. Since $J_0^b(t\bar{u}) \rightarrow -\infty$ as $t \rightarrow \infty$, there is a $t_0 > 0$ such that $J_h^b(u_{\min} + t\bar{u}) < J_h^b(u_0)$ for $t \geq t_0$. Hence,

$$\sup_{t \geq 0} J_h^b(u_{\min} + t\bar{u}) = \sup_{0 \leq t \leq t_0} J_h^b(u_{\min} + t\bar{u}). \quad (4.8)$$

Let $g_1(t) = J_h^b(u_{\min} + t\bar{u})$ for $t \geq 0$. By the continuity of $g_1(t)$, given

$$\varepsilon = \frac{1}{2} \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\} > 0, \quad (4.9)$$

there exists $t_1 \in (0, t_0)$ such that

$$g_1(t) < g_1(0) + \frac{1}{2} \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\} \quad \text{for } t \in [0, t_1]. \quad (4.10)$$

Then

$$\begin{aligned} \sup_{0 \leq t \leq t_1} J_h^b(u_{\min} + t\bar{u}) &\leq J_h^b(u_{\min}) + \frac{1}{2} \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\} \\ &< J_h^b(u_{\min}) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}. \end{aligned} \quad (4.11)$$

Now, we only need to show that

$$\sup_{t_1 \leq t \leq t_0} J_h^b(u_{\min} + t\bar{u}) < J_h^b(u_{\min}) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}. \quad (4.12)$$

Let $g_2(t) = J_0^b(t\bar{u})$ for $t \geq 0$. Then

$$\begin{aligned} g_2'(t) &= t \int_{\Omega} (|\nabla \bar{u}|^2 + \bar{u}^2) - t^{p-1} \int_{\Omega} b(x) \bar{u}^p, \\ g_2''(t) &= \int_{\Omega} (|\nabla \bar{u}|^2 + \bar{u}^2) - (p-1)t^{p-2} \int_{\Omega} b(x) \bar{u}^p. \end{aligned} \quad (4.13)$$

There is a unique $\bar{t} = [\int_{\Omega} (|\nabla \bar{u}|^2 + \bar{u}^2) / \int_{\Omega} b(x) \bar{u}^p]^{1/(p-2)}$ such that $g_2'(\bar{t}) = 0$ and $g_2''(\bar{t}) < 0$. Thus, $g_2(t)$ has an absolutely maximum at $\bar{t}$. Since

$$b(x) \geq b^\infty \left(\frac{\alpha_0^\infty(\Omega)}{\min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}} \right)^{(p-2)/2}, \quad (4.14)$$

we have

$$\bar{t} \leq \left(\frac{\min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}}{\alpha_0^\infty(\Omega)} \right)^{1/2}. \quad (4.15)$$

Therefore,

$$\begin{aligned} \sup_{t \geq 0} J_0^b(t\bar{u}) &= J_0^b(\bar{t}\bar{u}) = \left(\frac{1}{2} - \frac{1}{p} \right) \bar{t}^2 \int_{\Omega} (|\nabla \bar{u}|^2 + \bar{u}^2) \\ &\leq \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}. \end{aligned} \quad (4.16)$$

By (4.6), (4.7), and (4.16), we obtain

$$\begin{aligned} &\sup_{t_1 \leq t \leq t_0} J_h^b(u_{\min} + t\bar{u}) \\ &\leq J_h^b(u_{\min}) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\} \\ &\quad - \inf_{t_1 \leq t \leq t_0} \int_{\Omega} b(x) \left\{ \int_0^{t\bar{u}} [f(u_{\min} + s) - f(s) - f(u_{\min})] ds \right\} \\ &< J_h^b(u_{\min}) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}. \end{aligned} \quad (4.17)$$

Thus, $\sup_{t \geq 0} J_h^b(u_{\min} + t\bar{u}) < J_h^b(u_{\min}) + \min \{\alpha_0^\infty(\Omega_1), \alpha_0^\infty(\Omega_2), \dots, \alpha_0^\infty(\Omega_m)\}$. $\square$

LEMMA 4.6. Suppose that the domain Ω satisfies the condition $(\Omega 3')$ and the coefficient $b(x)$ satisfies the condition $(b2)$. Let $\bar{u}$ be a positive solution of (1.11) in Ω such that $J_0^\infty(\bar{u}) = \alpha_0^\infty(\Omega)$ and let $u_{\min}$ be the local minimum in Theorem 4.2. Then

$$\sup_{t \geq 0} J_h^b(u_{\min} + t\bar{u}) < J_h^b(u_{\min}) + \lim_{r \rightarrow \infty} \alpha_0^\infty(\Omega^c(r)). \quad (4.18)$$

The proof is similar to the proof of Lemma 4.5.

Now, we begin to show the proof of Theorem 1.6: for $u \in H_0^1(\Omega)$ with $\|u\|_{H^1} = 1$, by Lemma 2.7 there is a unique $t^-(u) > 0$ such that $t^-(u)$, $u \in \mathbf{M}_h^{b-}(\Omega)$ and

$$J_h^b(t^-(u)u) = \max_{t \geq t_{\max}} J_h^b(tu). \quad (4.19)$$

By Lemma 2.7(ii) and (iii), we have that $t^-(u)$ is a continuous function for nonzero u and

$$\mathbf{M}_h^{b-}(\Omega) = \left\{ u \in H_0^1(\Omega) \setminus \{0\} \mid \frac{1}{\|u\|_{H^1}} t^- \left(\frac{u}{\|u\|_{H^1}} \right) = 1 \right\}. \quad (4.20)$$

Let

$$\begin{aligned} A_1 &= \left\{ u \in H_0^1(\Omega) \setminus \{0\} \mid \frac{1}{\|u\|_{H^1}} t^- \left(\frac{u}{\|u\|_{H^1}} \right) > 1 \right\} \cup \{0\}, \\ A_2 &= \left\{ u \in H_0^1(\Omega) \setminus \{0\} \mid \frac{1}{\|u\|_{H^1}} t^- \left(\frac{u}{\|u\|_{H^1}} \right) < 1 \right\}. \end{aligned} \quad (4.21)$$

Then $\mathbf{M}_h^{b-}(\Omega)$ disconnects $H_0^1(\Omega)$ in two connected components A_1 and A_2 and $H_0^1(\Omega) \setminus \mathbf{M}_h^{b-}(\Omega) = A_1 \cup A_2$. For each $u \in \mathbf{M}_h^{b+}(\Omega)$, we have

$$1 < t_{\max}(u) < t^-(u). \quad (4.22)$$

Since $t^-(u) = (1/\|u\|_{H^1})t^-(u/\|u\|_{H^1})$, then $\mathbf{M}_h^{b+}(\Omega) \subset A_1$. In particular, $u_{\min} \in A_1$. We claim that there exists $t_0 > 0$ such that $u_{\min} + t_0\bar{u} \in A_2$. First, we find a constant $c > 0$ such that $0 < t^-((u_{\min} + t\bar{u})/\|u_{\min} + t\bar{u}\|_{H^1}) < c$ for each $t \geq 0$. Otherwise, there exists a sequence $\{t_n\}$ such that $t_n \rightarrow \infty$ and $t^-((u_{\min} + t_n\bar{u})/\|u_{\min} + t_n\bar{u}\|_{H^1}) \rightarrow \infty$ as $n \rightarrow \infty$. Let $v_n = (u_{\min} + t_n\bar{u})/\|u_{\min} + t_n\bar{u}\|_{H^1}$. Since $t^-(v_n)$, $v_n \in \mathbf{M}_h^{b-}(\Omega) \subset \mathbf{M}_h^b(\Omega)$, and by the Lebesgue dominated convergence theorem,

$$\begin{aligned} \int_{\Omega} b(x)v_n^p &= \frac{1}{\|u_{\min} + t_n\bar{u}\|_{H^1}^p} \int_{\Omega} b(x)(u_{\min} + t_n\bar{u})^p \\ &= \frac{1}{\|u_{\min}/t_n + \bar{u}\|_{H^1}^p} \int_{\Omega} b(x) \left(\frac{u_{\min}}{t_n} + \bar{u} \right)^p \rightarrow \frac{\int_{\Omega} b(x)\bar{u}^p}{\|\bar{u}\|_{H^1}^p} \quad \text{as } n \rightarrow \infty. \end{aligned} \quad (4.23)$$

We have

$$\begin{aligned} J_h^b(t^-(v_n)v_n) &= \frac{1}{2} [t^-(v_n)]^2 - \frac{1}{p} [t^-(v_n)]^p \int_{\Omega} b(x)v_n^p \\ &\quad - t^-(v_n) \int_{\Omega} h v_n \rightarrow -\infty \quad \text{as } n \rightarrow \infty. \end{aligned} \quad (4.24)$$

But J_h^b is bounded below on $\mathbf{M}_h^b(\Omega)$, a contradiction. Let

$$t_0 = \frac{|c^2 - \|u_{\min}\|_{H^1}^2|^{1/2}}{\|\bar{u}\|_{H^1}} + 1. \quad (4.25)$$

Then

$$\begin{aligned}
\|u_{\min} + t_0 \bar{u}\|_{H^1}^2 &= \|u_{\min}\|_{H^1}^2 + t_0^2 \|\bar{u}\|_{H^1}^2 + 2t_0 \langle u_{\min}, \bar{u} \rangle \\
&> \|u_{\min}\|_{H^1}^2 + |c^2 - \|u_{\min}\|_{H^1}^2| + 2 \int_{\Omega} b^{\infty} \bar{u}^{p-1} u_{\min} \\
&> c^2 > \left[t^- \left(\frac{u_{\min} + t_0 \bar{u}}{\|u_{\min} + t_0 \bar{u}\|_{H^1}} \right) \right]^2,
\end{aligned} \tag{4.26}$$

that is, $u_{\min} + t_0 \bar{u} \in A_2$. Define a path $\gamma(s) = u_{\min} + st_0 \bar{u}$ for $s \in [0, 1]$, then

$$\gamma(0) = u_{\min} \in A_1, \quad \gamma(1) = u_{\min} + t_0 \bar{u} \in A_2, \tag{4.27}$$

and there exists $s_0 \in (0, 1)$ such that $u_{\min} + s_0 t_0 \bar{u} \in \mathbf{M}_h^{b-}(\Omega)$. Thus, by Lemma 4.5,

$$\begin{aligned}
\alpha_h^-(\Omega) &\leq J_h^b(u_{\min} + s_0 t_0 \bar{u}) \leq \max_{s \in [0, 1]} J_h^b(\gamma(s)) \\
&< J_h^b(u_{\min}) + \min \{ \alpha_0^{\infty}(\Omega_1), \alpha_0^{\infty}(\Omega_2), \dots, \alpha_0^{\infty}(\Omega_m) \}.
\end{aligned} \tag{4.28}$$

By the Ekeland variational principle [25], there exists a sequence $\{u_n\}$ in $\mathbf{M}_h^{b-}(\Omega)$ such that

$$\begin{aligned}
J_h^b(u_n) &= \alpha_h^{b-}(\Omega) + o(1), \\
(J_h^b)'(u_n) &= o(1) \quad \text{strongly in } H^{-1}(\Omega).
\end{aligned} \tag{4.29}$$

Then by Proposition 4.3, there exist a subsequence $\{u_n\}$ and $u^0 \in \mathbf{M}_h^{b-}(\Omega)$ such that $u_n \rightarrow u^0$ strongly in $H_0^1(\Omega)$, u^0 is a solution of (1.1), and $J_h^b(u^0) = \alpha_h^{b-}(\Omega)$. By the Sobolev imbedding theorem, we have $u_n \rightarrow u^0$ strongly in $L^p(\Omega)$. Thus,

$$\|u^0\|_{H^1}^2 - (p-1) \int_{\Omega} b(x) |u^0|^p \leq 0. \tag{4.30}$$

Then $u^0 \in \mathbf{M}_h^{b-}(\Omega)$ and

$$J_h^b(u^0) = \alpha_h^{b-}(\Omega). \tag{4.31}$$

This implies that $u_{\min}$ and u^0 are distinct. Finally, since $h \geq 0$, by Lemma 2.7 there exists $t^-(|u^0|) > 0$ such that

$$\begin{aligned}
t^-(|u^0|) |u^0| &\in \mathbf{M}_h^{b-}(\Omega), \quad t^-(|u^0|) > t_{\max}(|u^0|) = t_{\max}(u^0), \\
\alpha_h^{b-}(\Omega) &\leq J_h^b(t^-(|u^0|) |u^0|) \leq J_h^b(t^-(|u^0|) u^0) \\
&\leq \max_{t \geq t_{\max}(u^0)} J_h^b(t u^0) = J_h^b(u^0) = \alpha_h^{b-}(\Omega).
\end{aligned} \tag{4.32}$$

Thus,

$$J_h^b(t^-(|u^0|) |u^0|) = J_h^b(t^-(|u^0|) u^0) = \alpha_h^{b-}(\Omega). \tag{4.33}$$

We concluded that $\int_{\Omega} hu^0 = \int_{\Omega} h|u^0|$. Let

$$u_+^0 = \max \{u^0, 0\}, \quad u_-^0 = \max \{-u^0, 0\}, \quad (4.34)$$

then $\int_{\Omega} hu_-^0 = 0$. Since $h \geq 0$ and $u_-^0 \geq 0$, we have $u_-^0 = 0$. Hence, u^0 is nonnegative. By the maximum principle, u^0 is positive. We complete the proof of Theorem 1.6.

Remark 4.7. The proof of Theorem 1.7 similar to Theorem 1.6.

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