Research Article

Traveling Wave Solutions in a Reaction-Diffusion Epidemic Model

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Received 4 February 2013; Revised 17 March 2013; Accepted 17 March 2013

Academic Editor: Anke Meyer-Baese

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We investigate the traveling wave solutions in a reaction-diffusion epidemic model. The existence of the wave solutions is derived through monotone iteration of a pair of classical upper and lower solutions. The traveling wave solutions are shown to be unique and strictly monotonic. Furthermore, we determine the critical minimal wave speed.

1. Introduction

Recently, great attention has been paid to the study of the traveling wave solutions in reaction-diffusion models [1–17]. In the sense of epidemiology, the traveling wave solutions describe the transition from a disease-free equilibrium to an endemic equilibrium; the existence and nonexistence of non-trivial traveling wave solutions indicate whether or not the disease can spread [11]. The results contribute to predicting the developing tendency of infectious diseases, to determining the key factors of the spread of infectious disease, and to seeking the optimum strategies of preventing and controlling the spread of the infectious diseases [18–21].

Some methods have been used to derive the existence of traveling wave solutions in reaction-diffusion models, and the monotone iteration method has been proved to be an effective one. Such a method reduces the existence of traveling wave solutions to that of an ordered pair of upper-lower solutions [6, 7, 9, 10, 14, 15].

In [22], Berezhovsky and coworkers introduced a simple epidemic model through the incorporation of variable population, disease-induced mortality, and emigration into the classical model of Kermack and McKendrick [23]. The total population (N) is divided into two groups of susceptible (S)

and infectious (I); that is to say, $N = S + I$. The model describing the relations between the state variables is

$$\begin{aligned}\frac{dS}{dt} &= rN \left(1 - \frac{N}{K}\right) - \beta \frac{SI}{N} - (\mu + m)S, \\ \frac{dI}{dt} &= \beta \frac{SI}{N} - (\mu + d)I,\end{aligned}\tag{1}$$

where the reproduction of susceptible follows a logistic equation with the intrinsic growth rate r and the carrying capacity K , β denotes the contact transmission rate (the infection rate constant), μ is the natural mortality; d denotes the disease-induced mortality, and m is the per-capita emigration rate of uninfected.

For model (1), the epidemic threshold, the so-called basic reproduction number R_0 , is then computed as $R_0 = \beta/(\mu + d)$. The disease will successfully invade when $R_0 > 1$ but will die out if $R_0 < 1$. $R_0 = 1$ is usually a threshold whether the disease goes to extinction or goes to an endemic. Large values of R_0 may indicate the possibility of a major epidemic [19]. In addition, the basic demographic reproductive number R_d is given by $R_d = r/(\mu + m)$. It can be shown that if $R_d > 1$ the population grows, while $R_d \leq 1$ implies that the population does not survive [22].

For simplicity, rescaling model (1) by letting $S \rightarrow S/K$, $I \rightarrow I/K$, and $t \rightarrow t/(\mu + d)$ leads to the following model:

$$\begin{aligned} \frac{dS}{dt} &= \nu R_d (S + I) [1 - (S + I)] - R_0 \frac{SI}{S + I} - \nu S, \\ \frac{dI}{dt} &= R_0 \frac{SI}{S + I} - I, \end{aligned} \quad (2)$$

where $\nu = (\mu + m)/(\mu + d)$ is defined by the ratio of the average life span of susceptibles to that of infectious.

For details, we refer the reader to [20, 22].

In this paper, we are interested in the existence of traveling wave solutions in the following reaction-diffusion epidemic model [20]:

$$\begin{aligned} \frac{\partial S}{\partial t} &= \nu R_d (S + I) [1 - (S + I)] - R_0 \frac{SI}{S + I} - \nu S + d \frac{\partial^2 S}{\partial x^2}, \\ \frac{\partial I}{\partial t} &= R_0 \frac{SI}{S + I} - I + d \frac{\partial^2 I}{\partial x^2}, \\ S(x, 0) &= S_0(x), \quad I(x, 0) = I_0(x), \end{aligned} \quad (3)$$

where ν, R_0, R_d are all positive constants, d is the diffusion coefficient, and $(x, t) \in \mathbb{R} \times \mathbb{R}^+$.

We are looking for the traveling wave solutions of model (3) with the following form:

$$S(x, t) = S(\xi), \quad I(x, t) = I(\xi), \quad \xi = x + ct, \quad (4)$$

satisfying the following boundary value conditions:

$$(S(-\infty), I(-\infty))^T = E_1, \quad (S(+\infty), I(+\infty))^T = E_2, \quad (5)$$

where E_1, E_2 are the equilibrium points of model (3).

This paper is arranged as follows. In Section 2, we construct a pair of ordered upper-lower solutions of model (3) and establish the uniqueness and strict monotonicity of the traveling wave solutions.

2. Existence and Asymptotic Decay Rates

In this section, we will establish the existence of traveling wave solutions of model (3) by constructing a pair of ordered upper-lower solutions. The definition of the upper solution and the lower solution is standard. We assume that the inequality between two vectors throughout this paper is componentwise.

Setting

$$\widehat{S} = \frac{R_d - 1}{R_d} - S, \quad \widehat{I} = I, \quad (6)$$

then model (3) can be written as

$$\begin{aligned} \frac{\partial \widehat{S}}{\partial t} &= -\nu R_d \left(\frac{R_d - 1}{R_d} - \widehat{S} + \widehat{I} \right) \left[1 - \left(\frac{R_d - 1}{R_d} - \widehat{S} + \widehat{I} \right) \right] \\ &\quad + R_0 \frac{((R_d - 1)/R_d - \widehat{S}) \widehat{I}}{(R_d - 1)/R_d - \widehat{S} + \widehat{I}} + \nu \left(\frac{R_d - 1}{R_d} - \widehat{S} \right) + d \frac{\partial^2 \widehat{S}}{\partial x^2}, \\ \frac{\partial \widehat{I}}{\partial t} &= R_0 \frac{((R_d - 1)/R_d - \widehat{S}) \widehat{I}}{(R_d - 1)/R_d - \widehat{S} + \widehat{I}} - \widehat{I} + d \frac{\partial^2 \widehat{I}}{\partial x^2}, \\ (\widehat{S}, \widehat{I})^T(-\infty) &= (0, 0)^T, \quad (\widehat{S}, \widehat{I})^T(+\infty) = (\widehat{S}^*, \widehat{I}^*)^T. \end{aligned} \quad (7)$$

For model (3), the equilibria are $E_1 = ((R_d - 1)/R_d, 0)$ and $E_2 = (S^*, I^*)$, where

$$S^* = \frac{\nu R_0 R_d - R_0 - \nu + 1}{\nu R_0^2 R_d}, \quad I^* = (R_0 - 1) S^*, \quad (8)$$

and for model (7), the equilibria are $\widehat{E}_1 = (0, 0)$ and $\widehat{E}_2 = (\widehat{S}^*, \widehat{I}^*)$, where

$$\begin{aligned} \widehat{S}^* &= \frac{(R_0 - 1)(\nu R_0 R_d - \nu R_0 - \nu + 1)}{\nu R_0^2 R_d}, \\ \widehat{I}^* &= \frac{(R_0 - 1)(\nu R_0 R_d - R_0 - \nu + 1)}{\nu R_0^2 R_d}. \end{aligned} \quad (9)$$

Obviously,

$$\begin{aligned} \widehat{I}^* - \widehat{S}^* &= \frac{(\nu - 1)(R_0 - 1)}{\nu R_0 R_d}, \\ \widehat{S}^* &= \frac{R_d - 1}{R_d} - S^*, \quad \widehat{I}^* = I^*. \end{aligned} \quad (10)$$

For simplicity, we define the following functions and constants:

$$\begin{aligned} \alpha_0 &= \frac{R_d - 1}{R_d}, \quad \phi(I) = \alpha_0 \widehat{I}^* + (\widehat{I}^* - \widehat{S}^*) I; \\ \beta_0 &= \alpha_0 (\widehat{I}^*)^2 \frac{R_0 - 1}{R_0} (R_0 - \nu + 1); \\ \gamma_0 &= 2\nu (R_d - 1) (\widehat{I}^* - \widehat{S}^*) - [\nu \widehat{I}^* + (R_0 - 1) \widehat{S}^*]; \\ \psi(I) &= \nu R_d (\widehat{I}^* - \widehat{S}^*)^2 I^2 + \gamma_0 \widehat{I}^* I + \beta_0; \\ \eta_0 &= -\frac{\gamma_0 \widehat{I}^*}{2\nu R_d (\widehat{I}^* - \widehat{S}^*)^2}; \\ \varphi(I) &= 1, \quad I > 0; \quad \varphi(I) = -1, \quad I \leq 0; \\ B &= \frac{\widehat{I}^*}{2} \left[1 + \varphi \left(\frac{\widehat{I}^*}{2} - \eta_0 \right) \right]. \end{aligned} \quad (11)$$

And we will always assume the following hypotheses throughout the rest of this paper:

[H1]

$$R_0 > 1, \quad 1 < R_d < \frac{2R_0^2 + 2R_0 - 2}{3R_0^2 - 2R_0},$$

$$\max \left\{ \frac{27R_0(R_d - 1)^2}{R_d^3}, \frac{R_0 - 1}{R_0 R_d - 1} \right\} < \nu < \frac{-1}{R_0 R_d - R_0 - 1}. \quad (12)$$

[H2]

$$\nu \geq \max \left\{ \frac{R_0}{2 - R_d}, \frac{R_0^3 - 2R_0^2 + 4R_0 - 2}{2R_0^2 + 2R_0 - 2 - (3R_0^2 - 2R_0)R_d} \right\}, \quad (13)$$

$$\psi(B) \leq 0.$$

Then we can obtain the following.

Lemma 1. If [H1] holds, then E_2 and $\widehat{E}_2$ are endemic points of model (3) and model (7), respectively.

Lemma 2. For model (7), if [H1] holds, then $\widehat{E}_1$ is unstable, and $\widehat{E}_2$ is stable.

For the sake of convenience, let $x = \sqrt{d} \tilde{x}$. For simplicity, we still use the variables S , I , and x instead of $\widehat{S}$, $\widehat{I}$, and $\tilde{x}$, respectively, then model (7) could be rewritten as

$$\begin{aligned} \frac{\partial S}{\partial t} &= -\nu R_d (\alpha_0 - S + I) [1 - (\alpha_0 - S + I)] \\ &\quad + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} + \nu (\alpha_0 - S) + \frac{\partial^2 S}{\partial x^2}, \\ \frac{\partial I}{\partial t} &= R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} - I + \frac{\partial^2 I}{\partial x^2}, \end{aligned} \quad (14)$$

$$(S, I)^T(-\infty) = (0, 0)^T, \quad (S, I)^T(+\infty) = (\widehat{S}^*, \widehat{I}^*)^T.$$

Following the definition of quasi-monotonicity [17], we can obtain the following results.

Lemma 3. Model (14) is a quasi-monotone decreasing system in $(S, I) \in [\widehat{E}_1, \widehat{E}_2]$.

Proof. Let

$$\begin{aligned} F_1(S, I) &= -\nu R_d (\alpha_0 - S + I) [1 - (\alpha_0 - S + I)] \\ &\quad + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} + \nu (\alpha_0 - S), \\ F_2(S, I) &= R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} - I. \end{aligned} \quad (15)$$

From [17], we can know that the functions $F_1(S, I)$ and $F_2(S, I)$ are said to possess a quasi-monotone nonincreasing

system, if the sign of $\partial F_1(S, I)/\partial I$ and $\partial F_2(S, I)/\partial S$ are both nonpositive.

Since

$$\begin{aligned} \frac{\partial F_2(S, I)}{\partial S} &= -R_0 \left(\frac{I}{\alpha_0 - S + I} \right)^2 \leq 0, \\ \frac{\partial F_1(S, I)}{\partial I} &= R_0 \left(\frac{\alpha_0 - S}{\alpha_0 - S + I} \right)^2 + 2\nu R_d (\alpha_0 - S + I) - \nu R_d, \\ \frac{\partial}{\partial S} \left(\frac{\partial F_1(S, I)}{\partial I} \right) &= -2\nu R_d - 2R_0 \frac{(\alpha_0 - S)I}{(\alpha_0 - S + I)^3} \leq -2\nu R_d < 0. \end{aligned} \quad (16)$$

Then,

$$\frac{\partial F_1(S, I)}{\partial I} \leq R_0 \left(\frac{\alpha_0}{\alpha_0 + I} \right)^2 + 2\nu R_d (\alpha_0 + I) - \nu R_d. \quad (17)$$

Let

$$G(z) = \frac{\alpha_0^2 R_0}{z^2} + 2\nu R_d z - \nu R_d, \quad z \in [\alpha_0, \alpha_0 + \widehat{I}^*], \quad (18)$$

then

$$G'(z) = 2\nu R_d - \frac{2\alpha_0^2 R_0}{z^3} = 0, \quad (19)$$

obviously, $z^* = \sqrt[3]{\alpha_0^2 R_0 / \nu R_d}$ is the unique real root of $G'(z)$.

Since $\nu > 27R_0(R_d - 1)^2/R_d^3$, consider $\alpha_0 = (R_d - 1)/R_d$, then we can get

$$G(z^*) = \frac{(\alpha_0^2 R_0)^{2/3} \left[(27\alpha_0^2 R_0)^{1/3} - (\nu R_d)^{1/3} \right]}{(z^*)^2} < 0. \quad (20)$$

And

$$\lim_{z \rightarrow 0^+} G(z) = \lim_{z \rightarrow +\infty} G(z) = +\infty; \quad (21)$$

hence, $G(z)$ has two positive roots.

Since $\nu \geq R_0/(2 - R_d)$, thus $G(\alpha_0) = R_0 + \nu R_d - 2\nu \leq 0$.

According to conditions [H1] and [H2], we can get

$$\begin{aligned} G(\alpha_0 + \widehat{I}^*) &= \nu R_d - 2\nu + 2\nu R_d \widehat{I}^* + R_0 \left(\frac{\alpha_0}{\alpha_0 + \widehat{I}^*} \right)^2 \\ &< \nu R_d - 2\nu + 2\nu R_d \widehat{I}^* + R_0 \\ &= \frac{(3R_0^2 R_d - 2R_0^2 - 2R_0 R_d - 2R_0 + 2)\nu}{R_0^2} \\ &\quad + \frac{(R_0^3 - 2R_0^2 + 4R_0 - 2)}{R_0^2} \\ &\leq 0. \end{aligned} \quad (22)$$

Then, $G([\alpha_0, \alpha_0 + \widehat{I}^*]) \leq 0$. Hence, $\partial F_1(S, I)/\partial I \leq 0$.

That is to say, model (14) is a quasi-monotone system in $(S, I) \in [\widehat{E}_1, \widehat{E}_2]$. $\square$

Since the traveling wave solution of model (14) has the following form

$$S(\xi) = S(x + ct), \quad I(\xi) = I(x + ct), \quad \xi = x + ct, \quad c > 0; \quad (23)$$

substituting (23) into model (14), we can get the following model:

$$\begin{aligned} S'' - cS' - \nu R_d(\alpha_0 - S + I)[1 - (\alpha_0 - S + I)] \\ + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} + \nu(\alpha_0 - S) &= 0, \\ I'' - cI' + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} - I &= 0, \end{aligned} \quad (24)$$

$$(S, I)^T(-\infty) = (0, 0)^T, \quad (S, I)^T(+\infty) = (\hat{S}^*, \hat{I}^*)^T.$$

Obviously, we can know the following.

Remark 4. Model (24) is also a quasi-monotone system in $(S, I) \in [\bar{E}_1, \bar{E}_2]$.

Now we establish the existence of traveling wave solutions of model (24) through monotone iteration of a pair of smooth upper and lower solutions. Following [17], we give the definitions of the upper and lower solutions of model (24) as follows, respectively.

Definition 5. A smooth function $(\bar{S}(\xi), \bar{I}(\xi))^T$ ($\xi \in \mathbb{R}$) is an upper solution of model (24) if its derivatives $(\bar{S}', \bar{I}')^T$ and $(\bar{S}'', \bar{I}'')$ are continuous on $\mathbb{R}$, and $(\bar{S}, \bar{I})^T$ satisfies

$$\begin{aligned} S'' - cS' - \nu R_d(\alpha_0 - S + I)[1 - (\alpha_0 - S + I)] \\ + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} + \nu(\alpha_0 - S) &\leq 0, \\ I'' - cI' + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} - I &\leq 0, \end{aligned} \quad (25)$$

with the following boundary value conditions

$$\begin{pmatrix} S \\ I \end{pmatrix}(-\infty) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad \begin{pmatrix} S \\ I \end{pmatrix}(+\infty) \geq \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix}. \quad (26)$$

Definition 6. A smooth function $(\underline{S}(\xi), \underline{I}(\xi))^T$ ($\xi \in \mathbb{R}$) is a lower solution of model (24) if its derivatives $(\underline{S}', \underline{I}')^T$ and $(\underline{S}'', \underline{I}'')$ are continuous on $\mathbb{R}$, and $(\underline{S}, \underline{I})^T$ satisfies

$$\begin{aligned} S'' - cS' - \nu R_d(\alpha_0 - S + I)[1 - (\alpha_0 - S + I)] \\ + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} + \nu(\alpha_0 - S) &\geq 0, \\ I'' - cI' + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} - I &\geq 0, \end{aligned} \quad (27)$$

with the following boundary value conditions

$$\begin{pmatrix} S \\ I \end{pmatrix}(-\infty) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad \begin{pmatrix} S \\ I \end{pmatrix}(+\infty) \leq \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix}. \quad (28)$$

The construction of the smooth upper-lower solution pair is based on the solution of the following KPP equation:

$$\begin{aligned} w'' - cw' + f(w) &= 0, \\ w(-\infty) &= 0, \quad w(+\infty) = b, \end{aligned} \quad (29)$$

where $f \in C^2([0, b])$ and $f > 0$ in the open interval $(0, b)$ with $f(0) = f(b) = 0$, $f'(0) = a_1 > 0$, and $f'(b) = -b_1 < 0$ [15]. First, let us recall the following result.

Lemma 7 (see [1, 15]). *Corresponding to every $c \geq 2\sqrt{a_1}$, model (29) has a unique (up to a translation of the origin) monotonically increasing traveling wave solution $w(\xi)$ for $\xi \in \mathbb{R}$. The traveling wave solution w has the following asymptotic behaviors.*

(i) *For the wave solution with noncritical speed $c > 2\sqrt{a_1}$, one has*

$$\begin{aligned} w(\xi) &= a_\omega e^{((c - \sqrt{c^2 - 4a_1})/2)\xi} \\ &+ o\left(e^{((c - \sqrt{c^2 - 4a_1})/2)\xi}\right) \quad \text{as } \xi \longrightarrow -\infty, \end{aligned} \quad (30)$$

$$\begin{aligned} w(\xi) &= b - b_\omega e^{((c - \sqrt{c^2 + 4b_1})/2)\xi} \\ &+ o\left(e^{((c - \sqrt{c^2 + 4b_1})/2)\xi}\right) \quad \text{as } \xi \longrightarrow +\infty, \end{aligned}$$

where a_ω and b_ω are positive constants.

(ii) *For the wave with critical speed $c = 2\sqrt{a_1}$, one has*

$$\begin{aligned} w(\xi) &= (a_c + d_c \xi) e^{\sqrt{a_1} \xi} + o\left(\xi e^{\sqrt{a_1} \xi}\right) \quad \text{as } \xi \longrightarrow -\infty, \\ w(\xi) &= b - b_c e^{(\sqrt{a_1} - \sqrt{a_1 + b_1})\xi} \\ &+ o\left(e^{(\sqrt{a_1} - \sqrt{a_1 + b_1})\xi}\right) \quad \text{as } \xi \longrightarrow +\infty, \end{aligned} \quad (31)$$

where the constant d_c is negative, b_c is positive, and $a_c \in \mathbb{R}$.

For constructing the upper solution of the model (24), we start with the following model:

$$\begin{aligned} I'' - cI' + I(\hat{I}^* - I) \frac{\alpha_0(R_0 - 1)}{\alpha_0 \hat{I}^* + (\hat{I}^* - \hat{S}^*)I} &= 0, \\ I(-\infty) &= 0, \quad I(+\infty) = \hat{I}^*. \end{aligned} \quad (32)$$

Define $f(I) = I(\hat{I}^* - I)(\alpha_0(R_0 - 1)/\phi(I))$, $I \in [0, \hat{I}^*]$, one can verify that all of the following conditions are satisfied:

(i) $f(I) = I(\hat{I}^* - I)(\alpha_0(R_0 - 1)/\phi(I)) \in C^2([0, \hat{I}^*])$;

- (ii) $f(I) > 0$, for all $I \in (0, \hat{I}^*)$ and $f(0) = f(\hat{I}^*) = 0$;
- (iii) $f'(0) = R_0 - 1 > 0$, $f'(\hat{I}^*) = -\alpha_0(R_0 - 1)^2/R_0\hat{I}^* < 0$.

From Lemma 7, we know that, for each $c \geq 2\sqrt{R_0 - 1}$, equation (32) has a unique traveling wave solution $\bar{I}(\xi)$ (up to a translation of the origin), satisfying the given boundary value conditions (26).

Define

$$\begin{pmatrix} \bar{S}(\xi) \\ \bar{I}(\xi) \end{pmatrix} = \begin{pmatrix} \frac{\hat{S}^*}{\hat{I}^*} \bar{I}(\xi) \\ \bar{I}(\xi) \end{pmatrix}, \quad \xi \in R, \quad (33)$$

then we can get the following result.

Lemma 8. For each $c \geq 2\sqrt{R_0 - 1}$, (33) is a smooth upper solution of model (24).

Proof. On the boundary,

$$\begin{pmatrix} \bar{S} \\ \bar{I} \end{pmatrix}(-\infty) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad \begin{pmatrix} \bar{S} \\ \bar{I} \end{pmatrix}(+\infty) \geq \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix}. \quad (34)$$

As for the I component, we have

$$\begin{aligned} & \bar{I}'' - c\bar{I}' + R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} - \bar{I} \\ &= -\bar{I}(\hat{I}^* - \bar{I}) \frac{\alpha_0(R_0 - 1)}{\phi(\bar{I})} + R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} - \bar{I} \\ &= -\bar{I}(\hat{I}^* - \bar{I}) \frac{\alpha_0(R_0 - 1)}{\phi(\bar{I})} + R_0 \frac{\alpha_0\hat{I}^* - \hat{S}^*\bar{I}}{\phi(\bar{I})} \bar{I} - \bar{I} \\ &= -\bar{I}(\hat{I}^* - \bar{I}) \frac{\alpha_0(R_0 - 1)}{\phi(\bar{I})} \\ & \quad + \bar{I} \frac{\alpha_0(R_0 - 1)\hat{I}^* - (R_0\hat{S}^* - \hat{S}^* + \hat{I}^*)\bar{I}}{\phi(\bar{I})} \\ &= -\bar{I}(\hat{I}^* - \bar{I}) \frac{\alpha_0(R_0 - 1)}{\phi(\bar{I})} \\ & \quad + \bar{I} \frac{\alpha_0(R_0 - 1)\hat{I}^* - \alpha_0(R_0 - 1)\bar{I}}{\phi(\bar{I})} \\ &= 0. \end{aligned} \quad (35)$$

As for the S component, since $\nu > 1$, then $\hat{I}^* - \hat{S}^* = (\nu - 1)(R_0 - 1)/\nu R_0 R_d > 0$. And

- (i) if $\eta_0 < \hat{I}^*/2$, then $\max_{\xi \in [0, \hat{I}^*]} \psi(\bar{I}) = \psi(\hat{I}^*) = \psi(B)$;
- (ii) if $\eta_0 \geq \hat{I}^*/2$, then $\max_{\xi \in [0, \hat{I}^*]} \psi(\bar{I}) = \psi(0) = \psi(B)$.

Thus we can get:

$$\begin{aligned} & \bar{S}'' - c\bar{S}' - \nu R_d (\alpha_0 - \bar{S} + \bar{I}) [1 - (\alpha_0 - \bar{S} + \bar{I})] \\ & \quad + R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} + \nu(\alpha_0 - \bar{S}) \\ &= \frac{\hat{S}^*}{\hat{I}^*} (\bar{I}'' - c\bar{I}') - \nu R_d (\alpha_0 - \bar{S} + \bar{I}) [1 - (\alpha_0 - \bar{S} + \bar{I})] \\ & \quad + R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} + \nu(\alpha_0 - \bar{S}) \\ &= \frac{\hat{S}^*}{\hat{I}^*} \left(-R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} + \bar{I} \right) - \nu R_d (\alpha_0 - \bar{S} + \bar{I}) \\ & \quad \times [1 - (\alpha_0 - \bar{S} + \bar{I})] + R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} + \nu(\alpha_0 - \bar{S}) \\ &= \left(1 - \frac{\hat{S}^*}{\hat{I}^*} \right) R_0 \frac{(\alpha_0 - \bar{S})\bar{I}}{\alpha_0 - \bar{S} + \bar{I}} + \frac{\hat{S}^*}{\hat{I}^*} \bar{I} - \nu R_d (\alpha_0 - \bar{S} + \bar{I}) \\ & \quad \times [1 - (\alpha_0 - \bar{S} + \bar{I})] + \nu(\alpha_0 - \bar{S}) \\ &= \frac{\hat{I}^* - \hat{S}^*}{\hat{I}^*} R_0 \frac{\alpha_0\hat{I}^* - \hat{S}^*\bar{I}}{\phi(\bar{I})} \bar{I} + \frac{\hat{S}^*}{\hat{I}^*} \bar{I} - \nu R_d \frac{\phi(\bar{I})}{\hat{I}^*} \left[1 - \frac{\phi(\bar{I})}{\hat{I}^*} \right] \\ & \quad + \nu \frac{\alpha_0\hat{I}^* - \hat{S}^*\bar{I}}{\hat{I}^*} \\ &= \frac{(\hat{I}^* - \hat{S}^*)\bar{I}\psi(\bar{I})}{(\hat{I}^*)^2\phi(\bar{I})} \leq 0. \end{aligned} \quad (36)$$

Hence, $(\bar{S}, \bar{I})$ forms a smooth upper solution for model (24). $\square$

For constructing the lower solution of the model (24), we start with the following model:

$$\begin{aligned} & I'' - cI' + I[\hat{I}^* - (1 + \varepsilon)I] \frac{\alpha_0(R_0 - 1)}{\alpha_0\hat{I}^* + (\hat{I}^* - \hat{S}^*)I} = 0, \\ & I(-\infty) = 0, \quad I(+\infty) = \frac{\hat{I}^*}{1 + \varepsilon}. \end{aligned} \quad (37)$$

Define $g(I) = I[\hat{I}^* - (1 + \varepsilon)I](\alpha_0(R_0 - 1)/(\alpha_0\hat{I}^* + (\hat{I}^* - \hat{S}^*)I))$, $I \in [0, \hat{I}^*/(1 + \varepsilon)]$. One can easily verify that all of the following conditions hold:

- (i) $g(I) = I[\hat{I}^* - (1 + \varepsilon)I](\alpha_0(R_0 - 1)/(\alpha_0\hat{I}^* + (\hat{I}^* - \hat{S}^*)I)) \in C^2([0, \hat{I}^*/(1 + \varepsilon)])$;
- (ii) $g(I) > 0$, for all $I \in (0, \hat{I}^*/(1 + \varepsilon))$ and $g(0) = g(\hat{I}^*/(1 + \varepsilon)) = 0$;

$$(iii) \ g'(0) = R_0 - 1 > 0, \ g'(\hat{I}^*/(1+\varepsilon)) = -(1+\varepsilon)\alpha_0(R_0 - 1)/(\varepsilon\alpha_0 + (R_0/(R_0 - 1))\hat{I}^*) < 0.$$

From Lemma 7, we know that, for each fixed $c \geq 2\sqrt{R_0 - 1}$, model (37) has a unique traveling wave solution $\underline{I}(\xi)$ (up to a translation of the origin), satisfying the given boundary value conditions (28).

Define

$$\begin{pmatrix} \underline{S}(\xi) \\ \underline{I}(\xi) \end{pmatrix} = \begin{pmatrix} \frac{\hat{S}^*}{\hat{I}^*} \underline{I}(\xi) \\ \underline{I}(\xi) \end{pmatrix}, \quad \xi \in R, \quad (38)$$

then we have the following result:

Lemma 9. For each fixed $c \geq 2\sqrt{R_0 - 1}$, (38) is a lower solution of model (24).

Proof. On the boundary,

$$\begin{pmatrix} \underline{S} \\ \underline{I} \end{pmatrix}(-\infty) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad \begin{pmatrix} \underline{S} \\ \underline{I} \end{pmatrix}(+\infty) = \begin{pmatrix} \frac{\hat{S}^*}{1+\varepsilon} \\ \frac{\hat{I}^*}{1+\varepsilon} \end{pmatrix} \leq \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix}. \quad (39)$$

As for the I component, we have

$$\begin{aligned} \underline{I}'' - c\underline{I}' + R_0 \frac{(\alpha_0 - \underline{S})\underline{I}}{\alpha_0 - \underline{S} + \underline{I}} - \underline{I} \\ = -\underline{I}[\hat{I}_* - (1+\varepsilon)] \frac{\alpha_0(R_0 - 1)}{\phi(\underline{I})} + R_0 \frac{(\alpha_0 - \underline{S})\underline{I}}{\alpha_0 - \underline{S} + \underline{I}} - \underline{I} \\ = -\underline{I}[\hat{I}_* - (1+\varepsilon)] \frac{\alpha_0(R_0 - 1)}{\phi(\underline{I})} + \underline{I}(\hat{I}^* - \underline{I}) \frac{\alpha_0(R_0 - 1)}{\phi(\underline{I})} \\ = \varepsilon(\underline{I})^2 \frac{\alpha_0(R_0 - 1)}{\phi(\underline{I})} \geq 0. \end{aligned} \quad (40)$$

As for the S component, we have

$$\begin{aligned} \underline{S}'' - c\underline{S}' - \nu R_d(\alpha_0 - \underline{S} + \underline{I})[1 - (\alpha_0 - \underline{S} + \underline{I})] \\ + R_0 \frac{(\alpha_0 - \underline{S})\underline{I}}{\alpha_0 - \underline{S} + \underline{I}} + \nu(\alpha_0 - \underline{S}) \\ = \frac{\hat{S}^*}{\hat{I}^*}(\underline{I}'' - c\underline{I}') - \nu R_d(\alpha_0 - \underline{S} + \underline{I})[1 - (\alpha_0 - \underline{S} + \underline{I})] \\ + R_0 \frac{(\alpha_0 - \underline{S})\underline{I}}{\alpha_0 - \underline{S} + \underline{I}} + \nu(\alpha_0 - \underline{S}) \end{aligned}$$

$$\begin{aligned} &= \frac{\hat{S}^*}{\hat{I}^*} \left\{ \left[\varepsilon(\underline{I})^2 \frac{\alpha_0(R_0 - 1)}{\phi(\underline{I})} \right] - \left[R_0 \frac{(\alpha_0 - \underline{S})\underline{I}}{\alpha_0 - \underline{S} + \underline{I}} - \underline{I} \right] \right\} \\ &\quad - \nu R_d(\alpha_0 - \underline{S} + \underline{I})[1 - (\alpha_0 - \underline{S} + \underline{I})] + R_0 \frac{(\alpha_0 - \underline{S})\underline{I}}{\alpha_0 - \underline{S} + \underline{I}} \\ &\quad + \nu(\alpha_0 - \underline{S}) \\ &= \varepsilon \frac{\hat{S}^*}{\hat{I}^*}(\underline{I})^2 \frac{\alpha_0(R_0 - 1)}{\phi(\underline{I})} + \frac{(\hat{I}^* - \hat{S}^*)\underline{I}\psi(\underline{I})}{(\hat{I}^*)^2 \phi(\underline{I})} \geq 0. \end{aligned} \quad (41)$$

Thus $(\underline{S}, \underline{I})$ forms a smooth lower solution for model (24). $\square$

Next, we show that, by shifting the upper solution far enough to the left, then the upper-lower solution in Lemmas 8 and 9 are ordered.

Lemma 10. Let $c \geq 2\sqrt{R_0 - 1}$, $(\bar{S}, \bar{I})^T$ and $(\underline{S}, \underline{I})^T$ be the upper solution and the lower solution defined in (33) and (38), then there exists a positive number r , such that $(\bar{S}, \bar{I})^T(\xi + r) \geq (\underline{S}, \underline{I})^T(\xi)$ for all $\xi \in R$.

Proof. Our proof is only for $c > 2\sqrt{R_0 - 1}$, and the proof for the case of $c = 2\sqrt{R_0 - 1}$ is similar to it.

First, we derive the asymptotic behaviors of the upper solution and the lower solution at infinities.

According to Lemma 7, when $\xi \rightarrow -\infty$, we can obtain:

$$\begin{aligned} \begin{pmatrix} \bar{S} \\ \bar{I} \end{pmatrix}(\xi) &= \begin{pmatrix} \frac{\hat{S}^*}{\hat{I}^*} A_1 \\ A_1 \end{pmatrix} e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}\right), \\ \begin{pmatrix} \underline{S} \\ \underline{I} \end{pmatrix}(\xi) &= \begin{pmatrix} \frac{\hat{S}^*}{\hat{I}^*} B_1 \\ B_1 \end{pmatrix} e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}\right). \end{aligned} \quad (42)$$

And let $\sigma_0 = (1/2)(c - \sqrt{c^2 + 4(\alpha_0(R_0 - 1)^2/R_0\hat{I}^*)}) < 0$, $\delta_0 = (1/2)(c - \sqrt{c^2 + 4((1+\varepsilon)\alpha_0(R_0 - 1)/(\varepsilon\alpha_0 + (R_0/(R_0 - 1))\hat{I}^*)}) < 0$, when $\xi \rightarrow +\infty$, we can get

$$\begin{aligned} \begin{pmatrix} \bar{S} \\ \bar{I} \end{pmatrix}(\xi) &= \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix} - \begin{pmatrix} \frac{\hat{S}^*}{\hat{I}^*} A_2 \\ A_2 \end{pmatrix} e^{\sigma_0 \xi} + o(e^{\sigma_0 \xi}), \\ \begin{pmatrix} \underline{S} \\ \underline{I} \end{pmatrix}(\xi) &= \frac{1}{1+\varepsilon} \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix} - \begin{pmatrix} \frac{\hat{S}^*}{\hat{I}^*} B_2 \\ B_2 \end{pmatrix} e^{\delta_0 \xi} + o(e^{\delta_0 \xi}), \end{aligned} \quad (43)$$

where, A_1, A_2, B_1, B_2 are all positive constants.

Since for any $\tilde{r} > 0$, $\bar{I}^{\tilde{r}}(\xi) \equiv \bar{I}(\xi + \tilde{r})$ is also a solution of model (32). Thus, $(\bar{S}^{\tilde{r}}, \bar{I}^{\tilde{r}})^T(\xi)$ is an upper solution of model (24). So, according to Lemma 7, when $\xi \rightarrow -\infty$, we can get:

$$\begin{pmatrix} \bar{S}^{\tilde{r}} \\ \bar{I}^{\tilde{r}} \end{pmatrix}(\xi) = \begin{pmatrix} \frac{\bar{S}^*}{\bar{I}^*} A_1 \\ A_1 \end{pmatrix} e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\tilde{r}} e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} + o\left(e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}\right). \quad (44)$$

Since $(c - \sqrt{c^2 - 4(R_0 - 1)})/2 > 0$, we can choose a large enough number $\tilde{r} \gg 0$, such that

$$\begin{pmatrix} \frac{\bar{S}^*}{\bar{I}^*} A_1 \\ A_1 \end{pmatrix} e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\tilde{r}} > \begin{pmatrix} \frac{\bar{S}^*}{\bar{I}^*} B_1 \\ B_1 \end{pmatrix}, \quad (45)$$

hence, there exists a large number $N_1 \gg 1$, such that

$$\begin{pmatrix} \bar{S}^{\tilde{r}}(\xi) \\ \bar{I}^{\tilde{r}}(\xi) \end{pmatrix} > \begin{pmatrix} \underline{S}(\xi) \\ \underline{I}(\xi) \end{pmatrix}, \quad \xi \in (-\infty, -N_1]. \quad (46)$$

By using a similar argument as above, there exists a large enough number $N_2 \gg 1$, such that

$$\begin{pmatrix} \bar{S}^{\tilde{r}}(\xi) \\ \bar{I}^{\tilde{r}}(\xi) \end{pmatrix} > \begin{pmatrix} \underline{S}(\xi) \\ \underline{I}(\xi) \end{pmatrix}, \quad \xi \in [N_2, +\infty). \quad (47)$$

Second, we show that

$$\begin{pmatrix} \bar{S}^{\tilde{r}}(\xi) \\ \bar{I}^{\tilde{r}}(\xi) \end{pmatrix} > \begin{pmatrix} \underline{S}(\xi) \\ \underline{I}(\xi) \end{pmatrix}, \quad \xi \in [-N_1, N_2]. \quad (48)$$

We deal with such two possible cases:

Case 1. If

$$\begin{pmatrix} \bar{S}^{\tilde{r}}(\xi) \\ \bar{I}^{\tilde{r}}(\xi) \end{pmatrix} > \begin{pmatrix} \underline{S}(\xi) \\ \underline{I}(\xi) \end{pmatrix}, \quad \xi \in [-N_1, N_2], \quad (49)$$

then, the proof is completed.

Case 2. If there exists a point $\xi_0 \in (-N_1, N_2)$, such that

$$\begin{pmatrix} \bar{S}^{\tilde{r}}(\xi_0) \\ \bar{I}^{\tilde{r}}(\xi_0) \end{pmatrix} \leq \begin{pmatrix} \underline{S}(\xi_0) \\ \underline{I}(\xi_0) \end{pmatrix} \quad (50)$$

satisfying $\bar{S}^{\tilde{r}}(\xi_0) < \underline{S}(\xi_0)$ or $\bar{I}^{\tilde{r}}(\xi_0) < \underline{I}(\xi_0)$.

In this case, we use the Sliding Domain method [15].

Step 1. we shift $(\bar{S}^{\tilde{r}}, \bar{I}^{\tilde{r}})^T$ to the left by increasing the number $\tilde{r}$ until finding a new number $r_1 > \tilde{r}$ such that $(\bar{S}^{r_1}, \bar{I}^{r_1})^T > (\underline{S}, \underline{I})^T$ on the smaller interval $[-N_1, N_2 - (r_1 - \tilde{r})]$.

Step 2. we shift $(\bar{S}^{r_1}, \bar{I}^{r_1})^T$ back to the right by decreasing r_1 to a smaller number $\tilde{r} < r_1$ such that one of the

branches of the upper solution touches its counterpart of the lower solution at some point ξ_1 in the interval $(-N_1 + (r_1 - r_2), N_2 - (r_1 - \tilde{r}))$. On the endpoints of the interval $(-N_1 + (r_1 - r_2), N_2 - (r_1 - \tilde{r}))$, we still have $(\bar{S}^{r_2}, \bar{I}^{r_2})^T > (\underline{S}, \underline{I})^T$.

Let $\tilde{W}(\xi) = (\bar{S}^{r_2}, \bar{I}^{r_2})^T - (\underline{S}, \underline{I})^T$ and $\tilde{F} = (F_1, F_2)^T$, where

$$\begin{aligned} F_1 &= -\nu R_d (\alpha_0 - S + I) [1 - (\alpha_0 - S + I)] \\ &\quad + R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} + \nu (\alpha_0 - S), \\ F_2 &= R_0 \frac{(\alpha_0 - S)I}{\alpha_0 - S + I} - I. \end{aligned} \quad (51)$$

For $\xi \in (-N_1 + (r_1 - r_2), N_2 - (r_1 - \tilde{r}))$, we get that

$$\begin{aligned} \tilde{W}'' - c\tilde{W}' + \left(\frac{\partial F_1}{\partial S} (\underline{S} + \zeta_1 \omega_1, \bar{I}^{r_2}) \frac{\partial F_1}{\partial I} (\bar{S}^{r_2}, \underline{I} + \zeta_2 \omega_2) \right. \\ \left. \frac{\partial F_2}{\partial S} (\underline{S} + \zeta_3 \omega_1, \bar{I}^{r_2}) \frac{\partial F_2}{\partial I} (\bar{S}^{r_2}, \underline{I} + \zeta_4 \omega_2) \right) \tilde{W} \\ = 0, \end{aligned} \quad (52)$$

where $\zeta_i \in [0, 1]$, $i = 1, 2, 3, 4$. Since the above model is monotone and the cube $[(0, 0), (\bar{S}^*, \bar{I}^*)]$ is convex, thus we can deduce by Maximum Principle that $\tilde{W} > 0$ for $\xi \in [-N_1 + (r_1 - r_2), N_2 - (r_1 - \tilde{r})]$. So ξ_1 does not exist and we can decrease r_2 further to $\tilde{r}$. It is calculated that the point ξ_0 does not exist either. The proof of this lemma is completed. $\square$

To ease the burden of notations, we still use $(\bar{S}, \bar{I})^T$ to denote the shifted upper solution as given in Lemma 8. Let

$$\begin{aligned} D_{11} &= -\frac{R_0^2 + \nu R_0 R_d + \nu R_0 - 4R_0 - 2\nu + 3}{R_0}, \\ D_{12} &= \frac{\nu R_0 R_d - 2R_0 - 2\nu + 3}{R_0}, \\ D_{21} &= -\frac{(R_0 - 1)^2}{R_0}, \\ D_{22} &= -\frac{R_0 - 1}{R_0}, \end{aligned} \quad (53)$$

$$\begin{aligned} \mu_1 &= \frac{-(D_{11} + D_{22}) + \sqrt{(D_{11} - D_{22})^2 + 4D_{12}D_{21}}}{2}, \\ \mu_2 &= \frac{-(D_{11} + D_{22}) - \sqrt{(D_{11} - D_{22})^2 + 4D_{12}D_{21}}}{2}. \end{aligned}$$

With such constructed ordered upper-lower solution pair, we can get the following.

Theorem 11. For $c \geq 2\sqrt{R_0 - 1}$, model (24) has a unique (up to a translation of the origin) traveling wave solution. The traveling wave solution is strictly increasing and has the following asymptotic properties:

(i) if $c > 2\sqrt{R_0 - 1}$, when $\xi \rightarrow -\infty$,

$$\begin{aligned} \begin{pmatrix} S \\ I \end{pmatrix}(\xi) &= \begin{pmatrix} A_1 \\ A_2 \end{pmatrix} e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ &+ o\left(e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}\right). \end{aligned} \quad (54)$$

when $\xi \rightarrow +\infty$, and if $\mu_1 \neq \mu_2$, then

$$\begin{aligned} \begin{pmatrix} S \\ I \end{pmatrix}(\xi) &= \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix} - \begin{pmatrix} \overline{A_1} \\ \overline{A_2} \end{pmatrix} e^{((c - \sqrt{c^2 + 4\mu})/2)\xi} \\ &+ o\left(e^{((c - \sqrt{c^2 + 4\mu})/2)\xi}\right), \end{aligned} \quad (55)$$

while $\mu_1 = \mu_2$:

$$\begin{aligned} \begin{pmatrix} S \\ I \end{pmatrix}(\xi) &= \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix} - \begin{pmatrix} \overline{A_{11}} + \overline{A_{12}}\xi \\ \overline{A_{21}} + \overline{A_{22}}\xi \end{pmatrix} e^{((c - \sqrt{c^2 + 4\mu})/2)\xi} \\ &+ o\left(e^{((c - \sqrt{c^2 + 4\mu})/2)\xi}\right), \end{aligned} \quad (56)$$

where, $\mu = \min\{\mu_1, \mu_2\} > 0$, $\overline{A_{11}}, \overline{A_{21}} \in \mathbb{R}$, $A_1, A_2, \overline{A_1}, \overline{A_2}, \overline{A_{12}}$ and $\overline{A_{22}}$ are all positive constants.

(ii) if $c = 2\sqrt{R_0 - 1}$, when $\xi \rightarrow -\infty$,

$$\begin{pmatrix} S \\ I \end{pmatrix}(\xi) = \begin{pmatrix} B_{11} + B_{12}\xi \\ B_{21} + B_{22}\xi \end{pmatrix} e^{\sqrt{R_0 - 1}\xi} + o\left(\xi e^{\sqrt{R_0 - 1}\xi}\right), \quad (57)$$

when $\xi \rightarrow +\infty$, and if $\mu_1 \neq \mu_2$, then

$$\begin{aligned} \begin{pmatrix} S \\ I \end{pmatrix}(\xi) &= \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix} - \begin{pmatrix} \overline{B_{11}} \\ \overline{B_{22}} \end{pmatrix} e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi} \\ &+ o\left(e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi}\right), \end{aligned} \quad (58)$$

while $\mu_1 = \mu_2$,

$$\begin{aligned} \begin{pmatrix} S \\ I \end{pmatrix}(\xi) &= \begin{pmatrix} \hat{S}^* \\ \hat{I}^* \end{pmatrix} - \begin{pmatrix} \overline{B_{11}} + \overline{B_{12}}\xi \\ \overline{B_{21}} + \overline{B_{22}}\xi \end{pmatrix} e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi} \\ &+ o\left(e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi}\right), \end{aligned} \quad (59)$$

where $\mu = \min\{\mu_1, \mu_2\} > 0$, $B_{12}, B_{22} < 0$, $B_{11}, B_{21}, \overline{B_{11}}, \overline{B_{21}} \in \mathbb{R}$, and $\overline{B_{11}}, \overline{B_{22}}, \overline{B_{12}}, \overline{B_{22}}$ are all positive constants.

Proof. From Lemma 3 and Remark 4, we know that model (24) is a quasi-monotone nonincreasing system in $(S, I) \in [\widehat{E_1}, \widehat{E_2}]$, and by using the monotone iteration scheme given in [3, 13], we can obtain the existence of the solution $(S, I)^T$ to

the first two equations in model (24) for every $c \geq 2\sqrt{R_0 - 1}$, which satisfies

$$\begin{pmatrix} S(\xi) \\ \underline{I}(\xi) \end{pmatrix} \leq \begin{pmatrix} S(\xi) \\ I(\xi) \end{pmatrix} \leq \begin{pmatrix} \overline{S}(\xi) \\ \overline{I}(\xi) \end{pmatrix}. \quad (60)$$

According to the above inequality, we can get that, on the boundary, the solution tends to $(0, 0)^T$ as $\xi \rightarrow -\infty$ and $(\hat{S}^*, \hat{I}^*)^T$ as $\xi \rightarrow +\infty$.

To derive the asymptotic decay rate of the traveling wave solutions as $\xi \rightarrow \pm\infty$, we just let $c > 2\sqrt{R_0 - 1}$ and

$$U(\xi) = (S(\xi), I(\xi))^T, \quad -\infty < \xi < +\infty \quad (61)$$

be the traveling wave solution of model (24) generated from the monotone iteration, since the case of (ii) $c = 2\sqrt{R_0 - 1}$ is similar to it.

We differentiate model (24) with respect to ξ , and note that $U'(\xi) = (\chi_1, \chi_2)^T(\xi)$ satisfies

$$\begin{aligned} \chi_1'' - c\chi_1' + C_{11}(S, I)\chi_1 + C_{12}(S, I)\chi_2 &= 0, \\ \chi_2'' - c\chi_2' + C_{21}(S, I)\chi_1 + C_{22}(S, I)\chi_2 &= 0, \end{aligned} \quad (62)$$

where

$$C_{11}(S, I) = \nu R_d [1 - (\alpha_0 - S + I)] - \nu R_d (\alpha_0 - S + I)$$

$$- \frac{R_0 I}{\alpha_0 - S + I} + \frac{R_0 (\alpha_0 - S) I}{(\alpha_0 - S + I)^2} - \nu,$$

$$C_{12}(S, I) = -\nu R_d [1 - (\alpha_0 - S + I)] + \nu R_d (\alpha_0 - S + I)$$

$$+ \frac{R_0 (\alpha_0 - S)}{\alpha_0 - S + I} - \frac{R_0 (\alpha_0 - S) I}{(\alpha_0 - S + I)^2},$$

$$C_{21}(S, I) = -\frac{R_0 I}{\alpha_0 - S + I} + \frac{R_0 (\alpha_0 - S) I}{(\alpha_0 - S + I)^2},$$

$$C_{22}(S, I) = \frac{R_0 (\alpha_0 - S)}{\alpha_0 - S + I} - \frac{R_0 (\alpha_0 - S) I}{(\alpha_0 - S + I)^2} - 1. \quad (63)$$

Now, we study the exponential decay rate of the traveling wave solution as $\xi \rightarrow -\infty$. The asymptotic model of model (62) as $\xi \rightarrow -\infty$ is

$$\begin{aligned} \lambda'' - c\lambda' + E_{11}\lambda + E_{12}\mu &= 0, \\ \mu'' - c\mu' + E_{21}\lambda + E_{22}\mu &= 0, \end{aligned} \quad (64)$$

where

$$\begin{aligned} E_{11} &= -\nu(R_d - 1), & E_{12} &= \nu R_d + R_0 - 2\nu, \\ E_{21} &= 0, & E_{22} &= R_0 - 1. \end{aligned} \quad (65)$$

The second equation of model (64) has two independent solutions with the following form:

$$\begin{aligned} \mu^{(1)}(\xi) &= e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}, \\ \mu^{(2)}(\xi) &= e^{((c + \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}. \end{aligned} \quad (66)$$

Relating the second equation of model (62) with the second equation of model (64), we can deduce that $\chi_2(\xi)$ has the following property as $\xi \rightarrow -\infty$:

$$\begin{aligned} \chi_2(\xi) = & \alpha [1 + o(1)] e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ & + \beta [1 + o(1)] e^{((c + \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \end{aligned} \quad (67)$$

for some constants α and β . Thus, we can obtain that

$$\begin{aligned} \chi_2(\xi) = & \alpha e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} + \beta e^{((c + \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ & + Y_1(\xi) + Y_2(\xi), \end{aligned} \quad (68)$$

where

$$\begin{aligned} \lim_{\xi \rightarrow -\infty} \frac{Y_1(\xi)}{e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}} &= 0, \\ \lim_{\xi \rightarrow -\infty} \frac{Y_2(\xi)}{e^{((c + \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}} &= 0. \end{aligned} \quad (69)$$

So we obtain that

$$\begin{aligned} \lim_{\xi \rightarrow -\infty} \frac{\chi_2(\xi) - \alpha e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}}{e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}} &= \lim_{\xi \rightarrow -\infty} \frac{Y_1(\xi) + Y_2(\xi) + \beta e^{((c + \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}}{e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}} \\ &= \lim_{\xi \rightarrow -\infty} \frac{Y_1(\xi)}{e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}} + \beta \lim_{\xi \rightarrow -\infty} e^{\sqrt{c^2 - 4(R_0 - 1)}\xi} \\ &+ \lim_{\xi \rightarrow -\infty} \frac{Y_2(\xi)}{e^{((c + \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}} \lim_{\xi \rightarrow -\infty} e^{\sqrt{c^2 - 4(R_0 - 1)}\xi} = 0. \end{aligned} \quad (70)$$

Thus, $\chi_2(\xi) = \alpha[1 + o(1)]e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}$.

Now, we consider the first equation of model (64). We rewrite it as

$$\lambda'' - c\lambda' - \nu(R_d - 1)\lambda = -(\nu R_d + R_0 - 2\nu)\mu. \quad (71)$$

One can verify that $(c - \sqrt{c^2 - 4(R_0 - 1)})/2$ is not a characteristic of

$$\lambda'' - c\lambda' - \nu(R_d - 1)\lambda = 0. \quad (72)$$

The above equation has two independent solutions of the following form:

$$\begin{aligned} \lambda^{(1)}(\xi) &= e^{((c - \sqrt{c^2 + 4\nu(R_d - 1)})/2)\xi}, \\ \lambda^{(2)}(\xi) &= e^{((c + \sqrt{c^2 + 4\nu(R_d - 1)})/2)\xi}. \end{aligned} \quad (73)$$

Thus, when $\xi \rightarrow -\infty$, $\chi_1(\xi)$ has the following property:

$$\begin{aligned} \chi_1(\xi) = & \bar{\alpha} [1 + o(1)] e^{((c + \sqrt{c^2 + 4\nu(R_d - 1)})/2)\xi} \\ & + \bar{\beta} [1 + o(1)] e^{((c - \sqrt{c^2 + 4\nu(R_d - 1)})/2)\xi} \\ & + \gamma [1 + o(1)] e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \end{aligned} \quad (74)$$

for some constants $\bar{\alpha}, \bar{\beta}; \gamma \neq 0$. Since $\chi_1(-\infty) = 0$, thus $\bar{\beta} = 0$. So, when $\xi \rightarrow -\infty$, we have the following formula:

$$\begin{pmatrix} \chi_1(\xi) \\ \chi_2(\xi) \end{pmatrix} = \begin{pmatrix} \gamma [1 + o(1)] e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ \alpha [1 + o(1)] e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \end{pmatrix}. \quad (75)$$

Then, we study the exponential decay rate of the traveling wave solution as $\xi \rightarrow +\infty$. The asymptotic model of model (62) as $\xi \rightarrow +\infty$ is

$$\begin{aligned} \psi_1'' - c\psi_1' + D_{11}\psi_1 + D_{12}\psi_2 &= 0, \\ \psi_2'' - c\psi_2' + D_{21}\psi_1 + D_{22}\psi_2 &= 0. \end{aligned} \quad (76)$$

By setting $(\psi_i)'(\xi) = \tilde{\psi}_i$, $i = 1, 2$, we rewrite model (76) as a first order model of ordinary differential equation in the four components $(\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T$:

$$\begin{aligned} \psi_1' &= \tilde{\psi}_1, \\ \tilde{\psi}_1' &= c\tilde{\psi}_1 - D_{11}\psi_1 - D_{12}\psi_2, \\ \psi_2' &= \tilde{\psi}_2, \\ \tilde{\psi}_2' &= c\tilde{\psi}_2 - D_{21}\psi_1 - D_{22}\psi_2. \end{aligned} \quad (77)$$

In the case of (i) $\mu_1 \neq \mu_2$, we can obtain that the solution of model (77) has the following form:

$$(\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T = \sum_{i=1}^4 c_i h_i e^{\lambda_i \xi}, \quad (78)$$

where

$$\begin{aligned} \lambda_1 &= \frac{c + \sqrt{c^2 + 4\mu_1}}{2}, & \lambda_2 &= \frac{c - \sqrt{c^2 + 4\mu_1}}{2}, \\ \lambda_3 &= \frac{c + \sqrt{c^2 + 4\mu_2}}{2}, & \lambda_4 &= \frac{c - \sqrt{c^2 + 4\mu_2}}{2}, \end{aligned} \quad (79)$$

and h_i ($i = 1, 2, 3, 4$) are the eigenvectors of the constant matrix with λ_i ($i = 1, 2, 3, 4$) as the corresponding eigenvalues, c_i ($i = 1, 2, 3, 4$) are arbitrary constants. Since

$$\lim_{\xi \rightarrow +\infty} (\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T = 0, \quad (80)$$

thus $(\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T = c_2 h_2 e^{\lambda_2 \xi} + c_4 h_4 e^{\lambda_4 \xi}$, so when $\xi \rightarrow +\infty$, we can get that

$$\begin{aligned} \begin{pmatrix} \chi_1(\xi) \\ \chi_2(\xi) \end{pmatrix} &= \begin{pmatrix} \kappa_1 [\Lambda_1 + o(1)] e^{((c - \sqrt{c^2 + 4\mu_1})/2)\xi} \\ \kappa_1 [\Gamma_1 + o(1)] e^{((c - \sqrt{c^2 + 4\mu_1})/2)\xi} \end{pmatrix} \\ &+ \begin{pmatrix} \kappa_2 [\Lambda_2 + o(1)] e^{((c - \sqrt{c^2 + 4\mu_2})/2)\xi} \\ \kappa_2 [\Gamma_2 + o(1)] e^{((c - \sqrt{c^2 + 4\mu_2})/2)\xi} \end{pmatrix}. \end{aligned} \quad (81)$$

Furthermore, we can obtain that

$$\begin{aligned}\chi_1(\xi) &= \kappa_1 \Lambda_1 e^{((c-\sqrt{c^2+4\mu_1})/2)\xi} + \kappa_2 \Lambda_2 e^{((c-\sqrt{c^2+4\mu_2})/2)\xi} \\ &\quad + \Omega_{11}(\xi) + \Omega_{12}(\xi), \\ \chi_2(\xi) &= \kappa_1 \Gamma_1 e^{((c-\sqrt{c^2+4\mu_1})/2)\xi} + \kappa_2 \Gamma_2 e^{((c-\sqrt{c^2+4\mu_2})/2)\xi} \\ &\quad + \Omega_{21}(\xi) + \Omega_{22}(\xi),\end{aligned}\quad (82)$$

where

$$\begin{aligned}\lim_{\xi \rightarrow +\infty} \frac{\Omega_{11}(\xi)}{e^{((c-\sqrt{c^2+4\mu_1})/2)\xi}} &= 0, & \lim_{\xi \rightarrow +\infty} \frac{\Omega_{12}(\xi)}{e^{((c-\sqrt{c^2+4\mu_2})/2)\xi}} &= 0, \\ \lim_{\xi \rightarrow +\infty} \frac{\Omega_{21}(\xi)}{e^{((c-\sqrt{c^2+4\mu_1})/2)\xi}} &= 0, & \lim_{\xi \rightarrow +\infty} \frac{\Omega_{22}(\xi)}{e^{((c-\sqrt{c^2+4\mu_2})/2)\xi}} &= 0.\end{aligned}\quad (83)$$

$\kappa_1, \kappa_2, \Lambda_1, \Lambda_2, \Gamma_1$, and Γ_2 are all constants.

Let $\mu = \min\{\mu_1, \mu_2\}$, then

$$\begin{aligned}\lim_{\xi \rightarrow +\infty} \frac{\chi_1(\xi)}{e^{((c-\sqrt{c^2+4\mu})/2)\xi}} &= \kappa_1 \Lambda_1 \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_1})/2)\xi} \\ &\quad + \kappa_2 \Lambda_2 \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_2})/2)\xi} \\ &\quad + \lim_{\xi \rightarrow +\infty} \frac{\Omega_{11}(\xi)}{e^{((c-\sqrt{c^2+4\mu_1})/2)\xi}} \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_1})/2)\xi} \\ &\quad + \lim_{\xi \rightarrow +\infty} \frac{\Omega_{12}(\xi)}{e^{((c-\sqrt{c^2+4\mu_2})/2)\xi}} \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_2})/2)\xi} \\ &= \Delta_1(\kappa_1, \kappa_2, \Lambda_1, \Lambda_2), \\ \lim_{\xi \rightarrow +\infty} \frac{\chi_2(\xi)}{e^{((c-\sqrt{c^2+4\mu})/2)\xi}} &= \kappa_1 \Gamma_1 \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_1})/2)\xi} \\ &\quad + \kappa_2 \Gamma_2 \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_2})/2)\xi} \\ &\quad + \lim_{\xi \rightarrow +\infty} \frac{\Omega_{21}(\xi)}{e^{((c-\sqrt{c^2+4\mu_1})/2)\xi}} \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_1})/2)\xi} \\ &\quad + \lim_{\xi \rightarrow +\infty} \frac{\Omega_{22}(\xi)}{e^{((c-\sqrt{c^2+4\mu_2})/2)\xi}} \lim_{\xi \rightarrow +\infty} e^{((\sqrt{c^2+4\mu}-\sqrt{c^2+4\mu_2})/2)\xi} \\ &= \Delta_2(\kappa_1, \kappa_2, \Gamma_1, \Gamma_2),\end{aligned}\quad (84)$$

where

$$\Delta_1(\kappa_1, \kappa_2, \Lambda_1, \Lambda_2) = \begin{cases} \kappa_1 \Lambda_1, & \mu_1 < \mu_2, \\ \kappa_2 \Lambda_2, & \mu_1 > \mu_2, \end{cases}\quad (85)$$

$$\Delta_2(\kappa_1, \kappa_2, \Gamma_1, \Gamma_2) = \begin{cases} \kappa_1 \Gamma_1, & \mu_1 < \mu_2, \\ \kappa_2 \Gamma_2, & \mu_1 > \mu_2; \end{cases}$$

thus, when $\xi \rightarrow +\infty$, we can get that

$$\begin{pmatrix} \chi_1(\xi) \\ \chi_2(\xi) \end{pmatrix} = \begin{pmatrix} \Delta_1(\kappa_1, \kappa_2, \Lambda_1, \Lambda_2) (1 + o(1)) e^{((c-\sqrt{c^2+4\mu})/2)\xi} \\ \Delta_2(\kappa_1, \kappa_2, \Gamma_1, \Gamma_2) (1 + o(1)) e^{((c-\sqrt{c^2+4\mu})/2)\xi} \end{pmatrix}.\quad (86)$$

In the case of (ii) $\mu_1 = \mu_2$, we can obtain that the solution of model (77) has the following form:

$$\begin{aligned}(\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T &= (G_1 + G_2 \xi) H_{1,2} e^{\bar{\lambda}_1 \xi} \\ &\quad + (G_3 + G_4 \xi) H_{3,4} e^{\bar{\lambda}_3 \xi},\end{aligned}\quad (87)$$

where $H_{1,2}$ is the eigenvector of the constant matrix with $\bar{\lambda}_1 = \bar{\lambda}_2$ as the corresponding eigenvalues, $H_{3,4}$ is the eigenvector of the constant matrix with $\bar{\lambda}_3 = \bar{\lambda}_4$ as the corresponding eigenvalues, G_i ($i = 1, 2, 3, 4$) are arbitrary constants.

Since $\lim_{\xi \rightarrow +\infty} (\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T = 0$, thus

$$(\psi_1, \tilde{\psi}_1, \psi_2, \tilde{\psi}_2)^T = (G_3 + G_4 \xi) H_{3,4} e^{\bar{\lambda}_3 \xi}.\quad (88)$$

So, when $\xi \rightarrow +\infty$, we can get that

$$\begin{pmatrix} \chi_1(\xi) \\ \chi_2(\xi) \end{pmatrix} = \begin{pmatrix} (G_{1,3} + G_{1,4} \xi) e^{\bar{\lambda}_3 \xi} + o(e^{\bar{\lambda}_3 \xi}) \\ (G_{2,3} + G_{2,4} \xi) e^{\bar{\lambda}_4 \xi} + o(e^{\bar{\lambda}_4 \xi}) \end{pmatrix}.\quad (89)$$

By comparing the upper solution and roughness of the exponential dichotomy [24], we obtain the asymptotic decay rate of the traveling wave solutions at $+\infty$ given in Theorem 11.

According to the monotone iteration process [3], the traveling wave solution $U(\xi)$ is increasing; thus $U'(\xi) \geq 0$ and $U'(\xi) = (\chi_1, \chi_2)^T(\xi)$ hold

$$\begin{aligned}\chi_1'' - c\chi_1' + \frac{\partial F_1}{\partial S} \chi_1 + \frac{\partial F_1}{\partial I} \chi_2 &= 0, \\ \chi_2'' - c\chi_2' + \frac{\partial F_2}{\partial S} \chi_1 + \frac{\partial F_2}{\partial I} \chi_2 &= 0,\end{aligned}\quad (90)$$

satisfying

$$(\chi_1, \chi_2)^T(\xi) \geq 0, \quad (\chi_1, \chi_2)^T(\pm\infty) = 0.\quad (91)$$

The strong Maximum Principle implies that $(\chi_1, \chi_2)^T(\xi) > 0$. So the strict monotonicity of the traveling wave solutions is concluded.

Now, we use the Sliding domain method to prove the uniqueness of the traveling wave solution. Let $U_1(\xi) = (S_1, I_1)^T(\xi)$ and $U_2(\xi) = (S_2, I_2)^T(\xi)$ be the traveling wave solution of model (24), with $c > 2\sqrt{R_0 - 1}$. Thus, there are some positive numbers A_i, B_i ($i = 1, 2, 3, 4$), such that for a big enough number $N \gg 1$, when $\xi < -N$, we have

$$\begin{aligned} \begin{pmatrix} S_1(\xi) \\ I_1(\xi) \end{pmatrix} &= \begin{pmatrix} A_1 e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi} \\ A_2 e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi} \end{pmatrix} \\ &\quad + o\left(e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi}\right), \\ \begin{pmatrix} S_2(\xi) \\ I_2(\xi) \end{pmatrix} &= \begin{pmatrix} A_3 e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi} \\ A_4 e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi} \end{pmatrix} \\ &\quad + o\left(e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi}\right), \end{aligned} \quad (92)$$

when $\xi > N$,

$$\begin{aligned} \begin{pmatrix} S_1(\xi) \\ I_1(\xi) \end{pmatrix} &= \begin{pmatrix} \hat{S}^* - B_1 e^{((c-\sqrt{c^2+4\mu})/2)\xi} \\ \hat{I}^* - B_2 e^{((c-\sqrt{c^2+4\mu})/2)\xi} \end{pmatrix} \\ &\quad + o\left(e^{((c-\sqrt{c^2+4\mu})/2)\xi}\right), \\ \begin{pmatrix} S_2(\xi) \\ I_2(\xi) \end{pmatrix} &= \begin{pmatrix} \hat{S}^* - B_3 e^{((c-\sqrt{c^2+4\mu})/2)\xi} \\ \hat{I}^* - B_4 e^{((c-\sqrt{c^2+4\mu})/2)\xi} \end{pmatrix} \\ &\quad + o\left(e^{((c-\sqrt{c^2+4\mu})/2)\xi}\right). \end{aligned} \quad (93)$$

Since the traveling wave solutions of model (24) are translation-invariant, then for any $\theta \in R$, $U_1^\theta(\xi) \equiv U_1(\xi + \theta) \equiv (S_1(\xi + \theta), I_1(\xi + \theta))^T$ is also a traveling wave solution of model (24). Thus, by using the same method as above, when $\xi < -N$, we can get

$$\begin{aligned} \begin{pmatrix} S_1(\xi + \theta) \\ I_1(\xi + \theta) \end{pmatrix} &= \begin{pmatrix} A_1 e^{((c-\sqrt{c^2-4(R_0-1)})/2)\theta} e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi} \\ A_2 e^{((c-\sqrt{c^2-4(R_0-1)})/2)\theta} e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi} \end{pmatrix} \\ &\quad + o\left(e^{((c-\sqrt{c^2-4(R_0-1)})/2)\xi}\right), \end{aligned} \quad (94)$$

when $\xi > N$,

$$\begin{aligned} \begin{pmatrix} S_1(\xi + \theta) \\ I_1(\xi + \theta) \end{pmatrix} &= \begin{pmatrix} \hat{S}^* - B_1 e^{((c-\sqrt{c^2+4\mu})/2)\theta} e^{((c-\sqrt{c^2+4\mu})/2)\xi} \\ \hat{I}^* - B_2 e^{((c-\sqrt{c^2+4\mu})/2)\theta} e^{((c-\sqrt{c^2+4\mu})/2)\xi} \end{pmatrix} \\ &\quad + o\left(e^{((c-\sqrt{c^2+4\mu})/2)\xi}\right). \end{aligned} \quad (95)$$

If θ is large enough, then we can obtain the following inequalities:

$$\begin{aligned} A_1 e^{((c+\sqrt{c^2-4(R_0-1)})/2)\theta} &> A_3, \\ A_2 e^{((c+\sqrt{c^2-4(R_0-1)})/2)\theta} &> A_4, \\ B_1 e^{((c-\sqrt{c^2+4\mu})/2)\theta} &< B_3, \\ B_2 e^{((c-\sqrt{c^2+4\mu})/2)\theta} &< B_4. \end{aligned} \quad (96)$$

Thus, if θ is large enough, then $U_1^\theta(\xi) > U_2(\xi)$, for all $\xi \in R \setminus [-N, N]$.

Now, we consider model (24) on the interval $[-N, N]$.

First, suppose that

$$W(\xi) \equiv U_1^\theta(\xi) - U_2(\xi) \geq 0, \quad \xi \in [-N, N], \quad (97)$$

then

$$\begin{aligned} W'' - cW' &+ \left(\frac{\partial F_1}{\partial S}(S_2 + \zeta_1 \omega_1, I_1) \frac{\partial F_1}{\partial I}(S_1, I_2 + \zeta_2 \omega_2) \right. \\ &\quad \left. + \left(\frac{\partial F_2}{\partial S}(S_2 + \zeta_3 \omega_1, I_1) \frac{\partial F_2}{\partial I}(S_1, I_2 + \zeta_4 \omega_2) \right) \right) W = 0, \\ W(-N) &> 0, \quad W(N) > 0, \end{aligned} \quad (98)$$

where, $\zeta_i \in (0, 1)$ ($i = 1, 2, 3, 4$), $\xi \in (-N, N)$. Since the above model is monotone, by the Maximum Principle, we can deduce that $W(\xi) > 0$, $\xi \in [-N, N]$. Consequently, we get that $U_1^\theta(\xi) > U_2(\xi)$, $\xi \in R$.

Second, we suppose that there exists a point $\xi_* \in (-N, N)$ such that

$$S_1^\theta(\xi_*) < S_2(\xi_*) \quad (99)$$

or

$$I_1^\theta(\xi_*) < I_2(\xi_*). \quad (100)$$

In this case, we increase θ , that is shifting U_1^θ to the left, so that $U_1^\theta(-N) > U_2(-N)$ and $U_1^\theta(N) > U_2(N)$. According to the monotonicity of U_1^θ and U_2 , we can find a number $\bar{\theta} \in (0, 2N)$ such that $U_1^\theta(\xi + \bar{\theta}) > U_2(\xi)$, $\xi \in (-N, N)$. Shifting $U_1^\theta(\xi + \bar{\theta})$ back until one component of $U_1^\theta(\xi + \bar{\theta})$ touches its counterpart of $U_2(\xi)$ at some point $\bar{\xi} \in (-N, N)$. Since $U_1^\theta(\xi + \bar{\theta})$ and $U_2(\xi)$ are strictly increasing, $\bar{\xi} \in (-N, N)$, thus, we get that $U_1^\theta(\xi + \bar{\theta}) > U_2(\xi)$, $\xi = \pm N$. However, by the Maximum Principle for that component again, we find that components of U_1^θ and U_2 are identically equal for all $\xi \in [-N, N]$ for a larger number θ . This is a contradiction, thus $U_1^\theta(\xi) > U_2(\xi)$, $\xi \in R$. Here, θ is a new number which is chosen by the above mean.

Now, decrease the θ until one of the following happens.

Case (a). There is a $\bar{\theta} \geq 0$, such that $U_1^{\bar{\theta}} = U_2(\xi)$, $\xi \in R$. In this case, we have finished the proof.

Case (b). There are a $\bar{\theta}$ and a point $\xi_1 \in R$, such that one of the components of $U_1^{\bar{\theta}}$ and U_2 are equal. And $U_1^{\bar{\theta}} \geq U_2$, $\xi \in R$. On R for that component, according to the Maximum Principle, we find that $U_1^{\bar{\theta}}$ and U_2 must be identical on that component. We can return to Case (a).

Consequently, in either situation, there exists a number $\bar{\theta} \geq 0$ such that

$$U_1^{\bar{\theta}}(\xi) = U_2(\xi), \quad \xi \in (-\infty, +\infty). \quad (101)$$

This ends of the proof. $\square$

By Theorem 11, we can get the following theorem:

Theorem 12. For each $c \geq 2\sqrt{R_0 - 1}$, model (3) has a unique (up to a translation of the origin) traveling wave solution. The traveling wave solution is strictly increasing and has the following asymptotic properties:

(i) $c > 2\sqrt{R_0 - 1}$: when $\xi \rightarrow -\infty$,

$$\begin{aligned} S(\xi) &= \frac{R_d - 1}{R_d} - A_1 e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}\right), \\ I(\xi) &= A_2 e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi} + o\left(e^{((c - \sqrt{c^2 - 4(R_0 - 1)})/2)\xi}\right). \end{aligned} \quad (102)$$

when $\xi \rightarrow +\infty$, and if $\mu_1 \neq \mu_2$, then

$$\begin{aligned} S(\xi) &= S^* + \overline{A_1} e^{((c - \sqrt{c^2 + 4\mu})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 + 4\mu})/2)\xi}\right), \\ I(\xi) &= I^* - \overline{A_2} e^{((c - \sqrt{c^2 + 4\mu})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 + 4\mu})/2)\xi}\right), \end{aligned} \quad (103)$$

if $\mu_1 = \mu_2$,

$$\begin{aligned} S(\xi) &= S^* + \left(\overline{A_{11}} + \overline{A_{12}}\xi\right) e^{((c - \sqrt{c^2 + 4\mu})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 + 4\mu})/2)\xi}\right), \\ I(\xi) &= I^* - \left(\overline{A_{21}} + \overline{A_{22}}\xi\right) e^{((c - \sqrt{c^2 + 4\mu})/2)\xi} \\ &\quad + o\left(e^{((c - \sqrt{c^2 + 4\mu})/2)\xi}\right), \end{aligned} \quad (104)$$

where $\mu = \min\{\mu_1, \mu_2\} > 0$, $\overline{A_{11}}, \overline{A_{21}} \in \mathbb{R}$, $A_1, A_2, \overline{A_1}, \overline{A_2}$, $\overline{A_{12}}$ and $\overline{A_{22}}$ are all positive constants.

(ii) $c = 2\sqrt{R_0 - 1}$: when $\xi \rightarrow -\infty$,

$$\begin{aligned} S(\xi) &= \frac{R_d - 1}{R_d} - (A_{11} + A_{12}\xi) e^{\sqrt{R_0 - 1}\xi} + o\left(\xi e^{\sqrt{R_0 - 1}\xi}\right), \\ I(\xi) &= (A_{11} + A_{12}\xi) e^{\sqrt{R_0 - 1}\xi} + o\left(\xi e^{\sqrt{R_0 - 1}\xi}\right). \end{aligned} \quad (105)$$

when $\xi \rightarrow +\infty$, and if $\mu_1 \neq \mu_2$, then

$$\begin{aligned} S(\xi) &= S^* + \overline{B_{11}} e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi} + o\left(e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi}\right), \\ I(\xi) &= I^* - \overline{B_{22}} e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi} + o\left(e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi}\right), \end{aligned} \quad (106)$$

if $\mu_1 = \mu_2$, then

$$\begin{aligned} S(\xi) &= S^* + \left(\overline{B_{11}} + \overline{B_{12}}\xi\right) e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi} \\ &\quad + o\left(e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi}\right), \\ I(\xi) &= I^* - \left(\overline{B_{21}} + \overline{B_{22}}\xi\right) e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi} \\ &\quad + o\left(e^{(\sqrt{R_0 - 1} - \sqrt{R_0 - 1 + \mu})\xi}\right), \end{aligned} \quad (107)$$

where $\mu = \min\{\mu_1, \mu_2\} > 0$, $B_{12}, B_{22} < 0$, $B_{11}, B_{21}, \overline{B_{11}}, \overline{B_{21}} \in \mathbb{R}$, $\overline{B_{11}}, \overline{B_{22}}, \overline{B_{12}}, \overline{B_{22}}$ are all positive constants.

Theorem 13. There is no monotone traveling wave solution of model (24) for any $0 < c < 2\sqrt{R_0 - 1}$. In other words, there is no monotone traveling wave solution of model (3) for any $0 < c < 2\sqrt{R_0 - 1}$.

Proof. Suppose there is a monotone traveling wave solution $\mathbb{L}(\xi) = (l_1(\xi), l_2(\xi))^T$ of model (24) with the wave speed c_0 , where $c_0 \in (0, 2\sqrt{R_0 - 1})$.

The asymptotic model of $\mathbb{L}(\xi) = (l_1(\xi), l_2(\xi))^T$ as $\xi \rightarrow -\infty$ is

$$\begin{aligned} \bar{\lambda}'' - c_0 \bar{\lambda}' - \nu(R_d - 1)\bar{\lambda} + (\nu R_d + R_0 - 2\nu)\bar{\mu} &= 0, \\ \bar{\mu}'' - c_0 \bar{\mu}' + (R_0 - 1)\bar{\mu} &= 0. \end{aligned} \quad (108)$$

The second function of (108) has two characteristics as the following ones: $(c_0 + \sqrt{4(R_0 - 1) - c_0^2})/2$, $(c_0 - \sqrt{4(R_0 - 1) - c_0^2})/2$. Thus it has two independent solutions of the following form:

$$\begin{aligned} \overline{\mu_{11}} &= e^{(c_0/2)\xi} \cos\left(\frac{\sqrt{4(R_0 - 1) - c_0^2}}{2}\xi\right), \\ \overline{\mu_{22}} &= e^{(c_0/2)\xi} \sin\left(\frac{\sqrt{4(R_0 - 1) - c_0^2}}{2}\xi\right). \end{aligned} \quad (109)$$

Similar to the proof of Theorem 11, we can get that, when $\xi \rightarrow -\infty$, $\chi_2(\xi)$ can be described as the following equation:

$$\begin{aligned}\chi_2(\xi) &= K_1 e^{(c_0/2)\xi} \cos\left(\frac{\sqrt{4(R_0-1)-c_0^2}}{2}\xi\right) \\ &\quad + K_2 e^{(c_0/2)\xi} \sin\left(\frac{\sqrt{4(R_0-1)-c_0^2}}{2}\xi\right) + \text{h.o.t.}, \\ &= \sqrt{K_1^2 + K_2^2} e^{(c_0/2)\xi} \sin\left(\frac{\sqrt{4(R_0-1)-c_0^2}}{2}\xi + \tau(\xi)\right) \\ &\quad + \text{h.o.t.},\end{aligned}\tag{110}$$

where $\tan(\tau(\xi)) = K_1/K_2$, and h.o.t is the short notation for the higher order terms.

That is to say, $l_2(\xi)$ is oscillating. Thus, any solution of model (24) with $0 < c < 2\sqrt{R_0-1}$ is not strictly monotone. $\square$

Theorems 12 and 13 indicate that $c = 2\sqrt{R_0-1}$ is the critical minimal wave speed.

Acknowledgments

The authors would like to thank the anonymous referee for the very helpful suggestions and comments which led to improvements of our original paper. This research was supported by the Natural Science Foundation of Zhejiang Province (LY12A01014, R1110261, and LQ12A01009), the National Science Foundation of China (61272018), and the National Basic Research Program of China (2012CB426510).

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