Research Article

On the Number of Limit Cycles of a Piecewise Quadratic Near-Hamiltonian System

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This paper is concerned with the problem for the maximal number of limit cycles for a quadratic piecewise near-Hamiltonian system. By using the method of the first order Melnikov function, we find that it can have 8 limit cycles.

1. Introduction and Main Results

Recently, piecewise smooth systems attracted many researchers' attention since many real processes and different modern devices can be modeled by them; see [1, 2] and references therein. Due to their nonsmoothness, these systems can have richer dynamical phenomena than the smooth ones (see [3–5] and the references cited therein). For instance, nonsmooth system can have the sliding phenomena and in [6, 7] Giannakopoulos and Pliete have studied the existence of sliding cycles and sliding homoclinic cycles for a planar relay control feedback systems.

As we have seen, studying the existence and number of limit cycles is one of the main problems for piecewise smooth systems; see [8–16]. Limit cycles of piecewise smooth linear differential systems defined on two half-planes separated by a straight line $x = 0$ or $y = 0$ have been studied recently in [8–11], from which one can find that 3 limit cycles can appear for piecewise smooth linear systems. From [12–14], we can know that piecewise linear and quadratic near-Hamiltonian systems can have 2 and 3 limit cycles, respectively. In fact, these papers investigated the problem for limit cycle bifurcations of piecewise linear Hamiltonian systems under piecewise polynomial perturbations of degree n , obtaining some new and interesting results. Furthermore, the authors of [16] studied a piecewise quadratic Hamiltonian system (one side linear and another side quadratic) perturbed inside the class of piecewise polynomial differential systems of degree n and achieved that it can have $2n + [(n + 1)/2]$

limit cycles, which implies that piecewise quadratic near-Hamiltonian systems can have 5 limit cycles. The authors of [15] also obtained this result by perturbed quadratic polynomial differential systems containing an isochronous center with piecewise quadratic polynomials of degree 2.

In this paper, we mainly study the problem for the maximal number of a quadratic near-Hamiltonian system. That is to say, we consider a system of the form

$$\dot{x} = y + \varepsilon p(x, y), \quad \dot{y} = -g(x) + \varepsilon q(x, y), \quad (1)$$

where $\varepsilon > 0$ is a small parameter,

$$g(x) = \begin{cases} a_2 x^2 + a_1 x + a_0, & x < 0, \\ b_2 x^2 + b_1 x + b_0, & x > 0 \end{cases} \quad (2)$$

with $a_2 b_2 \neq 0$, and

$$p(x, y) = \begin{cases} \sum_{i+j=0}^2 a_{ij}^- x^i y^j, & x < 0, \\ \sum_{i+j=0}^2 a_{ij}^+ x^i y^j, & x > 0, \end{cases} \quad (3)$$

$$q(x, y) = \begin{cases} \sum_{i+j=0}^2 b_{ij}^- x^i y^j, & x < 0, \\ \sum_{i+j=0}^2 b_{ij}^+ x^i y^j, & x > 0. \end{cases}$$

Then, our main results can be stated as follows.

FIGURE 1: Level curves of $(1/2)y^2 + (1/3)ax^3 + (1/2)bx^2 + cx = h$.

Theorem 1. *There exists a system of form (1) which has 8 limit cycles.*

In the next section, the proof of Theorem 1 is presented.

2. Proof of Theorem 1

In this section, we divide the proof into two subsections: one section is to present some preliminary knowledge, which is useful to verify our main result; another section is to give the process of the proof of Theorem 1.

2.1. Preliminaries. Clearly, system (1) has the following two subsystems:

$$\dot{x} = y + \varepsilon \sum_{i+j=0}^2 a_{ij}^- x^i y^j, \quad (4a)$$

$$\dot{y} = -(a_2 x^2 + a_1 x + a_0) + \varepsilon \sum_{i+j=0}^2 b_{ij}^- x^i y^j,$$

$$\dot{x} = y + \varepsilon \sum_{i+j=0}^2 a_{ij}^+ x^i y^j, \quad (4b)$$

$$\dot{y} = -(b_2 x^2 + b_1 x + b_0) + \varepsilon \sum_{i+j=0}^2 b_{ij}^+ x^i y^j,$$

which are called left and right subsystems, respectively. And systems $(4a)|_{\varepsilon=0}$ and $(4b)|_{\varepsilon=0}$ are Hamiltonian with Hamiltonian functions, respectively,

$$\begin{aligned} H^-(x, y) &= \frac{1}{2}y^2 + \frac{1}{3}a_2 x^3 + \frac{1}{2}a_1 x^2 + a_0 x, \\ H^+(x, y) &= \frac{1}{2}y^2 + \frac{1}{3}b_2 x^3 + \frac{1}{2}b_1 x^2 + b_0 x. \end{aligned} \quad (5)$$

For systems $(4a)|_{\varepsilon=0}$ and $(4b)|_{\varepsilon=0}$, we make the following assumption:

$$(H) \quad a_1^2 > 4a_2 a_0, \quad b_1^2 > 4b_2 b_0.$$

Then, under the condition (H), for $\varepsilon = 0$, system $(4a)$ (resp., $(4b)$) has a hyperbolic saddle $S_1 = (\alpha_{S_1}, 0)$ (resp., $S_2 = (\alpha_{S_2}, 0)$) and an elementary center $C_1 = (\alpha_{C_1}, 0)$ (resp., $C_2 = (\alpha_{C_2}, 0)$), where

$$\begin{aligned} \alpha_{S_1} &= \frac{-a_1 - \sqrt{a_1^2 - 4a_2 a_0}}{2a_2}, & \alpha_{S_2} &= \frac{-b_1 - \sqrt{b_1^2 - 4b_2 b_0}}{2b_2}, \\ \alpha_{C_1} &= \frac{-a_1 + \sqrt{a_1^2 - 4a_2 a_0}}{2a_2}, & \alpha_{C_2} &= \frac{-b_1 + \sqrt{b_1^2 - 4b_2 b_0}}{2b_2}. \end{aligned} \quad (6)$$

We will assume that

$$a_2 < 0, \quad (7)$$

since, for $a_2 > 0$, we can make a variable transformation $x \rightarrow -x$, $y \rightarrow y$, together with a time rescaling $t \rightarrow -t$ such that the resulting system satisfies (7). Further, note that there exist 2 possible cases for level curves of $H(x, y) = (1/2)y^2 + (1/3)ax^3 + (1/2)bx^2 + cx = h$ on the plane, where $a \neq 0$, $b^2 > 4ac$; see Figure 1.

We remark that, in Figure 1, $C = ((-b + \sqrt{b^2 - 4ac})/2a, 0)$, $S = ((-b - \sqrt{b^2 - 4ac})/2a, 0)$. Then, one can obtain 9 possible phase portraits of system $(1)|_{\varepsilon=0}$ under (H) and (7) when it has a generalized closed orbit and two hyperbolic saddles; see Figure 2. For system $(1)|_{\varepsilon=0}$, a generalized closed orbit passes through the y -axis two times.

Now, we study limit cycle bifurcations on system (1). For simplicity, we consider the case

$$\alpha_{S_1} < 0 < \alpha_{S_2} < \alpha_{C_2}, \quad H^-(S_1) = H^+(S_2) \triangleq h_1; \quad (8)$$

 FIGURE 2: Possible portraits of $(1)|_{\varepsilon=0}$ with two saddles and a generalized closed orbit.

that is,

$$\begin{aligned}
 a_1^2 &> 4a_2a_0, & b_1^2 &> 4b_2b_0, & a_2 &< 0, & b_2 &> 0, \\
 a_1 &< 0, & b_1 &< 0, & a_0 &< 0, & b_0 &> 0, \\
 \frac{a_1^3 - 6a_0a_1a_2 + (a_1^2 - 4a_0a_2)^{3/2}}{a_2^2} & & & & & & & \\
 &= \frac{b_1^3 - 6b_0b_1b_2 + (b_1^2 - 4b_0b_2)^{3/2}}{b_2^2}.
 \end{aligned}$$

Then, the phase portrait of system $(1)|_{\varepsilon=0}$ is Figure 2(h), and there are four families of periodic orbits given by

$$\begin{aligned}
 L_{1h} &= \{(x, y) \mid H^-(x, y) = h, x_0^- < x < \alpha_{S_1}\}, & h_0^- < h < h_1, \\
 L_{2h} &= \{(x, y) \mid H^+(x, y) = h, \alpha_{S_2} < x < x_0^+\}, & h_0^+ < h < h_1,
 \end{aligned}$$

$$\begin{aligned}
 L_{3h} &= L_{3h}^- \cup L_{3h}^+ = \{(x, y) \mid H^-(x, y) = h, \alpha_{S_1} < x < 0\} \\
 &\quad \cup \{(x, y) \mid H^+(x, y) = h, 0 < x < \alpha_{S_2}\}, \\
 &\quad 0 < h < h_1, \\
 L_h &= L_h^- \cup L_h^+ = \{(x, y) \mid H^-(x, y) = h, x < 0\} \\
 &\quad \cup \{(x, y) \mid H^+(x, y) = h, x > 0\}, \quad h > h_1,
 \end{aligned}$$

where $H^-(x_0^-, 0) = h_1$, $H^+(x_0^+, 0) = h_1$ and $H^-(C_1) = h_0^-$, $H^+(C_2) = h_0^+$. Further, we have

$$\begin{aligned}
 \lim_{h \rightarrow h_1} L_{1h} &= L_1, & \lim_{h \rightarrow h_0^-} L_{1h} &= C_1, & \lim_{h \rightarrow h_1} L_{2h} &= L_2, \\
 \lim_{h \rightarrow h_0^+} L_{1h} &= C_2, & \lim_{h \rightarrow h_1} L_{3h}^- &= L^-,
 \end{aligned}$$

FIGURE 3: Phase portrait of system (1)|_{ε=0} under (9).

$$\begin{aligned} \lim_{h \rightarrow h_1} L_{3h}^+ &= L^+, & \lim_{h \rightarrow 0} L_{3h}^- &= \lim_{h \rightarrow 0} L_{3h}^+ = (0, 0), \\ \lim_{h \rightarrow h_1} L_h^- &= L_1 \cup L^-, & \lim_{h \rightarrow h_1} L_h^+ &= L_2 \cup L^+; \end{aligned} \quad (11)$$

see Figure 3. From (11), one can know that, as h approaches h_1 , the limit of L_h is $L_1 \cup L^- \cup L_2 \cup L^+$ denoted by L , which is a generalized polycycle with two hyperbolic saddles.

Then, from Theorem 1.1 of [17], for system (1), we have four first order Melnikov functions as follows:

$$\begin{aligned} M_1(h) &= \oint_{L_{1h}} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy, \\ &h \in (h_0^-, h_1), \\ M_2(h) &= \oint_{L_{2h}} \sum_{i+j=0}^2 b_{ij}^+ x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^+ x^i y^j dy, \\ &h \in (h_0^+, h_1), \\ M_3(h) &= M_3^-(h) + M_3^+(h) = \int_{L_{3h}^-} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx \\ &- \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy + \int_{L_{3h}^+} \sum_{i+j=0}^2 b_{ij}^+ x^i y^j dx \\ &- \sum_{i+j=0}^2 a_{ij}^+ x^i y^j dy, \quad h \in (0, h_1), \end{aligned}$$

$$\begin{aligned} M(h) &= M^-(h) + M^+(h) = \int_{L_h^-} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx \\ &- \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy + \int_{L_h^+} \sum_{i+j=0}^2 b_{ij}^+ x^i y^j dx \\ &- \sum_{i+j=0}^2 a_{ij}^+ x^i y^j dy, \quad h \in (h_1, +\infty), \end{aligned} \quad (12)$$

where $L_{1h}, L_{2h}, L_{3h}^-, L_{3h}^+, L_h^-$, and L_h^+ are given in (10).

Now, we deduce the expansion of $M_i(h)$, $i = 1, 2, 3$, and $M(h)$ near $h = h_1$ in the above. First, we have the following.

Lemma 2. The function $M_3^-(h)$ in (12) has the following expression:

$$\begin{aligned} M_3^-(h) &= c_0 + c_1(h - h_1) \ln(h_1 - h) + c_2(h - h_1) \\ &+ c_3(h - h_1)^2 \ln(h_1 - h) + O((h - h_1)^2), \quad (13) \\ &0 < h_1 - h \ll 1, \end{aligned}$$

where

$$\begin{aligned} c_0 &= \int_{L^-} \left[(a_{10}^- + b_{01}^-) y + (b_{11}^- + 2a_{20}^-) xy \right. \\ &\quad \left. + \frac{1}{2} (a_{11}^- + 2b_{02}^-) y^2 \right] dx - 2a_{00}^- \sqrt{2h_1} \\ &\quad - \frac{4}{3} h_1 \sqrt{2h_1} a_{02}^-, \\ c_1 &= -\frac{1}{(a_1^2 - 4a_0a_2)^{1/4}} \left[(a_{10}^- + b_{01}^-) + \alpha_{S_1} (2a_{20}^- + b_{11}^-) \right], \\ c_2|_{c_1=0} &= \int_{L^-} [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-) x + (a_{11}^- + 2b_{02}^-) y] dt \\ &\quad - \frac{\sqrt{2h_1}}{h_1} a_{00}^- - 2\sqrt{2h_1} a_{02}^-, \\ c_3 &= \frac{a_2}{2(a_1^2 - 4a_0a_2)^{5/4}} (2a_{20}^- + b_{11}^-). \end{aligned} \quad (14)$$

Proof. Denote by U a disk of diameter $\epsilon_0 > 0$ with its center at S_1 . Then, for $\epsilon_0 > 0$ small, rewrite $M_3^-(h) = I_1(h) + I_2(h)$, where

$$\begin{aligned} I_1(h) &= \int_{\bar{L}_{3h}} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy, \\ I_2(h) &= \int_{\bar{L}_{3h}^-} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy \end{aligned} \quad (15)$$

with

$$\bar{L}_{3h} = L_{3h}^- \cap U, \quad \bar{L}_{3h}^- = L_{3h}^- \setminus \bar{L}_{3h}. \quad (16)$$

Note that $I_2(h)$ is a C^∞ function for $h_1 - h > 0$ small. Thus, it suffices to investigate the expansion of $I_1(h)$ for $0 < h_1 - h \ll 1$. To do this, we make a variable transformation

$$\begin{aligned} u &= (a_1^2 - 4a_0a_2)^{1/4} (x - \alpha_{S_1}), \quad v = y, \\ \tau &= (a_1^2 - 4a_0a_2)^{1/4} t. \end{aligned} \quad (17)$$

Then system (4a) becomes

$$\begin{aligned} \dot{u} &= v + \varepsilon \bar{p}^-(u, v), \\ \dot{v} &= u - \frac{a_2 u^2}{(a_1^2 - 4a_0a_2)^{3/4}} + \varepsilon \bar{q}^-(u, v), \end{aligned} \quad (18)$$

where

$$\bar{p}^-(u, v) = \sum_{i+j=0}^2 \bar{a}_{ij}^- u^i v^j = \sum_{i+j=0}^2 \bar{a}_{ij}^- \left[\frac{u}{(a_1^2 - 4a_0a_2)^{1/4}} + \alpha_{S_1} \right]^i v^j, \quad (19)$$

$$\begin{aligned} \bar{q}^-(u, v) &= \sum_{i+j=0}^2 \bar{b}_{ij}^- u^i v^j \\ &= \sum_{i+j=0}^2 \frac{\bar{b}_{ij}^-}{(a_1^2 - 4a_0a_2)^{1/4}} \left[\frac{u}{(a_1^2 - 4a_0a_2)^{1/4}} + \alpha_{S_1} \right]^i v^j. \end{aligned} \quad (20)$$

For $\varepsilon = 0$, the Hamiltonian function of (18) has the form

$$\bar{H}(u, v) = \frac{1}{2} (v^2 - u^2) + \frac{a_2}{3(a_1^2 - 4a_0a_2)^{3/4}} u^3 + h_1 \quad (21)$$

and the corresponding first order Melnikov function can be denoted by $\bar{M}(h)$. By Remark 3.1.4 of [18], we can know that $M_3^-(h) = \bar{M}(h)$, and then, from Corollary 3.2.2 and the Theorem 3.2.1 of [18], I_1 can be written as

$$\begin{aligned} I_1(h) &= c_1 (h - h_1) \ln (h_1 - h) + c_3 (h - h_1)^2 \ln (h_1 - h) \\ &\quad + O((h - h_1)^2), \quad 0 < h_1 - h \ll 1, \end{aligned} \quad (22)$$

where

$$\begin{aligned} c_1 &= -(\bar{a}_{10}^- + \bar{b}_{01}^-) = -\frac{1}{(a_1^2 - 4a_0a_2)^{1/4}} \\ &\quad \times [(a_{10}^- + b_{01}^-) + \alpha_{S_1} (2a_{20}^- + b_{11}^-)], \\ c_3 &= \frac{3}{2} \times \frac{a_2}{3(a_1^2 - 4a_0a_2)^{3/4}} \times \frac{1}{(a_1^2 - 4a_0a_2)^{2/4}} (2\bar{a}_{20}^- + \bar{b}_{11}^-) \\ &= \frac{a_2}{2(a_1^2 - 4a_0a_2)^{5/4}} (2a_{20}^- + b_{11}^-). \end{aligned} \quad (23)$$

Thus, by the above discussion, one can obtain the expression of (13) and the formulas for c_1 and c_3 . Now, we give the expressions of c_0 and c_2 in (13).

On the curve L_{3h} intersect y -axis with two points denoted by $A_1(h) = (0, -\sqrt{2h})$ and $A(h) = (0, \sqrt{2h})$. Then, from (12), we have

$$\begin{aligned} M_3^-(h) &= \int_{\widehat{A_1(h)A(h)}} \sum_{i+j=0}^2 \bar{b}_{ij}^- x^i y^j dx - \sum_{i+j=0}^2 \bar{a}_{ij}^- x^i y^j dy \\ &= N_1(h) + N_2(h), \end{aligned} \quad (24)$$

where

$$\begin{aligned} N_1(h) &= \int_{\widehat{A_1(h)A(h) \cup \overline{A(h)A_1(h)}}} \sum_{i+j=0}^2 \bar{b}_{ij}^- x^i y^j dx \\ &\quad - \sum_{i+j=0}^2 \bar{a}_{ij}^- x^i y^j dy, \end{aligned} \quad (25)$$

$$N_2(h) = \int_{\widehat{A(h)A_1(h)}} (a_{00}^- + a_{01}^- y + a_{02}^- y^2) dy.$$

By Green's formula two times to $N_1(h)$, we obtain

$$\begin{aligned} N_1(h) &= \int_{\text{Int}(\widehat{A_1(h)A(h) \cup \overline{A(h)A_1(h)}})} [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-) x \\ &\quad + (a_{11}^- + 2b_{02}^-) y] dx dy \\ &= \oint_{\widehat{A_1(h)A(h) \cup \overline{A(h)A_1(h)}}} [(a_{10}^- + b_{01}^-) y + (b_{11}^- + 2a_{20}^-) xy \\ &\quad + \frac{1}{2} (a_{11}^- + 2b_{02}^-) y^2] dx \\ &= \int_{L_{3h}^-} [(a_{10}^- + b_{01}^-) y + (b_{11}^- + 2a_{20}^-) xy \\ &\quad + \frac{1}{2} (a_{11}^- + 2b_{02}^-) y^2] dx. \end{aligned} \quad (26)$$

Further, one has

$$\begin{aligned} N_2(h) &= \int_{\widehat{A(h)A_1(h)}} (a_{00}^- + a_{01}^- y + a_{02}^- y^2) dy \\ &= \int_{\sqrt{2h}}^{-\sqrt{2h}} (a_{00}^- + a_{01}^- y + a_{02}^- y^2) dy \\ &= -2a_{00}^- \sqrt{2h} - \frac{4}{3} h \sqrt{2h} a_{02}^-. \end{aligned} \quad (27)$$

Thus, we have

$$\begin{aligned} M_3^-(h) &= \int_{L_{3h}^-} [(a_{10}^- + b_{01}^-) y + (b_{11}^- + 2a_{20}^-) xy \\ &\quad + \frac{1}{2} (a_{11}^- + 2b_{02}^-) y^2] dx \\ &\quad - 2a_{00}^- \sqrt{2h} - \frac{4}{3} h \sqrt{2h} a_{02}^-. \end{aligned} \quad (28)$$

Then, let $h \rightarrow h_1$. Then, by (11), we can obtain the expression of c_0 in (13). Further, differentiating (13) with respect to h , together with Lemma 3.1.2 of [18], we obtain

$$\begin{aligned} \frac{d}{dh} M_3^-(h) &= \int_{L_{3h}^-} (p_x + q_y)|_{\varepsilon=0} dt \\ &= \int_{L_{3h}^-} [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-)x \\ &\quad + (a_{11}^- + 2b_{02}^-)y] dt - \frac{\sqrt{2h}}{h} a_{00}^- - 2\sqrt{2h} a_{02}^- \\ &= c_1 \ln(h_1 - h) + c_1 + c_2 \\ &\quad + O((h - h_1) \ln(h_1 - h)), \end{aligned} \quad (29)$$

which, as $c_1 = 0$, gives that

$$\begin{aligned} &\int_{L_{3h}^-} [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-)x + (a_{11}^- + 2b_{02}^-)y] dt \\ &\quad - \frac{\sqrt{2h}}{h} a_{00}^- - 2\sqrt{2h} a_{02}^- \\ &= c_2 + O((h - h_1) \ln(h_1 - h)). \end{aligned} \quad (30)$$

Let $h \rightarrow h_1$ in the above. Then we can easily obtain the formula for c_2 in (13). This ends the proof. $\square$

In the same way, one can obtain the expressions of $M_3^+(h)$.

Lemma 3. For $M_3^-(h)$ in (12), one has

$$\begin{aligned} M_3^+(h) &= \tilde{c}_0 + \tilde{c}_1 (h - h_1) \ln(h_1 - h) + \tilde{c}_2 (h - h_1) \\ &\quad + \tilde{c}_3 (h - h_1)^2 \ln(h_1 - h) + O((h - h_1)^2), \\ &\quad 0 < h_1 - h \ll 1, \end{aligned} \quad (31)$$

where

$$\begin{aligned} \tilde{c}_0 &= \int_{L^+} \left[(a_{10}^+ + b_{01}^+)y + (b_{11}^+ + 2a_{20}^+)xy \right. \\ &\quad \left. + \frac{1}{2} (a_{11}^+ + 2b_{02}^+)y^2 \right] dx + 2\sqrt{2h_1} a_{00}^+ \\ &\quad + \frac{4}{3} h_1 \sqrt{2h_1} a_{02}^+, \\ \tilde{c}_1 &= -\frac{1}{(b_1^2 - 4b_0 b_2)^{1/4}} [(a_{10}^+ + b_{01}^+) + \alpha_{S_2} (2a_{20}^+ + b_{11}^+)], \\ \tilde{c}_2|_{\tilde{c}_1=0} &= \int_{L^+} [(a_{10}^+ + b_{01}^+) + (b_{11}^+ + 2a_{20}^+)x + (a_{11}^+ + 2b_{02}^+)y] dt \\ &\quad + \frac{\sqrt{2h_1}}{h_1} a_{00}^+ + 2\sqrt{2h_1} a_{02}^+, \\ \tilde{c}_3 &= \frac{b_2}{2(b_1^2 - 4b_0 b_2)^{5/4}} (2a_{20}^+ + b_{11}^+). \end{aligned} \quad (32)$$

By applying Theorem 3.2.3 of [18], together with the proof of Lemma 2, we have the following.

Lemma 4. The functions $M_i(h)$, $i = 1, 2$, given in (12) can be expressed as

$$\begin{aligned} M_1(h) &= c_{01} + c_1 (h - h_1) \ln(h_1 - h) + c_{21} (h - h_1) \\ &\quad + c_3 (h - h_1)^2 \ln(h_1 - h) \\ &\quad + O((h - h_1)^2), \quad 0 < h_1 - h \ll 1, \\ M_2(h) &= c_{02} + \tilde{c}_1 (h - h_1) \ln(h_1 - h) + c_{22} (h - h_1) \\ &\quad + \tilde{c}_3 (h - h_1)^2 \ln(h_1 - h) \\ &\quad + O((h - h_1)^2), \quad 0 < h_1 - h \ll 1, \end{aligned} \quad (33)$$

where $c_1, c_3, \tilde{c}_1, \tilde{c}_3$ are given in Lemmas 2 and 3, and

$$\begin{aligned} c_{01} &= \oint_{L_1} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy, \\ c_{02} &= \oint_{L_2} \sum_{i+j=0}^2 b_{ij}^+ x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^+ x^i y^j dy, \\ c_{21}|_{c_1=0} &= \oint_{L_1} [(a_{10}^- + b_{01}^-) + (2a_{20}^- + b_{11}^-)x \\ &\quad + (a_{11}^- + 2b_{02}^-)y] dt|_{c_1=0}, \\ c_{22}|_{\tilde{c}_1=0} &= \oint_{L_2} [(a_{10}^+ + b_{01}^+) + (2a_{20}^+ + b_{11}^+)x \\ &\quad + (a_{11}^+ + 2b_{02}^+)y] dt|_{\tilde{c}_1=0}. \end{aligned} \quad (34)$$

By above three Lemmas 2, 3, and 4, together with (12), one can derive the following.

Lemma 5. For $M_3(h)$ and $M(h)$, one has

$$\begin{aligned} M_3(h) &= (c_0 + \tilde{c}_0) + (c_1 + \tilde{c}_1) (h - h_1) \ln(h_1 - h) \\ &\quad + (c_2 + \tilde{c}_2) (h - h_1) + (c_3 + \tilde{c}_3) (h - h_1)^2 \ln(h_1 - h) \\ &\quad + (O(h_1 - h)^2), \quad 0 < h_1 - h \ll 1, \\ M(h) &= (c_0 + c_{01} + \tilde{c}_0 + c_{02}) + 2(c_1 + \tilde{c}_1) \\ &\quad \times (h - h_1) \ln(h - h_1) + (c_2 + c_{21} + \tilde{c}_2 + c_{22}) \\ &\quad \times (h - h_1) + 2(c_3 + \tilde{c}_3) (h - h_1)^2 \ln(h - h_1) \\ &\quad + (O(h_1 - h)^2), \quad 0 < h - h_1 \ll 1, \end{aligned} \quad (35)$$

where the corresponding coefficients are given in (14), (32), and (34).

Then, using Lemmas 4 and 5, we can have the following.

Theorem 6. Let (H) and (7) hold. Then if there exists δ_0 such that

$$\begin{aligned} c_{01}(\delta_0) &= c_{02}(\delta_0) = (c_0 + \tilde{c}_0)(\delta_0) = (c_1 + \tilde{c}_1)(\delta_0) \\ &= (c_2 + \tilde{c}_2)(\delta_0) = 0, \\ (c_3 + \tilde{c}_3)(\delta_0) &< 0, \quad c_1(\delta_0) < 0, \quad \tilde{c}_1(\delta_0) > 0, \\ (c_1 + \tilde{c}_1)(\delta_0) &= 0, \quad (c_{21} + \tilde{c}_{22})(\delta_0) = 0, \\ \text{rank } \frac{\partial(c_{01}, c_{02}, c_0 + \tilde{c}_0, c_1 + \tilde{c}_1, c_2 + \tilde{c}_2)}{\partial \delta} &= 5, \end{aligned} \quad (36)$$

then system (1) can have 7 limit cycles near L .

Proof. By the assumption and from the expansions in Lemmas 4 and 5 we have

$$\begin{aligned} M_1(h) &= c_1(\delta_0)(h - h_1) \ln(h_1 - h) + O(h_1 - h) < 0, \\ &\quad 0 < h_1 - h \ll 1, \\ M_2(h) &= \tilde{c}_1(\delta_0)(h - h_1) \ln(h_1 - h) + O(h - h_1) > 0, \\ &\quad 0 < h_1 - h \ll 1, \\ M_3(h) &= (c_3 + \tilde{c}_3)(\delta_0)(h - h_1)^2 \ln(h_1 - h) \\ &\quad + (O(h_1 - h)^2) > 0, \quad 0 < h_1 - h \ll 1, \\ M(h) &= 2(c_3 + \tilde{c}_3)(h - h_1)^2 \ln(h - h_1) + (O(h_1 - h)^2) > 0, \\ &\quad 0 < h - h_1 \ll 1. \end{aligned} \quad (37)$$

Note that $c_{01}, c_{02}, c_0 + \tilde{c}_0, c_1 + \tilde{c}_1$, and $c_2 + \tilde{c}_2$ can be taken as free parameters. Then we can vary them one by one to change the signs of M_1, M_2, M_3 , and M .

First, let $c_{01} = c_{02} = c_0 + \tilde{c}_0 = c_1 + \tilde{c}_1 = 0$ and $0 < -(c_2 + \tilde{c}_2) \ll 1$. Then, in this case, we have

$$\begin{aligned} M_1(h) &< 0, \quad M_2(h) > 0, \quad M_3(h) > 0, \\ &\quad \text{for } 0 < h_1 - h \ll 1, \quad (38) \\ M(h) &< 0, \quad \text{for } 0 < h - h_1 \ll 1, \end{aligned}$$

which means that there exists a zero $\tilde{h}_1$ with $0 < \tilde{h}_1 - h_1$ such that $M(\tilde{h}_1) = 0$.

Second, fix $c_{01} = c_{02} = c_0 + \tilde{c}_0 = 0$ and vary $0 \ll -(c_1 + \tilde{c}_1) \ll -(c_2 + \tilde{c}_2)$, so that

$$\begin{aligned} M_1(h) &< 0, \quad M_2(h) > 0, \quad M_3(h) < 0, \\ &\quad \text{for } 0 < h_1 - h \ll 1, \quad (39) \\ M(h) &> 0, \quad \text{for } 0 < h - h_1 \ll 1. \end{aligned}$$

Hence, there exist $h_1 - \tilde{h}_2 > 0, \tilde{h}_3 \in (h_1, \tilde{h}_1)$ such that $M_3(\tilde{h}_2) = 0, M(\tilde{h}_3) = 0$.

Third, let $c_{01} = c_0 + \tilde{c}_0 = 0, 0 < -c_{02} \ll -(c_1 + \tilde{c}_1)$, implying

$$\begin{aligned} M_1(h) &< 0, \quad M_2(h) < 0, \quad M_3(h) < 0, \\ &\quad \text{for } 0 < h_1 - h \ll 1, \quad (40) \\ M(h) &< 0, \quad \text{for } 0 < h - h_1 \ll 1. \end{aligned}$$

Then $M_2(\tilde{h}_4) = 0, M(\tilde{h}_5) = 0$ for some $\tilde{h}_4 < h_1, \tilde{h}_5 \in (h_1, \tilde{h}_3)$. Fourth, let $c_0 + \tilde{c}_0 = 0$ and $0 < c_{01} \ll -c_{02}$; we have

$$\begin{aligned} M_1(h) &> 0, \quad M_2(h) < 0, \quad M_3(h) < 0, \\ &\quad \text{for } 0 < h_1 - h \ll 1, \quad (41) \\ M(h) &< 0, \quad \text{for } 0 < h - h_1 \ll 1. \end{aligned}$$

So there exists $\tilde{h}_6 < h_1$ such that $M_1(\tilde{h}_6) = 0$.

Finally, vary $0 < c_0 + \tilde{c}_0 \ll c_{01}$. In this case,

$$\begin{aligned} M_1(h) &> 0, \quad M_2(h) < 0, \quad M_3(h) > 0, \\ &\quad \text{for } 0 < h_1 - h \ll 1, \quad (42) \\ M(h) &< 0, \quad \text{for } 0 < h - h_1 \ll 1. \end{aligned}$$

So there exists $\tilde{h}_7 \in (\tilde{h}_2, h_1)$ such that $M_3(\tilde{h}_7) = 0$.

In a word, we have proved that if

$$0 < c_0 + \tilde{c}_0 \ll c_{01} \ll -c_{02} \ll -(c_1 + \tilde{c}_1) \ll -(c_2 + \tilde{c}_2) \ll 1 \quad (43)$$

then M_1, M_2, M_3 , and M have one zero, one zero, two zeros, and three zeros which implies that seven limit cycles can appear in the neighbourhood of L . $\square$

2.2. Proofs. In this section, we will use Theorem 6 to prove our main result. For that purpose, let

$$\begin{aligned} a_2 &= -1, \quad a_1 = -\frac{3}{2}, \quad a_0 = -\frac{1}{2}, \\ b_2 &= 1, \quad b_1 = -\frac{11}{3}, \quad b_0 = \frac{41}{48}. \end{aligned} \quad (44)$$

Then, the Hamiltonian functions in (5) can be rewritten as

$$\begin{aligned} H^-(x, y) &= \frac{1}{2}y^2 - \frac{1}{3}x^3 - \frac{3}{4}x^2 - \frac{1}{2}x, \\ H^+(x, y) &= \frac{1}{2}y^2 + \frac{1}{3}x^3 - \frac{11}{6}x^2 + \frac{11}{48}x. \end{aligned} \quad (45)$$

And in this case

$$\begin{aligned} S_1 &= \left(-\frac{1}{2}, 0\right), \quad S_2 = \left(\frac{1}{4}, 0\right), \quad C_1 = (-1, 0), \\ C_2 &= \left(\frac{41}{12}, 0\right), \quad h_1 = \frac{5}{48}. \end{aligned} \quad (46)$$

Under (44), it is not hard to obtain the expressions for $c_1, \tilde{c}_1, c_3$, and $\tilde{c}_3$ in Lemmas 2 and 3.

Lemma 7. Under (44), for system (1), one has

$$\begin{aligned} c_1 &= -\sqrt{2} \left[(a_{10}^- + b_{01}^-) - \frac{1}{2} (2a_{20}^- + b_{11}^-) \right], \\ c_3 &= -\sqrt{2} (2a_{20}^- + b_{11}^-), \\ \tilde{c}_1 &= -\sqrt{\frac{6}{19}} \left[(a_{10}^+ + b_{01}^+) + \frac{1}{4} (2a_{20}^+ + b_{11}^+) \right], \\ \tilde{c}_3 &= \frac{1}{2} \left(\frac{6}{19} \right)^{5/2} (2a_{20}^+ + b_{11}^+). \end{aligned} \quad (47)$$

For the other expressions of coefficients in Lemmas 4 and 5, one has the following.

Lemma 8. Under (44), for system (1), one has

$$\begin{aligned} c_{01} &= \frac{3\sqrt{2}}{20} (a_{10}^- + b_{01}^-) - \frac{39\sqrt{2}}{280} (2a_{20}^- + b_{11}^-), \\ c_{21} &= -3\sqrt{2} (2a_{20}^- + b_{11}^-), \\ c_{02} &= \frac{361\sqrt{114}}{180} (a_{10}^+ + b_{01}^+) + \frac{29963\sqrt{114}}{5040} (2a_{20}^+ + b_{11}^+), \\ c_{22} &= \sqrt{114} (2a_{20}^+ + b_{11}^+), \\ c_0 &= \frac{3\sqrt{2}}{20} (a_{10}^- + b_{01}^-) - \left(\frac{39\sqrt{2}}{280} - \frac{5\sqrt{30}}{168} \right) (2a_{20}^- + b_{11}^-) \\ &\quad - \frac{\sqrt{30}}{6} a_{00}^- - \frac{5\sqrt{30}}{432} a_{02}^-, \\ \tilde{c}_0 &= \left(-\frac{35}{9}\sqrt{30} + \frac{361}{180}\sqrt{114} \right) (a_{10}^+ + b_{01}^+) \\ &\quad + \left(\frac{29963}{5040}\sqrt{114} - \frac{730}{63}\sqrt{30} \right) (2a_{20}^+ + b_{11}^+) + \frac{\sqrt{30}}{6} a_{00}^+ \\ &\quad + \frac{5\sqrt{30}}{432} a_{02}^+, \\ c_2 &= (\sqrt{30} - 3\sqrt{2}) (2a_{20}^- + b_{11}^-) - \frac{4\sqrt{30}}{5} a_{00}^- - \frac{\sqrt{30}}{6} a_{02}^-, \\ \tilde{c}_2 &= (-2\sqrt{30} + \sqrt{114}) (2a_{20}^+ + b_{11}^+) + \frac{4\sqrt{30}}{5} a_{00}^+ + \frac{\sqrt{30}}{6} a_{02}^+. \end{aligned} \quad (48)$$

Proof. Note that $c_{01} = \oint_{L_1} \sum_{i+j=0}^2 b_{ij}^- x^i y^j dx - \sum_{i+j=0}^2 a_{ij}^- x^i y^j dy$. Then, by Green's formula two times, we derive

$$\begin{aligned} c_{01} &= \int_{\text{Int}(L_1)} [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-) x \\ &\quad + (a_{11}^- + 2b_{02}^-) y] dx dy \\ &= \oint_{L_1} \left[(a_{10}^- + b_{01}^-) y + (b_{11}^- + 2a_{20}^-) xy \right. \\ &\quad \left. + \frac{1}{2} (a_{11}^- + 2b_{02}^-) y^2 \right] dx. \end{aligned} \quad (49)$$

On the other hand, along the curve L_1 , we have $(1/2)y^2 - (1/3)x^3 - (3/4)x^2 - (1/2)x = 5/48$, for $-(5/4) \leq x \leq -(1/2)$. Then, (49) can be represented as

$$\begin{aligned} c_{01} &= 2 \int_{-5/4}^{-1/2} [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-) x] \\ &\quad \times \sqrt{\frac{5}{24} + \frac{2}{3}x^3 + \frac{3}{2}x^2 + x} dx. \end{aligned} \quad (50)$$

By applying Maple, one can compute that

$$\begin{aligned} \int_{-5/4}^{-1/2} \sqrt{\frac{5}{24} + \frac{2}{3}x^3 + \frac{3}{2}x^2 + x} dx &= \frac{3\sqrt{2}}{40}, \\ \int_{-5/4}^{-1/2} x \sqrt{\frac{5}{24} + \frac{2}{3}x^3 + \frac{3}{2}x^2 + x} dx &= -\frac{39\sqrt{2}}{560}. \end{aligned} \quad (51)$$

Inserting the above into (50) gives the expression of c_{01} .

Recall that $c_1 = 0$ if and only if $a_{10}^- + b_{01}^- = (1/2)(2a_{20}^- + b_{11}^-)$. Thus, by (4a)| $_{\varepsilon=0}$, we rewrite $c_{21}|_{c_1=0}$ as

$$\begin{aligned} c_{21}|_{c_1=0} &= \oint_{L_1} \frac{[(1/2)(2a_{20}^- + b_{11}^-) + (2a_{20}^- + b_{11}^-)x + (a_{11}^- + 2b_{02}^-)y]}{y} dx \\ &= \frac{1}{2} (2a_{20}^- + b_{11}^-) \oint_{L_1} \frac{1+2x}{y} dx \\ &= (2a_{20}^- + b_{11}^-) \int_{-5/4}^{-1/2} \frac{1+2x}{\sqrt{(5/24) + (2/3)x^3 + (3/2)x^2 + x}} dx, \end{aligned} \quad (52)$$

which, by Maple, yields the conclusion for c_{02} . Using a similar way, one can derive the formulas for c_{02} and c_{22} . Now, we deduce the formulas for c_0 and c_2 .

Note that, on the curve $L^- : (1/2)y^2 - (1/3)x^3 - (3/4)x^2 - (1/2)x = (5/48)$, $-(1/2) < x < 0$, we have

$$\begin{aligned} c_0 &= 2 \int_{-1/2}^0 [(a_{10}^- + b_{01}^-) + (b_{11}^- + 2a_{20}^-) x] \\ &\quad \times \sqrt{\frac{5}{24} + \frac{2}{3}x^3 + \frac{3}{2}x^2 + x} dx \\ &\quad - \frac{\sqrt{30}}{6} a_{00}^- - \frac{5\sqrt{30}}{432} a_{02}^-. \end{aligned} \quad (53)$$

By Maple again, we obtain that

$$\begin{aligned} \int_{-1/2}^0 \sqrt{\frac{5}{24} + \frac{2}{3}x^3 + \frac{3}{2}x^2 + x} dx &= \frac{3\sqrt{2}}{40}, \\ \int_{-1/2}^0 x \sqrt{\frac{5}{24} + \frac{2}{3}x^3 + \frac{3}{2}x^2 + x} dx &= -\frac{39\sqrt{2}}{560} + \frac{5\sqrt{30}}{336}. \end{aligned} \quad (54)$$

Thus, we can obtain the conclusion for c_0 .

When $c_1 = 0$, we have $(a_{10}^- + b_{01}^-) = (1/2)(2a_{20}^- + b_{11}^-)$. Thus, c_2 in (14) can be rewritten as

$$\begin{aligned} c_2 &= \frac{1}{2} (2a_{20}^- + b_{11}^-) \int_{L^-} \frac{1+2x}{y} dx - \frac{2a_{00}^-}{\sqrt{2h_1}} \\ &\quad - \frac{4a_{02}^-}{3} \left(\sqrt{2h_1} + \frac{h_1}{\sqrt{2h_1}} \right) \Big|_{h_1=5/48} \\ &= (2a_{20}^- + b_{11}^-) \int_{-1/2}^0 \frac{1+2x}{\sqrt{(5/24) + (2/3)x^3 + (3/2)x^2 + x}} dx \\ &\quad - \frac{4}{5} \sqrt{6} a_{00}^- - \frac{\sqrt{30}}{6} a_{02}^-, \end{aligned} \quad (55)$$

which, by Maple, implies the conclusion. Further, we can prove $\tilde{c}_0$ and $\tilde{c}_2$ in the same way to c_0 and c_2 , respectively. This finishes the proof. $\square$

Proof of Theorem 1. Let

$$\begin{aligned} \delta_1 &= a_{10}^- + b_{01}^-, & \delta_2 &= 2a_{20}^- + b_{11}^-, & \delta_3 &= a_{10}^+ + b_{01}^+, \\ \delta_4 &= 2a_{20}^+ + b_{11}^+, & \delta_5 &= a_{00}^+ - a_{00}^-, & \delta_6 &= a_{02}^+ - a_{02}^-. \end{aligned} \quad (56)$$

Then, solve the equations

$$c_{01} = c_{02} = c_0 + \tilde{c}_0 = c_1 + \tilde{c}_1 = c_2 + \tilde{c}_2 = 0, \quad (57)$$

which, together with (56), gives that

$$\begin{aligned} \delta_1 &= \frac{13\sqrt{57}}{42} \delta_4, & \delta_2 &= \frac{\sqrt{57}}{3} \delta_4, & \delta_3 &= -\frac{83}{28} \delta_4, \\ \delta_5 &= \left(\frac{5\sqrt{57}}{42} - \frac{5}{7} \right) \delta_4, & \delta_6 &= \left(-\frac{18\sqrt{57}}{7} + \frac{108}{7} \right) \delta_4, \end{aligned} \quad (58)$$

where $\delta_4 > 0$. Then, under (58)

$$\begin{aligned} c_1 &= -\frac{\sqrt{114}}{7} \delta_4 < 0, & \tilde{c}_1 &= \frac{\sqrt{114}}{7} \delta_4 > 0, \\ c_3 + \tilde{c}_3 &= -\frac{6805}{20577} \sqrt{114} \delta_4 < 0, & c_1 + \tilde{c}_1 &= 0, \\ c_{21} + \tilde{c}_{22} &= 0. \end{aligned} \quad (59)$$

Further,

$$\text{rank} \frac{\partial (c_{01}, c_{02}, c_0 + \tilde{c}_0, c_1 + \tilde{c}_1, c_2 + \tilde{c}_2)}{\partial (\delta_1, \delta_2, \delta_3, \delta_5, \delta_6)} = 5. \quad (60)$$

Thus, by Theorem 6, there exist 7 limit cycles near L for $\varepsilon > 0$ small.

Under (58), we have for $0 < (5/48) - h \ll 1$

$$\begin{aligned} M_1(h) &= -\frac{\sqrt{114}}{7} \delta_4 \left(h - \frac{5}{48} \right) \ln \left(\frac{5}{48} - h \right) \\ &\quad + O \left(\left(h - \frac{5}{48} \right) \right) < 0, \\ M_2(h) &= \frac{\sqrt{114}}{7} \delta_4 \left(h - \frac{5}{48} \right) \ln \left(\frac{5}{48} - h \right) \\ &\quad + O \left(\left(h - \frac{5}{48} \right) \right) > 0. \end{aligned} \quad (61)$$

Now, we derive the expansion of $M_1(h)$. To achieve this, make the transformation

$$u = \frac{\sqrt{2}}{2} (x + 1), \quad v = y, \quad (62)$$

together with time scaling $\tau = (\sqrt{2}/2)t$. Then the system becomes

$$\begin{aligned} \dot{u} &= v + \varepsilon \bar{p}^-(u, v), \\ \dot{v} &= -u - 2\sqrt{2}u^2 + \varepsilon \bar{q}^-(u, v), \end{aligned} \quad (63)$$

where

$$\begin{aligned} \bar{p}^-(u, v) &= \sum_{i+j=0}^2 \bar{a}_{ij}^- u^i v^j = \sum_{i+j=0}^2 a_{ij}^- (\sqrt{2}u - 1)^i v^j, \\ \bar{q}^-(u, v) &= \sum_{i+j=0}^2 \bar{b}_{ij}^- u^i v^j = \sum_{i+j=0}^2 \sqrt{2} b_{ij}^- (\sqrt{2}u - 1)^i v^j. \end{aligned} \quad (64)$$

For $\varepsilon = 0$, the Hamiltonian function of (63) has the form

$$\bar{H}(u, v) = \frac{1}{2} (v^2 + u^2) + \frac{2\sqrt{2}}{3} u^3 + \frac{1}{12}. \quad (65)$$

Denote $\bar{M}_1(h)$ by the first Melnikov function of system (63) near C_1 , and note that

$$q^- dx - p^- dy = \left(\frac{\sqrt{2}}{2} \bar{q}^- \right) \frac{du}{(\sqrt{2}/2)} - \bar{p}^- dv = \bar{q}^- du - \bar{p}^- dv. \quad (66)$$

Then we have $\bar{M}_1(h) = M_1(h)$. From Theorem 2.4.1 of [18], we can obtain

$$\begin{aligned} M_1(h) &= \tilde{b}_0 \left(h - \frac{1}{12} \right) + O \left(\left(h - \frac{1}{12} \right)^2 \right) < 0, \\ 0 &< h - \frac{1}{12} \ll 1, \end{aligned} \quad (67)$$

where, under (58),

$$\begin{aligned} \tilde{b}_0 &= 2\pi (\bar{p}_u^- + \bar{q}_v^-) (0, 0) \\ &= 2\sqrt{2}\pi [(a_{10}^- + b_{01}^-) - (2a_{20}^- + b_{11}^-)] = 2\sqrt{2}\pi (\delta_1 - \delta_2) \\ &= -\frac{\sqrt{114}}{21} \pi \delta_4. \end{aligned} \quad (68)$$

Similarly, under (58) we have

$$M_2(h) = -\frac{\sqrt{114}}{228}\delta_4\left(h + \frac{1681}{324}\right) + O\left(\left(h - h + \frac{1681}{324}\right)^2\right) < 0,$$

$$0 < h + \frac{1681}{324} \ll 1. \quad (69)$$

This implies that there exists another zero between $h = -(1681/324)$ and $h = 5/48$ for $M_2(h)$. This ends the proof. $\square$

Conflict of Interests

The author declares that there is no conflict of interests regarding the publication of this paper.

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