

Research Article

Multiple Soliton Solutions of the Sawada-Kotera Equation with a Nonvanishing Boundary Condition and the Perturbed Korteweg de Vries Equation by Using the Multiple Exp-Function Scheme

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The Sawada-Kotera equation with a nonvanishing boundary condition, which models the evolution of steeper waves of shorter wavelength than those depicted by the Korteweg de Vries equation, is analyzed and also the perturbed Korteweg de Vries (pKdV) equation. For this goal, a capable method known as the multiple exp-function scheme (MEFS) is formally utilized to derive the multiple soliton solutions of the models. The MEFS as a generalization of Hirota's perturbation method actually suggests a systematic technique to handle nonlinear evolution equations (NLEEs).

1. Introduction

In the applied science, NLEEs are extensively used in theoretical studies to model a wide range of nonlinear phenomena. To comprehend the mechanisms of nonlinear phenomena, it is vital to investigate the solutions of NLEEs [1–15]. One specific tool that has recently achieved a special interest from academic researchers is the multiple exp-function scheme [16, 17]. The MEFS supposes that the multisoliton solutions of NLEEs can be presented as $u(x, t) = p/q$ in which p and q are polynomials of exponential functions. The SK equation with a nonvanishing boundary condition [18, 19]

$$u_t + a \left(\frac{a}{5b} u + 3u^2 + u_{xx} \right)_x + b \left(15u^3 + 15uu_{xx} + u_{xxxx} \right)_x = 0 \quad (1)$$

is one of NLEEs that models the evolution of steeper waves of shorter wavelength than those explained by the KdV equation

and its perturbed form. By using the binary-Bell-polynomial Hirota method and symbolic computation, the bilinear form and N-soliton solutions for this model were derived in [19]. The perturbed form of KdV equation [18, 20–22]

$$u_t + 6uu_x + u_{xxx} + \epsilon \left(30u^2u_x + 10uu_{xxx} + 20u_xu_{xx} + u_{xxxxx} \right) = 0, \quad (2)$$

is another kind of NLEEs describing some arrays of wave crests. The bilinear form, Bäcklund transformation, superposition formulae, and N-soliton solutions in terms of the Wronskian were done in [22].

For computational purposes, we use the transformation

$$u = v_x \quad (3)$$

to convert (1) and (2) into the following, respectively:

$$v_{xt} + a \left(\frac{a}{5b} v_{xx} + 6v_x v_{xx} + v_{xxx} \right) + b \left(45v_x^2 v_{xx} + 15v_x v_{xxx} + 15v_{xx} v_{xxx} + v_{xxxx} \right) = 0 \quad (4)$$

and

$$v_{xt} + 6v_x v_{xx} + v_{xxxx} + \epsilon \left(30v_x^2 v_{xx} + 10v_x v_{xxx} + 20v_{xx} v_{xxx} + v_{xxxx} \right) = 0. \quad (5)$$

The key goal of present work is applying the MEFS to generate the multiple soliton solutions of the models (1) and (2).

2. Multiple Exp-Function Method

The key steps of MEFS can be summarized as follows [16, 17].

Step 1. Let us consider the following (1 + 1)-dimensional NLEE:

$$P(x, t, u_x, u_t, \dots) = 0. \quad (6)$$

Step 2. Suppose the solution of above NLEE can be expressed as

$$u(x, t) = \frac{p(\eta_1, \eta_2, \dots, \eta_n)}{q(\eta_1, \eta_2, \dots, \eta_n)},$$

$$p = \sum_{r,s=1}^n \sum_{i,j=0}^M p_{rs,ij} \eta_r^i \eta_s^j, \quad (7)$$

$$q = \sum_{r,s=1}^n \sum_{i,j=0}^N q_{rs,ij} \eta_r^i \eta_s^j,$$

in which $p_{rs,ij}$ and $q_{rs,ij}$ are unknowns to be determined and

$$\eta_i = c_i e^{\xi_i},$$

$$\xi_i = k_i x - \omega_i t, \quad (8)$$

$$1 \leq i \leq n,$$

Step 3. Substituting (7) and its derivatives into (6) yields the following transformed equation:

$$Q(x, t, \eta_1, \eta_2, \dots, \eta_n) = 0. \quad (9)$$

Step 4. By setting the numerator of the function $Q(x, t, \eta_1, \eta_2, \dots, \eta_n)$ to zero, we will reach an algebraic system which its solution yields the multiple wave solution of (6) as

$$u(x, t) = \frac{p(c_1 e^{k_1 x - \omega_1 t}, \dots, c_n e^{k_n x - \omega_n t})}{q(c_1 e^{k_1 x - \omega_1 t}, \dots, c_n e^{k_n x - \omega_n t})}. \quad (10)$$

2.1. Multiple Soliton Solutions of SK Equation with a Nonvanishing Boundary Condition. In the current subsection, the multiple soliton solutions of SK equation with a nonvanishing boundary condition are derived through the MEFS.

2.1.1. One-Soliton Solution of (1). To obtain one-soliton solution, it is assumed

$$v(x, t) = \frac{p}{q},$$

$$p = A_1 e^{\theta_1}, \quad (11)$$

$$q = 1 + e^{\theta_1},$$

where A_1 is a constant and

$$\theta_1 = k_1 x - \omega_1 t \quad (12)$$

the dispersion is

$$\omega_1 = \frac{k_1 (5b^2 k_1^4 + 5abk_1^2 + a^2)}{5b}. \quad (13)$$

Now, applying the MEFS results in

$$A_1 = 2k_1, \quad (14)$$

one-soliton solution can be presented as

$$v(x, t) = \frac{2k_1 e^{k_1 x - \omega_1 t}}{1 + e^{k_1 x - \omega_1 t}} \quad (15)$$

where k_1 is arbitrary but ω_1 is defined by (13).

2.1.2. Two-Soliton Solution of (1). To seek two-soliton solution, the following ansatz is considered

$$v(x, t) = \frac{p}{q} \quad (16)$$

in which p and q are defined as

$$p = 2k_1 e^{k_1 x - \omega_1 t} + 2k_2 e^{k_2 x - \omega_2 t} + 2A_{12} (k_1 + k_2) e^{k_1 x - \omega_1 t} e^{k_2 x - \omega_2 t}, \quad (17)$$

$$q = 1 + e^{k_1 x - \omega_1 t} + e^{k_2 x - \omega_2 t} + A_{12} e^{k_1 x - \omega_1 t} e^{k_2 x - \omega_2 t}. \quad (18)$$

Now, by applying the MEFS, we acquire

$$A_{12} = \frac{5bk_1^4 - 15bk_1^3k_2 + 20bk_1^2k_2^2 - 15bk_1k_2^3 + 5bk_2^4 + 3ak_1^2 - 6ak_1k_2 + 3ak_2^2}{5bk_1^4 + 15bk_1^3k_2 + 20bk_1^2k_2^2 + 15bk_1k_2^3 + 5bk_2^4 + 3ak_1^2 + 6ak_1k_2 + 3ak_2^2},$$

$$\omega_1 = \frac{k_1(5b^2k_1^4 + 5abk_1^2 + a^2)}{5b},$$

$$\omega_2 = \frac{k_2(5b^2k_2^4 + 5abk_2^2 + a^2)}{5b}.$$
(19)

2.1.3. *Three-Soliton Solution of (1).* To derive three-soliton solution, it is assumed

$$v(x, t) = \frac{p}{q} \tag{20}$$

in which p and q are defined as

$$p = 2k_1e^{k_1x-\omega_1t} + 2k_2e^{k_2x-\omega_2t} + 2k_3e^{k_3x-\omega_3t}$$

$$+ 2A_{12}(k_1 + k_2)e^{-\omega_1t+k_1x}e^{-\omega_2t+k_2x}$$

$$+ 2A_{13}(k_1 + k_3)e^{k_1x-\omega_1t}e^{k_3x-\omega_3t} + 2A_{23}(k_2 + k_3)$$

$$\cdot e^{k_2x-\omega_2t}e^{k_3x-\omega_3t} + 2A_{12}A_{13}A_{23}(k_1 + k_2 + k_3)$$

$$\cdot e^{k_1x-\omega_1t}e^{k_2x-\omega_2t}e^{k_3x-\omega_3t},$$

$$q = 1 + e^{k_1x-\omega_1t} + e^{k_2x-\omega_2t} + e^{k_3x-\omega_3t}$$

$$+ A_{12}e^{k_1x-\omega_1t}e^{-\omega_2t+k_2x} + A_{13}e^{k_1x-\omega_1t}e^{k_3x-\omega_3t}$$

$$+ A_{23}e^{k_2x-\omega_2t}e^{k_3x-\omega_3t}$$

$$+ A_{12}A_{13}A_{23}e^{k_1x-\omega_1t}e^{k_2x-\omega_2t}e^{k_3x-\omega_3t}.$$
(22)

Now, applying the MEFS yields

$$A_{12} = \frac{5bk_1^4 - 15bk_1^3k_2 + 20bk_1^2k_2^2 - 15bk_1k_2^3 + 5bk_2^4 + 3ak_1^2 - 6ak_1k_2 + 3ak_2^2}{5bk_1^4 + 15bk_1^3k_2 + 20bk_1^2k_2^2 + 15bk_1k_2^3 + 5bk_2^4 + 3ak_1^2 + 6ak_1k_2 + 3ak_2^2},$$

$$A_{13} = \frac{5bk_1^4 - 15bk_1^3k_3 + 20bk_1^2k_3^2 - 15bk_1k_3^3 + 5bk_3^4 + 3ak_1^2 - 6ak_1k_3 + 3ak_3^2}{5bk_1^4 + 15bk_1^3k_3 + 20bk_1^2k_3^2 + 15bk_1k_3^3 + 5bk_3^4 + 3ak_1^2 + 6ak_1k_3 + 3ak_3^2},$$

$$A_{23} = \frac{5bk_2^4 - 15bk_2^3k_3 + 20bk_2^2k_3^2 - 15bk_2k_3^3 + 5bk_3^4 + 3ak_2^2 - 6ak_2k_3 + 3ak_3^2}{5bk_2^4 + 15bk_2^3k_3 + 20bk_2^2k_3^2 + 15bk_2k_3^3 + 5bk_3^4 + 3ak_2^2 + 6ak_2k_3 + 3ak_3^2},$$

$$\omega_1 = \frac{k_1(5b^2k_1^4 + 5abk_1^2 + a^2)}{5b},$$

$$\omega_2 = \frac{k_2(5b^2k_2^4 + 5abk_2^2 + a^2)}{5b},$$

$$\omega_3 = \frac{k_3(5b^2k_3^4 + 5abk_3^2 + a^2)}{5b}.$$
(23)

2.2. *Multiple Soliton Solutions of pKdV Equation (2).* In the present subsection, the multiple soliton solutions of pKdV equation are obtained through the MEFS.

2.2.1. *One-Soliton Solution of (2).* To obtain one-soliton solution, it is assumed

$$v(x, t) = \frac{p}{q},$$

$$p = A_1e^{\theta_1},$$

$$q = 1 + e^{\theta_1},$$
(24)

where A_1 is a constant and

$$\theta_1 = k_1x - \omega_1t$$
(25)

and the dispersion relation is

$$\omega_1 = \epsilon k_1^5 + k_1^3.$$
(26)

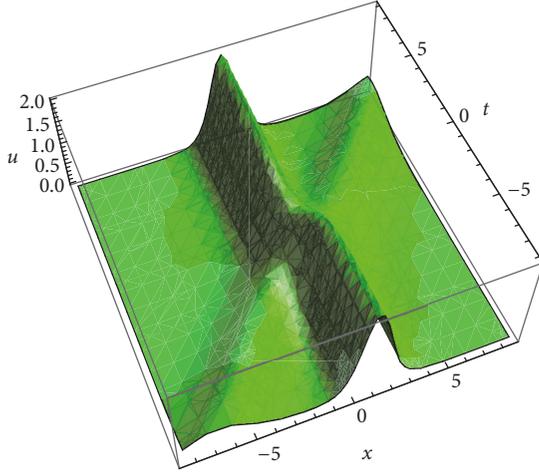


FIGURE 1: Evolution of the two-soliton solution (29).

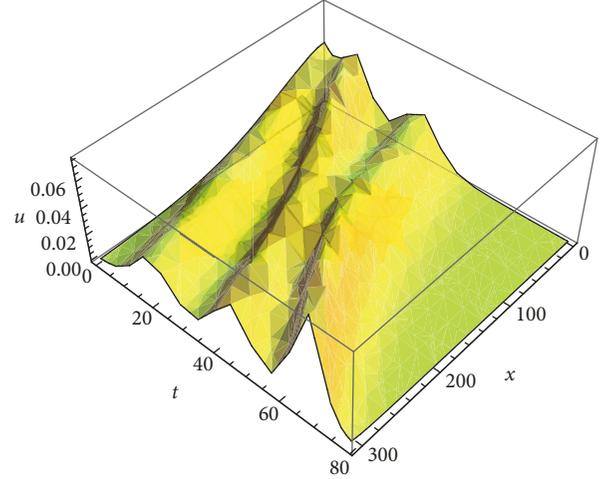


FIGURE 2: Evolution of the three-soliton solution (33).

Now, applying the MEFS results in

$$A_1 = 2k_1 \quad (27)$$

and, so, the resulting one-soliton solution can be presented as

$$v(x, t) = \frac{2k_1 e^{k_1 x - \omega_1 t}}{1 + e^{k_1 x - \omega_1 t}}, \quad (28)$$

where k_1 is arbitrary but ω_1 is defined by (26).

2.2.2. Two-Soliton Solution of (2). To seek two-soliton solution, the following ansatz is considered

$$v(x, t) = \frac{p}{q} \quad (29)$$

in which p and q are defined as

$$p = 2k_1 e^{k_1 x - \omega_1 t} + 2k_2 e^{k_2 x - \omega_2 t} + 2A_{12}(k_1 + k_2) e^{k_1 x - \omega_1 t} e^{k_2 x - \omega_2 t}, \quad (30)$$

$$q = 1 + e^{k_1 x - \omega_1 t} + e^{k_2 x - \omega_2 t} + A_{12} e^{k_1 x - \omega_1 t} e^{k_2 x - \omega_2 t}. \quad (31)$$

Now, by applying the MEFS, we acquire

$$A_{12} = \frac{k_1^2 - 2k_1 k_2 + k_2^2}{k_1^2 + 2k_1 k_2 + k_2^2}, \quad (32)$$

$$\omega_1 = \epsilon k_1^5 + k_1^3,$$

$$\omega_2 = \epsilon k_2^5 + k_2^3.$$

A profile of the evolution of the two-soliton solution (29) is given in Figure 1.

2.2.3. Three-Soliton Solution of (2). To derive three-soliton solution, it is assumed

$$v(x, t) = \frac{p}{q} \quad (33)$$

in which p and q are defined as

$$p = 2k_1 e^{k_1 x - \omega_1 t} + 2k_2 e^{k_2 x - \omega_2 t} + 2k_3 e^{k_3 x - \omega_3 t} + 2A_{12}(k_1 + k_2) e^{-\omega_1 t + k_1 x} e^{-\omega_2 t + k_2 x} + 2A_{13}(k_1 + k_3) e^{k_1 x - \omega_1 t} e^{k_3 x - \omega_3 t} + 2A_{23}(k_2 + k_3) e^{k_2 x - \omega_2 t} e^{k_3 x - \omega_3 t} + 2A_{12}A_{13}A_{23}(k_1 + k_2 + k_3) e^{k_1 x - \omega_1 t} e^{k_2 x - \omega_2 t} e^{k_3 x - \omega_3 t}, \quad (34)$$

$$q = 1 + e^{k_1 x - \omega_1 t} + e^{k_2 x - \omega_2 t} + e^{k_3 x - \omega_3 t} + A_{12} e^{k_1 x - \omega_1 t} e^{-\omega_2 t + k_2 x} + A_{13} e^{k_1 x - \omega_1 t} e^{k_3 x - \omega_3 t} + A_{23} e^{k_2 x - \omega_2 t} e^{k_3 x - \omega_3 t} + A_{12}A_{13}A_{23} e^{k_1 x - \omega_1 t} e^{k_2 x - \omega_2 t} e^{k_3 x - \omega_3 t}. \quad (35)$$

Now, applying the MEFS yields

$$A_{12} = \frac{k_1^2 - 2k_1 k_2 + k_2^2}{k_1^2 + 2k_1 k_2 + k_2^2},$$

$$A_{13} = \frac{k_1^2 - 2k_1 k_3 + k_3^2}{k_1^2 + 2k_1 k_3 + k_3^2},$$

$$A_{23} = \frac{k_2^2 - 2k_2 k_3 + k_3^2}{k_2^2 + 2k_2 k_3 + k_3^2}, \quad (36)$$

$$\omega_1 = \epsilon k_1^5 + k_1^3,$$

$$\omega_2 = \epsilon k_2^5 + k_2^3,$$

$$\omega_3 = \epsilon k_3^5 + k_3^3.$$

A profile of the evolution of the three-soliton solution (33) is given in Figure 2.

3. Concluding Remarks

We note that multiple soliton solutions of (1) and (2) are in agreement with [19, 21, 22]. Moreover, we should emphasize the approach employed here was independent of the bilinear forms, simplified Hirota method, Darboux transformation method, or Bell polynomial technique. Finally we can say the multiple exp-function algorithm is an elegant and versatile method that can be adopted to other NLEEs of mathematical physics.

Data Availability

The data used to support the findings of this study are included within the article.

Conflicts of Interest

The authors declare that they have no conflicts of interest.

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