Research Article

Application of Resonant Nuclear Reactions for Studying the Diffusion of Nitrogen and Silicon in Ti-Modified Stainless Steel

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1. Introduction

The structural materials like fuel clad and wrapper in sodium-cooled fast reactors are exposed to a hostile environment of intense fast neutron bombardment at elevated temperatures. One of the main problems which result on account of the radiation damage is the dimensional changes brought about by void swelling, leading to a decreased residence time of the fuel assemblies in the reactor core and lowering of fuel burnup. Therefore, resistance to void swelling is a major consideration in the choice of materials for the core components.

Improvement in the resistance to irradiation-induced void swelling can be achieved by the introduction of efficient traps for vacancies and helium such as dislocations and precipitate-matrix interfaces, by adjusting alloying elements and through thermomechanical treatments [1, 2]. The titanium-modified steel exhibits improved swelling resistance under breeder reactor conditions and, consequently, has become a prime candidate for structural applications [3]. The role of fast diffusing alloying elements, namely, Ti, Si, and P in suppressing void nucleation is known [4]. The fast diffusing species increases the effective vacancy diffusion coefficient, and the consequent reduction in the void nucleation rates plays a crucial role in imparting significant void swelling resistance to the alloy. In order to model such effects, knowledge of the diffusion behavior of constituent species is essential due to their considerable redistribution during irradiation. Therefore, thermal diffusion behavior of solutes like silicon and nitrogen in titanium-modified steels has been studied here.

Though some amount of studies on the effect of minor alloying elements on the self-diffusion of the constituent elements have been carried out [5], reports on the diffusion behavior of silicon in SS 316 are sparse, and to the best of our knowledge, there is hardly any literature describing the diffusion behavior of silicon in titanium-modified steels.

Also, studies concerning iron-based alloys and their variants with interstitial solute additions (C, N) continue to attract attention due to the rich technological interest in improving radiation resistance, mechanical and tribological properties. In nitriding processes involving ion implantation, an improvement in mechanical property is achieved due to the increased quantity of interstitial nitrogen and their interaction with crystal defects. There also exists an alternative conjecture that the trapping of vacancies with interstitial impurity atoms enhances the probability of recombination, hence aiding in reduction of void swelling. Thus, practical applications have necessitated a full understanding of the underlying physical processes involving interstitial solutes,
namely, the phase formation, diffusion, trapping with vacancies, dislocations, and other second-phase particles during irradiation. This paper portrays the application of the resonance nuclear reaction analysis (RNRA) technique in understanding the phenomena related to defect-impurity interaction (vacancy-nitrogen complexes), solute and impurity diffusion (Si and N thermal diffusion behaviour) in Ti-modified austenitic stainless steel.

2. Experimental Details

2.1. Sample Preparation and Ion Implantation. D9 is a titanium-modified austenitic stainless steel whose chemical composition is given in Table 1. The samples used for the experiments are spark cut from 20% cold-worked hexagonal wrappers. The samples were mechanically polished and subsequently subjected to solution annealing treatment which involves the heating of samples in vacuum at 1343 K for 30 minutes and subsequently at 1373 K for 5 minutes followed by furnace cooling. This solution annealing treatment ensures the samples to be free from the defect microstructure introduced during cold-work.

Solution-annealed D9 samples were implanted with 200 keV $^{30}$Si ions to a fluence of $3 \times 10^{16}$ atoms cm$^{-2}$ using a 1.7 MV tandem accelerator at IGCAR, Kalpakkam. For nitrogen isotope implantation, solution-annealed D9 samples were implanted with 30 keV $^{15}$N ions to a fluence of $5 \times 10^{15}$ ions cm$^{-2}$ using a low-energy ion implantation facility at GNS Science, New Zealand. Both silicon and nitrogen implantations were carried out at room temperature and at a vacuum level better than $\sim 1 \times 10^{-7}$ mbar.

2.2. Depth Profiling of Nitrogen and Silicon. The depth profiling of the implanted atoms has been carried out using narrow resonances of the nuclear reaction which are ideal for high depth resolution [6]. Depth profiling of $^{15}$N has been carried out using 429 keV resonance of the $^{15}$N($p,\alpha\gamma$) $^{12}$C reaction ($\Gamma \sim 120$ eV, $E_p = 4.43$ MeV) and that of $^{30}$Si using 620 keV resonance ($\Gamma \sim 68$ eV, $E_p = 7.9$ MeV) of the $^{30}$Si($p,\gamma$) $^{31}$P nuclear reaction. The applicability of these reactions in obtaining quantitative high-resolution depth profiles of nitrogen and silicon in a material is well known [7, 8]. The profiling experiments involve increasing incident proton energies beyond the resonance energy in steps of 1 keV, and collecting the gamma rays using a 3”$\times$3” NaI detector serves as the signal for depth profiling the implanted species. The gamma ray yield is proportional to the concentration of the implanted atoms at depths described by the incident energy. The relationship between the incident energy and the depth at which the reaction is taking place can be established by taking into account the energy loss of protons in the alloy. Since both reactions have a very narrow resonance width, the depth resolution for either reactions is estimated to be $\sim 3$ nm at the surface. The concentration of Si and N at the peak of the implantation depth profiles is $\sim 5$ at % and $\sim 3$ at %, respectively.

2.3. Diffusion of Nitrogen and Silicon. For nitrogen diffusion studies, the samples were isochronally annealed in vacuum ($\sim 1 \times 10^{-6}$ mbar) from the ambient to 973 K in steps of 50 K for a period of 30 min. At the end of each annealing step, the samples were quenched using a jet of cooled helium gas. Similarly for silicon diffusion studies, the samples were vacuum annealed up to 873 K in steps of 50 K isocrally (30 min). However, for obtaining diffusion profiles at temperatures greater than 873 K, separate sets of samples were implanted under similar irradiation condition and annealed in a tubular furnace under vacuum ($\sim 1 \times 10^{-6}$ mbar). These identical samples have been annealed from 873 K to 1073 K in steps of 50 K at different annealing time in order to obtain a measurable broadening of the depth profiles.

3. Results and Discussion

The experimental depth profile of $^{15}$N in D9 steel obtained by RNRA is shown in Figure 1(a). The depth profile exhibits an asymmetric Gaussian profile near the surface, and the experimental data has been fitted into two Gaussians with a variance of 0.95. SRIM simulation [9] of 30 keV nitrogen ions incident on steel shows that the peak damage and the projected range are located around 20 nm and 40 nm, respectively (Figure 1(b)). The peak position of the fitted Gaussians coincides with the peak damage region and the projected range of implantation, defined by the SRIM simulation. It is observed that a part of implanted nitrogen atoms is located around the peak damage region in addition to being present at the predicted range of implantation. The redistribution occurs in order to relieve the stress caused by the N$^+$ implantation possibly forming vacancy-nitrogen complexes. In order to explore this possibility, the depth profile of implantation-induced vacancy complexes has been probed using positron annihilation spectroscopy. From the changes observed in the defect-sensitive S-parameter, the evidence for the formation of vacancy-nitrogen complexes in the peak damage region was revealed [10].

The as-implanted sample was subjected to isochronal annealing in steps of 50 K, and it was found that there was no significant alteration of the nitrogen concentration profile up to 823 K as compared to the as implanted profile. However, beyond this temperature, the profile broadens and extends into the depth of the sample, while the nitrogen concentration remains almost constant near the surface.

<table>
<thead>
<tr>
<th>Table 1: D9 composition (in weight %).</th>
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<tbody>
<tr>
<td>Element</td>
</tr>
<tr>
<td>Cr</td>
</tr>
<tr>
<td>Mo</td>
</tr>
<tr>
<td>Si</td>
</tr>
<tr>
<td>C</td>
</tr>
<tr>
<td>Co</td>
</tr>
<tr>
<td>Al</td>
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<tr>
<td>P</td>
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<tr>
<td>Ta</td>
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<tr>
<td>As</td>
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<tr>
<td>Fe</td>
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</table>
Figure 1: (a) Experimental depth profile of $^{15}$N implanted in steel to a fluence of $5 \times 10^{15}$ N/cm$^2$ obtained using RNRA. The experimental data points (closed triangle) are fitted into two Gaussians whose peaks are centered at $\sim 20$ nm (dotted line) and $42$ nm (dashed line). (b) SRIM simulation of $30$ keV nitrogen in steel depicting the implanted ion and vacancy profiles.

Figure 2: Depth profile of implanted nitrogen after annealing the same sample at $823$ K, $873$ K, $923$ K, and $973$ K for $30$ min duration.

The profile broadening of the as-implanted sample (Gaussian distribution) can be used to deduce the diffusion coefficient $D$ by knowing the width of the implantation profile. In the present case, as the implantation profile is asymmetric and close to the surface, the implanted nitrogen was assumed as a thin source which diffuses into the bulk during annealing. Therefore, the following expression for the diffusion from an infinitely thin source into a semi-infinite medium is used [11, 12]. The concentration $C$ at the depth $x$ after annealing the sample for time $t$ is

$$C = A \exp \left( -\frac{x^2}{4Dt} \right),$$  

(1)

where $D$ is the diffusion coefficient. The plot of $\ln C$ versus $x^2$ gives the nitrogen penetration profiles in steel (Figure 3). The diffusion coefficient can be calculated from the slope of the penetration profiles [12]

$$D = \left( \frac{p - p_0}{4p_0^2t} \right),$$  

(2)

where $p_0$ and $p$ are the slopes of penetration profiles before and after annealing at a temperature $T$. The major uncertainty in the calculation of diffusion coefficients arises from obtaining straight line fits to the penetration profiles, and it is estimated to be $\sim 5\%$. Table 2 gives the diffusion coefficients of nitrogen in D9 at various temperatures. From the Arrhenius plot of $\ln D$ versus $(kT)^{-1}$, the effective activation energy ($E_a$) for nitrogen diffusion in steel is estimated to be $0.9$ eV (Figure 4). This value is comparable to the activation energy derived in austenitic stainless steels involving low-energy ion implantation ($\sim 1$k eV): $0.8$ eV [13], $1$ eV [14] and plasma nitriding: $1.07$ eV [15]. Although nitrogen diffusion in these cases was obtained at elevated substrate temperatures under irradiation conditions, the activation energies were calculated from the penetration of nitrogen from the surface into the bulk of the sample which is similar to the method employed here. Comparing the activation energy obtained in the present studies with other postimplantation annealing studies in similar austenitic steel [16], it is found that the present value is much lower. On studying the annealing behavior of implanted nitrogen ($40$ keV) in austenitic stainless steel (SS316), Hirvonen and Anttila determined the activation energy to be $1.87$ eV [16]. In their study [16], nitrogen was implanted to a fluence of $7.5 \times 10^{16}$ ions/cm$^2$, where there could be trapping of nitrogen due to the possible formation of V-N complexes in great extent, and also when nitrogen is in excess, precipitates can form hindering the movement of nitrogen. In the present case also, we inferred from the positron experiments that there is trapping of nitrogen at vacancies [10]. This could give rise to a decrease in the concentration of free nitrogen available for diffusion. We observe that there is no significant change in the near-surface concentration of nitrogen. This suggests that the V-N complexes remain immobile and the observed diffusion is due to the free interstitial nitrogen. At higher annealing temperatures, the V-N complexes might become mobile, or a breakup of the complexes is possible.
temperatures in D9 steel.

### Table 2: Diffusion coefficients of N in D9.

<table>
<thead>
<tr>
<th>Temperature (K)</th>
<th>Time ($10^3$ sec)</th>
<th>Diffusion coefficient (m$^2$/sec)</th>
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<tbody>
<tr>
<td>873</td>
<td>1.8</td>
<td>$6.1 \times 10^{-16}$</td>
</tr>
<tr>
<td>923</td>
<td>1.8</td>
<td>$1.2 \times 10^{-15}$</td>
</tr>
<tr>
<td>973</td>
<td>1.8</td>
<td>$2.1 \times 10^{-15}$</td>
</tr>
</tbody>
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### Table 3: Diffusion coefficients of Si in D9.

<table>
<thead>
<tr>
<th>Temperature (K)</th>
<th>Time ($10^3$ sec)</th>
<th>Diffusion coefficient D (m$^2$/sec)</th>
</tr>
</thead>
<tbody>
<tr>
<td>873</td>
<td>39.6</td>
<td>$4.6 \times 10^{-20}$</td>
</tr>
<tr>
<td>923</td>
<td>18</td>
<td>$1.3 \times 10^{-19}$</td>
</tr>
<tr>
<td>973</td>
<td>3.6</td>
<td>$8.8 \times 10^{-19}$</td>
</tr>
<tr>
<td>1023</td>
<td>3.6</td>
<td>$2.6 \times 10^{-18}$</td>
</tr>
<tr>
<td>1073</td>
<td>0.9</td>
<td>$6.5 \times 10^{-18}$</td>
</tr>
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In fact, small changes in the near surface composition of nitrogen is observed as a function of annealing temperature, and a more detailed study of this might throw light on the evolution of the nitrogen trapped at vacancies.

Proceeding to the next study in $^{30}$Si-implanted D9 samples, the experimentally obtained depth distribution of the as-implanted silicon is illustrated in Figure 5, which is a Gaussian profile with centre and full width at half-maximum (FWHM) as 120 nm and 95 nm, respectively. The experimentally determined depth profile is in good agreement with the theoretical profile computed from SRIM program [9]. It should be noted here that there is no redistribution of implanted silicon quite contrary to the observation in nitrogen implantation. The absence of such a redistribution of the implanted silicon may be due to the presence of silicon in substitutional position of the host lattice. This implanted silicon distribution serves as the marker layer, and its broadening is used for following the diffusion behavior of silicon in D9 samples.

The silicon implanted sample was annealed in steps of 50 K in vacuum for a period of 30 min at each step followed by RNRA experiments after each annealing. It is found that there is no significant change in the width of the implanted profile in the sample up to 873 K with the annealing time of half an hour. After prolonged annealing at 873 K for few hours followed by RNRA experiment at different increments of time, the implantation profile showed a discernible change after eleven hours (Figure 5). Further, four numbers of Si-implanted D9 samples were annealed independently at four different temperatures from 923 K to 1073 K in steps of 50 K for different durations. Figure 5 shows the diffusion profile of silicon in the temperature interval of 873 K to 1073 K, and the annealing time for each temperature is indicated in the figure. At higher annealing temperatures beyond 973 K, segregation of Si towards the surface is also observed. However, broadening of depth profiles in this temperature regime is quite evident. The method used for estimating N diffusion coefficients is employed for Si diffusion also. Figure 6 shows the penetration profiles of Si in D9 at various temperatures with the respective annealing time. The accuracy of the calculated $D$ values is limited by the linear fitting of the slopes (error less than ~10%). The diffusion coefficients at different annealing temperatures are tabulated in Table 3. From the Arrhenius plot of ln $D$ versus the inverse of $kT$ (Figure 7), activation energy of 2.1 eV is obtained for the silicon diffusion in D9 steel.

The activation energy for silicon diffusion in the present experiment may be compared to the self-diffusion energy of iron (2.84 eV) in fcc-Fe [17], and it is also comparable with the diffusion of alloying elements like nickel (2.72 eV), iron (2.60 eV), and chromium (2.51 eV) reported in austenitic stainless steel [5]. In addition, the activation energy deduced here for silicon diffusion is much greater than the experimentally determined effective vacancy migration energy (1.13 eV) in solution-annealed D9 steel [18]. Hence, silicon diffusion in the solution annealed D9 steel is regarded to be a vacancy-assisted mechanism. The present value of activation energy for silicon diffusion is lower than a previously reported value of 2.43 eV for silicon diffusion in fcc-Fe [17]. Addition of fast diffusing solutes can increase the vacancy migration and in turn the diffusion of solvent atoms. Evidences exist in the literature with respect to the increase
in the diffusivity of major alloying elements like Fe, Cr, and Ni due to silicon additions [5, 19]. It should be noted that another fast diffusing species, namely titanium is present in D9 steel. Titanium being an oversized element in the D9 alloy would diffuse via vacancy mechanism. In a work by Okita et al. [20] to understand the effect of titanium on void swelling, studies were carried out on solution-annealed austenitic Fe-15Cr-16Ni ternary alloys (similar wt% of Cr and Ni in D9 steel) without titanium and with the addition of 0.25 wt% of Ti (similar wt% of Ti in D9 steel). In that paper [20], reduction in void nucleation due to titanium addition was attributed to the enhancement of the effective vacancy diffusion coefficient. Hence, the reason for the relatively lower activation energy obtained here for silicon diffusion in solution-annealed D9 steel is ascribed to the increased vacancy migration due to the presence of both silicon and titanium in D9 steels. Therefore, the present observation strongly suggests that the synergetic effects due to the presence of various fast diffusing solutes contribute to enhanced vacancy mobility and the consequent void swelling suppression.

4. Conclusions

Resonant nuclear reaction analysis has been utilized to obtain the diffusion profiles of nitrogen and silicon in Ti-modified stainless steel. Results show that in the temperature range of annealing, a fraction of nitrogen remains trapped near
the surface possibly vacancy–nitrogen complexes, and the activation energy for diffusion of free nitrogen is obtained. The thermal diffusion behavior of silicon in titanium-modified steel is also determined. From the estimated activation energy, it can be concluded that in a thermal environment silicon diffuses by vacancy mechanism in steel.

Acknowledgment

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References

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