

Research Article

Optimal Modes for the Fabrication of Aluminum Nanopowders by the Electrical Explosion of Wires

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The paper is aimed at studying the impact of initial conditions of electrical explosion of wires on energy characteristics of the explosion and some other properties of the obtained aluminum powders. Explosion modes where the energy input into the wire has the maximal level were found. These modes are optimal for fabrication of powders with the best properties. The powders have the highest value of the specific surface of $14.5 \text{ m}^2/\text{g}$, a narrow histogram of the particle size distribution, and a narrow distribution histogram with a high polydispersity coefficient of 0.7.

1. Introduction

Metal nanopowders are used in the development of high-quality construction materials. Because of their structural peculiarities, the products of powder metallurgy are more thermally stable and resist better cyclic temperature gradients and deformation stress [1]. The transition from bulk metals to highly dispersed powders and nanoparticles is characterized by a change in a number of fundamental material properties, such as melting temperature, electron work function, and chemical reactivity [2]. This allows us to prepare novel metal and composite materials with improved mechanical, electrophysical, magnetic, and physicochemical characteristics. In addition, these nanopowders are used as highly active catalysts [3, 4], lubricants [5, 6], and activators of combustion and sintering [7, 8].

One of the fabrication methods for metal nanopowders is the electrical explosion of wires (EEW) [8]. EEW is characterized by low energy consumption related to the pulsed supply of energy to the wire, technology simplicity, and the availability of raw materials. EEW is carried out in an inert gas and experiment parameters are empirically found in order to produce pure metal nanopowders. So far, the level of the energy input into the wire during EEW has been seen to exercise the most significant influence over the

obtained powder dispersivity [9]. However, a great number of variables such as the capacity of capacitor battery, the initial voltage of its charge, the diameter and length of the exploded wire, circuit inductance, and gas pressure complicate the calculation of EEW in ensuring the desired mode of explosion for fabrication of powders with chosen wires.

The objective of the present work is to study the influence of the energy input into the wire and the initial parameters of the electrical explosion of aluminum wires on the properties of produced aluminum nanopowders.

2. Methods and Approach

Experiments on electrical explosion of aluminum wires were carried out in the setup described in [10]. To simplify the experiments, EEW was performed at constant capacity of $1.08 \mu\text{F}$, circuit inductance of $0.82 \mu\text{H}$, and argon pressure of $3 \cdot 10^5 \text{ Pa}$. Aluminum wire of 0.25 mm in diameter was used for the experiments.

The starting voltage was varied from 20 to 30 kV, and the exploded wire length was varied from 30 to 180 mm. The response on a change of parameters was controlled with waveforms recorded using a current shunt and an oscilloscope Tektronix 2014B. Additionally, the EEW current and voltage waveforms were treated in order to calculate the

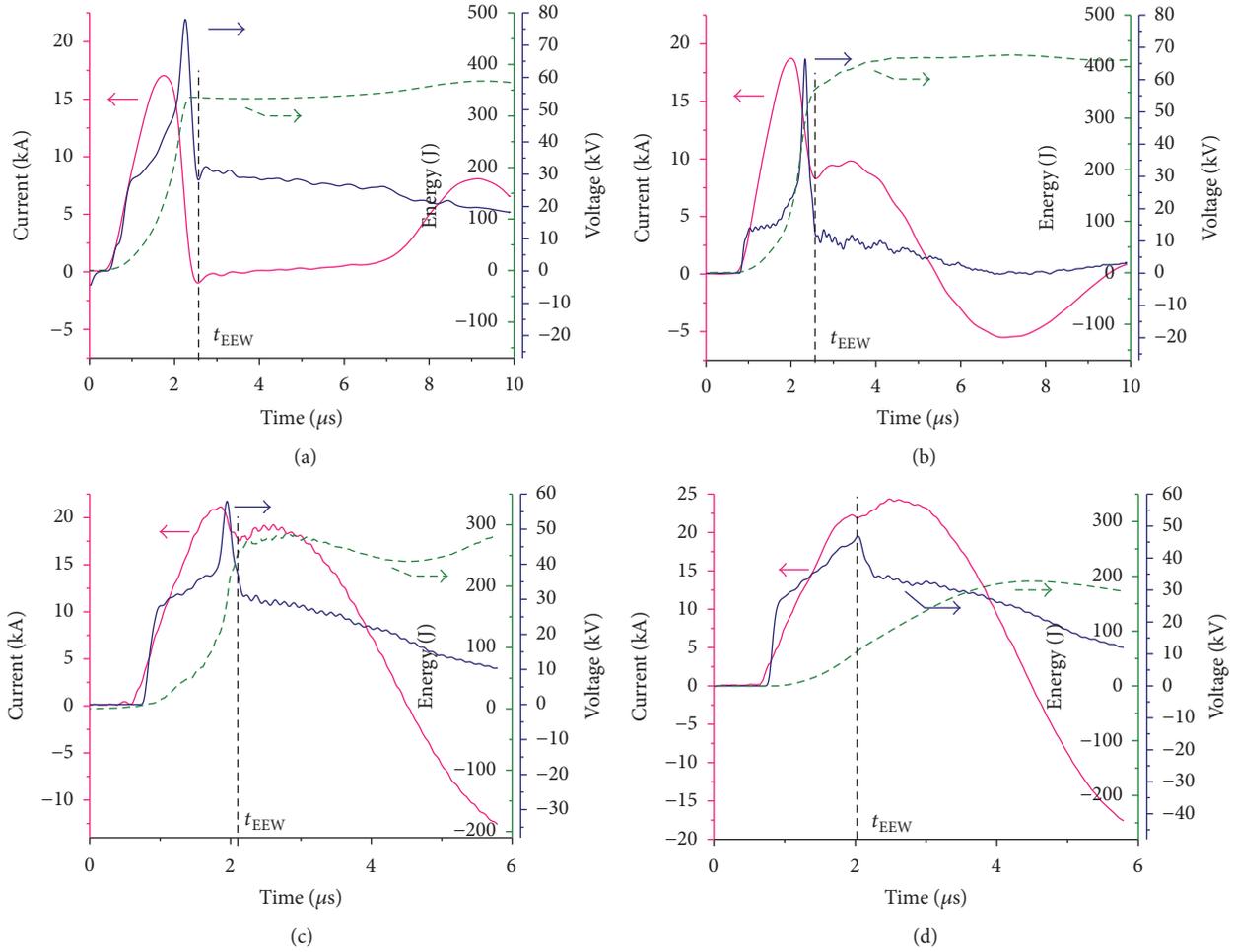


FIGURE 1: The EEW current and voltage waveforms: $U_0 = 30$ kV. (a) $l = 175$ mm, $e/e_s = 1.1$, and $e_{arc}/e_s = 0.2$; (b) $l = 135$ mm, $e/e_s = 1.6$, and $e_{arc}/e_s = 0.4$; (c) $l = 85$ mm, $e/e_s = 2.0$, and $e_{arc}/e_s = 1.2$; (d) $l = 60$ mm, $e/e_s = 1.7$, and $e_{arc}/e_s = 2.8$.

energy input into the wire and also the energy released during the charge arc stage. The energy released in the discharge gap in the range of time from the beginning of the current rise until the inflection moment of the current waveform (t_{EEW} in Figure 1) was considered as the energy input into the wire and the rest of the released energy as that released during the arc discharge stage.

The energy input into the wire was calculated using (1) offered by the authors of [11]:

$$W_{EEW} = \frac{CU_0^2}{2} - \frac{C(U_0 - (1/C) \int_0^{t_{EEW}} i(t) dt)^2}{2} - \frac{L \cdot i(t_{EEW})^2}{2} - R \int_0^{t_{EEW}} i(t)^2 dt, \quad (1)$$

where C is the capacitor capacity, U_0 is the voltage of capacitor charge, i is the EEW current, L is the inductance of setup circuit, R is the circuit resistance, and t_{EEW} is the duration of an EEW current pulse. According to the authors of [11], the accuracy of the calculations of the energy input into the wire is not less than 4–6%. The calculations did not take into

account the inductance of the exploded wire which was not less than 10%. In order to compare the various modes of electrical explosion of wires, the energy input into the wire was expressed in the form of a dimensionless value:

$$\frac{e}{e_s} = \frac{W_{EEW}}{V_W \cdot e_s}, \quad (2)$$

where e is the specific energy, V_W is the volume of the exploded wire, and e_s is the energy of aluminum sublimation (32.9 J/mm³). The energy released during the charge arc stage was also converted into a dimensionless value e_{arc}/e_s .

The obtained aluminum powders were passivated in air using the method described in [12]. The shape and the average size of the particles were studied using a transmission electron microscopy JEM-2100 F (JEOL, Japan). A particle size distribution histogram was additionally built. For this purpose, more than 15 microphotographs obtained at different magnitudes (totally more than 3000 particles) were treated. This histogram served for the calculations of the number-average (a_n), surface-average (a_s), and mass-average (a_m) particle size. According to a_n/a_m ratio, the particle

polydispersion coefficient (k) characterizing the particle size “dispersion” was calculated.

The specific surface area (S_{sp}) was measured according to the BET method using a Sorbtometer-M device. The specific surface area was used for the calculations of the surface-average particle diameter according to the following equation:

$$d_s = \frac{6}{\rho S_{sp}} \cdot 1000 \text{ [nm]}, \quad (3)$$

where ρ is the density of aluminum, 2.7 g/cm^3 .

The phase composition was determined using an X-ray diffractometer, Shimadzu XRD-7000 (Japan). According to the data of X-ray peak maximum enlargement, the size of coherent scattering regions (average size of crystallites) was calculated using the Debye-Scherrer equation:

$$d_{CSR} = \frac{k \cdot \lambda}{\beta \cdot \cos(2\theta/2)}, \quad (4)$$

where k is the dimensionless coefficient (for spherical particles it is 0.9), λ is the X-ray wavelength (0.1540593 nm), β is the full width at half maximum, and 2θ is the diffraction angle.

The content of metallic aluminum (Al^0) in samples was measured using the volumetric method during interaction of the metal with a 10% NaOH solution.

3. Results and Discussion

Figure 1 shows the EEW current and voltage waveforms. This figure also illustrates the dependence of the energy input during the discharge period versus time at a fixed voltage of 30 kV. In this experiment, the only exploded wire length was varied from 175 to 30 mm.

Variation in the exploded wire length leads to a rather high change in the EEW waveforms. At the longest length (Figure 1(a)), the EEW mode is characterized by two separately occurring processes. The first process is electrical explosion itself. The second one is the EEW arc stage. This explosion mode is characterized by a high coefficient of energy transfer from the capacitor to the wire: more than 80%. In the literature, this mode is called mode with a «current pause» [9, 13]. Nevertheless, the value of the energy input into the wire is not high and is equal to 330 J ($1.1e_s$). The decrease of the exploded wire length up to 85 mm leads to fusion of electrical explosion itself and the discharge arc stage: a transition to the mode without current pause [9, 13, 14]. The level of the energy input into the wire decreases down to 270 J . However, an increase in the specific energy (up to $2e_s$) into the wire $2e_s$ is observed due to an increase in the exploded wire volume. Moreover, the energy released during the arc discharge stage increases up to $1.2e_s$ (Figure 1(c)). Then, when decreasing the length (Figure 1(d)), a bigger part of energy is released during the arc stage, whereas the specific energy input into the wire decreases down to $1.7e_s$ (130 J). Probably, it is due to a rather low electrical strength of argon. When the wire length is less than 80 mm, an earlier start of the arc

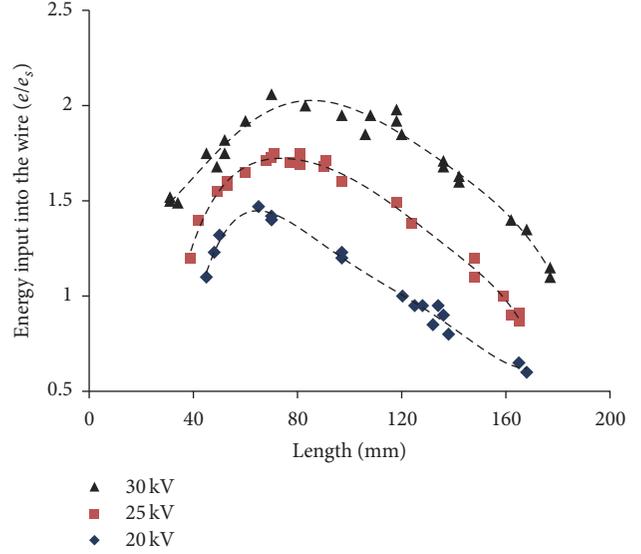


FIGURE 2: Energy input into the wire during explosion versus the exploded wire length and the voltage.

discharge stage is observed. It leads to shunting of the energy input into the wire (Figure 1(d)).

Figure 2 depicts the dependence of the specific energy input into the wire during the explosion on the exploded wire length for a series of experiments carried out at a voltage of 20, 25, and 30 kV.

The regularities of the change of e/e_s from l at different U_0 reveal similar character. The minimal e/e_s value is observed during electrical explosion of wires with the maximal length. A decrease in the exploded wire length leads to an increase in e/e_s . However, an increase in the specific energy input into the wire is observed up to a definite value of the working voltage. For the series of experiments at $U = 20 \text{ kV}$, the maximal e/e_s value (1.5) is observed at $l = 65 \text{ mm}$. For the series of experiments at $U = 30 \text{ kV}$, the maximal e/e_s value is 2 at $l \approx 80 \text{ mm}$. The further increase of the wire length leads to a fast decrease of the energy input into the wire.

Initial EEW conditions for fabrication of nanopowders were chosen taking into account the obtained data. The initial EEW conditions and properties of the powders produced under these conditions are given in Table 1.

Sample Al-1, obtained at the lowest level of the specific energy input into the wire during the explosion ($0.9e_s$), has the lowest specific surface area that equals $5.6 \text{ m}^2/\text{g}$. The particles have a spherical shape; there are both big particles of $4.5 \mu\text{m}$ and small particles of 25 to 400 nm. According to the literature data [15], the big particles form during the metal melting and sputtering, whereas the origin of the small particles is related to the metal evaporation during the explosion and the further condensation of the vapors. On the particles surface, we observe an approximately 5 nm thick well-defined layer which consists of aluminum oxide-hydroxide phases according to the literature data [2] formed during the passivation of the powders in the air. The presence of these phases is confirmed by the volumetric analysis,

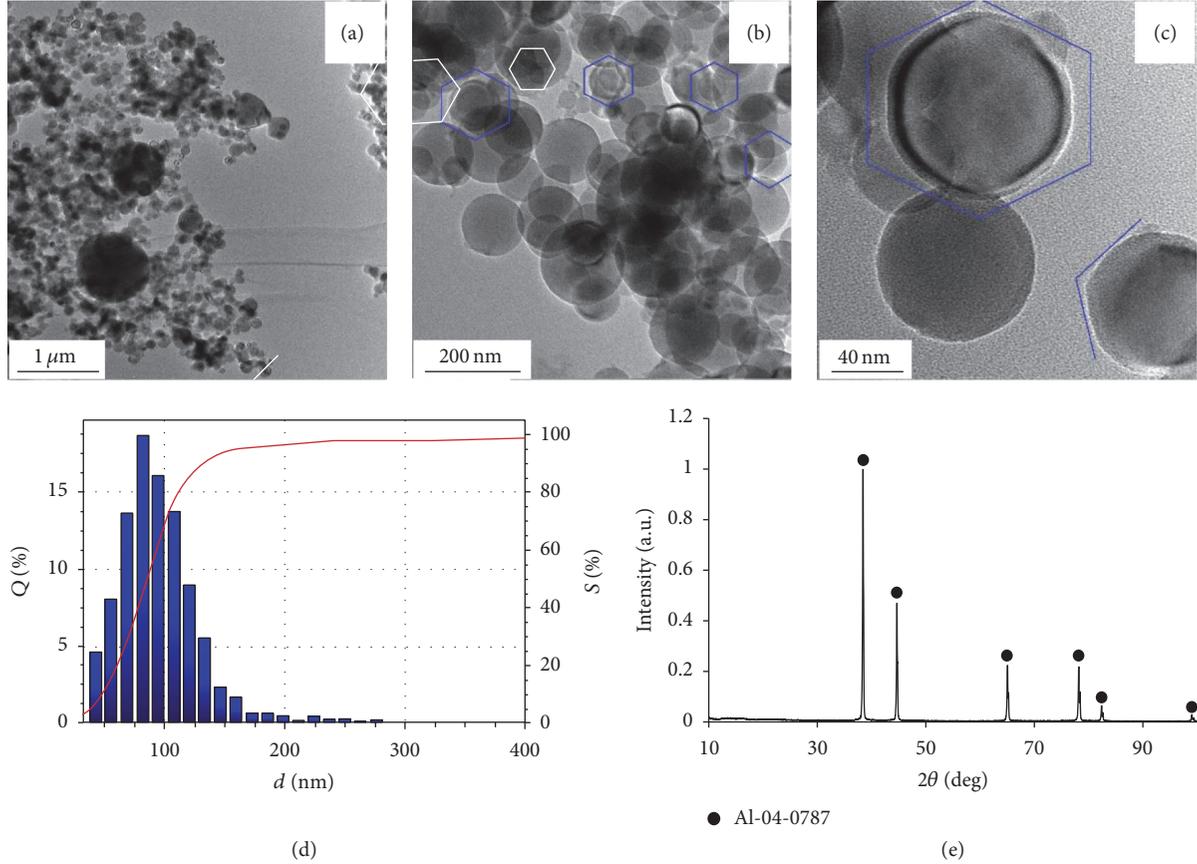


FIGURE 3: Microphotographs (a, b, c), histogram of the particle size distribution (d), and XRD pattern (e): sample Al-5: $a_n = 130$ nm; $a_s = 150$ nm; $a_m = 180$ nm; $k = 0.7$.

TABLE 1: Impact of EEW energetic parameters on the characteristics of Al powders.

Sample	U_0 , kV	l , mm	e/e_S	e_{ars}/e_S	S_{sp} , m^2/g	d_s , nm	Al ⁰ , wt.%
Al-1	20	120	0.9	0.1	5.6	400	96
Al-2	20	75	1.4	0.2	8.1	275	93
Al-3	25	120	1.5	0.1	9.2	240	92
Al-4	25	75	1.7	0.8	11.2	200	91
Al-5	30	85	2.0	1.2	14.5	150	89
Al-6	30	60	1.7	2.8	13.9	160	90
Al-7	30	30	1.5	7.5	13.1	170	90

according to which the content of metallic aluminum in the sample is 96 wt.%.

The size of 95% of the particles ranges from 25 to 400 nm; their maximum distribution corresponds to 115 nm. The presence of 5% of particles with a size from 400 nm to 4.5 μ m leads to an increase in the number of particles of average size up to 160 nm and up to 350 nm in the weight-average.

The sample is characterized by a rather low polydispersity coefficient: 0.4. When the particle size distribution is large, the surface-average particle diameter, calculated from the distribution histogram (a_s), is practically 2 times lower than the analogous value calculated from the specific surface area (d_s). The difference between a_s and d_s can be explained by the difference in techniques of sample preparation for

analysis. The BET measurements are more accurate because the surface of all particles is measured in contrast of TEM, where most of big particles, sedimented in the solution, are not analyzed.

The increase in e/e_S leads to an increase in the specific surface area of the obtained particles. Highest S_{sp} was found for the samples prepared under the modes with the maximal level e/e_S equal to 2.0 (Al-5, Table 1). The microphotographs of sample Al-5 are given in Figure 3. The sample is characterized by a narrow range of the particle size distribution and a rather high polydispersity coefficient: 0.7. As a result, the surface-average particle diameter, calculated from the distribution histogram, corresponds to the particle size calculated using the powder specific surface area. The size of most particles is

in the range of 25–200 nm, whereas particles of approx. 1 μm are observed but their part is not high.

Most particles of nanometer size have a spherical shape. However, in contrast to the previous sample, we can observe particles having the shape of a hexagon (Figures 3(b) and 3(c)), where the diameter of such particles is always less than 100 nm. According to the authors [2], the presence of the faceted particles (white hexagonal or polygonal lines in Figures 3(b) and 3(c)) is related to the mechanism of their crystallization. Droplets with the size of 100 or more nanometers crystallize from a large number of nuclei of which the mutual growth limitations contribute to the formation of particles with a spherical shape. Crystallization of relatively small droplets can occur during the formation and growth of one or a few nuclei when the particle forms in the form of a hexagon.

The increase in the specific surface area of the studied samples from 5.6 m^2/g to 14.5 m^2/g leads to a twofold decrease in the sizes of coherent scattering regions from 150 nm to 70 nm which can be explained by a decrease in crystallite sizes. However, the crystallite size for the studied samples calculated using the Debye-Scherrer formula is lower than the average particle size calculated using the powder specific surface area (Table 1). This allows us to assume that the powder particle produced at $2.0e_s$ has a polycrystalline structure.

The study of the modes with an increase of e/e_s and fast rise of e_{arc}/e_s showed that the powder specific surface area under these explosion modes decreases with a decrease in e/e_s (samples Al-6 and Al-7, Table 1). The powder particles obtained at the highest energy released during the arc discharge stage (sample Al-7) have both a spherical shape and polygon shape. The difference between this sample and the one considered above consists in the appearance of particles sintered with each other. Probably it is due to the release of large amount of discharge arc stage heat resulting in the sintering of particles.

4. Conclusions

An increase in the specific energy input into the wire is observed when decreasing the length of the exploded aluminum wire and keeping constant the other EEW conditions. However, this increase occurs until a definite moment. The dependence of e/e_s on U_0 and l has a maximum; the further decrease in the length leads to either significant increase in the working voltage or increase in the energy released at the discharge arc stage and decrease in the specific energy input into the wire during the explosion.

The optimal EEW modes for the fabrication of metal Al powders in the studied range are the explosion modes with the maximal level of the energy input into the wire ($U = 30 \text{ kV}$; $l = 85 \text{ mm}$). The powders have the highest value of the specific surface of 14.5 m^2/g , a narrow histogram of the particle size distribution, and a high polydispersity coefficient.

Under the modes with an increase in the specific energy input into the wire during the arc discharge stage, the specific surface area decreases; the samples were found to have a high amount of sintered particles which can be provoked by

the high thermal action of the energy released during the discharge arc stage.

Competing Interests

The authors declare that there are no competing interests regarding the publication of this paper.

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