Research Article

Representative Stress-Strain Curve by Spherical Indentation on Elastic-Plastic Materials

Chao Chang,1 M. A. Garrido,2 J. Ruiz-Hervias,3 Zhu Zhang,1 and Le-le Zhang4

1School of Applied Science, Taiyuan University of Science and Technology, Taiyuan 030024, China
2Departamento de Tecnología Química y Ambiental, Tecnología Química y Energética y Tecnología Mecánica, Escuela Superior de Ciencias Experimentales y Tecnología, Universidad Rey Juan Carlos, c/Tulipán s/n Móstoles, Madrid, Spain
3Departamento de Ciencia de Materiales, UPM, E.T.S.I. Caminos, Canales y Puertos, c/ Professor Aranguren s/n, 28040 Madrid, Spain
4School of Mechanical, Electronic and Control Engineering, Beijing Jiaotong University, Beijing 100044, China

Correspondence should be addressed to Chao Chang; chao_chang_tyust@163.com

Received 12 August 2017; Revised 4 December 2017; Accepted 25 December 2017; Published 26 February 2018

Abstract

Tensile stress-strain curve of metallic materials can be determined by the representative stress-strain curve from the spherical indentation. Tabor empirically determined the stress constraint factor (stress CF), \( \psi \), and strain constraint factor (strain CF), \( \beta \), but the choice of value for \( \psi \) and \( \beta \) is still under discussion. In this study, a new insight into the relationship between constraint factors of stress and strain is analytically described based on the formation of Tabor’s equation. Experiment tests were performed to evaluate these constraint factors. From the results, representative stress-strain curves using a proposed strain constraint factor can fit better with nominal stress-strain curve than those using Tabor’s constraint factors.

1. Introduction

The instrumented indentation technique consists of applying load to the sample by means of an indenter of known geometry, while the applied load and the penetration depth of the indenter are recorded simultaneously during a loading and unloading cycle. The load-penetration data can be used to determine the mechanical properties of the material without having to image the residual impression left on the material’s surface. Consequently, this technique can be applied at different scales, from macro to nano. The main mechanical properties measured by this technique are Young’s modulus, \( E \), and hardness, \( H \), by sharp indenter [1, 2]. This methodology could also measure material properties like yield stress, creep property, and fracture [3–6]. However, these properties are not sufficient to characterize a material. The present work focuses on the methodology to determine the stress-strain curve of metallic materials by the depth-sensing indentation technique using a spherical indenter. Sharp indenters like Berkovich or Vicker are characterized by inducing a constant strain on the indented materials. This deformation depends on the indenter angle. Consequently, if the complete stress-strain curve is needed, the sharp indenters are not the best option because they would give only a point of such curve. This is the reason why the interest for the spherical indenter has been recently grown. The experimental curve from spherical indenter provides more information, as this type of indenter has a smooth transition from elastic to elastic-plastic contact [7].

There were many methodologies for the extraction of flow stress with the spherical indentation technique. One category belongs to mathematical and numerical methods (e.g., dimensionless method and artificial neural network (ANN)). After running many finite element models within certain range of material properties, the indentation curves, contact stiffness, or the work of loading-unloading is obtained, which are verified by corresponding experimental indentation and establish a relation between the indentation characteristics and material properties [8–15]. However, these methods consume great computational effort. The unique solution of these methods is still under discussion [16, 17]. Moreover, the contamination in the experimental...
process like deficient indenter tip, roughness of specimen, or intrinsic noise of instrument machine is not considered in these methods.

Another category is to determine the material parameters from spherical indentation by defining the representative strain and stress. Regarding the experimental measurements, these methods are more convenient and can be used directly. Following the work of Meyer [18], Tabor defined the representative indentation strain, \( \varepsilon_r \), at contact edge of the spherical indenter as \( \varepsilon_r = 0.2 (a/R) \), where the values of \( \psi \) and \( \beta \) empirically were determined as 2.8 and 0.2, respectively, based on the quantity of tensile tests on common metals [19]. Most of the researchers agreed with Tabor's indentation strain but emerge with the controversy on the choice of values of \( \psi \) and \( \beta \) [5, 20–25]. Richmond et al. [21] predicted that the mean pressure was approximately equal to 3 times the yield stress and that the representative strain was approximately equal to 0.32 times the impression to ball diameter ratio. Herbert et al. [23] had used higher stress constraint factor of 3.7 to obtain the representative stress-strain curve for Al 6061-T6 by spherical indentation with a radius of 385nm. Based on the work of Matthews [22], Tirupataiah and Sundararajan [26] derived an expression for the indentation stress factor, \( \psi \), which was related with hardening exponent, \( n \). Similarly, Yetna N’Jock et al. [27] determined the tensile property by spherical indentation by simply using the expression of mean pressure to stress ratio as a function of hardening exponent which was derived from the FEM works of Taljat et al. [28]. Using the representative strain under spherical indentation, Ahn and Kwon [29] developed a shear strain definition by differentiating the displacement in the depth direction, in which the ratio of mean pressure to representative stress was equal to 3 in the fully plastic period by conducting instrumented spherical indentation on the steel specimen. By carrying out extensive forward analyses in FEM, Xu and Chen [24] using the indentation strain by Ahn and Kwon [29] found the indentation stress constraint factors, \( \psi \), depended almost linearly on hardening exponent, \( n \). Additionally, the corresponding indentation strain constraint factor \( \beta \) depended on both \( n \) and the ratio of Young’s modulus to yield stress. Milman et al. [30] assumed that the fully plastic zone beneath indenter was incompressible and proposed a new representative strain which was related with contact radius and contact depth. Fu et al. [31] used the Milman representative stress [30] and proposed a novel iterative process to determine the tensile stress-strain curve. Recently, Kalidindi et al. [32, 33] argued that the definition of indentation strain as \( a/R \) lacks any physical interpretation as a measure of strain and proposed a new definition of the indentation strain consistent with the Hertz theory, which was evaluated from several FEM simulations as well as from the analysis of experimental measurements.

From the literature, there is not uniform agreement among investigations concerning the representative stress and strain equations and also the values for the factors involved in the expressions. The main purpose of the present investigation is to systematically study Tabor’s indentation strain and propose the possibility to develop an analytical procedure to extract the stress-strain curve using experimental data from spherical indentations, which would be comparable to that obtained from a uniaxial test (i.e., tensile test).

2. Theoretical Background

2.1. The Analytical Relationship between Stress and Strain Factors. In 1908, Meyer had found that for many materials, the mean pressure increased with \( a/R \) according to the simple power law [18]

\[
P_m = k \left( \frac{a}{R} \right)^m,
\]

in which \( P_m \) is the mean pressure, \( a/R \) is the ratio of indentation radius to the ball radius, and \( m \) is the Meyer index.

In (1), \( k \) and \( m \) are constants. The Meyer equation was verified by further experimental studies [19, 34], and they suggested that the Meyer index, \( m \), was related with the hardening exponent, \( n \). Following Meyer’s work, Tabor had proposed the concept of representative strain or stress, by which the mean pressure in Meyer’s equation and the \( a/R \) ratio can be converted into the true stress-strain curve. He assumed that the mean pressure, \( P_m \), at the fully plastic regime was proportional to the representative stress, \( \sigma_r \), and the impression radius was proportional to the corresponding representative strain, \( \varepsilon_r \). Consequently, the representative stress and strain can be expressed as [19]

\[
\varepsilon_r = \beta \frac{a}{R},
\]

\[
\sigma_r = \frac{F}{\psi \pi a^2},
\]

where \( \beta \) is the indentation strain constraint factor, \( a \) is the contact radius, \( R \) is the radius of the indenter, \( \psi \) is the indentation stress constraint factor, and \( F \) is the indentation load.

Tabor determined the parameters \( \beta \) and \( \psi \) from empirically experimental data from spherical indentations under the fully plasticity regime. Generally, the indentation strain constraint factor \( \beta \) is considered to be equal to 0.2, and the indentation stress constraint factor \( \psi \) ranges from 2.8 to 3.2 [35]. It should be emphasized that the representative strain and stress defined by Tabor are an average value of the stress and strain states induced inside the material [36, 37]. The stress and strain constraint factors allow us to establish equivalence between the stress-strain indentation curve and the corresponding one obtained from a uniaxial test. The true stress-strain curve from uniaxial tensile or compression test can be expressed as

\[
\sigma = E \varepsilon \quad \text{for} \quad \sigma \leq \sigma_y \quad \text{elastic regime},
\]

\[
\sigma = K \varepsilon^n \quad \text{for} \quad \sigma \geq \sigma_y \quad \text{Hollomon’s equation for plastic regime},
\]

where \( E \) is the elastic modulus, \( \sigma_y \) is the yield stress, \( K \) is the strength coefficient, and \( n \) is the strain-hardening exponent.
Substituting (2) into the Hollomon equation:
\[
F = K \psi \pi a^2 \left( \frac{\beta}{\pi} \right)^n,
\]
which represents the relationship between load, \( F \), and the ratio of contact radius to indenter radius according to the power law at the fully plastic regime. Transforming (5) into the natural logarithm as
\[
\ln F = \ln \left( \frac{K \psi \pi a^2}{\pi / \beta} \right) + (n + 2) \ln a.
\]
(6)

According to Hertz equation under the elastic regime, the loading and contact radius can be expressed as
\[
F = \frac{4E_r a^3}{3R},
\]
(7)
where \( 1/E_r = (1 - \nu^2/E) + (1 - \nu^2/E_i) \), \( E, \nu \) and \( E_i \) and \( \nu_i \) are Young’s modulus and Poisson’s ratio for the bulk material to be measured and for the indenter, respectively.

Transforming (7) into the natural logarithm
\[
\ln F = \ln \left( \frac{4E_r}{3R} \right) + 3 \ln a.
\]
(8)

Most metals have an \( n \) value between 0.10 and 0.50, plotting (6) and (8) in the \( \ln a \) versus \( \ln F \) coordinate, as shown in Figure 1. Inversely extending the two lines, there should be an intersectional point “\( e \)” for different linear slopes.

Therefore, the representative strain and stress at the point “\( e \)” can be expressed as
\[
\varepsilon_r = E \beta \left( \frac{a_e}{R} \right),
\]
(9)

Combining (9) and (10), we can obtain
\[
\psi \beta = \frac{R}{\pi E} \cdot \frac{E_r \sigma_r}{a_e^3}
\]
(11)

Comparing (10) with Hertz equation in (7), we can obtain
\[
\phi = \psi \beta = \frac{4}{3 \pi} \cdot \frac{E_r \sigma_r}{E}
\]
(12)

In (12), the product, \( \phi \), of stress constraint factor \( \psi \) and strain constraint factor \( \beta \) is constant, which is dependent on the ratio of reduced Young’s modulus \( E_r \) to Young’s modulus \( E \). According to (12), the constraint factors can be determined if one of them is known.

It should be noted that Tabor’s representative method is restricted within the fully plastic regime. In this study, the procedure assumes that elastic-plastic transition is negligible, and the relationship between stress and strain constraint factors is affirmed from Hertz’s theory. In Kalidindi et al.’s [32, 33] work, they utilized indentation strain derived from the Hertz theory for extracting the full stress-strain curve of metallic materials. Similarly, we utilize the Hertz equation to study the relationship between the stress and strain constraint factors to extract the stress-strain curve based on Tabor’s representative method. For the soft metallic materials with high \( E/\sigma_y \) which has very short elastic-plastic transition, the intersection point “\( e \)” can be considered as a simplistic interpretation of the elastic-plastic transition regime. Especially, in experimental practice, when performing the nanoindentation test on the soft metallic

![Figure 1: Equations (6) and (8) are in the same ln a-ln F coordinate, and "e" is the intersectional point.](image_url)
Figure 2: Scanning electron microscope image of a worn diamond conical-spherical indenter.

Figure 3: The load versus penetration depth curves for the three indented materials performed with spherical tip. (a) Al-1050, (b) brass, and (c) steel.
materials, the penetration depth is very shallow to turn into the fully plastic regime. Consequently, this assumption is reasonable only for certain soft metallic materials. The experimental verification will be discussed later.

In Tabor’s equation, the value of stress constraint factor is selected as 3 for common engineering metals in literature [5, 31, 38]. However, there are few studies on the analytical solution for the stress constraint factor. The indentation pressure is found to be proportional to the yield strength of the material by the constraint factor, C. The magnitude of the constraint factor has been found to be depended on the material properties, particularly on the ratio of the modulus to the yield stress. For most metals, \( E/\sigma_y > 100 \), and the constraint factor is approximately equal to 3 \( (C \approx 3) \) [19]. For polymers, the ratio of \( E/\sigma_y \) is less than 3 [39]. An approach to large strain plasticity problems in which the material is considered to behave in a plastic-elastic fashion, instead of as a plastic-rigid body, is applied to the axisymmetric blunt indenter. The ratio of the mean stress on the punch face to the uniaxial flow stress of the material (constraint factor C) is found to be 2.82 for an extensive specimen. However, it is shown that a small part of the punch face is elastically loaded, and if the loaded punch area is assumed to be equal to the size of the plastic impression, then the constraint factor to be used is 3 instead of 2.82. This is the value to be used in interpreting the ordinary Brinell test [40]. Additionally, Shield found the exact solution for the axisymmetric flat punch. He found a maximum pressure under the punch of 2.8\( \sigma_y \) at the centre [41]. Von Mises yield conditions are more representative of engineering materials than Tresca. Although the slip line field theory can be applied to both types of yield conditions, Tresca is generally selected as it leads to a simpler equation that can be solved analytically. However, Von Mises yield can give up to 15.5% higher limit load value than Tresca, which would lead to \( C = 3.2 \) [42]. Unlike the punch indenter, where the contact area is constant, the spherical one offers a gradual transition of the contact area. It could be assumed that the spherical indenter is corresponding to a punch indenter at a given penetration depth. As a result, the value of stress constraint factor is taken as 3.2, which is consistent with the analytical solution by Shield [41]. As the product of stress and strain constraint factors is constant, according to (12), the strain constraint factor can be determined.

2.2. Estimation of Contact Radius. In order to calculate the representative stress and strain, an accurate measurement of the contact radius at defined indentation load would be obtained. In the case of a flat punch indenter, Sneddon’s analysis proposed an estimation expression of contact radius, \( a \):

\[
   a = \frac{S_u}{2F_r},
\]

where \( S_u \) is the contact stiffness.

From Eq. (13), the contact radius and Young’s modulus can be extracted from the initial unloading slope [43]. According to (13), the contact radius could be determined when reduced Young’s modulus and initial unloading slope of the unloading are known. The procedure to estimate the contact radius is consistent with Kalidindi et al.’s [32, 33] work.

3. Experimental Procedure

Depth-sensing indentation (DSI) tests were performed on a Nanoindentor XP (NanoInstruments Innovation Center, MTS systems, TN, USA) by using the continuous stiffness measurement (CSM) methodology. A conical (angle 60°) spherical diamond indenter with a tip radius of 10 \( \mu m \) was selected. To verify the radius of the indenter, scanning electron microscope was used to measure the radius of the tip before the test. Continuous loading and unloading cycles were conducted during the loading branch by imposing a small dynamic oscillation of 2 nm and 45 Hz on the displacement signal and measuring the amplitude and phase of the corresponding force [44, 45]. Consequently, the contact stiffness was continuously measured as a function of the penetration depth during the experiment. A tip calibration procedure was carried out using the fused silica in accordance with the CSM methodology [46]. Prior to the indentation tests, the samples were ground and polished in two steps, using a mechanical polisher (Labopol-5, Struers, Copenhagen, Denmark) and finished with colloidal suspension of silica of 0.05 \( \mu m \). A total of 25 indentations were performed in displacement control on three samples of commercial metallic alloys: a carbon steel (F-114), a brass alloy, and an aluminum alloy (Al-1050). The penetration depth was selected according to the grain size of indented materials: 3000 nm for Al-1050 and brass and 2000 nm for steel. The average of these curves is selected to determine the representative stress-strain curve by spherical indentation which can be comparable with the true stress-strain curve by tensile tests [3]. Tensile tests were also conducted.
on carbon steel (F-114), aluminum (Al-1050), and brass alloy samples. They were performed on a universal testing machine INSTRON 8501 with a 10 kN load cell at room temperature under displacement control by using a displacement rate of 0.5 mm/min. Cylindrical specimens were machined according to the ASTM E8 standard [47]. The gage length was 12.5 mm for the samples.

4. Results and Discussion

In the tensile test, according to Hollomon’s equation, the mechanical properties can be obtained from the stress-strain curve. Young’s modulus, yield stress which is determined by means of the 0.2% strain rule, and the hardening exponent are shown in Table 1.

In order to check the radius of the spherical indenter, as shown in Figure 2, the best fitting spherical radius is found to be about 9.2 µm from SEM measurement which is smaller than the nominal value. From the zoomed-in figure, the spherical indenter is worn at shallow penetration depth. In addition, inherent noise of continuous stiffness measurement and the roughness of specimen could scatter the representative stress-strain curve. Consequently, penetration depths less than 10 nm will not be utilized for the extraction of representative stress-strain curve.

Figure 3 shows the load versus penetration depth curves for the three indented materials performed with spherical tip. The indenter curves of brass exhibit considerable scatter, as the diameter of the indenter is comparative to the grain size of brass listed in Table 1. On the
contrary, the indenter curves of steel and Al-1050 are much more reproducible for the multiple grain indents. The average of these curves is selected to determine the representative stress-strain curve. The procedure is consistent with the previous work [3].

When Young’s modulus was confirmed by the tensile test and the contact stiffness can be measured by CSM, the contact radius can be estimated by Sneddon’s equation [48]. It should be noted that, when using (13) to estimate the contact radius, Young’s modulus is assumed as a constant during the indentation process [32, 49]. Tabor had concluded that the slope of the linear regression is Meyer’s index $m$ which is equal to $n + 2$, where $n$ is the hardening exponent in the Hollomon equation. Figure 4 shows the linear regression of $\log(P)$ versus $\log(a)$ for the experimental materials. The values of hardening exponent from the slope are in good agreement with those obtained from the tensile experiments in Table 1.

In order to check the reasonable determination of constraint factors in Tabor’s equation, the tensile stress-strain curves are compared with representative stress-strain curves from a different constraint factor (CF): stress CF = 3.2 in this study and stress CF = 3 according to Tabor’s equation. In the experimental analysis, the value of stress constraint factor is taken as 3.2, and the strain constraint factor is confirmed from (12). In order to verify the reasonability of the confirmation of constraint factors from experimental data, the representative stress-strain curve extracted from this study will be compared with the curve extracted from Tabor’s method.

Figure 5 shows the plots for the indentation stress-strain curve using the constraint factors in the study, the indentation stress-strain curve using the constraint factors in Tabor’s equation, and the true stress-strain curve extracted from the tensile test. The values of $E/\sigma_y$ for studied materials range from 250 to 320. Although the initial part of the true stress-strain curve and the representative stress-strain curves shows little derivation between them, the representative stress-strain curve using the constraint factors in the study shows well agreement with the true stress-strain curve from the tensile test in the fully plastic regime. In this study, the product $\phi$ is attained from the intersection between the elastic regime and the fully plastic regime, in which the effect of the transition from the elastic to plastic regime is ignored. This assumption is only suitable for the soft metal which has short transition of the elastic-plastic regime, and the fully plastic regime is developed at shallow penetration depth. For metals with longer transition of the elastic-plastic regime, it can be seen that there is significant derivation in the initial part of the representative stress-strain curve regarding the true stress strain curve obtained from tensile test. As noted in literature [32, 33, 50, 51], it is very difficult to make a comparable stress-strain curve with the tensile test in the elastic-plastic transition by Tabor’s representative method for the metallic materials with low $E/\sigma_y$. The FEM simulations also showed that representative stress-strain from Kalidindi’s equation led to large deviation for some materials with large transition of the elastic-plastic regime [52]. However, in the study, utilizing the proposed procedure to determine the constraint factors could make a comparable stress-strain curve in the fully plastic regime.

5. Conclusion

The main conclusion in the paper can be summarized as follows:

(1) In this study, a new insight into the constraint factors in the formation of Tabor’s equation is analytically described. The product $\phi$ of strain constraint factor $\beta$ and stress constraint factor $\psi$ is constant at the transition position between the elastic regime and fully plastic regime, which is related to Young’s modulus of contact materials.

(2) An experimental procedure is performed to evaluate the analytical analysis for extracting the representative stress-strain curve from the spherical indentation. The representative stress-strain curve using the proposed constraint factors in the study shows well agreement with the nominal true stress-strain curve.

(3) Due to the drawback of the formation of Tabor’s equation, the representative stress-strain curve derived from the Tabor’s formation could have large deviation at the initial strain period for the materials with longer elastic-plastic transition. In future, formation of new representative strain for spherical indentation needs more study.

Nomenclature

- $p_m$: Mean pressure
- $\alpha$: Contact radius
- $R$: Indenter radius
- $m$: Meyer’s index
- $k$: Coefficient in Meyer’s equation
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\beta$: Indentation strain constraint factor
- $\psi$: Indentation stress constraint factor
- $E$: Elastic modulus
- $\sigma_r$: Yield stress
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress
- $E$: Elastic modulus
- $K$: Strength coefficient in Hollomon’s equation
- $n$: Strain coefficient in Hollomon’s equation
- $\psi$: Indentation stress constraint factor
- $\beta$: Indentation strain constraint factor
- $\phi$: Product of stress constraint factor $\psi$ and strain constraint factor $\beta$
- $\varepsilon_r$: Representative strain
- $\sigma_r$: Representative stress
- $\sigma_y$: Yield stress

Conflicts of Interest

The authors declare that they have no conflicts of interest regarding the publication of this paper.

Acknowledgments

This work was supported by the National Key R&D Program of China (2016YFB1200403-A-02), the Natural Science Foundation of Shanxi Province (no. 201701D221008), the
References


