

Research Article

Dynamics and Patterns of a Diffusive Prey-Predator System with a Group Defense for Prey

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We study a diffusive prey-predator system with a group defense for prey. Under Neumann boundary condition, we analyze local and stability of nonnegative constant steady states and the existence and nonexistence of nonconstant steady states. These results also exhibit the critical role of the system parameters leading to the formation of spatiotemporal patterns.

1. Introduction

The predator-prey system first proposed by [1, 2] is one of the fundamental ecological systems in both ecology and mathematical ecology. Based on different settings, various types of predator-prey models described by differential systems have been proposed and the dynamics of these systems are studied [3–6]. The basic form of these models is as follows:

$$\begin{aligned} \frac{dx}{dt} &= rx \left(1 - \frac{x}{K} \right) - P(x, y), \\ \frac{dy}{dt} &= -sy + cP(x, y), \end{aligned} \quad (1)$$

where r is the intrinsic growth rate and K is the environmental carrying capacity of prey population, and the function $P(x)$ is the functional response; the constant $c(>0)$ is the ratio of biomass conversion and s is the natural death rate of predator species. The simplest functional response is Lotka-Volterra function which is described as

$$P(x, y) = \begin{cases} axy, & 0 \leq x \leq \frac{k}{a}, \\ ky, & x \geq \frac{k}{a}, \end{cases} \quad (2)$$

which is also called Holling type I function. However, the curve defined by the Lotka-Volterra response function is a

straight line through the origin and is unbounded. Thus, more reasonable response functions should be nonlinear and bounded. In 1913, Michaelis and Menten proposed the response function

$$P(x, y) = \frac{mxy}{a + x}, \quad (3)$$

where $m > 0$ denotes the maximal growth rate of the species and $a > 0$ is the half-saturation constant. It is now referred to as a Michaelis-Menten function or a Holling type II function. Another class of response function is

$$P(x, y) = \frac{mx^2y}{a + bx + x^2} \quad (4)$$

which is called a sigmoidal response function, while the simplification

$$P(x, y) = \frac{mx^2y}{a + x^2} \quad (5)$$

is known as a Holling type III function. Some authors [7, 8] considered system (1) with following response function:

$$P(x, y) = \frac{mxy}{a + x^2} \quad (6)$$

which is called Holling type IV function. Besides, Beddington-DeAngelis type $P(x, y) = \beta xy/(\alpha + x + my)$ and more complicated functional response $P(x, y) = \beta x^2 y/(x^2 + my^2)$ are also considered by some researchers [9, 10].

Recently, some works consider the case when animals join together in herds in order to provide a self-defense from predators. In [11], the authors argued that it is more appropriate to model the response functions of prey that exhibit herd behavior in terms of the square root of the prey population. Inspired by this thought, the authors in [12] choose response function $P(x) = \sqrt{x}$ to reflect this fact. When motion is allowed, [13] considered the spatiotemporal behavior of a prey-predator system with a group defense for prey by means of extensive computer simulations. The proposed model is as follows:

$$\begin{aligned} \frac{\partial u}{\partial t} &= d_1 \Delta u + r_1 u \left(1 - \frac{u}{K}\right) - \rho u^\alpha v, \\ (x, t) &\in \Omega \times (0, +\infty), \\ \frac{\partial v}{\partial t} &= d_2 \Delta v + \beta u^\alpha v - r_2 v, \quad (x, t) \in \Omega \times (0, +\infty) \end{aligned} \quad (7)$$

$$u(x, 0) = u_0(x),$$

$$v(x, 0) = v_0(x),$$

$$x \in \Omega,$$

where u and v denote, respectively, the densities of prey and predator species. r is the growth rate of prey species, K is its carrying capacity, r_2 is the mortality rate of predator species, ρ is the search efficiency of predator for prey, β is the biomass conversion coefficient, and $\alpha \in (0, 1)$ represents a kind of aggregation efficiency. The local dynamics for nonspatial model was studied, such as Hopf bifurcation and existence of extinction domain. For model (7), the authors only give some numerical simulations to find some spatiotemporal features. Reference [14] considers the direction and the stability of the bifurcating periodic solutions for model (7) with $\alpha = 1/2$ under Neumann boundary conditions. Reference [15] investigated the global dynamics of nonspatial model including the nonexistence of periodic orbits and the existence and uniqueness of limit cycles. We refer readers to [16–21] as some other related works on predator-prey model with herd behavior.

It is noted that up to now no one has studied the existence and nonexistence of positive steady state solutions of (7). Therefore, the main aim of this article is to study the existence and nonexistence of nonconstant positive solutions of the following elliptic system:

$$\begin{aligned} -d_1 \Delta u &= r_1 u \left(1 - \frac{u}{K}\right) - \rho u^\alpha v, \quad \text{in } \Omega, \\ -d_2 \Delta v &= \beta u^\alpha v - r_2 v, \quad \text{in } \Omega, \\ \frac{\partial u}{\partial \nu} &= \frac{\partial v}{\partial \nu} = 0, \quad \text{on } \partial \Omega, \end{aligned} \quad (8)$$

where ν is the outward unit normal vector on $\partial \Omega$, and we impose a homogeneous Neumann type boundary condition,

which implies that (8) is a closed system and has no flux across the boundary $\partial \Omega$.

The structure of this paper is arranged as follows. In Section 2, we estimate the a priori bounds of positive solutions of (7). In Section 3, the local and global stabilities of nonnegative constant steady states of (7) are discussed. In Section 4, we give a priori estimate for the positive solutions of (8) by using maximum principle and Harnack inequality. In Section 5, we give a nonexistence result of nonconstant solutions of (8). In Section 6, we consider the existence of nonconstant positive solutions of (8). Finally, to support our theoretical predictions, some numerical simulations are given.

2. Basic Dynamics and a Priori Bound

Theorem 1. *For system (7), one has the following.*

- (a) *If $u_0(x) \geq 0$, $v_0(x) \geq 0$, then system (7) has a unique solution $(u(t, x), v(t, x))$ such that $u(t, x) > 0$, $v(t, x) > 0$ for $t \in (0, +\infty)$ and $x \in \bar{\Omega}$.*
- (b) *Any solution (u, v) of (7) satisfies*

$$\begin{aligned} \limsup_{t \rightarrow +\infty} u(t, x) &\leq K, \\ \int_{\Omega} v(x, t) dx &\leq \frac{\beta}{\rho} \left(1 + \frac{r_1}{r_2}\right) K |\Omega|. \end{aligned} \quad (9)$$

Proof. (a) Define

$$\begin{aligned} f(u, v) &= r_1 u \left(1 - \frac{u}{K}\right) - \rho u^\alpha v, \\ g(u, v) &= \beta u^\alpha v - r_2 v. \end{aligned} \quad (10)$$

Then $f_v = -\rho u^\alpha \leq 0$ and $g_u = \beta \alpha u^{\alpha-1} v \geq 0$ in $\overline{\mathbb{R}_+^2} = \{u \geq 0, v \geq 0\}$. Hence, (7) is a mixed quasi-monotone system. Consider following system:

$$\begin{aligned} \frac{du}{dt} &= r_1 u \left(1 - \frac{u}{K}\right) - \rho u^\alpha v, \\ \frac{dv}{dt} &= \beta u^\alpha v - r_2 v, \\ u(0) &= u_0, \\ v(0) &= v_0. \end{aligned} \quad (11)$$

Assume $u(t; u_0, v_0)$, $v(t; u_0, v_0)$ are the unique solution to system (11). Let

$$\begin{aligned} \max_{\bar{\Omega}} u_0(x) &= \phi_M, \\ \max_{\bar{\Omega}} v_0(x) &= \psi_M. \end{aligned} \quad (12)$$

Obviously, $(\underline{u}(t, x), \underline{v}(t, x)) = (0, 0)$ and $(\bar{u}(t), \bar{v}(t)) = (u(t; \phi_M, \psi_M), v(t; \phi_M, \psi_M))$ are a pair of lower-solution and upper-solution to system (7). Therefore, according to the Theorem 8.3.3 in [22] or Theorem 5.3.2 in [23], system (7)

has a unique globally defined solution $(u(x, t), v(x, t))$ which satisfies

$$\begin{aligned} 0 &\leq u(x, t) \leq u(t; \phi_M, \psi_M), \\ 0 &\leq v(x, t) \leq v(t; \phi_M, \psi_M). \end{aligned} \tag{13}$$

The strong maximum principle implies that $u(x, t), v(x, t) > 0$ when $t > 0$ for all $x \in \bar{\Omega}$.

(b) By the first equation of (7), we easily obtain the fact that $r_1 u(1-u/K) - \rho u^\alpha v \leq r_1 u(1-u/K)$ in $[0, +\infty) \times \Omega$; the first result follows easily from the simple comparison argument for parabolic problems, and thus there exists $T \in (0, +\infty)$ such that $u(t, x) \leq K + \epsilon$ in $[T, +\infty) \times \Omega$ for an arbitrary constant $\epsilon > 0$.

For the estimate of $v(x, t)$, let $U(t) = \int_{\Omega} u(x, t) dx$, $V(t) = \int_{\Omega} v(x, t) dx$; then

$$\begin{aligned} \frac{dU}{dt} &= \int_{\Omega} u_t dx \\ &= d_1 \int_{\Omega} \Delta u dx + \int_{\Omega} \left[r_1 u \left(u - \frac{u}{K} \right) - \rho u^\alpha v \right] dx \\ &= \int_{\Omega} \left[r_1 u \left(u - \frac{u}{K} \right) - \rho u^\alpha v \right] dx, \\ \frac{dV}{dt} &= \int_{\Omega} v_t dx = d_2 \int_{\Omega} \Delta v dx + \int_{\Omega} (\beta u^\alpha v - r_2 v) dx \\ &= \int_{\Omega} (\beta u^\alpha v - r_2 v) dx. \end{aligned} \tag{14}$$

Multiplying (14) by β/ρ and adding it to (15), we have

$$\begin{aligned} \left(\frac{\beta}{\rho} U + V \right)_t &= -r_2 V + \frac{\beta r_1}{\rho} \int_{\Omega} u \left(u - \frac{u}{K} \right) dx \\ &\leq -r_2 \left(\frac{\beta}{\rho} U + V \right) \\ &\quad + \left(\frac{\beta r_2}{\rho} + \frac{\beta r_1}{\rho} \right) \rho (K + \epsilon) |\Omega|. \end{aligned} \tag{16}$$

Integration of the inequality leads to

$$\begin{aligned} \int_{\Omega} v(x, t) dx = V(t) &< \frac{\beta}{\rho} U(t) + V(t) \\ &\leq \frac{\beta}{\rho} \left(1 + \frac{r_1}{r_2} \right) (K + \epsilon) |\Omega|. \end{aligned} \tag{17}$$

□

3. Stability of the Nonnegative Constant Steady States of (7)

In this section, we will analyze the stability of nonnegative constant steady states of (7). By the direct computation, we

see that the possible nonnegative constant steady states of (7) are

$$\begin{aligned} E_0 &= (0, 0), \\ E_1 &= (K, 0), \\ E^* &= (u^*, v^*), \end{aligned} \tag{18}$$

where $u^* = (r_2/\beta)^{1/\alpha}$, $v^* = r_1 u^{*(1-\alpha)} (1-u^*/K)/\rho$. Obviously, the positive constant steady state E^* exists if $r_2 < \beta K^\alpha$ holds.

Notation 1. Let $0 = \mu_0 < \mu_1 < \mu_2 < \dots < \mu_n < \dots \rightarrow \infty$ be the eigenvalues of $-\Delta$ on Ω under homogeneous Neumann boundary condition. We define the following space decomposition:

- (i) $S(\mu_n)$ is the space of eigenfunctions corresponding to μ_n for $n = 0, 1, 2, \dots$
- (ii) $X_{ij} := \{c \cdot \phi_{ij} : c \in \mathbb{R}^2\}$, where $\{\phi_{ij}\}$ are orthonormal basis of $S(\mu_n)$ for $j = 1, 2, \dots, \dim[S(\mu_n)]$.
- (iii) $\mathbf{X} := \{u = (u, v) \in [C^1(\bar{\Omega})]^2 : \partial u/\partial \mathbf{n} = \partial v/\partial \mathbf{n} = 0\}$, and so $\mathbf{X} = \oplus_{i=1}^{\infty} \mathbf{X}_i$, where $\mathbf{X}_i = \oplus_{j=1}^{\dim[S(\mu_i)]} \mathbf{X}_{ij}$.

Let \tilde{E} be a nonnegative constant steady state of (7); then the linearization of (7) at a constant solution \tilde{E} can be expressed by

$$\mathbf{u}_t = (D\Delta + J) \mathbf{u}, \tag{19}$$

where $D = \text{diag}(d_1, d_2)$, $\mathbf{u} = (u(x, t), v(x, t))^T$, and

$$J = \begin{pmatrix} r_1 - \frac{2r_1 u^*}{K} - \rho \alpha u^{*(\alpha-1)} v^* & -\rho u^{*\alpha} \\ -\beta v^* u^{*(\alpha-1)} & \beta u^{*\alpha} - r_2 \end{pmatrix}. \tag{20}$$

In view of Notation 1, we can induce the eigenvalues of system (19) confined on the subspace \mathbf{X}_i . If λ is an eigenvalue of (19) on \mathbf{X}_i , it must be an eigenvalue of the matrix $-\mu_n D + J$ for each $n \geq 0$. It is easy to see that λ satisfies the characteristic equation

$$\begin{aligned} \lambda^2 &+ \left(d_1 \mu_n + d_2 \mu_n - r_1 + \frac{2r_1 u^*}{K} + \rho \alpha u^{*(\alpha-1)} v^* \right) \lambda \\ &+ d_2 \mu_n \left(d_1 \mu_n - r_1 + \frac{2r_1 u^*}{K} + \rho \alpha u^{*(\alpha-1)} v^* \right) \\ &+ \rho \beta u^{*(2\alpha-1)} v^* = 0. \end{aligned} \tag{21}$$

- Theorem 2.** (i) The trivial equilibrium $E_0 = (0, 0)$ is unstable.
 (ii) If $\beta K^\alpha < r_2$, then $E_1 = (K, 0)$ is globally asymptotically stable.
 (iii) If $1 - \alpha - ((2 - \alpha)/K)u^* < 0$, then E^* is locally asymptotically stable.

Proof. (i) For $E_0 = (0, 0)$, the eigenvalues are

$$\begin{aligned} \lambda_{1n} &= r_1 - d_1 \mu_n, \\ \lambda_{2n} &= -d_2 \mu_n. \end{aligned} \tag{22}$$

Obviously, E_0 is unstable.

(ii) For $E_1 = (K, 0)$, the eigenvalues are

$$\begin{aligned}\lambda_{1n} &= -r_1 - d_1\mu_n, \\ \lambda_{2n} &= \beta K^\alpha - r_2 - d_2\mu_n.\end{aligned}\quad (23)$$

If $\beta K^\alpha < r_2$, then λ_{1n} and λ_{2n} are all negative. Therefore E_1 is locally asymptotically stable. Indeed, E_1 is globally asymptotically stable.

On account of Theorem 1, we have $\limsup_{t \rightarrow +\infty} \max_{\bar{\Omega}} u(\cdot, t) \leq K$, and thus there exists $T_1 \in (0, +\infty)$ such that, for an arbitrary constant $0 < \epsilon < (r_2/\beta)^{1/\alpha} - K$,

$$u(\cdot, t) \leq K + \epsilon, \quad t \geq T_1. \quad (24)$$

It follows from the second equation of (7) that

$$v_t - d_2\Delta v \leq v(\beta(K + \epsilon)^\alpha - r_2), \quad t \geq T_1. \quad (25)$$

Therefore, $\limsup_{t \rightarrow +\infty} \max_{\bar{\Omega}} v(\cdot, t) \leq 0$, and there exists $T_2 > T_1$ such that

$$v(\cdot, t) \leq \epsilon, \quad t \geq T_2. \quad (26)$$

It follows from the first equation of (7) that

$$u_t - d_1\Delta u \geq r_1u \left(1 - \frac{u}{K}\right) - \rho u^\alpha \epsilon, \quad t > T_2, \quad x \in \Omega. \quad (27)$$

On account of and the arbitrariness of $\epsilon > 0$, we have $\liminf_{t \rightarrow +\infty} \min_{\bar{\Omega}} u(\cdot, t) \geq K$. This combined with $\limsup_{t \rightarrow +\infty} \max_{\bar{\Omega}} u(\cdot, t) \leq K$ allows us to derive

$$\lim_{t \rightarrow +\infty} \max_{\bar{\Omega}} |u(\cdot, t) - K| = 0. \quad (28)$$

Hence, E_1 is globally asymptotically stable when $r_2 > K^\alpha \beta$.

(iii) When $E^* = (u^*, v^*)$ exists, the corresponding characteristic equation is as follows:

$$\begin{aligned}\lambda^2 &+ \left(d_1\mu_n + d_2\mu_n - r_1 \left(1 - \alpha - \frac{2-\alpha}{K}u^*\right)\right)\lambda \\ &+ d_2\mu_n \left(d_1\mu_n - r_1 \left(1 - \alpha - \frac{2-\alpha}{K}u^*\right)\right) \\ &+ \rho\beta u^{*(2\alpha-1)}v^* = 0.\end{aligned}\quad (29)$$

Obviously, we have

$$\begin{aligned}\lambda_{1n} + \lambda_{2n} &= r_1 \left(1 - \alpha - \frac{2-\alpha}{K}u^*\right) - d_1\mu_n - d_2\mu_n, \\ \lambda_{1n}\lambda_{2n} &= d_2\mu_n \left(d_1\mu_n - r_1 \left(1 - \alpha - \frac{2-\alpha}{K}u^*\right)\right) \\ &+ \rho\beta u^{*(2\alpha-1)}v^*.\end{aligned}\quad (30)$$

If $1 - \alpha - ((2-\alpha)/K)u^* < 0$, then $\lambda_{1n} + \lambda_{2n} < 0$ and $\lambda_{1n}\lambda_{2n} > 0$. Hence, all the roots of (29) have negative real part which means that E^* is locally asymptotically stable when $1 - \alpha - ((2-\alpha)/K)u^* < 0$. \square

4. The Prior Estimate

In this section, we will give some a priori estimates of positive solutions to (8). Firstly, we give two known lemmas.

Lemma 3 (Harnack inequality (cf. [24])). *Let $\omega \in C^2(\Omega) \cap C^1(\bar{\Omega})$ be a positive classical solution to*

$$\begin{aligned}\Delta\omega(x) + c(x)\omega(x) &= 0 \quad \text{in } \Omega, \\ \frac{\partial\omega}{\partial n} &= 0 \quad \text{on } \partial\Omega.\end{aligned}\quad (31)$$

Then there exists a positive constant C such that

$$\max_{\bar{\Omega}} \omega \geq C \min_{\bar{\Omega}} \omega. \quad (32)$$

Lemma 4 (maximum principle (cf. [25])). *Suppose that $g \in C(\bar{\Omega} \times \mathbb{R})$.*

(i) *Assume that $\omega \in C^2(\Omega) \cap C^1(\bar{\Omega})$ satisfies*

$$\begin{aligned}\Delta\omega(x) + c(x)\omega(x) &\geq 0 \quad \text{in } \Omega, \\ \frac{\partial\omega}{\partial n} &\leq 0 \quad \text{on } \partial\Omega.\end{aligned}\quad (33)$$

If $\omega(x_0) = \max_{\bar{\Omega}} \omega$, then $g(x_0, \omega(x_0)) \geq 0$.

(ii) *Assume that $\omega \in C^2(\Omega) \cap C^1(\bar{\Omega})$ satisfies*

$$\begin{aligned}\Delta\omega(x) + c(x)\omega(x) &\leq 0 \quad \text{in } \Omega, \\ \frac{\partial\omega}{\partial n} &\geq 0 \quad \text{on } \partial\Omega.\end{aligned}\quad (34)$$

If $\omega(x_0) = \min_{\bar{\Omega}} \omega$, then $g(x_0, \omega(x_0)) \leq 0$.

Lemma 5. *For any positive solution (u, v) of system (8),*

$$\begin{aligned}0 &< u(x) < K, \\ 0 &< v(x) < \frac{\beta}{\rho} \left(\frac{d_1}{d_2} + \frac{r_1}{r_2}\right) K\end{aligned}\quad (35)$$

for any $x \in \bar{\Omega}$.

Proof. Form Lemma 4, $u(x) \leq K$ and from the strong maximum principle $u(x) < K$ for all $x \in \bar{\Omega}$. Multiplying the first equation of (8) by β/ρ and adding it to the second equation, we have

$$\begin{aligned}-\left(\frac{\beta}{\rho}d_1\Delta u + d_2\Delta v\right) &= \frac{\beta}{\rho}r_1u \left(1 - \frac{u}{K}\right) - r_2v \\ &\leq \frac{\beta}{\rho}r_1K + \frac{d_1r_2\beta}{\rho d_2}K \\ &\quad - \frac{r_2}{d_2} \left(\frac{\beta}{\rho}d_1u + d_2v\right).\end{aligned}\quad (36)$$

Then the maximum principle implies that

$$\frac{\beta}{\rho}d_1u + d_2v \leq \frac{\beta}{\rho} \left(\frac{d_2r_1}{r_2} + d_1\right) K. \quad (37)$$

Hence, $v(x) < (\beta/\rho)(d_1/d_2 + r_1/r_2)K$. \square

In the following, we estimate the positive lower bound of positive solution of (8).

Theorem 6. *Let Ω be a bounded smooth domain in R^n . There exist two positive constants $\underline{C} < \overline{C}$ depending possibly on $d_1, d_2, K, \beta, \alpha, \rho$, and Ω , such that such that any positive solution $(u(x), v(x))$ of system (8) satisfies*

$$\begin{aligned} \underline{C} &\leq u(x) \leq \overline{C}, \\ \underline{C} &\leq v(x) \leq \overline{C} \end{aligned} \quad (38)$$

for any $x \in \overline{\Omega}$.

Proof. From Lemma 5, we obtain

$$u(x), v(x) \leq \overline{C} := \max \left\{ K, \frac{\beta}{\rho} \left(\frac{d_1}{d_2} + \frac{r_1}{r_2} \right) K \right\}, \quad (39)$$

where \overline{C} depends on $d_1, d_2, K, \beta, \alpha$, and ρ .

From Lemma 3, we obtain the fact that there exists a positive constant C_2 such that

$$\begin{aligned} \sup_{\overline{\Omega}} u(x) &\leq C_2 \inf_{\overline{\Omega}} u(x), \\ \sup_{\overline{\Omega}} v(x) &\leq C_2 \inf_{\overline{\Omega}} v(x). \end{aligned} \quad (40)$$

On the contrary, suppose the result is false. Then there exists a sequence $\{(u_n, v_n)\}$ of positive solutions to system (8) such that

$$\begin{aligned} \sup_{\overline{\Omega}} u_n &\longrightarrow 0 \\ \text{or } \sup_{\overline{\Omega}} v_n &\longrightarrow 0 \end{aligned} \quad (41)$$

as $n \longrightarrow +\infty$.

By the regularity theory for elliptic equations, there exists a subsequence of $\{(u_n, v_n)\}$, which will be denoted again by $\{(u_n, v_n)\}$, such that $\{(u_n, v_n)\} \rightarrow (u_0, v_0)$ in $C^2(\overline{\Omega})$ as $n \rightarrow +\infty$. Observe that $u_0 \leq K$ and, from (41), either $u_0 \equiv 0$ or $v_0 \equiv 0$. Therefore, we have the following two cases:

- (i) $u_0 \equiv 0, v_0 \neq 0$; or $u_0 \equiv 0, v_0 \equiv 0$.
- (ii) $u_0 \neq 0, v_0 \equiv 0$.

Since $\{(u_n, v_n)\}$ is a positive solution of (8), one can obtain the following integral equation by integrating (8) for u_n and v_n over Ω :

$$\begin{aligned} \int_{\Omega} \left(r_1 u_n \left(1 - \frac{u_n}{K} \right) - \rho u_n^\alpha v_n \right) dx &= 0, \\ \int_{\Omega} (\beta u_n^\alpha v_n - r_2 v_n) dx &= 0. \end{aligned} \quad (42)$$

- (i) In this case, $u_0 \equiv 0$; then

$$\beta u_n^\alpha - r_2 \longrightarrow -r_2 < 0 \quad (43)$$

uniformly as $n \rightarrow \infty$ and $v_n > 0$; then for sufficiently large n , we have

$$\int_{\Omega} (\beta u_n^\alpha v_n - r_2 v_n) dx < 0, \quad (44)$$

which is a contradiction.

- (ii) If $u_0 \neq 0, v_0 \equiv 0$, then this implies that u_0 satisfies (8). So $u_0 \equiv K$ for large n . Thus

$$\beta u_n^\alpha - r_2 \longrightarrow \beta K^\alpha - r_2 > 0 \quad (45)$$

for large n since $\beta K^\alpha < r_2$, which derives a contradiction again to the second integral equation of (42). This completes the proof. \square

5. Nonexistence of Nonconstant Positive Steady States

In this section, we can show the nonexistence of nonconstant positive solutions to system (8) when the diffusion coefficients d_1 and d_2 are large.

Theorem 7. *There exists a positive constant d^* such that elliptic problem (8) has no nonconstant positive solution if $\min\{d_1, d_2\} > d^*$.*

Proof. Suppose that $(u(x), v(x))$ is a nonconstant positive solution of system (8). Denote $\bar{u} = |\Omega|^{-1} \int_{\Omega} u(x) dx \geq 0$, $\bar{v} = |\Omega|^{-1} \int_{\Omega} v(x) dx \geq 0$. Then

$$\begin{aligned} \int_{\Omega} (u - \bar{u}) dx &= 0, \\ \int_{\Omega} (v - \bar{v}) dx &= 0. \end{aligned} \quad (46)$$

Define

$$h(u) = \frac{\bar{u} - u}{\bar{u}^\alpha - u^\alpha} \quad (47)$$

for $\bar{u} \neq u$. Indeed, we can prove that $h(u) > 0$ and $h'(u) > 0$. In fact, notice

$$h'(u) = \frac{(1 - \alpha) u^\alpha + \alpha \bar{u} u^{(\alpha-1)} - \bar{u}^\alpha}{(\bar{u}^\alpha - u^\alpha)^2}. \quad (48)$$

Let

$$h_1(u) = (1 - \alpha) u^\alpha + \alpha \bar{u} u^{(\alpha-1)} - \bar{u}^\alpha, \quad (49)$$

and we have

$$h_1'(u) = \alpha(1 - \alpha) u^{(\alpha-2)} (u - \bar{u}), \quad (50)$$

which implies that $h_1(u) > \min h_1(u) = h_1(\bar{u}) = 0$ for $u \neq \bar{u}$. Therefore, we obtain the fact that $h(u) > 0$.

Furthermore, multiplying the first equation of (8) by β/ρ , adding it to the second equation, and integrating over Ω , we get

$$\begin{aligned} - \int_{\Omega} \left(\frac{\beta}{\rho} d_1 \Delta u + d_2 \Delta v \right) dx \\ = \int_{\Omega} \left(\frac{\beta}{\rho} r_1 u \left(1 - \frac{u}{K} \right) - r_2 v \right) dx, \end{aligned} \quad (51)$$

and then the Neumann boundary conditions lead to

$$r_2 \int_{\Omega} v dx = \int_{\Omega} \frac{\beta}{\rho} r_1 u \left(1 - \frac{u}{K}\right) dx \leq \frac{K\beta r_1}{4\rho} |\Omega|. \quad (52)$$

Thus

$$\bar{v} = \frac{1}{\Omega} \int_{\Omega} v dx \leq \frac{K\beta r_1}{4r_2\rho}. \quad (53)$$

Multiplying the first equation in (8) by $u - \bar{u}$, we have

$$\begin{aligned} d_1 \int_{\Omega} |\nabla(u - \bar{u})|^2 dx &= \int_{\Omega} (u - \bar{u}) \left(r_1 u \left(1 - \frac{u}{K}\right) \right. \\ &\quad \left. - \rho u^\alpha v \right) dx = \int_{\Omega} (u - \bar{u}) \left(r_1 (u - \bar{u}) \right. \\ &\quad \left. - \frac{r_1}{K} (u^2 - \bar{u}^2) - \rho u^\alpha v + \rho \bar{u}^\alpha \bar{v} \right) dx \\ &\leq 3r_1 \int_{\Omega} (u - \bar{u})^2 dx \\ &\quad + \rho K^\alpha \int_{\Omega} |(u - \bar{u})(v - \bar{v})| dx + \rho \int_{\Omega} \frac{\bar{v}}{h(u)} (u \\ &\quad - \bar{u})^2 dx \leq \left(3r_1 + \frac{\rho K^\alpha}{2} + \frac{K\beta r_1}{4r_2 h(0)} \right) \int_{\Omega} (u \\ &\quad - \bar{u})^2 dx + \frac{\rho K^\alpha}{2} \int_{\Omega} (v - \bar{v})^2 dx. \end{aligned} \quad (54)$$

Multiplying the second equation in (8) by $v - \bar{v}$, we have

$$\begin{aligned} d_2 \int_{\Omega} |\nabla(v - \bar{v})|^2 dx &= \int_{\Omega} (v - \bar{v}) (\beta u^\alpha v - r_2 v) dx \\ &= \int_{\Omega} (v - \bar{v}) (-r_2 (v - \bar{v}) + \beta u^\alpha v - \beta \bar{u}^\alpha \bar{v}) dx \\ &\leq \left(r_2 + \beta K^\alpha + \frac{K\beta^2 r_1}{8r_2 \rho h(0)} \right) \int_{\Omega} (v - \bar{v})^2 dx \\ &\quad + \frac{K\beta^2 r_1}{8r_2 \rho h(0)} \int_{\Omega} (u - \bar{u})^2 dx. \end{aligned} \quad (55)$$

From (54) and (55) and the Poincaré inequality, we obtain

$$\begin{aligned} d_1 \int_{\Omega} |\nabla(u - \bar{u})|^2 dx + d_2 \int_{\Omega} |\nabla(v - \bar{v})|^2 dx \\ \leq A \int_{\Omega} (u - \bar{u})^2 dx + B \int_{\Omega} (v - \bar{v})^2 dx \\ \leq \frac{1}{\mu_1} \left(A \int_{\Omega} |\nabla(u - \bar{u})|^2 dx + B \int_{\Omega} |\nabla(v - \bar{v})|^2 dx \right), \end{aligned} \quad (56)$$

where

$$\begin{aligned} A &= 3r_1 + \frac{\rho K^\alpha}{2} + \frac{K\beta r_1}{4r_2 h(0)} + \frac{K\beta^2 r_1}{8r_2 \rho h(0)}, \\ B &= r_2 + \beta K^\alpha + \frac{K\beta^2 r_1}{8r_2 \rho h(0)} + \frac{\rho K^\alpha}{2}. \end{aligned} \quad (57)$$

Hence, if

$$\min \{d_1, d_2\} > d^* := \frac{1}{\mu_1} \max \{A, B\}, \quad (58)$$

then

$$\nabla(u - \bar{u}) = \nabla(v - \bar{v}) = 0, \quad (59)$$

and (u, v) must be a constant solution. \square

6. Existence of Nonconstant Positive Steady States

In this subsection, we discuss the existence of nonconstant positive solutions to system (8) when the diffusion coefficients d_1 and d_2 vary while the parameters r_1 , K , α , ρ , β , and r_2 are fixed by using the Leray-Schauder degree theory. Throughout this section, we assume that the positive constant steady state $E^* = (u^*, v^*)$ exists.

For simplicity, denote $\mathbf{u} = (u, v)$ and

$$\begin{aligned} a_1 &= r_1 - \frac{2r_1 u^*}{K} - \rho \alpha (u^*)^{(\alpha-1)} v^*, \\ a_2 &= \rho (u^*)^\alpha, \\ a_3 &= \beta v^* (u^*)^{(\alpha-1)}, \\ \Phi(\mathbf{u}) &= \begin{pmatrix} r_1 u \left(1 - \frac{u}{K}\right) - \rho u^\alpha v \\ \beta u^\alpha v - r_2 v \end{pmatrix}, \\ \Phi_{\mathbf{u}}(E^*) &= \begin{pmatrix} a_1 & -a_2 \\ a_3 & 0 \end{pmatrix}, \\ D &= \text{diag} \{d_1, d_2\}, \\ X &= \left\{ (u, v) \in [C^1(\bar{\Omega})]^2 : \frac{\partial u}{\partial n} = \frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega \right\}, \\ X^+ &= \{(u, v) : u, v \geq 0, (u, v) \in X\}, \\ \Lambda &= \{(u, v) \in X : C^{-1} < u, v < C \text{ on } \bar{\Omega}\}. \end{aligned} \quad (60)$$

Thus, (8) can be written as

$$-D\Delta \mathbf{u} = \Phi(\mathbf{u}) \quad \text{in } \frac{\partial \mathbf{u}}{\partial n} = 0 \text{ on } \Omega, \quad (61)$$

and, obviously, \mathbf{u} is a positive solution of (61) if and only if

$$\mathcal{F}(\mathbf{u}) := \mathbf{u} - (I - \Delta)^{-1} (D^{-1} \Phi(\mathbf{u}) + \mathbf{u}) \quad \text{in } X^+, \quad (62)$$

where $(I - \Delta)^{-1}$ is the inverse of $I - \Delta$ with the homogeneous Neumann boundary condition. As $\mathcal{F}(\cdot)$ is a compact perturbation of the identity operator, the Leray-Schauder degree $\text{deg}(\mathcal{F}(\cdot), \Lambda, 0)$ is well-defined from Theorem 6. By direct computation, we have

$$\mathcal{F}_{\mathbf{u}}(E^*) = I - (I - \Delta)^{-1} (D^{-1} \Phi_{\mathbf{u}}(E^*) + I). \quad (63)$$

If $\mathcal{F}_{\mathbf{u}}(E^*)$ is invertible, the index of \mathcal{F} is defined as

$$\text{index}(\mathcal{F}(\cdot), E^*) = (-1)^\gamma, \quad (64)$$

where γ is the number of negative eigenvalues of $\mathcal{F}_{\mathbf{u}}(E^*)$. Note that λ is an eigenvalue of $\mathcal{F}_{\mathbf{u}}(E^*)$ on X_j if and only if it is an eigenvalue of the matrix

$$\begin{aligned} B_j &= I - \frac{1}{1 + \mu_j} [D^{-1}\Phi_{\mathbf{u}}(E^*) + I] \\ &= \frac{1}{1 + \mu_j} [\mu_j I - D^{-1}\Phi_{\mathbf{u}}(E^*)]. \end{aligned} \quad (65)$$

Thus $\mathcal{F}_{\mathbf{u}}(E^*)$ is invertible if and only if, for all $j \geq 0$, the matrix B_j is nonsingular. Writing

$$\begin{aligned} H(d_1, d_2; \mu) &= \det[\mu I - D^{-1}\Phi_{\mathbf{u}}(E^*)] \\ &= \frac{1}{d_1 d_2} \det[\mu D - \Phi_{\mathbf{u}}(E^*)], \end{aligned} \quad (66)$$

we have that if $H(d_1, d_2; \mu) \neq 0$, then $H(d_1, d_2; \mu) < 0$ if and only if the number of negative eigenvalues of $\mathcal{F}_{\mathbf{u}}(E^*)$ in X_j is odd. The following lemma gives the explicit formula of calculating the index.

Lemma 8. *If $H(d_1, d_2; \mu_i) \neq 0$ for all $i \geq 0$, then*

$$\begin{aligned} \text{index}(\mathcal{F}(\cdot), E^*) &= (-1)^\gamma, \\ \gamma &= \sum_{i \geq 0, H(d_1, d_2; \mu_i) < 0} m(\mu_i), \end{aligned} \quad (67)$$

where $m(\mu_i)$ is the algebraic multiplicity of μ_i .

To facilitate our computation of $\text{deg}(\mathcal{F}(\cdot), E^*)$, we only need consider the sign of $\det[\mu D - \Phi_{\mathbf{u}}(E^*)]$. The direct calculation gives

$$\det[\mu D - \Phi_{\mathbf{u}}(E^*)] = d_1 d_2 \mu^2 - a_1 d_2 \mu + a_2 a_3. \quad (68)$$

Obviously, nonnegative roots of (68) exist if and only if $a_1^2 d_2 - 4d_1 a_2 a_3 > 0$ and $a_1 > 0$. Assume that μ^+ and μ^- are the two roots of (68), we have the following conclusion.

Theorem 9. *Assuming that $\beta K^\alpha > r_2$,*

$$\frac{d_2}{d_1} > \frac{4a_2 a_3}{a_1^2}, \quad a_1 > 0, \quad (69)$$

and there exist $i, j \in N$, such that $0 \leq \mu_j < \mu^- < \mu_{j+1} \leq \mu_i < \mu^+ < \mu_{i+1}$ and $\sum_{k=j+1}^i m(\mu_k)$ is odd, then (8) has at least one nonconstant positive solution.

Proof. For $t \in [0, 1]$, we define

$$\mathcal{A}_t(\mathbf{u}) \triangleq (-\Delta + I)^{-1} \begin{pmatrix} u + \left(\frac{1-t}{d^*} + \frac{t}{d_1} \right) f(u, v) \\ v + \left(\frac{1-t}{d^*} + \frac{t}{d_2} \right) g(u, v) \end{pmatrix}, \quad (70)$$

where d^* is defined in Theorem 7.

The positive solutions of the problem

$$\begin{aligned} \mathcal{A}_t(\mathbf{u}) &= \mathbf{u} \quad \text{in } \Omega, \\ \frac{\partial \mathbf{u}}{\partial n} &= 0 \quad \text{on } \partial\Omega \end{aligned} \quad (71)$$

are contained in Λ . Note that \mathbf{u} is a positive solution of system (8) if and only if it is a positive solution of (71) with $t = 1$. \mathbf{u}^* is the unique positive constant solution of (71) for any $t \in [0, 1]$. According to the choice of d^* in Theorem 7, we have E^* which is the only fixed point of \mathcal{A}_0 .

$$\text{deg}(I - \mathcal{A}_0, \Lambda, 0) = \text{index}(I - \mathcal{A}_0, \Lambda, E^*) = 1. \quad (72)$$

Since $\mathcal{F} = I - H(\cdot, 1)$ and if (8) has no other solutions except the constant one E^* , then we have

$$\begin{aligned} \text{deg}(I - \mathcal{A}_1, \Lambda, (0, 0)) &= \text{index}(\mathcal{F}, E^*) \\ &= (-1)^{\sum_{k=j+1}^i m(\mu_k)} = -1. \end{aligned} \quad (73)$$

On the other hand, by the homotopy invariance of the topological degree,

$$\text{deg}(I - \mathcal{A}_0, \Lambda, 0) = \text{deg}(I - \mathcal{A}_1, \Lambda, 0), \quad (74)$$

which is a contradiction. Therefore, there exists at least one nonconstant solution of (8). \square

7. Numerical Simulation

7.1. Global Stability of Equilibrium E_1 . Consider system (7) with following parameters: $d_1 = 0.8$, $d_2 = 0.9$, $r_1 = 0.9$, $\beta = 0.1$, $r_2 = 0.2$, $K = 2$, $\rho = 0.1$, and $\alpha = 2/3$. According to the discussions in Section 3, the steady state E_1 is globally asymptotically stable; see Figure 1.

7.2. Stability of Positive Steady State E^* . Consider system (7) with following parameters: $d_1 = 0.8$, $d_2 = 0.9$, $r_1 = 0.9$, $\beta = 0.3$, $r_2 = 0.2$, $K = 2$, $\rho = 0.1$, and $\alpha = 2/3$. According to the discussions in Section 3, the positive steady state E^* is locally asymptotically stable; see Figure 2.

8. Conclusions

In this paper, we have investigated the existence/nonexistence of nonconstant positive steady states for a diffusive predator-prey system with a group defense for prey under Neumann boundary conditions. The existence results provide a theoretical support for pattern formation caused by diffusion. We also study the stability of nonnegative equilibria and obtain the fact that E_1 is globally asymptotically stable when $\beta K^\alpha < r_2$. In fact, the positive steady state does not exist at this time. If $\beta K^\alpha > r_2$ and $1 - \alpha - ((2 - \alpha)/K)u^* < 0$, then the positive steady state E^* is locally asymptotically stable. It is easily obtained that when $1 - \alpha - ((2 - \alpha)/K)u^* = 0$, characteristic equation (29) has a pair of purely imaginary roots. Therefore, system (7) occurs with Hopf bifurcation, as shown in Figure 3.

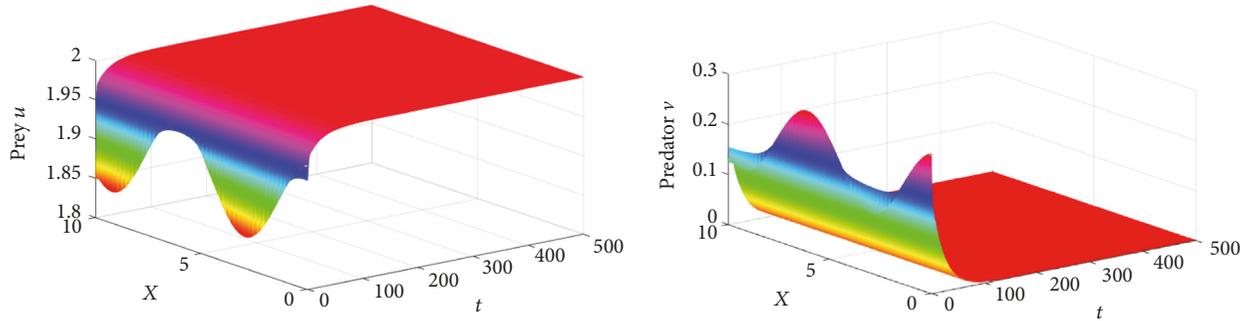


FIGURE 1: The steady state E_1 is globally asymptotically stable.

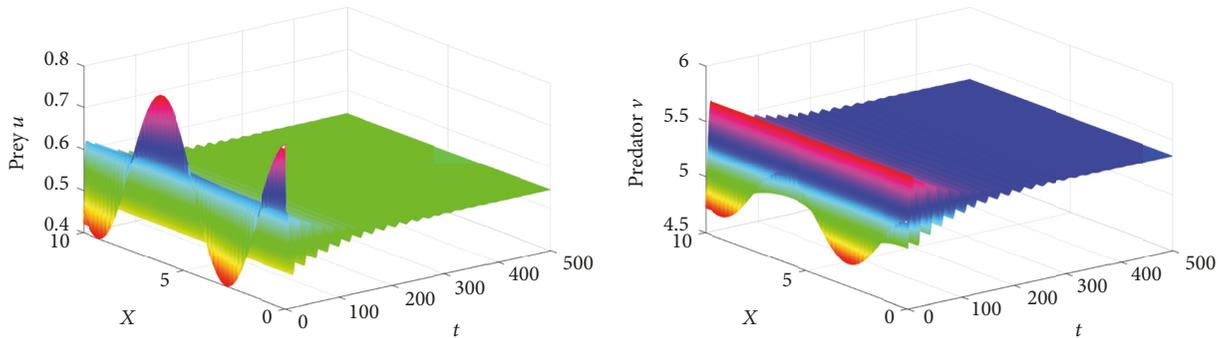


FIGURE 2: The positive steady state E^* is locally asymptotically stable.

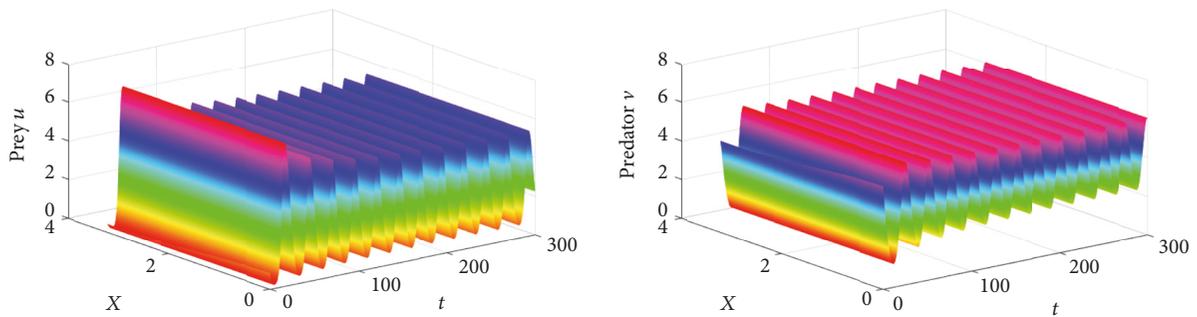


FIGURE 3: Hopf bifurcation occurs with parameters $r_1 = 0.9$, $\alpha = 2/3$, $\beta = 0.1$, $r_2 = 0.18$, $K = 10$; $\rho = 0.2$.

Conflicts of Interest

The authors declare that they have no conflicts of interest.

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