

LINEAR ALGEBRA AND DIFFERENTIAL GEOMETRY ON ABSTRACT HILBERT SPACE

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Isomorphisms of separable Hilbert spaces are analogous to isomorphisms of n -dimensional vector spaces. However, while n -dimensional spaces in applications are always realized as the Euclidean space R^n , Hilbert spaces admit various useful realizations as spaces of functions. In the paper this simple observation is used to construct a fruitful formalism of local coordinates on Hilbert manifolds. Images of charts on manifolds in the formalism are allowed to belong to arbitrary Hilbert spaces of functions including spaces of generalized functions. Tensor equations then describe families of functional equations on various spaces of functions. The formalism itself and its applications in linear algebra, differential equations, and differential geometry are carefully analyzed.

1. Introduction

Various integral transforms are known to be extremely useful in analysis and applications. One example is the Fourier transform in analysis and applied problems; another example is the Segal-Bargmann transform [1] in quantum theory. An important common property of integral transforms is that they relate various spaces of functions and various operators on these spaces and allow one to “transplant” a problem from one space to another. Because of that, the problem at hand may become easier to solve. A somewhat similar situation arises when working with tensor equations in the finite-dimensional setting. Namely, by an appropriate choice of coordinates one can significantly simplify a given tensor equation. Although the analogy is obvious, an infinite-dimensional setting offers a significantly larger variety of situations. In particular, by using the Segal-Bargmann transform, one can relate problems on spaces of ordinary or even generalized functions to problems on spaces of holomorphic functions.

In the paper we attempt to build a systematic approach to functional transformations on Hilbert spaces based on the above mentioned analogy between integral transforms and changes of coordinates on an n -dimensional manifold. Initial results in this direction were announced in [5, 6]. Various applications of the formalism to quantum theory, especially to the problem of emergence of the classical space-time, were considered in [5, 7, 8]. The goal here is to approach the subject in a more formal mathematical way and to justify the

previously obtained results. The reader is referred to [9] for an introduction to Hilbert spaces and applications.

We begin by describing a method of building various Hilbert spaces of functions including spaces of smooth and generalized functions. Then a simple formalism based on isomorphisms of these spaces is constructed. The formalism allows one to “move” in a systematic way between Hilbert spaces of functions and at the same time to formulate problems in a way independent of any particular functional realization. The formalism is then applied to building linear algebra on Hilbert spaces, which deals with the ordinary and the generalized functions on an equal footing. After discussing various special transformations preserving properties of differential operators, we concentrate on differential geometry of Hilbert spaces of functions. Namely, using the developed formalism we find a natural isometric embedding of finite dimensional Riemannian manifolds into Hilbert spaces of functions. Finally, the formalism is applied to construct a Riemannian metric on the unit sphere in a Hilbert space in such a way that solutions of Schrödinger equation are geodesics on the sphere.

2. Hilbert spaces of C^∞ -functions and their duals

We first discuss a general method of constructing various Hilbert spaces of smooth and generalized functions. Consider the convolution

$$(f * \rho)(x) = \int f(y)e^{-(x-y)^2} dy \tag{2.1}$$

of a function $f \in L_2(R)$ and the Gaussian function $\rho(x) = e^{-x^2}$. It is the standard result that such a convolution is in $C^\infty \cap L_2(R)$ and

$$D^p(f * \rho)(x) = (f * D^p\rho)(x), \tag{2.2}$$

for any order p of the derivative D .

THEOREM 2.1. *The linear operator $\rho : L_2(R) \rightarrow L_2(R)$ defined by*

$$\rho f = f * \rho \tag{2.3}$$

is a bounded invertible operator.

Proof. The operator ρ is bounded because $\|\rho * f\|_{L_2} \leq \|\rho\|_{L_1} \|f\|_{L_2}$. To check that it is invertible, assume that $(\rho * f)(x) = 0$ and differentiate both sides of this equation an arbitrary number of times p . We then conclude that all Fourier coefficients c_p of f in the orthonormal in $L_2(R)$ basis of Hermite functions vanish. □

The operator ρ induces the Hilbert metric on the image $H = \rho(L_2(R))$ by

$$(\varphi, \psi)_H = (\rho^{-1}\varphi, \rho^{-1}\psi)_{L_2}. \tag{2.4}$$

In particular, the space H with this metric is Hilbert.

THEOREM 2.2. *The embedding $H \subset L_2(R)$ is continuous and H is dense in $L_2(R)$.*

Proof. We have

$$\|\varphi\|_{L_2} = \|\rho f\|_{L_2} \leq \|\rho\|_{L_1} \|f\|_{L_2} = M\|\varphi\|_H, \tag{2.5}$$

where M is a constant. In addition, the functions $e^{-(x-a)^2}$, $a \in R$, being in H , form a complete system in $L_2(R)$. Therefore, H is dense in $L_2(R)$. \square

Provided we identify those linear functionals on H and $L_2(R)$ which are equal on H , we then have

$$L_2^*(R) \subset H^*. \tag{2.6}$$

Indeed, any functional continuous on L_2 will be continuous on H .

THEOREM 2.3. *When using the standard norm*

$$\|f\|_{H^*} = \sup_{\|\varphi\|_H \leq 1} |(f, \varphi)| \tag{2.7}$$

on H^* and similarly on L_2^* , we have a continuous embedding $L_2^*(R) \hookrightarrow H^*$.

Proof. Because $\|\varphi\|_{L_2} \leq M\|\varphi\|_H$, for any $f \in L_2^*$ we have

$$\|f\|_{H^*} \leq N\|f\|_{L_2^*}. \tag{2.8}$$

That is, we have a continuous embedding $L_2^*(R) \hookrightarrow H^*$. \square

By Riesz theorem, the norm on H^* is induced by the inner product

$$(f, g)_{H^*} = (\varphi, \psi)_H, \tag{2.9}$$

where $f = (\varphi, \cdot)_H = \widehat{G}\varphi$ and $g = (\psi, \cdot)_H = \widehat{G}\psi$ and $\widehat{G} : H \rightarrow H^*$ is the Riesz isomorphism. Similarly for L_2^* .

On the other hand, since $\rho : L_2 \rightarrow H$ is an isomorphism of Hilbert spaces, the adjoint operator $\rho^* : H^* \rightarrow L_2^*$, defined for any $f \in H^*$, $\varphi \in L_2$ by

$$(\rho^* f, \varphi) = (f, \rho\varphi), \tag{2.10}$$

is continuous and invertible (with $(\rho^*)^{-1} = (\rho^{-1})^*$).

THEOREM 2.4. *The operator ρ^* is an isomorphism of Hilbert spaces H^* and L_2^* with the above inner product (2.9). The metric on H^* can be defined by the isomorphism $\widehat{G}^{-1} : H^* \rightarrow H$, $\widehat{G}^{-1} = \rho\rho^*$.*

Proof. For any $f, g \in H^*$ the inner product $(f, g)_{H^*}$ induced by ρ^* is given by

$$(f, g)_{H^*} = (\rho^* f, \rho^* g)_{L_2^*} = (f, \rho\rho^* g), \tag{2.11}$$

where $\rho\rho^* : H^* \rightarrow H$ defines the metric on H^* . On the other hand, if $f = (\varphi, \cdot)_H$ and $g = (\psi, \cdot)_H$, then $f = \widehat{G}\varphi$, and $g = \widehat{G}\psi$, where $\widehat{G} : H \rightarrow H^*$ is the Riesz isomorphism.

Therefore, metric (2.9) on H^* can be written as follows:

$$(f, g)_{H^*} = (\varphi, \psi)_H = (\widehat{G}\varphi, \psi) = (\widehat{G}\widehat{G}^{-1}f, \widehat{G}^{-1}g) = (f, \widehat{G}^{-1}g). \tag{2.12}$$

From (2.4) we conclude that $\widehat{G}^{-1} = \rho\rho^*$. Also, because the inner products (2.11) and (2.12) coincide, ρ^* is an isomorphism of Hilbert spaces H^* and L_2^* . \square

Because the kernel of ρ is given by $\rho(x, y) = e^{-(x-y)^2}$, the kernel $k(x, y)$ of \widehat{G}^{-1} , computed up to a constant factor, is equal to

$$k(x, y) = e^{-(x-y)^2/2}. \tag{2.13}$$

Therefore, the inner product on H^* can be written as

$$(\varphi, \psi)_{H^*} = \int e^{-(x-y)^2/2} \varphi(x)\psi(y) dx dy. \tag{2.14}$$

Note that the integral symbol here denotes the action of the bilinear functional G^{-1} on $H^* \times H^*$. Only in special cases does this action coincide with Lebesgue integration with respect to x, y .

THEOREM 2.5. *The space H^* contains the pointwise evaluation functional (the delta-function) $\delta_a : \varphi \rightarrow \varphi(a)$.*

Proof. The sequence

$$f_L = \frac{L}{\sqrt{\pi}} e^{-L^2 x^2} \tag{2.15}$$

is fundamental in H^* and so it converges in H^* . Also, the sequence f_L is a delta-converging sequence, so that $(f_L, \varphi) \rightarrow (\delta_a, \varphi)$ as $L \rightarrow \infty$ for any $\varphi \in H$. In particular, f_L strongly converges to the evaluation functional δ_a in H^* . \square

Formally,

$$(\delta_a, \delta_b)_{H^*} = \int e^{-(x-y)^2/2} \delta(x-a)\delta(y-b) dx dy = e^{-(a-b)^2/2}. \tag{2.16}$$

Note that the inner product on $L_2(R)$ can be written in the form

$$(\varphi, \psi)_{L_2} = \int \delta(x-y)\varphi(x)\psi(y) dx dy, \tag{2.17}$$

where $\int \delta(x-y)\psi(y) dy$ denotes the convolution $(\delta * \psi)(x) = \psi(x)$. The main difference between metrics (2.14) and (2.17) is that the metric on $L_2(R)$ is “diagonal,” while the metric on H^* is not.

Consider now another Hilbert space obtained in a similar fashion. For this, note that since $H \hookrightarrow L_2(R)$, the Fourier transform σ on H is defined. Moreover, σ induces a Hilbert metric on the image $\tilde{H} = \sigma(H)$ by

$$(\sigma(\varphi), \sigma(\psi))_{\tilde{H}} = (\varphi, \psi)_H, \tag{2.18}$$

for any φ, ψ in H .

We can then define the Fourier transform on H^* in the standard way by

$$(\sigma(f), \sigma(\varphi)) = 2\pi(f, \varphi), \tag{2.19}$$

for any $f \in H^*$ and $\varphi \in H$. The image $\sigma(H^*)$ is the Hilbert space \tilde{H}^* dual to \tilde{H} . The metric on this space is induced by the map $\sigma\rho : L_2 \rightarrow \tilde{H}$:

$$(f, g)_{\tilde{H}^*} = ((\sigma\rho)^* f, (\sigma\rho)^* g)_{L_2^*}. \tag{2.20}$$

In other words, the metric on \tilde{H}^* is given by $\sigma\rho\rho^*\sigma^* : \tilde{H}^* \rightarrow \tilde{H}$. The kernel of this metric is then equal to

$$\tilde{k}(x, y) = \frac{1}{\sqrt{2\pi}} e^{-x^2/2} \delta(x - y), \tag{2.21}$$

where $\delta(x - y)$ is the kernel of the operator δ^* of convolution with the delta function. Note that because of the weight $e^{-x^2/2}$, the space \tilde{H}^* contains the plane wave functions e^{ikx} .

So, we have the Hilbert space H of C^∞ (in fact, analytic) functions, its dual H^* containing singular generalized functions, in particular, $\delta, D^p\delta$. We also have the Fourier transformed space \tilde{H} and its dual \tilde{H}^* , which contains, in particular, the plane wave function e^{ipx} .

The method used to construct the above spaces of generalized functions is very general. All we need is a linear injective map ρ on a space L_2 of square-integrable functions. By changing ρ , we change the metric \hat{G} on the Hilbert space $H = \rho(L_2)$ and therefore the metric \hat{G}^{-1} on the dual space H^* . Roughly speaking, by “smoothing” the kernel of the metric \hat{G}^{-1} , we extend the class of (generalized) functions for which the norm defined by this metric is finite. The same is true if we “improve” the behavior of the metric for large $|x|$. Conversely, by “spoiling the metric” we make the corresponding Hilbert space “poor” in terms of a variety of the elements of the space. This is further illustrated by the following theorem.

THEOREM 2.6. *If ρ is the map on $L_2(\mathbb{R})$ with the kernel $\rho(x, y) = e^{-(x-y)^2-x^2}$, then the Hilbert space $\rho(L_2(\mathbb{R}))$ with the induced metric is continuously embedded into the Schwartz space \mathcal{W} of C^∞ rapidly decreasing functions on \mathbb{R} .*

Proof. From Theorem 2.1 we immediately conclude that the map ρ is injective and the image $\rho(L_2(\mathbb{R}))$ consists of C^∞ functions. Moreover, we have

$$\left| \int f(y) e^{-(x-y)^2-x^2} dy \right| \leq M \|f\|_{L_2} e^{-x^2}, \tag{2.22}$$

where M is a constant. In particular, the functions in $\rho(L_2(\mathbb{R}))$ decrease faster than any power of $1/|x|$. Now, the topology on \mathcal{W} may be defined by the countable system of norms

$$\|\varphi\|_p = \sup_{x \in \mathbb{R}; k, q \leq p} |x^k \varphi^{(q)}(x)|, \tag{2.23}$$

where k, q, p are nonnegative integers and $\varphi^{(q)}$ is the derivative of φ of order q . For any $\varphi \in H$, by the Schwarz inequality, we have

$$\begin{aligned} |x^k \varphi^{(q)}(x)| &= \left| \int x^k f(y) \frac{d^q}{dx^q} e^{-(x-y)^2 - x^2} dy \right| \\ &\leq \left(\int x^{2k} P_q^2(x) e^{-2(x-y)^2 - 2x^2} dy \right)^{1/2} \|f\|_{L_2} = M_{k,q} \|f\|_{L_2}, \end{aligned} \tag{2.24}$$

where $\rho f = \varphi$, P_q is a polynomial of degree q and $M_{k,q}$ are constants depending only on k, q . This proves that topology on H is stronger than topology on W , that is, $H \subset W$ is a continuous embedding. \square

It follows from the theorem that $W^* \subset H^*$ as a set. In fact, any functional continuous on W will be continuous on H . Moreover, by choosing the strong topology on W^* one can show that the embedding of W^* into H^* is continuous. Notice also that one could choose the weak topology on W^* instead as the weak and strong topologies on W^* are equivalent [2].

Remarks. (1) The considered Hilbert spaces of generalized functions could have also been obtained by completing spaces of ordinary functions with respect to the discussed inner products. More generally, in many cases the expression

$$(f, g)_k = \int k(x, y) f(x) g(y) dx dy, \tag{2.25}$$

where k is a continuous function on a region D in R^{2n} , defines an inner product on $L_2(R^n)$. Then the completion with respect to the corresponding metric yields a Hilbert space containing singular generalized functions.

(2) The isomorphism ρ defined on $L_2(R)$ by $(\rho f)(x) = \int e^{-(x-y)^2} f(y) dy$ is closely related to the Segal and Bargmann transform $C_t : L_2(R^d) \rightarrow H(C^d)$, where $H(C^d)$ is the space of holomorphic functions on C^d , given by

$$(C_t f)(z) = \int (2\pi t)^{-d/2} e^{-(z-y)^2/2t} f(y) dy. \tag{2.26}$$

The theorem by Segal and Bargmann [1] (see also [4]) says the following.

THEOREM 2.7. *For each t the map C_t is an isomorphism of $L_2(R^d)$ onto $HL_2(C^d, \nu_t)$. Here $HL_2(C^d, \nu_t)$ denotes the space of holomorphic functions that are square-integrable with respect to the measure $\nu_t(z) dz$, where dz is the $2d$ -dimensional Lebesgue measure on C^d and the density ν_t is given by*

$$\nu_t(x + iy) = (\pi t)^{-d/2} e^{-y^2/t}, \quad x, y \in R^d. \tag{2.27}$$

(3) We also remark that Hilbert spaces on which the action of pointwise evaluation functional is continuous are often called *reproducing kernel Hilbert spaces*.

3. Generalized eigenvalue problem and functional tensor equations

Consider the *generalized eigenvalue problem* for a linear operator A on a Hilbert space H :

$$f(\hat{A}\varphi) = \lambda f(\varphi). \tag{3.1}$$

The problem consists in finding all functionals $f \in H^*$ and the corresponding numbers λ for which (3.1) is satisfied for all functions φ in a Hilbert space H of functions.

For instance, let \tilde{H} be the Hilbert space with the metric (2.18). As we know, the metric on the dual space \tilde{H}^* is given by the kernel $e^{-x^2/2}\delta(x - y)$. The generalized eigenvalue problem for the operator $\hat{A} = -i(d/dx)$ is given by

$$f\left(-i\frac{d}{dx}\varphi\right) = pf(\varphi). \tag{3.2}$$

Equation (3.2) must be satisfied for every φ in \tilde{H} . The functionals

$$f(x) = e^{ipx} \tag{3.3}$$

are the eigenvectors of A . Note that these eigenvectors belong to \tilde{H}^* so that the generalized eigenvalue problem has a solution.

As before, the Fourier transform $\sigma : \tilde{H} \rightarrow H$,

$$\psi(k) = (\sigma\varphi)(k) = \int \varphi(x)e^{ikx} dx, \tag{3.4}$$

induces a Hilbert structure on the space $H = \sigma(\tilde{H})$. Relative to this structure σ is an isomorphism of the Hilbert spaces \tilde{H} and H . The inverse transform $\omega = \sigma^{-1}$ is given by

$$(\omega\psi)(x) = \frac{1}{2\pi} \int \psi(k)e^{-ikx} dk. \tag{3.5}$$

Notice that the Fourier transform of e^{ipx} is $\delta(k - p)$ which is to say again that the space H^* dual to H contains delta-functions. As we know from (2.13), the kernel of the metric on H^* is proportional to $e^{-(1/2)(x-y)^2}$.

Using transformations σ, ω , we can rewrite the generalized eigenvalue problem (3.2) as

$$\omega^* f(\sigma\hat{A}\omega\psi) = p\omega^* f(\psi), \tag{3.6}$$

where $\omega\psi = \varphi$. We have

$$\hat{A}\omega\psi = -i\frac{d}{dx}\frac{1}{2\pi} \int \psi(k)e^{-ikx} dk = \frac{1}{2\pi} \int k\psi(k)e^{-ikx} dk. \tag{3.7}$$

Therefore,

$$(\sigma\hat{A}\omega\psi)(k) = k\psi(k). \tag{3.8}$$

So, the new eigenvalue problem is as follows:

$$g(k\psi) = pg(\psi). \quad (3.9)$$

Thus, we have the eigenvalue problem for the operator of multiplication by the variable. The eigenfunctions here are the elements of H^* given by

$$g(k) = \delta(p - k), \quad (3.10)$$

where $g = \omega^* f$. Indeed,

$$(\omega^* f)(k) = \frac{1}{2\pi} \int f(x)e^{-ikx} dx = \frac{1}{2\pi} \int e^{ipx} e^{-ikx} dx = \delta(p - k). \quad (3.11)$$

Two important things are happening here. First of all, for the given operator \hat{A} we have been able to find an appropriate Hilbert space of generalized functions so that the generalized eigenvectors of the operator are its elements. We also conclude that each eigenvalue problem actually defines a whole family of “unitary equivalent” problems obtained via isomorphisms of Hilbert spaces.

This is analogous to the case of a single tensor equation on R^n considered in different linear coordinates (i.e., in different bases). Instead of a change in *numeric* components of vectors and other tensors, we now have a change in functions, operators on spaces of functions and so forth.

The idea is then to consider different equivalence classes of functional objects related by isomorphisms of Hilbert spaces as realizations of invariant objects, which themselves are independent of any particular functional realization. In particular, the eigenvalue problems (3.2) and (3.9) can be considered as two coordinate expressions of a single equation which is itself independent of a chosen realization. In applications, such realizations, although mathematically equivalent, may represent different physical situations. In the next section this approach will be formalized.

4. Functional coordinate formalism on Hilbert manifolds

Here are the main definitions.

- (i) A *string space* \mathbf{S} is an abstract topological vector space topologically linearly isomorphic to a separable Hilbert space.
- (ii) The elements Φ, Ψ, \dots of \mathbf{S} are called strings.
- (iii) A *Hilbert space of functions* (or a *coordinate space*) is either a Hilbert space H , elements of which are equivalence classes of maps between two given subsets of R^n , or the Hilbert space H^* dual to H .
- (iv) A linear isomorphism e_H from a Hilbert space H of functions onto \mathbf{S} is called a *string basis* (or a *functional basis*) on \mathbf{S} .
- (v) The inverse map $e_H^{-1} : \mathbf{S} \rightarrow H$ is called a *linear coordinate system on \mathbf{S}* (or a *linear functional coordinate system*).

Analogy. \mathbf{S} is analogous to the abstract n -dimensional vector space V . A coordinate space H is analogous to a particular realization R^n of V as a space of column vectors. A string

basis e_H is analogous to the ordinary basis $\{e_k\}$ on V . Namely, the string basis identifies a string with a function: if $\Phi \in \mathbf{S}$, then $\Phi = e_H(\varphi)$ for a unique $\varphi \in H$. The ordinary basis $\{e_k\}$ identifies a vector with a column of numbers. It can be therefore identified with an isomorphism $e_k : R^n \rightarrow V$. The inverse map e_k^{-1} is then a global coordinate chart on V or a linear coordinate system on V .

Remark 4.1. Of course, any separable Hilbert space \mathbf{S} can be realized as a space l_2 of sequences by means of an orthonormal basis on \mathbf{S} . So it seems that we already have a clear infinite-dimensional analogue of the n -dimensional vector space. However, such a realization is too restrictive. In fact, unlike the case of a finite number of dimensions, there exist various realizations of \mathbf{S} as a space of functions. These realizations are also “numeric” and useful in applications and they should not be discarded. Moreover, we will see that they can be put on an equal footing with the space l_2 .

The basis $e_{H^*} : H^* \rightarrow \mathbf{S}^*$ is called *dual* to the basis e_H if for any string $\Phi = e_H(\varphi)$ and for any functional $F = e_{H^*}(f)$ in \mathbf{S}^* the following is true:

$$F(\Phi) = f(\varphi). \tag{4.1}$$

Analogy. The dual string basis e_{H^*} is analogous to the ordinary basis $\{e^m\}$ dual to a given basis $\{e_k\}$. Indeed, if Φ is a vector in the finite-dimensional vector space V , and $F \in V^*$ is a linear functional on V , then

$$F(\Phi) = f^k \varphi_k, \tag{4.2}$$

where f^k and φ_k are components of F and Φ in the bases $\{e^m\}$ and $\{e_k\}$, respectively.

By definition, the string space \mathbf{S} is isomorphic to a separable Hilbert space. Assume that \mathbf{S} itself is a Hilbert space. Assume further that the string bases e_H are isomorphisms of Hilbert spaces. That is, the Hilbert metric on any coordinate space H is determined by the Hilbert metric on \mathbf{S} and the choice of a string basis. Conversely, we have the following proposition.

PROPOSITION 4.2. *The choice of a coordinate Hilbert space determines the corresponding string basis e_H up to a unitary transformation.*

Proof. Indeed, with H fixed, any two bases e_H, \tilde{e}_H can only differ by an automorphism of H , that is, by a unitary transformation. □

Assume for simplicity that H is a real Hilbert space (generalization to the case of a complex Hilbert space will be obvious). We have

$$(\Phi, \Psi)_S = \mathbf{G}(\Phi, \Psi) = G(\varphi, \psi) = \int g(x, y)\varphi(x)\psi(y)dx dy = g_{xy}\varphi^x\psi^y, \tag{4.3}$$

where $\mathbf{G} : \mathbf{S} \times \mathbf{S} \rightarrow R$ is a bilinear functional defining the inner product on \mathbf{S} and $G : H \times H \rightarrow R$ is the induced bilinear functional. As before, the integral sign is a symbol of action of G on $H \times H$ and the expression on the right is a convenient form of writing this action.

A string basis e_H in \mathbf{S} will be called *orthogonal* if for any $\Phi, \Psi \in \mathbf{S}$, we have

$$(\Phi, \Psi)_S = f_\varphi(\psi), \tag{4.4}$$

where f_φ is a *regular* functional and $\Phi = e_H\varphi, \Psi = e_H\psi$ as before. That is,

$$(\Phi, \Psi)_S = f_\varphi(\psi) = \int \varphi(x)\psi(x)d\mu(x), \tag{4.5}$$

where \int here denotes an actual integral over a μ -measurable set $D \in R^n$ which is the domain of definition of functions in H .

If the integral in (4.5) is the usual Lebesgue integral and/or a sum over a discrete index x , the corresponding coordinate space will be called an L_2 -space. In this case we will also say that the basis e_H is *orthonormal*. If the integral is a more general Lebesgue-Stieltjes integral, the coordinate space defined by (4.5) will be called an L_2 -space with the weight μ and the basis e_H will be called *orthogonal*.

Analogy. An orthonormal string basis is analogous to the usual orthonormal basis. In fact, if φ^k, ψ^k are components of vectors Φ, Ψ in an orthonormal basis e_k on Euclidean space V , then $(\Phi, \Psi)_V = \delta_{kl}\varphi^k\psi^l = \varphi^k\psi^k$, which is similar to (4.5).

Roughly speaking, the metric on Hilbert spaces defined by orthogonal string bases has a “diagonal” kernel. In particular, the kernel may be proportional to the delta-function (as in $L_2(R)$) or to the Kronecker symbol (as in l_2). More general coordinate Hilbert spaces have a “nondiagonal” metric. The metrics (2.14) and (2.14) provide examples.

It is important to distinguish clearly the notion of a string basis from the notion of an ordinary basis on a Hilbert space of functions. Namely, a string basis permits us to represent invariant objects in string space (strings) in terms of functions, which are elements of a Hilbert space of functions. A basis on the space of functions then allows us to represent functions in terms of numbers; that is, in terms of the components of the functions in the basis. On the other hand, we have the following proposition.

PROPOSITION 4.3. *The ordinary orthonormal basis on \mathbf{S} can be identified with the string basis $e_{l_2} : l_2 \rightarrow \mathbf{S}$.*

Proof. The ordinary orthonormal basis on \mathbf{S} provides an isomorphism ω from \mathbf{S} onto l_2 which associates with each element Φ in \mathbf{S} the sequence of its Fourier components in the basis. Conversely, given an isomorphism $\omega : \mathbf{S} \rightarrow l_2$, we can associate with each $\Phi \in \mathbf{S}$ the sequence $\omega(\Phi)$ of its Fourier components in a unique (for all Φ) orthonormal basis. Therefore, the ordinary orthonormal basis on \mathbf{S} can be identified with the string basis $e_{l_2} = \omega^{-1}$. □

A *linear coordinate transformation on \mathbf{S}* is an isomorphism $\omega : \tilde{H} \rightarrow H$ of Hilbert spaces which defines a new string basis $e_{\tilde{H}} : \tilde{H} \rightarrow \mathbf{S}$ by $e_{\tilde{H}} = e_H \circ \omega$.

PROPOSITION 4.4. *Let $\varphi = e_H^{-1}\Phi, \hat{A} = e_H^{-1}\hat{A}e_H$, and $\hat{G} = (e_H^{-1})^*\hat{G}e_H^{-1}$ be the coordinate expressions of a string Φ , an operator $\hat{A} : \mathbf{S} \rightarrow \mathbf{S}$, and the metric $\hat{G} : \mathbf{S} \rightarrow \mathbf{S}^*$ in a basis e_H . Let*

Table 4.1. Classical versus string bases.

Basis	$e_H : H \rightarrow \mathbf{S}$	$e_k : R^n \rightarrow V$
Inner product	$(\Phi, \Psi)_S = \int g(x, y)\varphi(x)\psi(y)dx dy$	$(\Phi, \Psi)_V = g_{kl}\varphi^k\psi^l$
Dual basis	$e_{H^*}, F(\Phi) = f(\varphi)$	$e^k, F(\Phi) = f_k\varphi^k$
Ortho basis	$e_{L_2}, (\Phi, \Psi)_S = \int \varphi(x)\psi(x)dx$	$e_k, (\Phi, \Psi)_V = \delta_{kl}\varphi^k\psi^l$
Change of basis	$e_{\tilde{H}} = e_H \circ \omega, \hat{G}_{\tilde{H}} = \omega^* \hat{G}\omega, \dots$	$e_{\tilde{k}} = e_k \omega_{\tilde{k}}^k, \hat{G}_{\tilde{kl}} = \dots$

$\omega : \tilde{H} \rightarrow H$ be a linear coordinate transformation on \mathbf{S} . Then we have the following transformation laws:

$$\begin{aligned} \varphi &= \omega\tilde{\varphi}, \\ \hat{G}_{\tilde{H}} &= \omega^* \hat{G}\omega, \\ \hat{A}_{\tilde{H}} &= \omega^{-1} \hat{A}\omega, \end{aligned} \tag{4.6}$$

where $\tilde{\varphi}, \hat{A}_{\tilde{H}}$, and $\hat{G}_{\tilde{H}}$ are coordinate functions of Φ, \hat{A} , and \hat{G} in the basis $e_{\tilde{H}}$.

Proof. This follows immediately from

$$(\hat{G}\varphi, \hat{A}\psi) = (\omega^* \hat{G}\omega\tilde{\varphi}, \omega^{-1} \hat{A}\omega\tilde{\psi}) = (\hat{G}_{\tilde{H}}\tilde{\varphi}, \hat{A}_{\tilde{H}}\tilde{\psi}). \tag{4.7}$$

□

The correspondence between the classical and “string” definitions is summarized in Table 4.1.

More generally, consider an arbitrary Hilbert manifold S modeled on \mathbf{S} . Let (U_α, π_α) be an atlas on S . A collection of quadruples $(U_\alpha, \pi_\alpha, \omega_\alpha, H_\alpha)$, where each H_α is a Hilbert space of functions and ω_α is an isomorphism of \mathbf{S} onto H_α , is called a *functional atlas* on S . A collection of all compatible (i.e., related by diffeomorphisms) functional atlases on S is called a *coordinate structure* on S . A Hilbert manifold S with the above coordinate structure is called a *string manifold* or a *functional manifold*.

Let (U_α, π_α) be a chart on S . If $p \in U_\alpha$, then $\omega_\alpha \circ \pi_\alpha(p)$ is called the *coordinate* of p . The map $\omega_\alpha \circ \pi_\alpha : U_\alpha \rightarrow H_\alpha$ is called a *coordinate system*. The diffeomorphisms $\omega_\beta \circ \pi_\beta \circ (\omega_\alpha \circ \pi_\alpha)^{-1} : \omega_\alpha \circ \pi_\alpha(U_\alpha \cap U_\beta) \rightarrow \omega_\beta \circ \pi_\beta(U_\alpha \cap U_\beta)$ are called *string (or functional) coordinate transformations*.

As S is a differentiable manifold one can also introduce the tangent bundle structure $\tau : TS \rightarrow S$ and the bundle $\tau_s^r : T_s^r S \rightarrow S$ of tensors of rank (r, s) . Whenever necessary to distinguish tensors (tensor fields) on ordinary Hilbert manifolds from tensors on manifolds with a coordinate structure, we will call the latter tensors the *string tensors* or the *functional tensors*. Accordingly, the equations invariant under string coordinate transformations will be called the *string tensor* or the *functional tensor equations*.

A coordinate structure on a Hilbert manifold permits one to obtain a functional description of any string tensor. Namely, let $\mathbf{G}_p(F_1, \dots, F_r, \Phi_1, \dots, \Phi_s)$ be an (r, s) -tensor on S . The coordinate map $\omega_\alpha \circ \pi_\alpha : U_\alpha \rightarrow H_\alpha$ for each $p \in U_\alpha$ yields the linear map of tangent spaces $d\rho_\alpha : T_{\omega_\alpha \circ \pi_\alpha(p)} H_\alpha \rightarrow T_p S$, where $\rho_\alpha = \pi_\alpha^{-1} \circ \omega_\alpha^{-1}$. This map is called a *local coordinate string basis* on S . Notice that for each p the map $e_{H_\alpha} \equiv e_{H_\alpha}(p)$ is a string basis as defined

earlier. Therefore, the local dual basis $e_{H_\alpha^*} = e_{H_\alpha^*}(p)$ is defined for each p as before and is a function of p .

We now have $F_i = e_{H_\alpha^*} f_i$, and $\Phi_j = e_{H_\alpha} \varphi_j$ for any $F_i \in T_p^*S$, $\Phi_j \in T_pS$ and some $f_i \in H_\alpha^*$, $\varphi_j \in H_\alpha$. Therefore, the equation

$$\mathbf{G}_p(F_1, \dots, F_r, \Phi_1, \dots, \Phi_s) = G_p(f_1, \dots, f_r, \varphi_1, \dots, \varphi_s) \tag{4.8}$$

defines component functions of the (r, s) -tensor \mathbf{G}_p in the local coordinate basis e_{H_α} .

There are two important aspects associated with the formalism. First of all, by changing the string basis e_H , we can reformulate a given problem on Hilbert space in any way we want. In particular, we can go back and forth between spaces of smooth and generalized functions and choose the one in which the problem in hand is simpler to deal with. In this respect the theory of generalized functions is now a part of the above formalism on Hilbert spaces.

The second aspect is that we can now work directly with families of “unitary equivalent” problems rather than with particular realizations only. In fact, a single string tensor equation describes infinitely many equations on various spaces of functions.

In the following sections various examples and related propositions and theorems will illustrate the importance of these two aspects.

5. The generalized eigenvalue problem as a functional tensor equation

We begin with a *string basis independent* formulation of the discussed earlier generalized eigenvalue problem. Namely, consider the generalized eigenvalue problem

$$F(\hat{\mathbf{A}}\Phi) = \lambda F(\Phi), \tag{5.1}$$

for a linear operator $\hat{\mathbf{A}}$ on \mathbf{S} . The problem consists in finding all functionals $F \in \mathbf{S}^*$ and the corresponding numbers λ for which the string tensor equation (5.1) is satisfied for all $\Phi \in \mathbf{S}$.

Assume that the pair F, λ is a solution of (5.1) and e_H is a string basis on \mathbf{S} . Then we have

$$e_H^* F(e_H^{-1} \hat{\mathbf{A}} e_H \varphi) = \lambda e_H^* F(\varphi), \tag{5.2}$$

where $e_H \varphi = \Phi$ and $e_H^{-1} \hat{\mathbf{A}} e_H$ is the representation of $\hat{\mathbf{A}}$ in the basis e_H . By defining $e_H^* F = f$ and $\hat{A} = e_H^{-1} \hat{\mathbf{A}} e_H$, we have

$$f(\hat{A}\varphi) = \lambda f(\varphi). \tag{5.3}$$

The latter equation describes not just one eigenvalue problem, but a family of such problems, one for each string basis e_H . As we change e_H , the operator \hat{A} in general changes as well, as do the eigenfunctions f .

Assume that in a particular string basis e_H the problem (5.3) is the already-discussed generalized eigenvalue problem for the operator of differentiation $-i(d/dx)$. Then (5.1) is nothing but the corresponding “realization-independent” generalized eigenvalue problem given by a functional tensor equation.

6. The spectral theorem

Here we apply the coordinate formalism to reformulate the standard results of linear algebra in Hilbert spaces.

Definition 6.1. A string basis e_H is called the *proper basis* of a linear operator \mathbf{A} on \mathbf{S} with eigenvalues (eigenvalue function) $\lambda = \lambda(k)$ if

$$\mathbf{A}e_H(\varphi) = e_H(\lambda\varphi) \quad (6.1)$$

for any $\varphi \in H$.

As any string basis, the proper basis of \mathbf{A} is a linear map from H onto \mathbf{S} and a numeric function λ is defined on the same set as functions $\varphi \in H$.

By rewriting (6.1) as

$$e_H^{-1}\mathbf{A}e_H(\varphi) = \lambda\varphi, \quad (6.2)$$

we see that the problem of finding a proper basis of \mathbf{A} is equivalent to the problem of finding such a string basis e_H in which the action of \mathbf{A} reduces to multiplication by a function λ .

Because the expression “the adjoint of an operator” has at least two different meanings, we accept the following.

Definition 6.2. Let \hat{A} be a continuous linear operator which maps a space H into a space \hat{H} . Then the *adjoint* \hat{A}^* of operator \hat{A} maps the space \hat{H}^* into the space H^* according to

$$(\hat{A}^*f, \varphi) = (f, \hat{A}\varphi) \quad (6.3)$$

for any $\varphi \in H$, $f \in \hat{H}^*$. Assume further that H is continuously embedded into an L_2 -space, and under the identification $L_2^* = L_2$, the action of \hat{A}^* and \hat{A} on their common domain H coincide. Then the operator \hat{A} will be called *self-adjoint*.

Definition 6.3. Let A be a continuous linear operator on a Hilbert space H . Then the *Hermitian conjugate* operator A^+ of A is defined on H by

$$(A^+\varphi, \psi)_H = (\varphi, A\psi)_H, \quad (6.4)$$

for any $\varphi, \psi \in H$. If $A = A^+$, the operator is called Hermitian.

Let now A be given on H and let $\hat{G}: H \rightarrow H^*$ define a metric on H . Then

$$(\hat{G}\varphi, A\psi) = (A^*\hat{G}\varphi, \psi) = (\hat{G}A^+\varphi, \psi) \quad (6.5)$$

and the relationship between the operators is as follows:

$$A^+ = \hat{G}^{-1}A^*\hat{G}. \quad (6.6)$$

THEOREM 6.4. Any Hermitian operator $\hat{\mathbf{A}}$ on \mathbf{S} possesses a proper basis with the eigenvalue function $\lambda(x) = x$.

Proof. The theorem simply says that any Hermitian operator is unitary equivalent to the operator of multiplication by the variable, which is the standard result of spectral theorem. \square

THEOREM 6.5. Let \mathbf{A} be a Hermitian operator on \mathbf{S} and let e_H be a proper basis of $\hat{\mathbf{A}}$. Assume that H here is a Hilbert space of numeric-valued functions on a (possibly infinite) interval (a, b) and the eigenvalue function is given by $\lambda(x) = x, x \in (a, b)$. Assume further that the fundamental space K of infinitely differentiable functions of bounded support in (a, b) is continuously embedded into H as a dense subset. Then the proper basis of $\hat{\mathbf{A}}$ is orthogonal.

Proof. In the proper basis e_H of $\hat{\mathbf{A}}$ we have

$$(\Phi, \mathbf{A}\Psi)_S = (e_H\varphi, \mathbf{A}e_H\psi)_S = (\varphi, \lambda\psi)_H = \int g(x, y)\varphi(x)y\psi(y)dx dy. \tag{6.7}$$

In agreement with Section 4, the integral symbol is used here for the action of the bilinear metric functional G with the kernel g . Hermiticity gives then

$$\int g(x, y)(x - y)\varphi(x)\psi(y)dx dy = 0 \tag{6.8}$$

for any $\varphi, \psi \in H$.

We show that $g(x, y) = a(x)\delta(x - y)$, that is, H must coincide with the space $L_2(a, b)$ with weight $a(x)$. Note first of all that because $K \hookrightarrow H$, any bilinear functional on H is also a bilinear functional on K . Also, by the kernel theorem [3], any bilinear functional $G(\varphi, \psi) = \int g(x, y)\varphi(x)\psi(y)dx dy$ on the space K of infinitely differentiable functions of bounded support has the form

$$G(\varphi, \psi) = (f, \varphi(x)\psi(y)), \tag{6.9}$$

where f is a linear functional on the space K_2 of infinitely differentiable functions $\varphi(x, y)$ of bounded support in $(a, b) \times (a, b) \subset R^2$. If $\Omega \subset R^2$ is a bounded domain, $\varphi, \psi \in K$ are arbitrary with support in Ω and $x \neq y$ on Ω , then linear combinations of the functions $(y - x)\varphi(x)\psi(y)$ form a dense subset in the space $K_2(\Omega)$ of infinitely differentiable functions with support in Ω . From (6.8) it follows then that the functional f in (6.9) is zero on $K_2(\Omega)$ for any such Ω . Therefore, the bilinear functional G is concentrated on the diagonal $x = y$ in R^2 . That is, G is equal to zero on all pairs of functions $\varphi, \psi \in H$ such that $\varphi(x)\psi(y)$ is equal to zero in a neighborhood of the diagonal. Therefore, G reduces to a finite sum

$$G(\varphi, \psi) = \int \sum_{k=0}^n a_k(x)D^k\delta(y - x)(y - x)\varphi(x)\psi(y)dx dy, \tag{6.10}$$

where D^k is the derivative of order k with respect to y (see [2]).

Notice that $a_k(x) \neq 0$ almost everywhere on R as otherwise G would not be positive-definite. “Integration by parts” in (6.10) gives zero for $k = 0$ term and

$$\pm \int p\psi^{(p-1)}(x)a_p(x)\varphi(x)dx \tag{6.11}$$

for $k = p$ term and any $0 < p \leq n$. By choosing ψ so that the functions $\psi^{(p)}$ are linearly independent in $L_2(a, b)$, we conclude that the corresponding components of $a_p(x)\varphi(x)$ are equal to zero which is not true for a general $\varphi \in K$. Therefore, n must be equal to zero and the functional G is given by the kernel $g(x, y) = a(x)\delta(x - y)$. As K is dense in H , this kernel also defines the bilinear functional on H . We conclude that the proper basis of \hat{A} is orthogonal (as defined in Section 4). \square

Note that the obtained result can be generalized to the case of functions of several variables on a domain $D \subset R^n$. Note also that this theorem is nothing but the generalization to the case of a continuous spectrum of the well-known theorem on orthogonality of eigenvectors of a Hermitian operator corresponding to different eigenvalues.

The conclusion of the theorem seems to favor the use of L_2 -spaces when working with Hermitian operators. Notice however that in an orthogonal proper basis of \hat{A} with $H = L_2(a, b)$ we have

$$(\Phi, \hat{A}\Psi)_S = (f, \lambda g)_{L_2} = \int f(x)\lambda(x)g(x)dx, \tag{6.12}$$

where $f, g \in L_2(a, b)$. Therefore, it is impossible for $f(x)$ to be a (generalized) eigenvector of the operator \hat{A} of multiplication by the function $\lambda(x)$. Indeed, the eigenvectors of such an operator would be δ -functionals and the latter ones do not belong to $L_2(a, b)$.

To include such functionals in the formalism we need to consider Hilbert spaces containing “more” functions than $L_2(a, b)$. By the above, this in general requires consideration of non-Hermitian operators.

Example 6.6. The operator of multiplication by the variable $\hat{A} = x$ has no eigenvectors in $L_2(0, 1)$. Consider then a different Hilbert space H^* of functions of a real variable $x \in (0, 1)$ with the metric G given by a smooth kernel $g(k, m)$ and with the dual space H consisting of continuous functions only. We have

$$(G\varphi, \hat{A}\psi) = \int g(k, m)\varphi(k)m\psi(m)dk dm. \tag{6.13}$$

Although the operator $\hat{A} = x$ is not Hermitian on H , it is self-adjoint (see Definition 6.2). In fact, if φ, ψ are continuous functions, then

$$\int g(k, m)\varphi(k)m\psi(m)dk dm = \int mg(k, m)\varphi(k)\psi(m)dk dm. \tag{6.14}$$

So on the functionals

$$F(m) = \int g(k, m)\varphi(k)dk \tag{6.15}$$

which are functions in H , we have

$$F(m\psi) = (mF)(\psi). \tag{6.16}$$

Note also that the eigenfunctions of $\hat{A} = x$ are in H^* .

The example shows that at least in some cases there is a possibility to accommodate the Hermiticity of an operator in L_2 , the self-adjointness of its restriction \hat{A} onto a Hilbert subspace $H \subset L_2$, and the inclusion of generalized eigenvectors of \hat{A} into the conjugate space H^* . Before formalizing and generalizing this statement we need the following definition (see [3]).

Definition 6.7. Let H_λ^* be the *eigenspace* of \hat{A} consisting of all the generalized eigenvectors f_λ of \hat{A} whose eigenvalue is λ . Associate with each element $\varphi \in H$ and each number λ a linear functional $\tilde{\varphi}_\lambda$ on H_λ^* which takes the value $f_\lambda(\varphi)$ on the element f_λ . This gives vector functions of λ whose values are linear functionals on H_λ^* . The correspondence $\varphi \rightarrow \tilde{\varphi}_\lambda$ is called the *spectral decomposition* of φ corresponding to the operator \hat{A} . The set of generalized eigenvectors of \hat{A} is called *complete* if $\tilde{\varphi}_\lambda = 0$ implies $\varphi = 0$.

We have now the following.

THEOREM 6.8. *Let \hat{B} be a Hermitian operator on a Hilbert space L_2 . Then there exists a topological Hilbert subspace H of L_2 which is dense in L_2 and such that the restriction \hat{A} of \hat{B} onto H is a self-adjoint operator and the conjugate space H^* contains the complete set of eigenvectors of \hat{A} . Moreover, there exists a coordinate transformation $\rho : H \rightarrow \tilde{H}$ such that the transformed operator $\tilde{\hat{A}} = \rho\hat{A}\rho^{-1}$ is the operator of multiplication by x .*

Proof. Notice first of all that the property of an operator \hat{A} to be Hermitian is invariant under functional coordinate transformations.

Let $\omega : L_2 \rightarrow L_2(R)$ be an isomorphism of Hilbert spaces. By the above, $\hat{B}_1 = \omega\hat{B}\omega^{-1}$ is an Hermitian operator in $L_2(R)$.

Let W be the Schwartz space of infinitely differentiable functions on R . It is known that W is a nuclear space and the triple $W \subset L_2(R) \subset W^*$ is a rigged Hilbert space [3]. The operator \hat{B}_1 is a Hermitian operator in the rigged space. From [3] we know that the operator \hat{B}_1 has a complete set of generalized eigenvectors belonging to W^* .

We choose a Hilbert space H_1 such that H_1 is a topological subspace of W which is also dense in $L_2(R)$. We know from Theorem 2.6 that such a Hilbert space H_1 exists. As $W^* \subset H_1^*$, the space H_1^* contains the complete set of eigenvectors of the restriction \hat{A}_1 of \hat{B}_1 onto H_1 . The image $H = \omega^{-1}(H_1)$ is a Hilbert subspace of L_2 with the induced Hilbert structure. It is dense in L_2 and on this subspace the operator $\hat{A} = \omega^{-1}\hat{A}_1\omega$ is self-adjoint and coincides with the original operator \hat{B} . This proves the first part of the theorem.

Now, by the abstract theorem on spectral decomposition of Hermitian operators, there exists a realization $\tau : L_2(R) \rightarrow \tilde{L}_2$ such that the operator \hat{B}_1 is given by multiplication by x . Here \tilde{L}_2 denotes in general a direct integral of Hilbert spaces of the L_2 -type.

Consider the restriction of τ onto H_1 . As τ is an isomorphism, it induces a Hilbert structure on the image $\tau(H_1) = \tilde{H}$.

Consider then the isomorphism $\kappa = \tau \circ \omega : H \rightarrow \tilde{H}$ as a coordinate transformation. This transformation takes the operator \hat{A} in H into the operator of multiplication by x in \tilde{H} . This completes the proof. \square

Assume now that \hat{A} is the operator constructed in the theorem. That is, \hat{A} is a self-adjoint operator in a Hilbert space H such that the conjugate space H^* contains the complete set of generalized eigenvectors f_λ of \hat{A} . Let $\varphi \rightarrow \tilde{\varphi}_\lambda$ be the spectral decomposition of the element $\varphi \in H$ corresponding to the operator \hat{A} . Let e_H be a string basis on \mathbf{S} and $\Phi = e_H \varphi$, $F_\lambda = (e_H^*)^{-1} f_\lambda$, $\hat{\mathbf{A}} = e_H \hat{A} e_H^{-1}$ as always.

The spectral decomposition $\varphi \rightarrow \tilde{\varphi}_\lambda$ establishes the correspondence $\Phi \rightarrow \tilde{\varphi}_\lambda$ which is the spectral decomposition of the string Φ corresponding to the operator $\hat{\mathbf{A}}$.

More directly, to construct $\Phi \rightarrow \tilde{\varphi}_\lambda$ we introduce the eigenspace \mathbf{S}_λ^* of $\hat{\mathbf{A}}$ which consists of all eigenvectors F_λ of $\hat{\mathbf{A}}$ whose eigenvalue is λ . Then we associate with each string $\Phi \in \mathbf{S}$ and each number λ a linear functional $\tilde{\varphi}_\lambda$ on \mathbf{S}_λ^* which takes the value $F_\lambda(\Phi)$ on the element F_λ .

As the set of generalized eigenvectors f_λ of \hat{A} is complete, so is the set of eigenvectors F_λ of the operator $\hat{\mathbf{A}}$. In fact, whenever $\tilde{\varphi}_\lambda = 0$ we have $\varphi = 0$ by completeness of the set of eigenvectors f_λ . But then $\Phi = 0$ as well, that is, the set of eigenvectors F_λ of $\hat{\mathbf{A}}$ is complete.

The correspondence $\rho_{\tilde{H}} : \Phi \rightarrow \tilde{\varphi}_\lambda$ is an isomorphism of linear spaces \mathbf{S} and $\tilde{H} = \rho_{\tilde{H}}(\mathbf{S})$. In fact, assume that to strings Φ_1, Φ_2 and to a number λ it corresponds respectively linear functionals $\tilde{\varphi}_{\lambda 1}, \tilde{\varphi}_{\lambda 2}$ on the eigenspace \mathbf{S}_λ^* of $\hat{\mathbf{A}}$ taking values $F_\lambda(\Phi_1), F_\lambda(\Phi_2)$ on the element F_λ . Then to the string $\Phi = \alpha\Phi_1 + \beta\Phi_2$ and to the same number λ corresponds the functional $\tilde{\varphi}_\lambda$ on \mathbf{S}_λ^* taking the value $F_\lambda(\alpha\Phi_1 + \beta\Phi_2) = \alpha F_\lambda(\Phi_1) + \beta F_\lambda(\Phi_2)$. This means that $\rho_{\tilde{H}}$ is linear. The fact that $\rho_{\tilde{H}}$ is injective follows from completeness of the set of eigenvectors F_λ of $\hat{\mathbf{A}}$. In fact, if $\rho_{\tilde{H}}(\Phi) = 0$, then $\tilde{\varphi}_\lambda = 0$, thus, $\Phi = 0$.

The isomorphism $\rho_{\tilde{H}}$ induces a Hilbert structure on $\tilde{H} = \rho_{\tilde{H}}(\mathbf{S})$. Therefore it becomes an isomorphism of Hilbert spaces \mathbf{S} and \tilde{H} .

THEOREM 6.9. *The string basis $e_{\tilde{H}} = \rho_{\tilde{H}}^{-1}$ is proper.*

Proof. To verify, notice that if $\Phi \rightarrow \tilde{\varphi}_\lambda$ is the spectral decomposition of Φ , then the spectral decomposition of $\Psi = \hat{\mathbf{A}}\Phi$ is $\tilde{\psi}_\lambda = \lambda\tilde{\varphi}_\lambda$. In fact, for any functional $F_\lambda \in \mathbf{S}_\lambda^*$ we have

$$F_\lambda(\Psi) = F_\lambda(\hat{\mathbf{A}}\Phi) = \lambda F_\lambda(\Phi), \tag{6.17}$$

so that

$$\tilde{\psi}_\lambda = \lambda\tilde{\varphi}_\lambda. \tag{6.18}$$

We therefore conclude that

$$F(\hat{\mathbf{A}}\Phi) = F(e_{\tilde{H}}^{-1} e_{\tilde{H}} \hat{\mathbf{A}} \Phi) = (e_{\tilde{H}}^{-1})^* F(\lambda\tilde{\varphi}_\lambda) = \tilde{f}(\lambda\tilde{\varphi}_\lambda), \tag{6.19}$$

where $\tilde{f} = (e_{\tilde{H}}^{-1})^* F$. By (6.1) this means that the basis $e_{\tilde{H}}$ is proper. \square

Moreover, in the considered case the complete set of eigenvectors F_λ of \hat{A} belongs to \mathbf{S}^* . Therefore, the complete set of eigenvectors of the operator of multiplication by λ on \tilde{H} belongs to \tilde{H}^* .

Definition 6.10. If a proper basis e_H is such that the complete set of eigenvectors of the operator \hat{A} in this basis belongs to H^* (alternatively, the complete set of eigenvectors of \hat{A} belongs to \mathbf{S}^*), then the basis e_H is called the *string basis of eigenvectors of \hat{A}* or the *string eigenbasis of \hat{A}* .

We have thus proven the following theorem.

THEOREM 6.11. *If a complete set of eigenvectors F_λ of an operator \hat{A} is contained in \mathbf{S}^* , then there exists a string eigenbasis of \mathbf{A} . That is, assume $F_\lambda(\hat{A}\Phi) = \lambda F_\lambda(\Phi)$, where the eigenvectors F_λ of \hat{A} form a complete set in \mathbf{S}^* . Then there exists a string basis $e_{\tilde{H}} : \tilde{H} \rightarrow \mathbf{S}$ such that for any functional $F \in \mathbf{S}^*$ and any string Φ we have $F(\mathbf{A}\Phi) = \tilde{f}(\lambda\tilde{\varphi}_\lambda)$, where $\tilde{\varphi}_\lambda$ and \tilde{f} are coordinates of Φ and F in the basis $e_{\tilde{H}}$ and its dual, respectively.*

It is well known that many operators useful in applications are not bounded. However, by an appropriate choice of metric on the image of a linear operator \hat{A} , one can ensure the continuity of \hat{A} . In other words, one can consider an unbounded operator \hat{A} on a Hilbert space H as a bounded operator which maps H into another Hilbert space \tilde{H} . In particular, we have the following.

THEOREM 6.12. *Let \hat{A} be an invertible (possibly unbounded) operator on a dense subset $D(\hat{A}) \subset L_2$ having a range $R(\hat{A}) \supset D(\hat{A})$. Then there exists a Hilbert metric on $R(\hat{A})$, in which \hat{A} is a bounded operator from $D(\hat{A})$ onto $R(\hat{A})$.*

Proof. For any $f, g \in R(\hat{A})$ define the inner product $(f, g)_H$ by

$$(f, g)_H = (\hat{A}^{-1}f, \hat{A}^{-1}g)_{L_2}. \tag{6.20}$$

Then $\|\hat{A}f\|_H = \|f\|_{L_2}$ for any $f \in D(\hat{A})$. That is, \hat{A} is bounded as a map from $D(\hat{A})$ with the metric L_2 onto $R(\hat{A})$ with the metric (6.20). Note that when \hat{A} is extended to the entire L_2 it becomes an isomorphism from L_2 onto the Hilbert space H which is a completion of $R(\hat{A})$ in the metric (6.20). □

We also remark that if e_H and $e_{\tilde{H}}$ are string bases, the realization $\hat{A} = e_{\tilde{H}}^{-1}\hat{A}e_H$ of a continuous operator $\hat{A} : \mathbf{S} \rightarrow \mathbf{S}$ may not be continuous when considered as a map into H . In other words, the unbounded operators in a Hilbert space H may be realizations of continuous operators on \mathbf{S} .

Notice finally that when a realization \hat{A} of operator $\hat{A} : \mathbf{S} \rightarrow \mathbf{S}$ is a map between two different Hilbert spaces, it becomes useful to generalize Definition 6.1 of the proper basis. Namely, for any $F \in \mathbf{S}^*$, $\Phi \in \mathbf{S}$, and $\hat{A} : \mathbf{S} \rightarrow \mathbf{S}$, we have

$$F(\hat{A}\Phi) = F(\hat{A}e_H\varphi) = F(e_{\tilde{H}}e_{\tilde{H}}^{-1}\hat{A}e_H\varphi) = e_{\tilde{H}}^*F(e_{\tilde{H}}^{-1}\hat{A}e_H\varphi) = f(\hat{A}\varphi), \tag{6.21}$$

where $f = e_{\tilde{H}}^*F \in \tilde{H}^*$ and $\hat{A} : H \rightarrow \tilde{H}$ is given by $\hat{A} = e_{\tilde{H}}^{-1}\hat{A}e_H$. We then have the following.

Definition 6.13. The realization $\hat{A} = e_{\tilde{H}}^{-1} \hat{A} e_H$ of \hat{A} is called *proper* if for any $F \in \mathbf{S}^*$, $\Phi \in \mathbf{S}$,

$$F(\hat{\mathbf{A}}\Phi) = f(\hat{A}\varphi) = f(\lambda\varphi), \tag{6.22}$$

where $f = e_{\tilde{H}}^* F$, $\varphi = e_{\tilde{H}}^{-1} \Phi$ and λ is a function such that the operator of multiplication by λ maps H into \tilde{H} .

If $H = \tilde{H}$ this is equivalent to Definition 6.1.

7. Isomorphisms of Hilbert spaces preserving locality of operators

We saw that a single eigenvalue problem for an operator on the string space leads to a family of eigenvalue problems in particular string bases. Obviously, it is a general feature of tensor equations in the formalism: a specific functional form of an equation depends on the choice of functional coordinates.

In the process of changing coordinates, one may desire to preserve some specific properties of the equation. It becomes then important to describe all coordinate transformations preserving these properties.

Definition 7.1. An operator \hat{A} on a Hilbert space H is called *differential* or *local* if

$$\hat{A} = \sum_{|q| \leq r} a_q D^q. \tag{7.1}$$

Here $x \in R^n$, r is a nonnegative integer, $q = (q_1, \dots, q_n)$ is a set of nonnegative integers, $|q| = q_1 + \dots + q_n$, the coefficients a_q are functions of $x = (x_1, \dots, x_n)$, and $D^q = \partial^{|q|} / (\partial x_1^{q_1} \dots \partial x_n^{q_n})$.

It is easy to see that locality of an operator is not a coordinate invariant property. That is, the operator which is local in one system of functional coordinates does not have to be local in a different system of coordinates. At the same time, the property of being local is extremely important in applications.

We describe the coordinate transformations $\omega : \tilde{H} \rightarrow H$ mapping differential operators on H into differential operators on a Hilbert space \tilde{H} . We assume here that H, \tilde{H} are either spaces of sufficiently smooth functions of bounded support (or sufficiently fast decreasing at infinity) or dual to such spaces.

All such transformations can be found by solving the equation

$$\omega^{-1} \hat{A} \omega = \hat{B}, \tag{7.2}$$

where \hat{A}, \hat{B} are appropriate differential operators. In expanded form we have

$$\omega^{-1} \sum_{|q| \leq r} a_q D^q \omega = \sum_{|q| \leq s} b_q D^q. \tag{7.3}$$

The latter equation can also be written in terms of the kernels of \hat{A} and \hat{B} :

$$\int \sum_{|q| \leq r} a_q(x) D^q \delta(z-x) \omega(z, y) dz = \int \sum_{|q| \leq s} \omega(x, z) b_q(z) D^q \delta(y-z) dz. \quad (7.4)$$

Example 7.2. Let H, \tilde{H} are spaces of functions of a single variable, $\hat{A} = aD, \hat{B} = b$, where a, b are functions. Then (7.3) reduces to

$$aD\omega = \omega b. \quad (7.5)$$

By solving this differential equation we see that the kernel $\omega(x, y)$ of ω has the form

$$\omega(x, y) = g(y) e^{c(x)b(y)}, \quad (7.6)$$

where $c(x) = -\int(dx/a(x))$ and g is an arbitrary smooth function. To be a coordinate transformation, ω must be an isomorphism as well. In particular, the Fourier transform is a solution of (7.5) with

$$\omega(x, y) = e^{ixy}. \quad (7.7)$$

More generally, if \hat{A}, \hat{B} contain only one term each, (7.4) yields

$$a(x) \frac{\partial^n \omega(x, y)}{\partial x^n} = (-1)^m \frac{\partial^m (\omega(x, y) b(y))}{\partial y^m}. \quad (7.8)$$

Example 7.3. Consider (7.8) with $n = m = 1$ assuming $a(x)$ and $b(y)$ are functions. In this case the equation reads

$$a(x) \frac{\partial \omega(x, y)}{\partial x} + \frac{\partial (\omega(x, y) b(y))}{\partial y} = 0. \quad (7.9)$$

We look for a solution in the form

$$\omega(x, y) = e^{f(x)g(y)}. \quad (7.10)$$

Then (7.9) yields

$$a(x) f'(x) g(y) + b(y) f(x) g'(y) + b'(y) = 0. \quad (7.11)$$

If $b(y) = 1$, (7.11) is a separable equation and we have

$$\frac{a(x) f'(x)}{f(x)} = -\frac{g'(y)}{g(y)} = C, \quad (7.12)$$

where C is a constant. Solving this, we obtain

$$\omega(x, y) = e^{C \int (C_1/a(x)) dx e^{-C_1 y}}. \quad (7.13)$$

Taking for example $C = C_1 = 1$ and $a(x) = x$, we have

$$\omega(x, y) = e^{xe^{-y}}. \tag{7.14}$$

The corresponding transformation changes the operator xD into the operator D .

8. Isomorphisms of Hilbert spaces preserving the derivative operator

Among solutions of (7.8) those preserving the derivative operator D are of particular interest. To find them consider (7.8) with $n = m = 1$ and with the constant coefficients $a = b = 1$. Then (7.8) yields the following equation:

$$\frac{\partial\omega(x, y)}{\partial x} + \frac{\partial\omega(x, y)}{\partial y} = 0. \tag{8.1}$$

The smooth solutions of (8.1) are given by

$$\omega(x, y) = f(x - y), \tag{8.2}$$

where f is an arbitrary infinitely differentiable function on R . In particular, the function

$$\omega(x, y) = e^{-(x-y)^2} \tag{8.3}$$

satisfies (8.1). Also, we saw in Section 2 that the corresponding transformation considered on an appropriate Hilbert space \tilde{H} is injective and induces a Hilbert structure on the image H . Therefore, it provides an example of a coordinate transformation that preserves the derivative operator D and, more generally, D^q .

When H, \tilde{H} are spaces of functions of n variables, a similar role is played by the function

$$\omega(x, y) = e^{-(x-y)^2} \tag{8.4}$$

with $x = (x_1, \dots, x_n), y = (y_1, \dots, y_n)$, and the standard Euclidean metric on the space of variables.

THEOREM 8.1. *Let L be a polynomial function of n variables. Let $u, v \in \tilde{H}$ be functionals on the space K of functions of n variables which are infinitely differentiable and have bounded supports. Assume that u is a generalized solution of*

$$L\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)u = v. \tag{8.5}$$

Then there exists a smooth solution φ of

$$L\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)\varphi = \psi, \tag{8.6}$$

where $\varphi = \omega u, \psi = \omega v$ and ω is as in (8.4).

Proof. Consider first the case of the ordinary differential equation

$$\frac{d}{dx}u(x) = v(x). \tag{8.7}$$

Assume u is a generalized solution of (8.7). Define $\varphi = \omega u$ and $\psi = \omega v$, where ω is as in (8.3). Notice that φ, ψ are infinitely differentiable. In fact, any functional on the space K of infinitely differentiable functions of bounded support acts as follows (see [2]):

$$(f, \varphi) = \int F(x)\varphi^{(m)}(x)dx, \tag{8.8}$$

where F is a continuous function on R . Applying ω to f shows that the result is a smooth function.

As $\omega^{-1}(d/dx)\omega = d/dx$, we have

$$\omega^{-1}\frac{d}{dx}\omega u = v. \tag{8.9}$$

That is,

$$\frac{d}{dx}\varphi(x) = \psi(x), \tag{8.10}$$

proving the theorem in this case. The higher-order derivatives can be treated similarly as

$$\omega^{-1}\frac{d^n}{dx^n}\omega = \omega^{-1}\frac{d}{dx}\omega\omega^{-1}\frac{d}{dx}\omega \cdots \omega^{-1}\frac{d}{dx}\omega. \tag{8.11}$$

That is, transformation ω preserves derivatives of any order. Generalization to the case of several variables is straightforward. □

9. Transformations preserving a product of functions

Consider a simple algebraic equation

$$a(x)f(x) = h(x), \tag{9.1}$$

where f is an unknown (generalized) function of a single variable x and h is an element of a Hilbert space H of functions on a set $D \subset R$. To investigate transformation properties of this equation we need to interpret it as a tensor equation on the string space \mathbf{S} . The right-hand side is a function. Therefore, this must be a “vector equation” (i.e., both sides must be $(1,0)$ -tensors on the string space). If f is to be a function as well, a must be a $(1,1)$ -tensor.

To preserve the product-like form of the equation, we need such a coordinate transformation $\omega : \tilde{H} \rightarrow H$ that

$$\omega^{-1}a\omega = b. \tag{9.2}$$

In this case (9.1) in new coordinates is

$$b(x)\varphi(x) = \psi(x), \tag{9.3}$$

where $h = \omega\psi$, $f = \omega\varphi$, and $\varphi, \psi \in \tilde{H}$.

Equation (9.2) is clearly satisfied whenever $a = b = C$, where C is a constant function. On the other hand, whenever $b' \neq 0$ and H contains sufficiently many functions, we deduce as in the Theorem 6.5 that $\omega(x, y) = d(x)\delta(x - y)$ for some function d .

The obtained result then says that the operator of multiplication can be preserved only in trivial cases when $a(x) = C$ or ω itself is an operator of multiplication by a function.

In particular, the product of nonconstant functions of one and the same variable is not an invariant operation under a general transformation of functional coordinates.

10. Coordinate transformations of nonlinear functional tensor equations

It is known that the theory of generalized functions has been mainly successful with linear problems. The difficulty of course lies in defining the product of generalized functions. We saw in the previous section that multiplication of functions of the same variable is not an invariant operation. The idea is then to define an invariant operation which reduces to multiplication in particular coordinates.

More generally, the developed functional coordinate formalism offers a systematic approach for dealing with nonlinear equations in generalized functions. Namely, the terms in a nonlinear functional tensor equation represent tensors. Because of that, the change of coordinates is meaningful and can be used to extend nonlinear operations to generalized functions.

Example 10.1. Consider a nonlinear equation with the term $(\varphi(x))^2$. To interpret this term as a functional tensor, we write

$$\varphi(x) \cdot \varphi(x) = \int \delta(x - u)\delta(x - v)\varphi(u)\varphi(v)du dv. \tag{10.1}$$

Therefore, this term can be considered to be the convolution of the (1, 2)-tensor

$$c_{uv}^x = \delta(x - u)\delta(x - v) \tag{10.2}$$

with the pair of strings $\varphi^u = \varphi(u)$:

$$\varphi(x) \cdot \varphi(x) = c_{uv}^x \varphi^u \varphi^v. \tag{10.3}$$

With this identification we can now apply a coordinate transformation to make φ singular at the expense of smoothing c_{uv}^x . In particular, we can transform φ into the delta-function.

Example 10.2. Consider the equation

$$\int k(x - y) \frac{d\varphi_t(x)}{dt} \frac{d\varphi_t(y)}{dt} dx dy = 0, \tag{10.4}$$

where $\varphi_t(x)$ is an unknown function which depends on the parameter t and $k(x, y)$ is a smooth function on R^{2n} .

We look for a solution in the form $\varphi_t(x) = \delta(x - a(t))$. As

$$\frac{d\varphi_t(x)}{dt} = -\frac{\partial\delta(x - a(t))}{\partial x^\mu} \frac{da^\mu}{dt}, \tag{10.5}$$

we have after “integration by parts” the following equation:

$$\frac{\partial^2 k(x, y)}{\partial x^\mu \partial y^\nu} \Big|_{x=y=a(t)} \frac{da^\mu(t)}{dt} \frac{da^\nu(t)}{dt} = 0. \tag{10.6}$$

We remark here that, as explained in Section 11, formula (10.5) and the method of “integration by parts” in (10.4) are valid. If the tensor field

$$g_{\mu\nu}(a) \equiv \frac{\partial^2 k(x, y)}{\partial x^\mu \partial y^\nu} \Big|_{x=y=a} \tag{10.7}$$

is symmetric and positive definite, then (10.6) has only the trivial solution. However, if $g_{\mu\nu}(a)$ is nondegenerate and indefinite, then there is a nontrivial solution. In particular, we can choose $g_{\mu\nu}$ to be the Minkowski tensor $\eta_{\mu\nu}$ on space-time. For this assume that x, y are space-time points and take

$$k(x, y) = e^{-(x-y)^2/2}, \tag{10.8}$$

where $(x - y)^2 = \eta_{\mu\nu}(x - y)^\mu(x - y)^\nu$. Then we immediately conclude that $g_{\mu\nu}$ is the Minkowski tensor $\eta_{\mu\nu}$. Solutions to (10.6) are then given by the null lines $a(t)$. Therefore, the original problem (10.4) has solutions of the form $\varphi_t(x) = \delta(x - a(t))$ where $a(t)$ is a null line.

Notice that the obtained generalized functions φ_t are singular generalized solutions to the nonlinear equation (10.4). Once again, these solutions become possible because the kernel $k(x, y)$ in (10.4) is a smooth function, so the convolution $k_{xy}\hat{\varphi}^x\hat{\varphi}^y$ is meaningful.

Example 10.3. Consider the equation

$$\int e^{-(x-y)^2/2} \frac{d\varphi_t(x)}{dt} \frac{d\varphi_t(y)}{dt} dx dy = 1, \tag{10.9}$$

where the metric on the space of variables is Euclidean: $(x - y)^2 = \delta_{\mu\nu}(x - y)^\mu(x - y)^\nu$. Looking for solution in the form $\varphi_t(x) = \delta(x - a(t))$ and performing “integration by parts,” we obtain

$$\delta_{\mu\nu} \frac{da^\mu(t)}{dt} \frac{da^\nu(t)}{dt} = 1. \tag{10.10}$$

That is, $\|da(t)/dt\|_{R^n} = 1$. Therefore, $\varphi_t(x) = \delta(x - a(t))$, where the path $a(t)$ has a unit velocity vector at any t .

We apply a transformation ρ with kernel $\rho(x, y) = e^{-(x-y)^2}$ to the equation

$$\int e^{-(x-y)^2/2} \frac{d\delta(x - a(t))}{dt} \frac{d\delta(y - a(t))}{dt} dx dy = 1. \tag{10.11}$$

This yields

$$M \int \delta(x - y) \frac{de^{-(x-a(t))^2}}{dt} \frac{de^{-(y-a(t))^2}}{dt} dx dy = 1, \tag{10.12}$$

where M is a constant. This is the same equation in the coordinates in which the kernel of the metric is the delta function, while the solution given originally by the delta-function becomes a smooth exponential function.

11. Embeddings of n -dimensional manifolds into Hilbert spaces of functions

In this section, we will use various dual Hilbert spaces H, H^* , where the elements of H^* are smooth functions and the metric on the dual space H is given by a smooth-on- $R^n \times R^n$ kernel $k(x, y)$. Examples of such spaces were given in Section 2.

THEOREM 11.1. *The space H described above contains the evaluation functionals $\delta(x - a)$ for all a in R^n .*

Proof. Because $k(x, y)$ is smooth, the sequence $f_L = (L/\sqrt{\pi})^n e^{-L^2(x-a)^2}$ is fundamental in H . As H is complete, f_L converges to an element f in H . Therefore, f_L converges to f weakly in H as well. By Riesz theorem, this is equivalent to saying that $(f_L, \varphi) \rightarrow (f, \varphi)$ for all $\varphi \in H^*$. At the same time, f_L is a delta-converging sequence. In particular, $(f_L, \varphi) \rightarrow (\delta_a, \varphi)$ for all $\varphi \in H^*$. The uniqueness of the weak limit signifies then that f_L converges to δ_a in H . □

THEOREM 11.2. *If the space H^* contains sufficiently many functions, the subset M of all delta-functions in H forms an n -dimensional submanifold of H diffeomorphic to R^n .*

Proof. The map $\omega : a \rightarrow \delta(x - a)$ is smooth from R^n into H . Assume that for any two points $a, b \in R^n$ there is a function $\varphi \in H^*$ such that $\delta_a(\varphi) \neq \delta_b(\varphi)$. This is true in particular if H^* contains all C^∞ -functions of bounded support. In this case the map ω is injective. The smooth injective map ω parametrizes the set M of all delta-functions in H identifying M with a submanifold of H diffeomorphic to R^n . □

Note that, although M is not a linear subspace of H , the diffeomorphism $\omega : R^n \rightarrow M$ induces a linear structure on M . In fact, we can define linear operations \oplus, \odot on M by $\omega(x + y) = \omega(x) \oplus \omega(y)$ and $\omega(kx) = k \odot \omega(x)$ for any vectors $x, y \in R^n$ and any number k .

THEOREM 11.3. *These operations are continuous in topology of H .*

Proof. $\int k(x, y)(\lambda\delta(x - a) - \lambda_k\delta(x - a_k))(\lambda\delta(y - a) - \lambda_k\delta(y - a_k))dx dy = \lambda^2k(a, a) - \lambda\lambda_k(k(a, a_k) + k(a_k, a)) + \lambda_k^2k(a_k, a_k) \rightarrow 0$ provided $a_k \rightarrow a$ and $\lambda_k \rightarrow \lambda$. The continuity of addition is verified in a similar way. □

By using the same method, one can also derive topologically nontrivial spaces M . For example, let H^* be the Hilbert space of smooth functions on the interval $[0, 2\pi]$ which contains all smooth functions of bounded support in $(0, 2\pi)$. Assume further that

$$\varphi^{(n)}(0) = \varphi^{(n)}(2\pi) \tag{11.1}$$

for all φ in H^* and for all orders n of (one-sided) derivatives of φ . Consider the dual space H of functionals in H^* and assume that the kernel of the metric on H is a smooth function.

THEOREM 11.4. *The subset M of delta-functions in H forms a submanifold diffeomorphic to the circle S^1 .*

Proof. The map $\omega : a \rightarrow \delta(\theta - a)$ from $(0, 2\pi)$ into H is C^∞ . It is also injective, because H^* contains sufficiently many functions to distinguish any two delta-functions $\delta(x - a)$, $\delta(x - b)$ with $a, b \in (0, 2\pi)$, $a \neq b$. Finally, the functionals $\delta(\theta)$, $\delta(\theta - 2\pi)$ are identical on H^* as $\int \delta(\theta - 2\pi)\varphi(\theta)d\theta = \varphi(2\pi) = \varphi(0) = \int \delta(\theta)\varphi(\theta)d\theta$ for all φ in H^* . \square

More generally, a Hilbert space H^* of functions on an n -dimensional manifold can be identified with the space of functions on a subset of R^n . In fact, the manifold itself is a collection of nonintersecting “pieces” of R^n “glued” together. Functions on the manifold can be then identified with functions defined on the disjoint union of all pieces and taking equal values at the points identified under “gluing.” As a result, the dual space H of generalized functions “on” the manifold can be also identified with the corresponding space of generalized functions “on” a subset of R^n .

This fact allows us to conclude that topologically different manifolds M can be obtained by choosing an appropriate Hilbert space H of functions “on” a subset of R^n and identifying M with the submanifold of H consisting of delta-functions. The manifold structure on M is then induced by the inclusion of M into H and does not have to be defined in advance. In particular, the “gluing conditions” like (11.1) imposed on functions result in the corresponding “gluing” of the delta-functions and of the appropriate subsets of R^n .

Moreover, the tangent bundle structure and the Riemannian structure on M can also be induced by the embedding $i : M \rightarrow H$. This embedding is *natural* in a sense that the formalism of local coordinates on M turns out to be a “restriction” to M of the developed functional coordinate formalism.

In fact, given vector X tangent to a path Φ_t in \mathbf{S} at the point Φ_0 and a differentiable functional F on a neighborhood of Φ_0 in \mathbf{S} , the directional derivative of F at Φ_0 along X is defined by

$$XF = \left. \frac{dF(\Phi_t)}{dt} \right|_{t=0}. \tag{11.2}$$

By applying the chain rule, we have

$$XF = F'(\Phi) \Big|_{\Phi=\Phi_0} \Phi'_t \Big|_{t=0}, \tag{11.3}$$

where $F'(\Phi) \Big|_{\Phi=\Phi_0} : \mathbf{S} \rightarrow R$ is the derivative functional at $\Phi = \Phi_0$ and $\Phi'_t \Big|_{t=0} \in \mathbf{S}$ is the derivative of Φ_t at $t = 0$. Let e_H be a functional basis on \mathbf{S} . Writing (11.3) in coordinates (\mathbf{S}, e_H^{-1}) yields

$$XF = \int \left. \frac{\delta f(\varphi)}{\delta \varphi(x)} \right|_{\varphi=\varphi_0} \xi(x) dx, \tag{11.4}$$

where $\xi = \varphi'_t|_{t=0}$, $\varphi_t = e_H^{-1}\Phi_t$, and the linear functional $\delta f(\varphi)/\delta\varphi(x)|_{\varphi=\varphi_0} \in H^*$ is the derivative functional $F'(\Phi_0)$ in the dual basis e_H^* .

We select from all paths in H the paths with values in M . In the chosen coordinates any such path $\varphi_t : [a, b] \rightarrow M$ has the form

$$\varphi_t(x) = \delta(x - a(t)) \tag{11.5}$$

for some function $a(t)$ taking values in R^n .

THEOREM 11.5. *The expression (11.2) evaluated on appropriate functionals and on a path (11.5) yields the standard expression $(\partial f(a)/\partial a^\mu)|_{a=a(0)}(da^\mu/dt)|_{t=0}$ for action of the vector tangent to the path $a(t) : [a, b] \rightarrow R^n$ at $t = 0$ on differentiable functions f .*

Proof. Assume that f is an analytic functional represented on a neighborhood of $\varphi_0 = \varphi_t|_{t=0} = \delta_{a(0)}$ in H by a convergent in H power series

$$f(\varphi) = f_0 + \int f_1(x)\varphi(x)dx + \iint f_2(x, y)\varphi(x)\varphi(y)dx dy + \dots \tag{11.6}$$

Assume further that the function

$$f(a) = f(\delta(x - a)) = f_0 + f_1(a) + f_2(a, a) + \dots \tag{11.7}$$

is smooth on a neighborhood of $a_0 = a(0)$ (e.g., f_1 is a smooth function and $f_k = 0$ for $k \geq 2$). Then on the path $\varphi_t = \delta(x - a(t))$ we have

$$\left. \frac{df(\varphi_t)}{dt} \right|_{t=0} = \left. \frac{\partial f(a)}{\partial a^\mu} \right|_{a=a(0)} \left. \frac{da^\mu}{dt} \right|_{t=0}. \tag{11.8}$$

In particular, the expression on the right of (11.8) is the action of the n -vector $(da^\mu/dt)(\partial/\partial a^\mu)$ tangent to the path $a(t)$ on the smooth function $f(a)$. □

Assume once again that H is a real Hilbert space and let $K : H \times H \rightarrow R$ be the metric on H given by a smooth kernel $k(x, y)$. If $\varphi = \varphi_t(x) = \delta(x - a(t))$ is a path in M , then for the vector $\delta\varphi(x)$ tangent to the path at φ_0 we have

$$\delta\varphi(x) \equiv \left. \frac{d\varphi_t(x)}{dt} \right|_{t=0} = -\nabla_\mu \delta(x - a) \left. \frac{da^\mu}{dt} \right|_{t=0}. \tag{11.9}$$

Here $a = a(0)$ and derivatives are understood in a generalized sense. Therefore,

$$\|\delta\varphi\|_H^2 = \int k(x, y) \nabla_\mu \delta(x - a) \left. \frac{da^\mu}{dt} \right|_{t=0} \nabla_\nu \delta(y - a) \left. \frac{da^\nu}{dt} \right|_{t=0} dx dy. \tag{11.10}$$

“Integration by parts” in the last expression gives

$$\int k(x, y)\delta\varphi(x)\delta\varphi(y)dx dy = \frac{\partial^2 k(x, y)}{\partial x^\mu \partial y^\nu} \Big|_{x=y=a} \frac{da^\mu}{dt} \Big|_{t=0} \frac{da^\nu}{dt} \Big|_{t=0}. \tag{11.11}$$

Remark 11.6. Although the above manipulations with generalized functions are somewhat formal, they can be easily justified. In particular, from Theorem 11.1 we know that as $L \rightarrow \infty$, the sequence $f_L(x - a) = (L/\sqrt{\pi})^n e^{-L^2(x-a)^2}$ converges in norm in H to $\delta(x - a)$. Similarly, the sequence of the derivatives $\nabla_\mu f_L(x - a)$ converges to $\nabla_\mu \delta(x - a)$. Performing now the ordinary integration by parts in

$$\int k(x, y)\nabla_\mu f_L(x - a)\frac{da^\mu}{dt} \Big|_{t=0} \nabla_\nu f_L(y - a)\frac{da^\nu}{dt} \Big|_{t=0} dx dy \tag{11.12}$$

and taking L to infinity, we obtain the result (11.11).

By defining $(da^\mu/dt)|_{t=0} = da^\mu$, we now have

$$\int k(x, y)\delta\varphi(x)\delta\varphi(y)dx dy = g_{\mu\nu}(a)da^\mu da^\nu, \tag{11.13}$$

where

$$g_{\mu\nu}(a) = \frac{\partial^2 k(x, y)}{\partial x^\mu \partial y^\nu} \Big|_{x=y=a}. \tag{11.14}$$

As the functional K is symmetric, the tensor $g_{\mu\nu}(a)$ can be assumed to be symmetric as well. If in addition $(\partial^2 k(x, y)/\partial x^\mu \partial y^\nu)|_{x=y=a}$ is positive definite at every a , the tensor $g_{\mu\nu}(a)$ can be identified with the Riemannian metric on an n -dimensional manifold M given in local coordinates a^μ .

Example 11.7. Consider the Hilbert space H with metric given by the kernel $k(\mathbf{x}, \mathbf{y}) = e^{-(1/2)(\mathbf{x}-\mathbf{y})^2}$ for all $\mathbf{x}, \mathbf{y} \in R^3$. Using (11.14) and assuming $(\mathbf{x} - \mathbf{y})^2 = \delta_{\mu\nu}(x^\mu - y^\mu)(x^\nu - y^\nu)$ with $\mu, \nu = 1, 2, 3$, we immediately conclude that $g_{\mu\nu}(a) = \delta_{\mu\nu}$.

The resulting isometric embedding of the Euclidean space $N = R^3$ into H is illustrated in Figure 11.1. The cones in the figure represent delta-functions forming the manifold M which we denote in this case by M_3 .

The following theorem clarifies the embedding of M_3 into the space H with the metric given by $e^{-(1/2)(x-y)^2}$.

THEOREM 11.8. *The manifold M_3 is a submanifold of the unit sphere S^H in H . Moreover, the set M_3 forms a complete system in H . The elements of any finite subset of M_3 are linearly independent.*

Proof. Observe that the norm of any element $\delta(\mathbf{x} - \mathbf{a})$ in H is equal to 1. Therefore, the three-dimensional manifold M_3 is a submanifold of the unit sphere S^H in H .

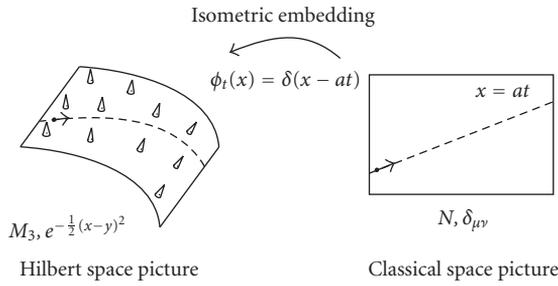


Figure 11.1

To show that the set M_3 forms a *complete system* in H we need to verify that there is no nontrivial element of H orthogonal to every element of M_3 . Assume that f is a functional in H such that $\int e^{-(1/2)(x-y)^2} f(x)\delta(y - \mathbf{u})d\mathbf{x}d\mathbf{y} = 0$ for all $\mathbf{u} \in R^3$. Then $\int e^{-(1/2)(x-\mathbf{u})^2} f(x)d\mathbf{x} = 0$ for all $\mathbf{u} \in R^3$. Since the metric $\hat{G} : H \rightarrow H^*$ given by the kernel $e^{-(1/2)(x-y)^2}$ is an isomorphism, we conclude that $f = 0$.

Assume now that $\sum_{k=1}^n c_k \delta(\mathbf{x} - \mathbf{a}_k)$ is the zero functional in H^* and the vectors $\mathbf{a}_k \in R^3$ are all different. For any such finite set of vectors, the space H^* contains functions φ_k , $k = 1, \dots, n$, with supports containing one and only one of the points \mathbf{a}_k each and so that no two supports contain the same point. Therefore the coefficients c_k must be all equal to zero, that is, the elements of any finite subset of M_3 are linearly independent. \square

Note that the set M_3 is uncountable and that no two elements of M_3 are orthogonal (although, provided $|\mathbf{a} - \mathbf{b}| \gg 1$, the elements $\delta(\mathbf{x} - \mathbf{a})$, $\delta(\mathbf{x} - \mathbf{b})$ are “almost” orthogonal).

Figures 11.2 and 11.3 help “visualizing” the embedding of M_3 into H . Under the embedding any straight line $\mathbf{x} = \mathbf{a}_0 + at$ in R^3 becomes a “spiral” $\varphi_t(\mathbf{x}) = \delta(\mathbf{x} - \mathbf{a}_0 - at)$ on the sphere S^H through dimensions of H . One such spiral is shown in Figure 11.2. The curve in Figure 11.2 goes through the tips of three shown linearly independent unit vectors. Imagine that each point on the curve is the tip of a unit vector and that any n of these vectors are linearly independent.

The manifold M_3 is spanned by such spirals. Figure 11.3 illustrates the embedding of M_3 into H in light of this result.

12. Riemannian metric on the unit sphere in L_2 and on the complex projective space CP^{L_2}

In this section we apply the developed coordinate formalism to demonstrate that solutions of the Schrödinger equation $d\varphi_t/dt = -i\hat{h}\varphi_t$ for a quantum system described by the Hamiltonian \hat{h} are geodesics in the appropriate Riemannian metric on the space of states of the system.

For this, it will turn out to be convenient to use the index notation introduced in Section 4. Thus, a string-tensor T or rank (r, s) in the index notation will be written as $t_{b_1 \dots b_s}^{a_1 \dots a_r}$. Assume that $\hat{K} : H \rightarrow H^*$ defines a Hermitian inner product $K(\xi, \eta) = (\hat{K}\xi, \eta)$ on

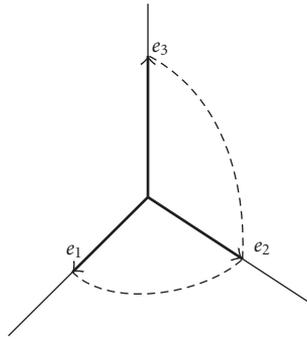


Figure 11.2

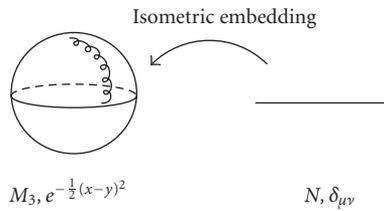


Figure 11.3

a complex Hilbert space H of complex-valued functions ξ . Let H_R be the real Hilbert space which is the realization of H . Namely, H_R is the space of pairs of vectors $X = (\xi, \bar{\xi})$ with multiplication by real numbers.

Since the inner product on H is Hermitian, it defines a real-valued Hilbert metric on H_R by

$$K_R(X, Y) = 2 \operatorname{Re} K(\xi, \eta), \tag{12.1}$$

for all $X = (\xi, \bar{\xi})$, $Y = (\eta, \bar{\eta})$ with $\xi, \eta \in H$. We will also use the “matrix” representation of the corresponding operator $\hat{K}_R : H_R \rightarrow H_R^*$:

$$\hat{K}_R = \begin{bmatrix} 0 & \hat{K} \\ \overline{\hat{K}} & 0 \end{bmatrix}. \tag{12.2}$$

In particular, we have

$$K_R(X, Y) = (\hat{K}_R X, Y) = [\xi, \bar{\xi}] \hat{K}_R \begin{bmatrix} \eta \\ \bar{\eta} \end{bmatrix} = 2 \operatorname{Re}(\hat{K} \xi, \eta), \tag{12.3}$$

where $\xi \hat{K} \bar{\eta}$ stands for the inner product $(\hat{K} \xi, \eta)$ and $\bar{\xi} \overline{\hat{K}} \eta$ stands for its conjugate.

We agree to use the capital Latin letters A, B, C, \dots as indices of tensors defined on direct products of copies of the real Hilbert space H_R and its dual. The small Latin letters

a, b, c, \dots and the corresponding overlined letters $\bar{a}, \bar{b}, \bar{c}, \dots$ will be reserved for tensors defined on direct products of copies of the complex Hilbert space H , its conjugate, dual, and dual conjugate. A single capital Latin index replaces a pair of lower Latin indices. For example, if $X \in H_R$, then $X^A = (X^a, X^{\bar{a}})$, with X^a representing an element of H and $X^{\bar{a}} = \overline{X^a}$.

Consider now the tangent bundle over a complex string space S which we identify here with a Hilbert space L_2 of square-integrable functions. We identify all fibers of the tangent bundle over L_2 (i.e., all tangent spaces $T_\varphi L_2, \varphi \in L_2$) with the complex Hilbert space H described above. We introduce a Hermitian $(0,2)$ tensor field G on the space L_2 without the origin as follows:

$$G(\xi, \eta) = \frac{(\widehat{K}\xi, \eta)}{(\varphi, \varphi)_{L_2}}, \tag{12.4}$$

for all ξ, η in the tangent space $T_\varphi L_2$ and all points $\varphi \in L_{2*}$. Here L_{2*} stands for the space L_2 without the origin.

The corresponding (strong) Riemannian metric G_R on L_2 is defined by

$$G_R(X, Y) = 2 \operatorname{Re} G(\xi, \eta), \tag{12.5}$$

where as before $X = (\xi, \bar{\xi})$ and $Y = (\eta, \bar{\eta})$. In the matrix notation of (12.2) we have, for the operator $\widehat{G}_R : H_R \rightarrow H_R^*$ defining the metric G_R ,

$$\widehat{G}_R = \begin{bmatrix} 0 & \widehat{G} \\ \widehat{G} & 0 \end{bmatrix}, \tag{12.6}$$

where $\widehat{G} : H \rightarrow H^*$ defines the metric G .

In our index notation the kernel of the operator \widehat{G} will be denoted by $g_{a\bar{b}}$, so that

$$g_{a\bar{b}} = \frac{k_{a\bar{b}}}{\|\varphi\|_{L_2}^2}, \tag{12.7}$$

where $k_{a\bar{b}}$ is the kernel of \widehat{K} . From (12.6) we have, for the components $(\widehat{G}_R)_{AB}$ of the metric \widehat{G}_R ,

$$\begin{aligned} (\widehat{G}_R)_{ab} &= (\widehat{G}_R)_{\bar{a}\bar{b}} = 0, \\ (\widehat{G}_R)_{a\bar{b}} &= g_{a\bar{b}}, \quad (\widehat{G}_R)_{\bar{a}b} = \bar{g}_{\bar{a}b}. \end{aligned} \tag{12.8}$$

For this reason and with the agreement that $g_{\bar{a}b}$ stands for $\bar{g}_{\bar{a}b}$, we can denote the kernel of \widehat{G}_R by g_{AB} . For the inverse metric, we have

$$\widehat{G}_R^{-1} = \begin{bmatrix} 0 & \widehat{G}^{-1} \\ \widehat{G}^{-1} & 0 \end{bmatrix}. \tag{12.9}$$

Let the notation $g^{\bar{a}b}$ stand for the kernel of the inverse operator \widehat{G}^{-1} and let $g^{a\bar{b}}$ stand for its conjugate $\overline{g^{\bar{a}b}}$. Then

$$\begin{aligned} (\widehat{G}_R)^{ab} &= (\widehat{G}_R)^{\bar{a}\bar{b}} = 0, \\ (\widehat{G}_R)^{\bar{a}b} &= g^{\bar{a}b}, \quad (\widehat{G}_R)^{a\bar{b}} = \overline{g^{\bar{a}b}}. \end{aligned} \tag{12.10}$$

Accordingly, without danger of confusion, we can denote the kernel of \widehat{G}_R^{-1} by g^{AB} .

Having the Riemannian metric G_R on L_2 we can define the compatible (*Riemannian*, or *Levi-Civita*) connection Γ by

$$2G_R(\Gamma(X, Y), Z) = dG_RX(Y, Z) + dG_RY(Z, X) - dG_RZ(X, Y), \tag{12.11}$$

for all vector fields X, Y, Z in H_R . Here, for example, the term $dG_RX(Y, Z)$ denotes the derivative of the inner product $G_R(Y, Z)$ evaluated on the vector field X . In the given realization of the tangent bundle, for any $\varphi \in L_2$, the connection Γ is an element of the space $L(H_R, H_R; H_R)$. The latter notation means that Γ is an H_R -valued 2-form on $H_R \times H_R$. In our index notation (12.11) can be written as

$$2g_{AB}\Gamma_{CD}^B = \frac{\delta g_{AD}}{\delta \varphi^C} + \frac{\delta g_{CA}}{\delta \varphi^D} - \frac{\delta g_{CD}}{\delta \varphi^A}. \tag{12.12}$$

Here for any $\varphi \in L_2$ the expression $g_{AB}\Gamma_{CD}^B$ is an element of $L(H_R, H_R, H_R; R)$, that is, it is an R -valued 3-form defined by

$$g_{AB}\Gamma_{CD}^B X^C Y^D Z^A = G_R(\Gamma(X, Y), Z) \tag{12.13}$$

for all $X, Y, Z \in H_R$. Similarly, for any $\varphi \in L_2$, the functional derivative $\delta g_{AD}/\delta \varphi^C$ is an element of $L(H_R, H_R, H_R; R)$ defined by

$$\frac{\delta g_{AD}}{\delta \varphi^C} X^C Y^D Z^A = dG_RX(Y, Z). \tag{12.14}$$

For any $\varphi \in L_2$, by leaving vector Z out, we can treat both sides of (12.11) as elements of H^* . Recall now that G_R is a strong Riemannian metric. That is, for any $\varphi \in L_2$ the operator $\widehat{G}_R : H_R \rightarrow H_R^*$ is an isomorphism, that is, \widehat{G}_R^{-1} exists. By applying \widehat{G}_R^{-1} to both sides of (12.11) without Z , we have in the index notation

$$2\Gamma_{CD}^B = g^{BA} \left(\frac{\delta g_{AD}}{\delta \varphi^C} + \frac{\delta g_{CA}}{\delta \varphi^D} - \frac{\delta g_{CD}}{\delta \varphi^A} \right), \tag{12.15}$$

where

$$\Gamma_{CD}^B X^C Y^D \Omega_B = (\widehat{G}_R^{-1}(\widehat{G}_R \Gamma(X, Y)), \Omega). \tag{12.16}$$

Formula (12.15) defines the connection “coefficients” (Christoffel symbols) of the Levi-Civita connection. From the matrix form of \widehat{G}_R and \widehat{G}_R^{-1} we can now easily obtain

$$\begin{aligned} \Gamma_{cd}^b &= \overline{\Gamma}_{\bar{c}\bar{d}}^{\bar{b}} = \frac{1}{2}g^{\bar{a}b} \left(\frac{\delta g_{d\bar{a}}}{\delta \varphi^c} + \frac{\delta g_{c\bar{a}}}{\delta \varphi^d} \right), \\ \Gamma_{\bar{c}\bar{d}}^{\bar{b}} &= \overline{\Gamma}_{cd}^{\bar{b}} = \frac{1}{2}g^{\bar{a}b} \left(\frac{\delta g_{c\bar{a}}}{\delta \overline{\varphi}^d} - \frac{\delta g_{d\bar{a}}}{\delta \overline{\varphi}^c} \right), \\ \Gamma_{\bar{c}\bar{d}}^b &= \overline{\Gamma}_{cd}^{\bar{b}} = \frac{1}{2}g^{\bar{a}b} \left(\frac{\delta g_{d\bar{a}}}{\delta \overline{\varphi}^c} - \frac{\delta g_{c\bar{a}}}{\delta \overline{\varphi}^d} \right), \end{aligned} \tag{12.17}$$

while the remaining components vanish. To compute the coefficients, we write the metric (12.7) in the form

$$g_{\bar{a}\bar{b}} = \frac{k_{\bar{a}\bar{b}}}{\delta_{u\bar{v}}\varphi^u\overline{\varphi}^v}, \tag{12.18}$$

where $\delta_{u\bar{v}} \equiv \delta(u - v)$ is the L_2 metric in the index notation. We then have for the derivatives

$$\begin{aligned} \frac{\delta g_{\bar{a}\bar{b}}}{\delta \varphi^c} &= -\frac{k_{\bar{a}\bar{b}}\delta_{c\bar{v}}\overline{\varphi}^v}{\|\varphi\|_{L_2}^4}, \\ \frac{\delta g_{\bar{a}\bar{b}}}{\delta \overline{\varphi}^c} &= -\frac{k_{\bar{a}\bar{b}}\delta_{uc}\varphi^u}{\|\varphi\|_{L_2}^4}. \end{aligned} \tag{12.19}$$

Using (12.17), we can now find the nonvanishing connection coefficients

$$\begin{aligned} \Gamma_{cd}^b &= \overline{\Gamma}_{\bar{c}\bar{d}}^{\bar{b}} = -\frac{(\delta_d^b\delta_{c\bar{v}} + \delta_c^b\delta_{d\bar{v}})\overline{\varphi}^v}{2\|\varphi\|_{L_2}^2}, \\ \Gamma_{\bar{c}\bar{d}}^{\bar{b}} &= \overline{\Gamma}_{cd}^{\bar{b}} = -\frac{(\delta_c^b\delta_{u\bar{d}} - k^{\bar{a}b}k_{c\bar{d}}\delta_{u\bar{a}})\varphi^u}{2\|\varphi\|_{L_2}^2}, \\ \Gamma_{\bar{c}\bar{d}}^b &= \overline{\Gamma}_{cd}^{\bar{b}} = -\frac{(\delta_d^b\delta_{u\bar{c}} - k^{\bar{a}b}k_{d\bar{c}}\delta_{u\bar{a}})\varphi^u}{2\|\varphi\|_{L_2}^2}. \end{aligned} \tag{12.20}$$

Consider now the unit sphere $S^{L_2} : \|\varphi\|_{L_2} = 1$ in the space L_2 . Let \widehat{A} be a (possibly unbounded) injective Hermitian operator defined on a set $D(\widehat{A})$ and with the image $\widehat{A}(D) \equiv R(\widehat{A})$. Here we assume for simplicity that $D(\widehat{A}) \subset R(\widehat{A})$ and that both $D(\widehat{A})$ and $R(\widehat{A})$ are dense subsets of L_2 . We define the inner product $(f, g)_H$ of any two elements f, g in $R(\widehat{A})$ by the formula $(f, g)_H \equiv (\widehat{A}^{-1}f, \widehat{A}^{-1}g)_{L_2} = ((\widehat{A}\widehat{A}^*)^{-1}f, g)$. By completing $R(\widehat{A})$ with respect to this inner product we obtain a Hilbert space H . Notice that \widehat{A} is bounded in this norm and can be therefore extended to the entire space L_2 and becomes an isomorphism from L_2 onto H . We will denote such an extension by the same symbol \widehat{A} . Let $\widehat{K} = (\widehat{A}\widehat{A}^*)^{-1}$, $\widehat{K} : H \rightarrow H^*$ be the metric operator on H . As before, we define the Riemannian metric on L_{2*} by

$$G_R(X, Y) = \frac{2 \operatorname{Re}(\widehat{K}\xi, \eta)}{(\varphi, \varphi)_{L_2}}, \tag{12.21}$$

where $X = (\xi, \bar{\xi})$, $Y = (\eta, \bar{\eta})$. Assume that the sphere $S^{L_2} \subset L_{2*}$ is furnished with the induced Riemannian metric. Consider now the vector field $A_\varphi = -i\hat{A}\varphi$ associated with the operator \hat{A} . The integral curves of this vector field are solutions of the equation $d\varphi_t/dt = -i\hat{A}\varphi_t$. These solutions are given by $\varphi_\tau = e^{-i\hat{A}\tau}\varphi_0$. Since $e^{-i\hat{A}\tau}$ denotes a one-parameter group of unitary operators, the integral curve φ_τ through a point $\varphi_0 \in S^{L_2}$ stays on S^{L_2} . In particular, the vector field A_φ is tangent to the sphere. In other words, the operator $-i\hat{A}$ maps points on the sphere into vectors tangent to the sphere.

We claim now that the curves $\varphi_\tau = e^{-i\hat{A}\tau}\varphi_0$ are geodesics on the sphere in the induced metric. That is, they satisfy the equation

$$\frac{d^2\varphi_\tau}{d\tau^2} + \Gamma\left(\frac{d\varphi_\tau}{d\tau}, \frac{d\varphi_\tau}{d\tau}\right) = 0. \tag{12.22}$$

In fact, using (12.20) and collecting terms, we obtain

$$\Gamma_{CD}^b \frac{d\varphi_\tau^C}{d\tau} \frac{d\varphi_\tau^D}{d\tau} = \frac{(\hat{K}(d\varphi_\tau/d\tau), (d\varphi_\tau/d\tau))\hat{A}^2\varphi_\tau}{\|\varphi_\tau\|_{L_2}^2}. \tag{12.23}$$

The expression for $\Gamma_{CD}^b(d\varphi_\tau^C/d\tau)(d\varphi_\tau^D/d\tau)$ turns out to be the complex conjugate of (12.23). Now, the substitution of $\varphi_\tau = e^{i\hat{A}\tau}\varphi_0$ and $\hat{K} = (\hat{A}\hat{A}^*)^{-1}$ into the right-hand side of (12.23) yields $\hat{A}^2\varphi_\tau$. At the same time, $d^2\varphi_\tau/d\tau^2 = -\hat{A}^2\varphi_\tau$ and therefore (12.22) is satisfied. That is, the curves $\varphi_\tau = e^{-i\hat{A}\tau}\varphi_0$ are geodesics in the metric (12.21) on L_{2*} . Since these curves also belong to the sphere S^{L_2} and the metric on the sphere is induced by the embedding $S^{L_2} \rightarrow L_{2*}$, we conclude that the curves φ_τ are geodesics on S^{L_2} .

Assume in particular that \hat{A} is the Hamiltonian \hat{h} of a quantum system. Then the above model demonstrates that, in the appropriate Riemannian metric on the unit sphere S^{L_2} , the Schrödinger evolution of the system is a motion along a geodesic of S^{L_2} . For a closely related metric on S^{L_2} this result was obtained earlier in [7] by means of variational principle.

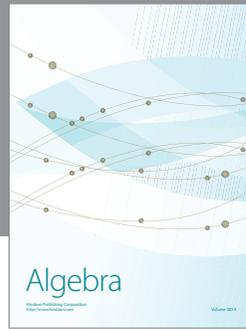
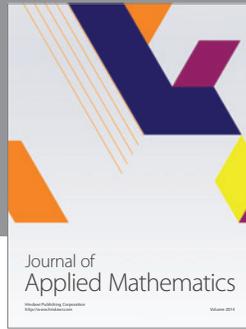
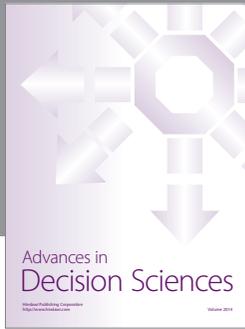
References

- [1] V. Bargmann, *On a Hilbert space of analytic functions and an associated integral transform*, Comm. Pure Appl. Math. **14** (1961), 187–214.
- [2] I. M. Gel’fand and G. E. Shilov, *Generalized Functions. Vol. 2. Spaces of Fundamental and Generalized Functions*, Academic Press, New York, 1968.
- [3] I. M. Gel’fand and N. Ya. Vilenkin, *Generalized Functions. Vol. 4. Applications of Harmonic Analysis*, Academic Press, New York, 1964.
- [4] B. C. Hall, *Holomorphic methods in analysis and mathematical physics*, First Summer School in Analysis and Mathematical Physics (Cuernavaca Morelos, 1998), Contemp. Math., vol. 260, American Mathematical Society, Rhode Island, 2000, pp. 1–59.
- [5] A. A. Kryukov, *Coordinate formalism on abstract Hilbert space: kinematics of a quantum measurement*, Found. Phys. **33** (2003), no. 3, 407–443.
- [6] ———, *Coordinate formalism on Hilbert manifolds*, Mathematical Physics Research at the Cutting Edge, Nova Science Publishers, New York, 2004, pp. 1–33.
- [7] ———, *On the problem of emergence of classical space-time: the quantum-mechanical approach*, Found. Phys. **34** (2004), no. 8, 1225–1248.

- [8] ———, *Quantum mechanics on Hilbert manifolds: the principle of functional relativity*, to appear.
- [9] D. Lokenath and P. Mikusinski, *Introduction to Hilbert Spaces with Applications*, 3rd ed., Elsevier, Boston, 2005.

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