

Research Article

Processing and Performance of Polymeric Transparent Conductive Composites

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Recent advances in microelectronic and optoelectronic industries have spurred interest in the development of reticulate doped polymer films containing “metallic” charge transfer complexes. In this study, such reticulate doped polymer films were prepared by exposing solid solutions of bis(ethylenedioxy) tetrathiafulvalene (BEDO-TTF) in polycarbonate (PC) to iodine, forming conductive charge transfer complexes. The resulting films exhibited room temperature conductivities ranging from 6.33 to $90.4 \times 10^{-5} \text{ S} \cdot \text{cm}^{-1}$. The colored iodine complexes in the film were reduced by cyclic voltammetry yielding conductive, colorless, transparent films. We were intrigued to examine the dielectric properties of BEDO-TTF in solid solution in PC prior to formation of the charge transfer complex as no such studies appear in the literature. Dielectric analysis (DEA) was used to probe relaxations in neat PC and BEDO-TTF/PC. BEDO-TTF plasticized the PC and decreased the glass transition temperature. Two secondary relaxations appeared in PC films, whereas the transitions merged in the BEDO-TTF/PC film. DEA also evidenced conductivity relaxations above 180°C which are characterized via electric modulus formalism and revealed that BEDO-TTF increased AC conductivity in PC.

1. Introduction

In the last few decades, there has been growing interest in conductive polymers and composites due to an array of potential applications in biological and chemical sensors, separations, microelectronics circuit boards, biomedical, coatings, and optical displays [1–5]. The ability to prepare thin, flexible, transparent films is an asset in applications requiring nanostructuring and miniaturization [6]. Conductivity is a general property of metals, but some polymers in combination with organic charge moieties exhibit metallic behavior [7]. After the discovery of the first organic superconductor, $(\text{TMTSeF})_2\text{PF}_6$, a growing number of organic charge transfer complexes and conductive salts have been investigated [8]. These salts exhibit a charge transfer between a donors and acceptors in the solid state; the donors and acceptors molecules form segregated stacked sheets of cations

and anions, respectively [9–11]. These materials also have low critical temperatures, T_c s (the temperature below which material is superconducting). $(\text{TMTSeF})_2\text{PF}_6$ was subsequently replaced by another class of organic superconductors with BEDT-TTF (bis(ethylenedithio)tetrathiafulvalene) donors (Figure 1(a)) containing sulfur heterocycles [9]. The peripheral sulfur atoms allow better orbital overlap between donor stacks, forming a two-dimensional network. Suzuki et al. proposed that the substitution of sulfur or selenium atoms by lighter atoms, such as oxygen, would increase the T_c of organic superconductors. They synthesized an oxygen analogue of BEDT-TTF, (bis(ethylenedioxy)tetrathiafulvalene) (BEDO-TTF) (Figure 1(b)) in which outer sulfur atoms are replaced by oxygen atoms [12, 13]. BEDO-TTF was synthesized postulating that π donation from oxygen atoms into tetrathiafulvalene ring would lower the first ionization energy and make it an easily oxidized donor molecule [14]. In its

exhibit color ranging from shiny purple to dark green. Light purple to olive green films with metallic sheen are conductive. The metallic sheen increased as the optimum conductivity was reached. BEDO-TTF/PC films exposed for a longer period of time to iodine vapors exhibited dull green to dark purple color with no metallic sheen, thereby resulting in nonconductive films. Long exposure time led to changes in the oxidation state of BEDO-TTF, resulting films with nonconductive surfaces [26].

2.4. Characterization

2.4.1. Differential Scanning Calorimetry. Calorimetric experiments were carried out using TA Instruments DSC 2920. The amount of sample used was about 5 mgs, and the heating rate was maintained at 10°C/min. Samples were scanned in the range from 25 to 300°C and were encapsulated in hermetically sealed aluminum pans under a nitrogen purge rate of 70–80 mL/min. Temperature calibrations were performed with indium as a standard. All data analysis was performed using the TA instruments universal analysis program, version 3.9A. Two scans were performed for films: the first scan erased thermal history, and the second scan was used to determine the glass transition temperature.

2.4.2. Dielectric Analysis. Note that dielectric analysis was conducted on neat PC and BEDO-TTF/PC films. The iodine complexes conduct electricity and are not suitable for DEA. Dielectric measurements were performed with the TA instruments DEA 2970 using single surface electrodes. The samples were heated to 225°C to embed the sample into the channels of the single surface and then cooled to –120°C with liquid nitrogen. A maximum force of 250 N was exerted on the samples to achieve a minimum spacing of 0.25 mm, which ensures good contact between the samples and the sensors. The experiments were performed with the temperature range from –120 to 225°C, with 4°C increments through a frequency range of 1 Hz to 100 kHz. A dry helium atmospheric purge of 500 mL/min was used to create an inert atmosphere. Capacitance and conductance were measured as a function of time, temperature, and frequency to obtain the dielectric constant or permittivity (ϵ'), dielectric loss (ϵ''), and loss tangent ($\tan \delta = \epsilon''/\epsilon'$).

2.4.3. UV-Vis Spectroscopy. UV-Vis films and solutions measurements were performed using a Perkin Elmer Lambda 40 UV-Vis-NIR double beam spectrophotometer with scan range from 190 nm to 800 nm. The UV-vis spectra of PC, BEDO-TTF, and BEDO-TTF/PC solutions in dichloromethane were recorded using quartz cell with 1 cm path length and dichloromethane as reference solution. PC and BEDO-TTF/PC and BEDO-TTF/PC-Iodine doped films were scanned with quartz slide.

2.4.4. Optical Images. Surface morphology of the BEDO-TTF/PC-Iodine films was studied using a Leica optical microscope in the reflection mode.

2.4.5. Four-Point Probe. The four-point probe method was used for the resistivity measurements on BEDO-TTF/PC-Iodine films. The probe consists of four linearly arranged and equally spaced electrodes, which remain in contact with a sample. The current, I , was supplied to the material through two outside probes with the help of a Keithley 6221 DC and AC current source, and steady voltage across the other two inside probes, V , was determined by Keithley 6514 system electrometer. Resistivity of BEDO-TTF/PC-Iodine films was measured randomly at different locations on the film's surface. Voltages were measured in volts and current in milliamperes. Electrical conductivity, σ , was obtained by simply inverting the corresponding values of the surface resistivity.

2.4.6. Electrochemistry. The electrochemical modifications on BEDO-TTF/PC-Iodine film (0.095 mol · L⁻¹, exposed for 2 minutes) were carried out using a CH 760 electrochemistry workstation (CH Instrument, TX, USA). An Ag/AgCl electrode (BASi) was used as the reference electrode, and a platinum wire was used as the counter electrode. The electrolyte was 0.1 M KI solution in water.

3. Results and Discussion

3.1. UV-Vis Spectroscopy. UV-vis spectra for solutions and films are shown in Figure 2. In solution, BEDO-TTF exhibits strong absorption peaks at 314 nm and 334 nm and a weaker absorption peak around 518 nm. These absorption bands have been observed in the molecular absorption of neutral BEDO-TTF (undoped) earlier [26]. PC in DCM exhibits an absorption peak at 235 nm and another at 265 nm with a shoulder around 271 nm. These absorptions are characteristic of the bisphenol A unit which exhibits a main absorbance band at 226 nm with a secondary band at 276 nm [29].

PC films with and without BEDO-TTF were scanned as well. The purpose was to note any scatter at wavelengths above the electronic transition bands. The films were approximately 50 microns thick. At this thickness the absorbance was over 3.0 at wavelengths below 390 nm. However, the baseline at wavelengths above the absorption maxima indicates a lack of any appreciable scattering, and this indicates that the BEDO-TTF did not fall out of solution in the PC matrix. BEDO-TTF/PC-Iodine doped films were not transparent and became more metallic in appearance as the conductivity increased.

3.2. Optical Images. Figure 3 shows surface morphology of BEDO-TTF/PC-Iodine films at different iodine concentrations, where the conductivity is maximum. The surface exhibits fine regular grooves on optimized BEDO-TTF/PC-Iodine films. Longer exposure times and lower iodine concentrations results yield more pronounced surface topology, which, as shown in the next section, correlates with increased conductivity.

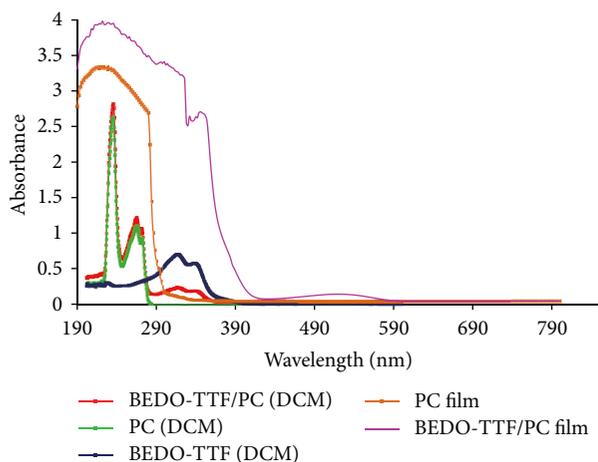


FIGURE 2: UV-Vis spectra of BEDO-TTF/PC, PC and BEDO-TTF dissolved in DCM and UV-Vis spectra of PC film and BEDO-TTF/PC film.

3.3. Electrical Properties

3.3.1. Four-Point Probe. Table 1 shows, as expected, that the surface of the BEDO-TTF/PC-Iodine films becomes conductive as a result of complex formation between acceptor iodine and donor species BEDO-TTF. This correlates with the surface topology viewed via light microscopy. The opposite side of the films are, of course, nonconductive [26]. It has been shown that there is a window of optimum exposure time, wherein the conductivity maximizes [26]. Beyond this time, films are nonconductive due to a change in the oxidation state of BEDO-TTF. Table 1 shows the calculated average resistivity and conductivity of these BEDO-TTF/PC-Iodine films. The exposure time required to reach minimum resistivity for each BEDO-TTF/PC-Iodine film increases with a decrease in iodine concentration. The surface resistivity values measured at three locations on the surface of the film were clustered closely around one central value. Standard deviations of surface resistivity for films when exposed for 2, 4, and 10 min are ± 1.0 K Ω /sq, ± 0.02 K Ω /sq, and ± 0.005 K Ω /sq, respectively.

3.3.2. Electrochemical Properties. Figure 4 shows the cyclic voltammogram of BEDO-TTF/PC-Iodine film in 0.1 M KI solution. The oxidation peak at 0.5 V in the anodic scan corresponds to the oxidation of iodide to iodine, and the reduction peak at -0.15 V in the reverse scan is due to reduction of iodine to iodide [30, 31]. The oxidation and reduction of iodine at these potentials is also confirmed by the visual observation of the coloration of the film at high potentials and discoloration at low potentials. It is interesting to note that while the reversible oxidation/reduction peak of BEDO or TTF was reported ca. 0.4–0.6 V in organic solvents [14, 32], no distinguishable peak for BEDO-TTF was observed on the film in aqueous solution. This is due to the presence of iodine in the film which can react with BEDO-TTF [33]. Furthermore, the electrochemical method can be used to control the oxidation state of iodine species in the film. For

instance, by holding the potential at -0.4 V, the iodine species in the film was reduced to iodide. Then, the conductivity of the reduced film was measured to ascertain the effect of iodine on the conductivity of the film. The discoloration gives rise to transparent film with 10% loss in conductivity. This finding encourages further studies on optimizing these systems.

3.3.3. Differential Scanning Calorimetry. The DSC thermograms for PC and BEDO-TTF/PC films are shown in Figure 5. The glass transition temperature for PC and BEDO-TTF/PC films are 148.5°C and 142.7°C, respectively [34–36]. The glass transition temperature of the BEDO-TTF/PC film decreased by about 6°C as compared to the neat film, indicating that the BEDO-TTF dissolves in and plasticizes the film.

3.3.4. Dielectric Analysis. The summary of dielectric analysis is provided for readers interested in a brief review. DEA is a thermal analysis technique well suited to the study of relaxations in polar polymers. In dielectric experiments, a sample is exposed to an alternating electric field, which generates an alternating electric polarization. The polarization causes the output current to lag behind the applied electric field by a phase shift angle, θ . DEA helps in determining the capacitance and conductance as $f(t, T, f)$ [37]. The capacitance and conductance of a material are measured over a range of temperatures and frequencies and are related to dielectric permittivity, ϵ' , and dielectric loss factor, ϵ'' , respectively. The dielectric permittivity, ϵ' , represents the amount of dipole alignment (both induced and permanent), and the dielectric loss factor, ϵ'' , measures the amount of energy required to align the dipoles or move ions. The complex permittivity, ϵ^* , is defined as follows [38]:

$$\epsilon^* = \epsilon' - i\epsilon''$$

$$\epsilon' = \epsilon'_{\text{induced dipole}} + \epsilon'_{\text{alignment of dipole}} \quad (1)$$

$$\epsilon'' = \epsilon''_{\text{dipole loss factor}} + \epsilon''_{\text{ionic conductance}}$$

where ϵ' and ϵ'' are the real and imaginary components of the dielectric complex permittivity ϵ^* . ϵ' , dielectric constant or permittivity, represents the amount of dipole alignment both induced and permanent in the sample. The permittivity of a dielectric material is measured relative to that of a vacuum ($\epsilon_0 = 8.85 \times 10^{-12}$ F·m $^{-1}$) [39]. Plots of ϵ'' versus temperature maximize at temperatures that increase with frequency, f . Plots of $\ln f$ versus $1/T_{\text{max}}$ are used to characterize transitions in polymers. Linear plots indicate secondary relaxations due to small scale motion such as side group rotation in the backbone. The large scale segmental motion accompanying the glass transition results in nonlinear plots discussed later.

However, at high temperatures conductivity often obscures the maxima noted in ϵ'' versus temperature. The Maxwell-Wagner-Sillars (MWS) effect created in a heterogeneous environment due to accumulation of charge, electrode polarization, impurities, and so forth masks relaxation behavior [40–43].

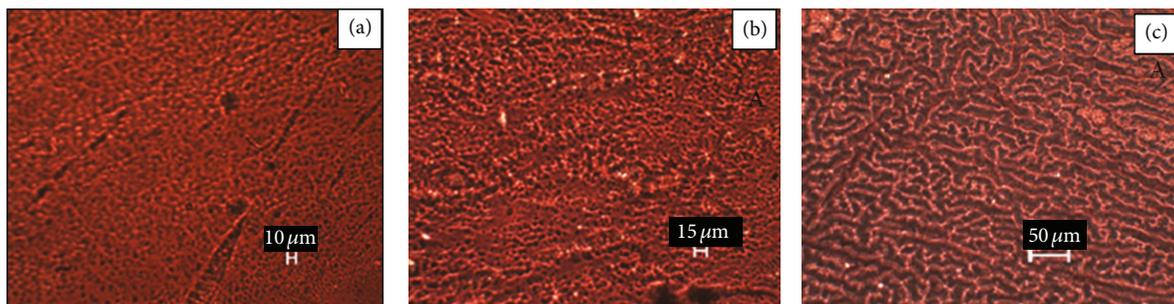


FIGURE 3: Optical images of BEDO-TTF/PC-Iodine film surface obtained using the transmission mode of a Leica microscope at (20X) with different exposure times and I_2 concentration (a) 0.095 mol/L-2 min, (b) 0.063 mol/L-4 min, and (c) 0.033 mol/L-10 min.

TABLE 1: Surface resistivity and conductivity of BEDO-TTF/PC-Iodine films.

BEDO-TTF (wt%) in PC	Iodine conc. in DCM (mol L ⁻¹)	Exposure time (mins)	Surface resistivity (KΩ/sq)	Conductivity (S cm ⁻¹)
2	0.095	2	15.82 ± 1.0	6.33E - 05
2	0.063	4	1.47 ± 0.02	6.80E - 04
2	0.033	10	1.10 ± 0.005	9.04E - 04

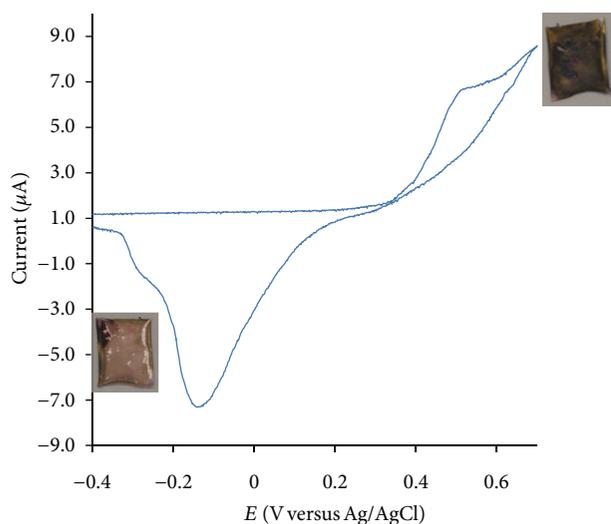


FIGURE 4: Cyclic voltammogram of BEDO-TTF/Iodine film with 0.095 mol/L (2 mins) in 0.1 M KI aqueous solution. Starting potential is 0.2 V, and scan rate is 5 mV/s.

McCrum et al. formulated a mathematical treatment of the complex permittivity, ϵ^* , and defined electric modulus (M^*) as the inverse of the complex permittivity (ϵ^*) as follows [38]:

$$M^* = \frac{1}{\epsilon^*} = M' + iM'' = \frac{\epsilon'}{\epsilon'^2 + \epsilon''^2} + i \frac{\epsilon''}{\epsilon'^2 + \epsilon''^2}, \quad (2)$$

where M^* is the complex electric modulus, M' is the electric storage modulus, and M'' is the electric loss modulus.

The dielectric loss spectra (ϵ'') versus temperature provide information on different relaxations, whereas loss modulus (M'') versus temperature reveal viscoelastic and ionic conductivity region. Figures 6(a) and 7(a) are plots of ϵ'' versus temperature for PC and BEDO-TTF/PC. Figures 6(b) and

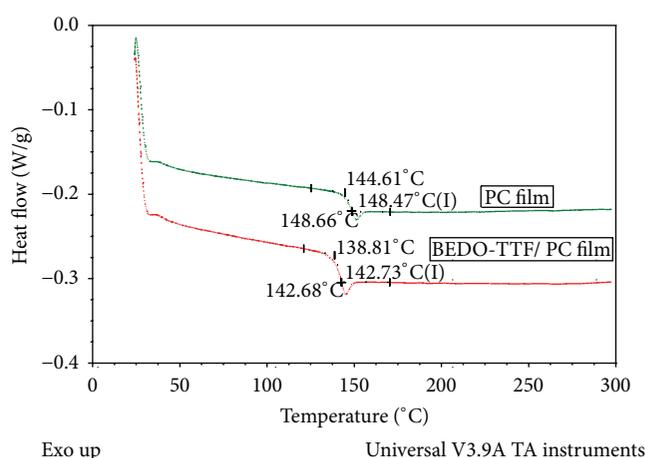


FIGURE 5: DSC thermograms for (a) PC film and (b) BEDO-TTF/PC film.

7(b) are plots for M'' versus temperature for PC and BEDO-TTF/PC. The secondary β relaxation is enlarged in the insets of the ϵ'' plots. The glass transition region also exhibits clear maxima in ϵ'' versus temperature plots. Many other polymers exhibit conductivity effects that obscure the glass transition region. However, when electric modulus formalism is used, the glass transition region is again visible, but in addition, a high temperature relaxation, the conductivity relaxation, is noted. This region is important because it provides information for calculating AC and DC conductivities, as well as activation energies for ion translation [44]. Each relaxation region is further discussed in detail.

(1) *Secondary Relaxation*. The secondary relaxation, β , in PC has been widely studied by a host of techniques such as NMR [45–48], neutron scattering [49], molecular dynamic simulations [50], dynamic mechanical spectroscopy [51–54], dielectric analysis [55, 56], depolarized Rayleigh scattering

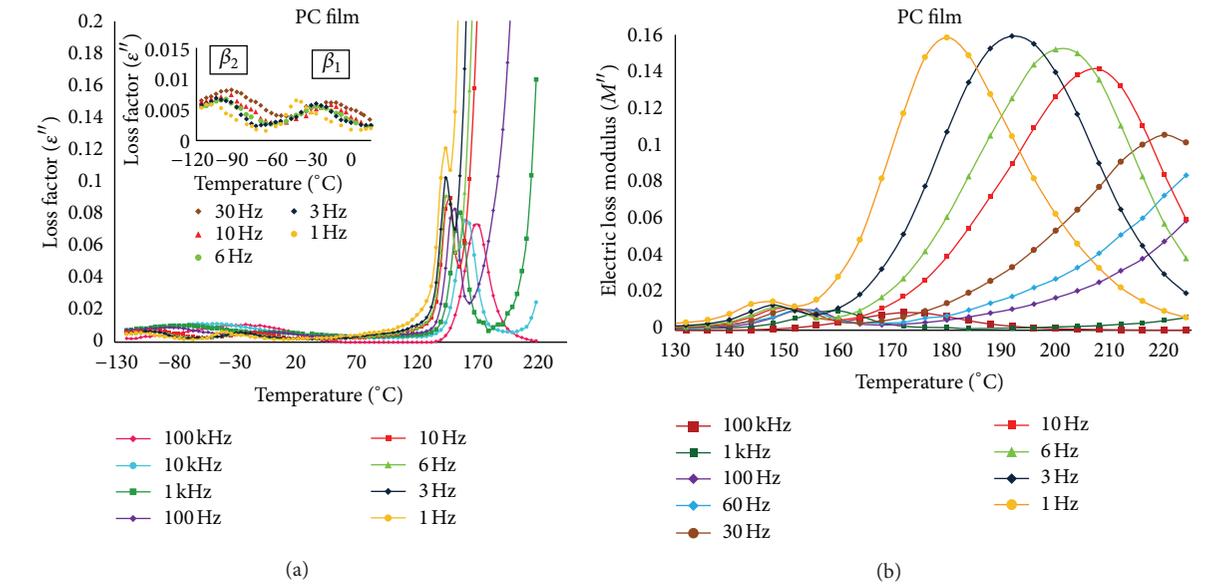


FIGURE 6: (a) ϵ'' versus temperature ($^{\circ}\text{C}$) and (b) M'' versus temperature ($^{\circ}\text{C}$) plots for PC film.

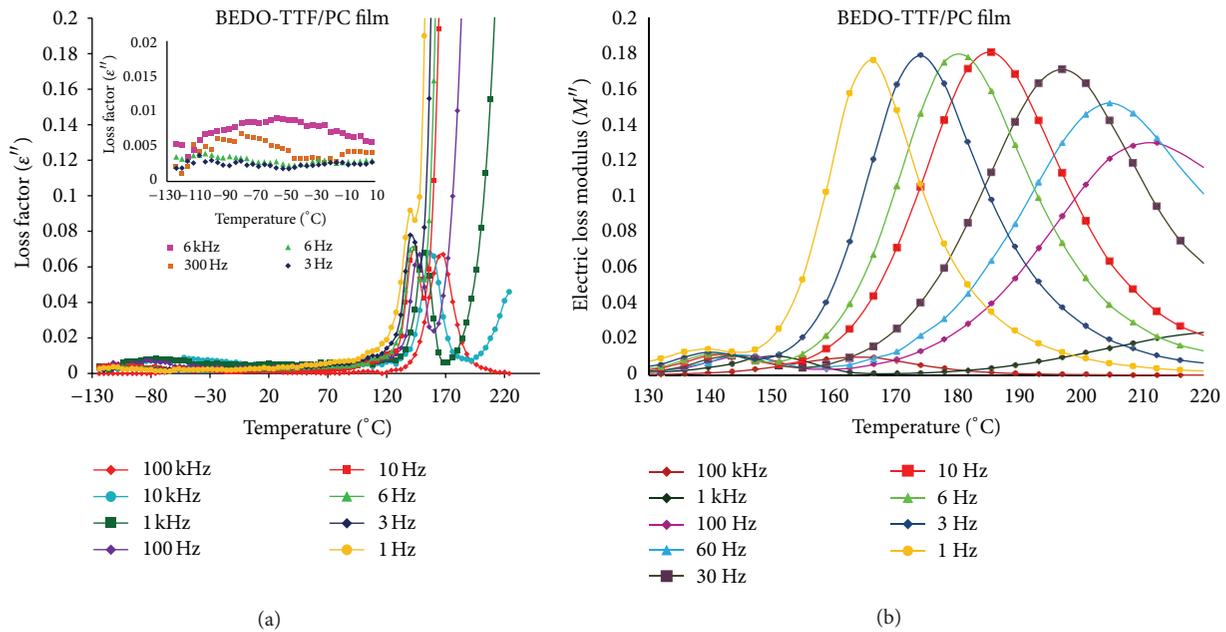


FIGURE 7: (a) ϵ'' versus temperature ($^{\circ}\text{C}$) and (b) M'' versus temperature ($^{\circ}\text{C}$) plots for BEDO-TTF/PC film.

[57], light scattering [58, 59], and thermally stimulated discharge current (TSC) [60]. Different explanations of the origin of beta relaxation have been proposed over the last few decades. Earlier investigations on PC ascertained that beta relaxation is a single, unresolved relaxation that involves rotation of carbonate groups [51, 55]. Later several studies were performed on beta relaxation over a broad frequency range and at the wide temperatures where the secondary transition separated into two or three regions resulting from phenyl ring motion, carbonate group motion, and coupled

phenyl ring and carbonate motion [61–66]. Our DEA studies for PC reveal a broad β relaxation observed at higher frequencies (1 kHz–100 kHz). The inset in Figure 6(a) shows that the broad β relaxation process was further resolved prominently into two prominent components at lower frequencies (1 Hz–30 Hz). The relaxation occurring at higher temperature is termed as β_1 relaxation, whereas other peaks at lower temperatures are termed as β_2 relaxation. In comparison to PC, the BEDO-TTF/PC film exhibits only one secondary relaxation in the frequency range 1 Hz–100 kHz.

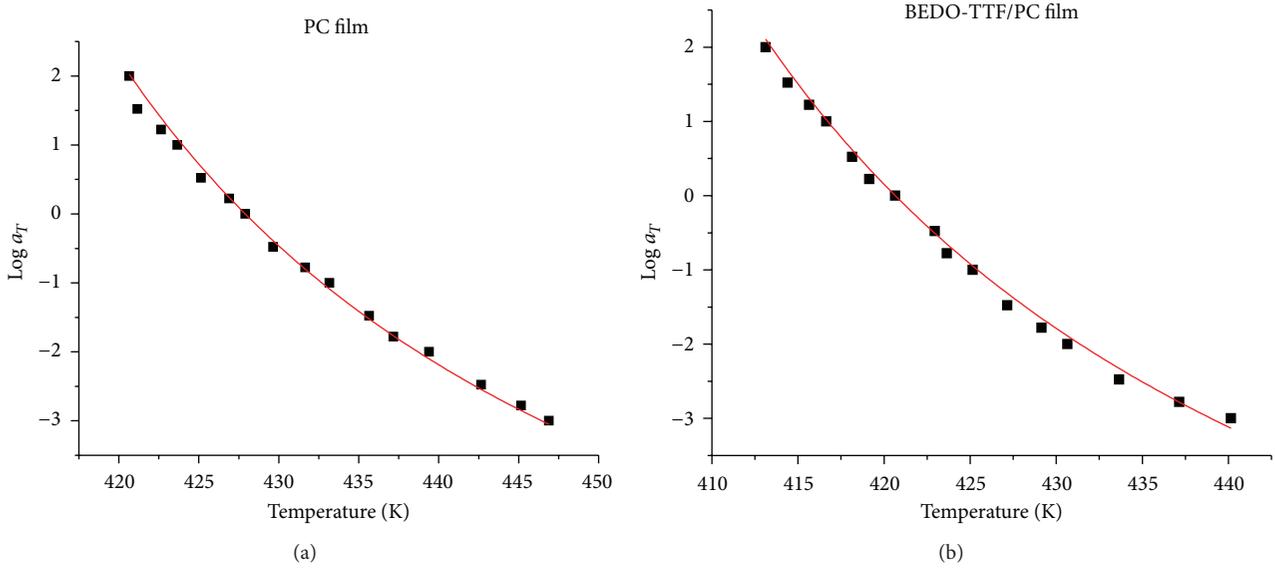


FIGURE 8: WLF plot for glass transition of (a) PC and (b) BEDO-TTF/PC films.

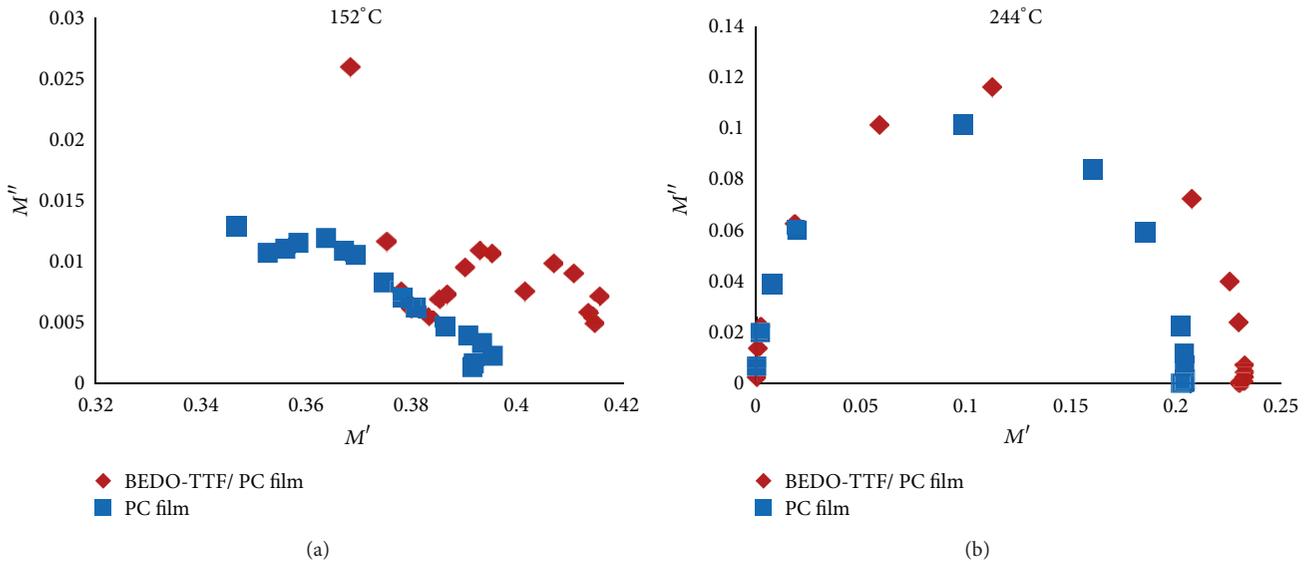


FIGURE 9: Argand plots derived from the conductivity relaxation region at different temperatures (152 and 224°C).

Peak maxima for each of the relaxation processes were assigned, and the slope of line obtained from Arrhenius plots of \ln frequency versus reciprocal of temperature was used to calculate activation energies. The activation energies for β_1 and β_2 processes in PC in frequency range 1 Hz–30 Hz were found to be 44 and 54 kJ/mol, whereas for PC, it was determined to be 47 kJ/mole in frequency range 1 kHz–100 kHz. The activation energy for BEDO-TTF/PC in frequency range 1 kHz–100 kHz is 38 kJ/mole, respectively.

(2) *Primary Relaxation.* In PC, the alpha transition or glass transition involves micro-Brownian motion in the main chain and conformational changes in the phenyl groups [67,

68]. The α relaxation is discernible in ϵ'' versus temperature plots. The maxima in ϵ'' for Tg occur from 148 to 174°C for PC and from 140 to 168°C for BEDO-TTF/PC. The Williams-Landel-Ferry (WLF) equation (3) was used to characterized relaxation behavior [69, 70]. Plots of $\log a_T$ versus temperature were constructed as shown in Figure 8 Consider

$$\ln a_T = \log \frac{\tau(T)}{\tau(T_0)} = -\log \frac{\omega}{\omega_0} = \frac{-C_1(T - T_0)}{C_2 + (T - T_0)}, \quad (3)$$

where a_T is the shift factor that corresponds to frequency, T is a given temperature, T_0 is the reference temperature, τ is the relaxation time, and ω_0 is the angular frequency at

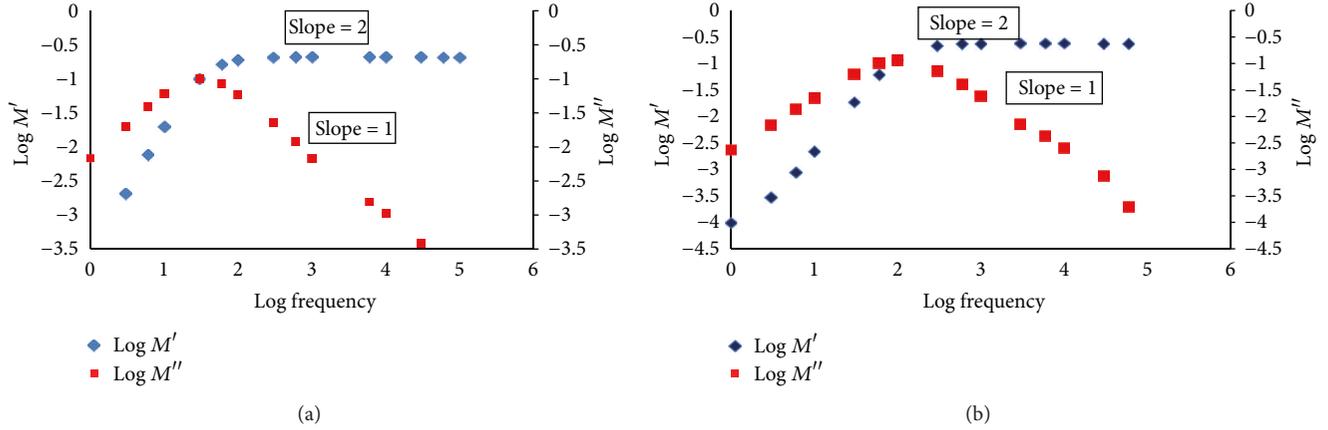


FIGURE 10: (a) Dependence of the real (M') and imaginary (M'') parts of the electric modulus on frequency in the region of the conductivity for PC and (b) BEDO-TTF/PC films, respectively.

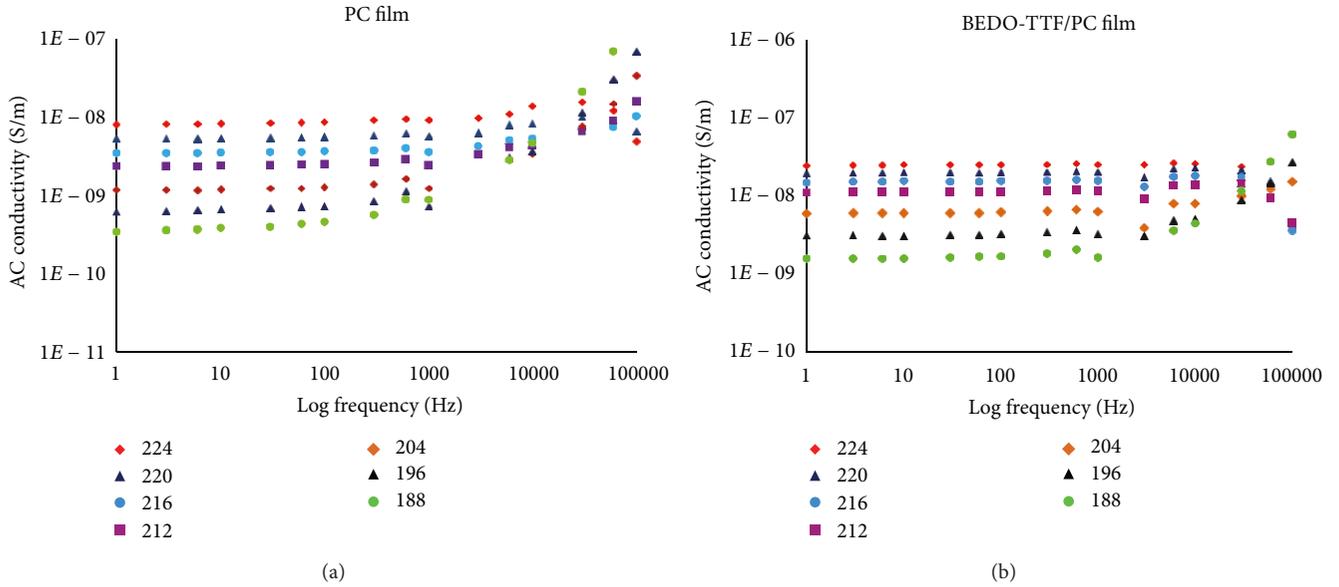


FIGURE 11: Frequency dependence of AC conductivity σ_{AC} (S/m) for (a) PC and (b) BEDO-TTF/PC film (2 wt%), respectively, above T_g temperature.

T_0 , C_1 , C_2 , and T_0 are WLF constants that are determined by curve fitting the data to the WLF equation using Origin software as shown in Figures 8(a) and 8(b). The values of C_1 , C_2 , and T_0 were found to be 9.8, 42.7 K, and 427.9 K within a temperature range of 148–174°C for PC and 10.1, 43.8 K, and 420.6 K within a temperature range of 140–167°C for BEDO-TTF/PC, respectively. The WLF constants C_1 and C_2 reveal information related to fractional free volume (f_g) and thermal expansion (α_f) parameters through.

$$f_g = \frac{B}{(2.303 * C_1)}, \quad (4)$$

$$\alpha_f = \frac{f_g}{C_2},$$

where f_g is an unoccupied (free) volume in the structure and B is taken as unity according to Doolittle equation. The f_g for PC and BEDO-TTF/PC are 0.044 and 0.042, respectively, whereas the thermal expansion for PC and BEDO-TTF/PC is 1.03×10^{-3} and 0.958×10^{-3} , respectively. The apparent activation energy (ΔH) values for PC and BEDO-TTF/PC were calculated accordingly to the Catsiff and Tobolsky equation [71]:

$$\Delta H = 2.303 \left(\frac{C_1}{C_2} \right) RT_g^2. \quad (5)$$

The values of activation energies for PC and BEDO-TTF/PC films were found to be 813 kJ/mol and 785 kJ/mol, respectively. These values for PC are in the range of values of 480 and 835 kJ/mol as reported in the literature [45, 55, 72].

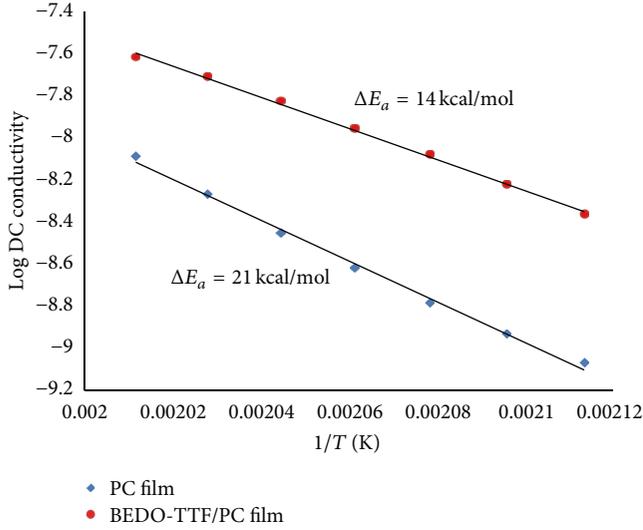


FIGURE 12: Log σ_{DC} versus inverse temperature plots of PC (red) and BEDO-TTF/PC film (blue).

WLF behavior for PC had been previously reported by others [39, 73–78].

Work herein demonstrates that BEDO-TTF plasticizes the PC matrix and reduces the glass transition temperature. Main chain motion occurs at lower temperatures in the BEDO-TTF plasticized film.

(3) *Ionic Conductivity Relaxation.* The high temperature region of the relaxation spectra was analyzed to confirm that the relaxation behavior was due to conductivity and not obstructed by any viscoelastic effects. Three proofs for this statement were undertaken, and these are described in the “the appendix” [28, 35]. This part of the research demonstrated that conductivity was enhanced in the PC containing BEDO-TTF, and this is likely a result of plasticization of the matrix.

4. Conclusion

This study first summarizes some of the excellent background work done by researchers in this field. The initial work reported herein reveals that films processed via the reticulate doped polymer method can be reduced by cyclic voltammetry to yield conductive, colorless, transparent films. These features make it possible to design optimum sensors. DSC studies of PC and BEDO-TTF/PC films reveal that the dye is soluble in and plasticized the PC matrix. Dielectric measurements carried out on PC and BEDO-TTF/PC confirm multiple transitions within the temperature range of -120 to 225°C for the frequency range 1–100 kHz. Different parameters such as relaxation broadening, activation energy, and conductivity (AC or DC) were found to be composition dependent. The AC conductivity becomes higher with the addition of BEDO-TTF. This is most likely due to enhanced mobility due to the plasticizing effect of the dye.

Appendix

Proof (Argand plot). The dielectric permittivity and loss factor for a relaxation with a single relaxation can be described by (A.1)

$$\epsilon' = \epsilon_U + \frac{(\epsilon_R - \epsilon_U)}{1 + \omega^2 \tau_E^2}, \quad (\text{A.1})$$

$$\epsilon'' = (\epsilon_R - \epsilon_U) \frac{\omega \tau_E}{1 + \omega^2 \tau_E^2},$$

where τ_E is the dielectric relaxation time and ω is the angular frequency. ϵ_U represents the high frequency, unrelaxed state, while ϵ_R represents the low frequency, relaxed state. By arranging (A.1), (A.2) is derived as

$$\left\{ \epsilon' - \frac{(\epsilon_R - \epsilon_U)}{2} \right\}^2 + (\epsilon'')^2 = \left(\frac{\epsilon_R - \epsilon_U}{2} \right)^2. \quad (\text{A.2})$$

Cole-Cole proposed that by plotting dielectric loss (ϵ'') against permittivity (ϵ') at a particular temperature, and a semicircle of radius $((\epsilon_R - \epsilon_U)/2)$ is obtained [38]. In Cole-Cole plots the high frequency region is on left side, while the low frequency exists on the right side of the plot. The Argand plots identify viscoelastic and conductivity effects [38, 79–81]. Argand plots were plotted between imaginary part (M'') of the complex modulus against real part (M') of the complex modulus as functions of the frequency at fixed temperatures. In an Argand plot, low frequency measurements exist on the left side of the semicircle, while the high frequency is on the right side. Argand plots can be obtained from

$$\left(M' - \left\{ \frac{M_U - M_R}{2} \right\} \right)^2 + (M'')^2 = \left(\frac{M_U - M_R}{2} \right)^2. \quad (\text{A.3})$$

Semicircular behavior is characteristic of Debye behavior for small rigid molecules and molecular liquids. This indicates a single relaxation time due to ionic conduction in the absence of any viscoelastic relaxation behavior [43]. By using the electric modulus, the space charge effects are reduced, and ionic conductivity peaks appear [82, 83]. Figure 9 shows semicircular behavior at 224°C . \square

Proof ($\log M''$, M' versus frequency). Ambrus et al. derived the electric modulus in terms of time, frequency, and modulus (see (A.4) and (A.5)) [84]. Starkweather et al. show that plots of $\log M''$ and $\log M'$ versus \log frequency at low frequencies will yield slopes of 1 and 2, respectively, when the electric modulus arises purely from ion conduction without contributions from viscoelastic relaxations [44]. Above T_g , the conduction is pure due to the diffusion of ions and independent of viscoelastic relaxation [14, 42, 81] as follows:

$$\begin{aligned} M &= M_s \left(\frac{i\omega\tau_\sigma}{1 + i\omega\tau_\sigma} \right) \\ &= M_s \left[\frac{(\omega\tau_\sigma)^2}{1 + (\omega\tau_\sigma)^2} \right] + iM_s \left[\frac{\omega\tau_\sigma}{1 + (\omega\tau_\sigma)^2} \right], \end{aligned} \quad (\text{A.4})$$

where

$$M_s = \frac{1}{\varepsilon_s}. \quad (\text{A.5})$$

Figure 10 shows that above T_g, PC, and BEDO-TTF/PC films samples approach ideal value of 2 or 1, respectively, confirming ionic conductivity. □

Proof (AC and DC conductivities). When viscoelastic effects are negligible, the loss factor is described by (A.6). Upon rearrangement, AC conductivity can be obtained by (A.7)

$$\varepsilon'' = \frac{\sigma_{AC}}{\omega\varepsilon_0}, \quad (\text{A.6})$$

where σ is the ionic conductivity, ω is the angular frequency ($2\pi f$), and ε_0 is the absolute permittivity of free space (8.854×10^{-14}) as follows:

$$\sigma_{AC} = \varepsilon'' \omega\varepsilon_0. \quad (\text{A.7})$$

As shown in Figures 11(a) and 11(b), the loss factor increases in the conductivity region at high temperatures and low frequencies.

Figure 11 shows frequency dependency of σ_{ac} at temperatures where conductivity effects dominate viscoelastic behavior. As temperature increases, σ_{ac} conductivity develops a plateau (ca.184°C) from 1 Hz to 10^2 Hz, signifying the beginning of the conductivity relaxation region. The σ_{ac} plateau then expands to higher frequencies 10^4 Hz as the temperature is increased, thus showing a frequency independent conductivity relaxation region. The absence of a frequency dependent region signifies negligible viscoelastic effects. This study reveals that the conductivity relaxation region exists between 180 and 224°C. At a lower temperature, where dielectric relaxation is dominant, the apparent conductivities (AC and DC) are strongly dependent on frequency. As temperature increases, conductivity becomes almost independent of the frequency. The BEDO-TTF/PC film has higher σ_{ac} values than that of the PC. The addition of BEDO-TTF presumably increases the amorphous content of the polymer. It speeds up segmental motion by increasing available free volume, thereby facilitating easy ion migration.

The DC conductivity values (σ_{DC}) have been obtained through the AC conductivity measurement [85–87] as shown by

$$\sigma_{AC}(\omega) = \omega\varepsilon_0\varepsilon''(\omega) = \sigma_{DC} + A\omega^s, \quad (\text{A.8})$$

where ω is the angular frequency, ε_0 is the vacuum permittivity, ε'' is the imaginary part of the complex permittivity, A is the temperature dependent parameter, and s is the frequency dependent exponent of the “universal” power law. The exponent s lies between $0 \leq s \leq 1$ [81, 83]. Frequency independent AC conductivity has been observed at higher temperatures, which signifies the long-range movement of mobile charge carriers. At higher temperatures, the graph is seen as linear on a logarithmic scale. The plateau values give DC conductivity. The Arrhenius relationship is expressed by

$$\log \sigma_{DC} = \log \sigma_0 \exp\left(\frac{-E_A}{kT}\right), \quad (\text{A.9})$$

where E_A is the apparent activation energy, σ_{DC} is the DC conductivity, σ_0 is the preexponential factor (conductivity at infinite temperature), and k is the Boltzmann constant.

Plots of \ln DC conductivity versus $1/\text{temperature}$ yield ionic conductivity and activation energies from the slopes. At higher temperatures, this plot exhibits a linear response. Figure 12 shows higher activation energy for PC (21 kcal/mole) due to lack of plasticization as compared to the plasticized BEDO-TTF/PC film which has activation energy of 14 kcal/mole. □

Conflict of Interests

None of the authors have conflict of interests with any of the aspects of this paper.

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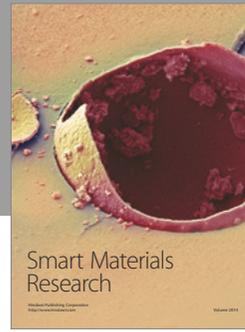
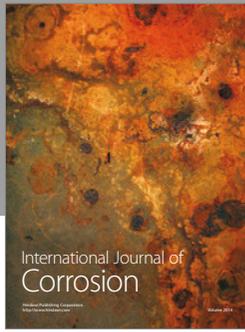
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