Influence of Rotation on Heat Transfer from a Two-Pass Channel with Periodically Placed Turbulence and Secondary Flow Promoters

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With the increase of operating costs and sensitivity to environmental pollution, there is a growing necessity to increase the thermal efficiency of turbines. One way to do this is to increase the inlet gas temperature. Therefore, turbine blades that can withstand higher thermal loads are in greater demand. Standard metal blades with special cooling techniques are employed in turbines since no other material now available can simultaneously withstand such high levels of heat load and thermomechanical stress. Turbine blades require sophisticated external and internal cooling techniques to sustain the stress-handling capabilities of the blades in a very hot operating environment.

Figure 1 shows the sectional view of a typical turbine blade. A two-pass rotating channel simulates the conditions of internal serpentine cooling passages of turbine blades. Earlier studies on cooling passages were based on stationary models, which were easier to construct and work on. It was observed that surface heat transfer in stationary ducts can be augmented by casting rib turbulators (see Han [1984]; Han et al. [1985, 1991]). Unlike a stationary duct flow situation, rotation introduces Coriolis and centrifugal buoyancy effects in the flow resulting in asymmetric velocity and turbulence intensity profiles (Johnston et al. [1972]). Therefore, to understand and improve internal cooling of turbine blades it is necessary to investigate the heat transfer patterns in an actual rotating frame.

There are fewer experimental studies on rotating channels than on stationary ducts with similar geometrical features. Han and Zhang [1992] covered most of the important related publications on rotating channels to that date. Therefore, a concise discussion on the previous works is given here. Initial studies were conducted on the effects of rotation on heat transfer characteristics in smooth channels with radial outward flow by Mori et al. [1971], Clifford et al. [1984], Harasgama and Morris [1988], Guidez [1989], Morris and Ghavami-Nasr [1991], and Soong et al. [1991]. After the initial investigations, more complicated test cases were reported with roughened heat transfer surfaces and with multipass serpentine...
Impingement cooling

Turbulence promoters

Film cooling

Suction surface

Shaped internal channel

Pressure surface

Trailing edge ejection

FIGURE 1  Schematic of the turbine blade coolant passages (from Parsons et al. [1994]).

Han and coworkers have been engaged in ribbed-duct heat transfer experiments for more than a decade. After successfully completing a series of rib configurations on a stationary duct, the emphasis is now shifted to a rotating channel with a two-pass flow passage. Han et al. [1994] reported the effects of different surface heating conditions (uniform wall temperature and uniform wall heat flux) in rotating flow with smooth walls. The first-pass trailing and second-pass leading surfaces showed enhancement in Nusselt numbers with rotation. Unlike the first-pass trailing and the second-pass leading surfaces, the first-pass leading and the second-pass trailing surfaces showed a decrease in Nusselt numbers with an increase in the rotation number. The reduction in Nusselt numbers on the first-pass leading and the second-pass trailing surfaces make them the critical surfaces for possible burn out. Ribs were introduced in the leading and trailing sides of the two-pass channel to increase heat transfer. Parsons et al. [1994] and Zhang et al. [1993b] reported the results. It was seen that heat transfer increases with the addition of ribs, but the increase from the smooth wall condition differs in the leading and trailing sides of the channel. The rotation number was the key influence in developing the difference between leading and trailing side Nusselt numbers.

In one stationary test case, Han et al. [1991] studied the heat transfer augmentation in stationary square channel with parallel, crossed, and V-shaped ribs. It was shown that V-shaped ribs can increase the heat transfer rate threefold over the heat transfer rate from the smooth wall case. The heat transfer for V-shaped and 60° parallel ribs were better than 90° transverse ribs for a wide range of Reynolds numbers. This paper tests the effect of rotation on different rib configurations that, in a stationary frame, showed significant improvements over standard 90° rib (Han et al. [1991]).

The present experimental work uses two types of rib configurations: 60° parallel rib and 60° staggered rib. These rib configurations are of special interest because ribs placed at an angle to the coolant flow induce secondary flow patterns that interact with the secondary flow generated by rotation. The performance of 60° staggered ribs in a rotating channel are compared with the standard 90° transverse and parallel ribs placed at a 60° angle to the flow. Figure 2 shows the different rib configurations used and referred to in this experimental work. The smooth duct results of Han et al. [1993], the 90° transverse rib of Parsons et al. [1994], and the 60° parallel rib of Zhang et al. [1993b] are included to compare the performance of the 60° parallel rib and the 60° staggered rib of the present work. Note that the layout of the 60° parallel rib of this work and the 60° parallel rib of Zhang et al. [1993b] are different (Figures 2c and 2d). The ribs at the leading and trailing walls are aligned to be on top of each other.

Figure 3 shows the schematic of the vortex structures due to skewed rib configuration (no rotation) and rotation (smooth wall). From the rib orientations it is postulated that the 60° parallel rib would induce a two-vortex structure and the 60° staggered rib would generate a four-vortex structure. The staggered orientation would cause the vortex generated by one pair of leading–trailing 60° staggered ribs to dominate the upstream pair’s vortices. For clarity, Figure 3a shows the stronger vortices in solid lines and the weaker vortices in broken lines. The vortex structure for rotation in a smooth channel is taken from Dutta et al. [1994]. In their numerical work, Dutta et al. [1994] observed two types of vortices for radially outward flow (Figure 3b). The second-pass (radially inward flow) vortices for rotation are extrapolated from the first-pass (radially outward flow) simulations. Figure 3 shows that the vortex structures of ribs and rotation are not aligned and would interact to give new complicated vortex structures. The new structure would be dependent on the relative strengths of the interacting vortices and would differ with Reynolds number and rotation number.

Local heat transfer results from the leading and trailing sides of the channel are presented in this paper. The surface averaged Nusselt number ratios are also included. The duct is rotated at an angular speed of 800 rpm and the Reynolds number of the flow altered to get four different rotation numbers varying from 0.0352 to 0.352. The four walls of the square channel are maintained at uniform temperatures at all axial locations.

**EXPERIMENTAL SET-UP**

Han et al. [1993] provide a detailed illustrated description of the experimental set-up. Therefore, only a brief description of the experimental facility is given here. Figure 4 shows the schematic of the test rig. Regulated compressed air is used as the working fluid. An orifice meter measures the mean flow rate passing through the test section, which rotates in a horizontal plane. There are two sets of slip ring units. One maintains a connection be-
(a) Smooth - no rib
   (Han et al. [1993])

(b) 90° Transverse rib
   (Parsons et al. [1994])

(c) 60° Parallel rib
   (Zhang et al. [1993b])

(d) 60° Parallel rib

(e) 60° Staggered rib

FIGURE 2 Rib configurations used and compared with.
(a) Vortex structures due to the rib orientations at no rotation.

(b) Vortex structures due to rotation in a smooth duct.

FIGURE 3 Schematic of the secondary flow vortex structures generated by rib orientations and by rotation.
tween thermocouples and the data logger interfaced to a personal computer. The other connects the power transformers to the heater units built in to the test section. A digital photo tachometer measures the rotation speed.

Each pass of the two-pass duct is 12.7 mm by 12.7 mm in cross-section. The flow in the first pass is radially outward, while the flow in the second pass is radially inward. Figure 5 shows a cut-away view of the two-pass test section with 60° staggered ribs. To get local heat transfer coefficients, the duct is divided into twelve sections (Figure 5) with six in the first pass and six in the second pass. The heater sections are two-hydraulic-diameters-long copper plates insulated from each other. Each of the nine walls (eight walls of the two passes and the turn) is insulated from each other and have independent heater coils powered by VARIAC transformers. The power input and thus the temperature of any surface can be adjusted to a desired level. The power to the heaters is adjusted so that the wall temperatures remain close to 40°C above ambient. Typical wall and bulk temperature profiles were given in Han et al. [1993].

The channel length to hydraulic diameter ratio is 24, while each pass is 12 hydraulic diameters long connected by a sharp 180° turn. The mean rotating arm length is 30 hydraulic diameters. The ribs are placed on the leading and trailing sides of the heated duct. Brass ribs of square cross-section are glued to the heater surfaces. The glue thickness is kept to less than 0.3 mm. The rib height to hydraulic diameter ratio is 0.125 and the pitch to height ratio is 10. The angle between the coolant flow direction and the rib axis is 60°. The ribs on both the leading and trailing surfaces are exactly on top of each other and not staggered. Figures 2d and 2e show the two rib configurations used in this paper. The arrows show bulk flow directions.

The experiment uses four Reynolds numbers: 2,500, 5,000, 10,000, and 25,000. Variation of the mean flow velocity through the test section changes the Reynolds number. The duct, rotates at 800 rpm, resulting in rotation numbers of 0.352, 0.176, 0.088, and 0.0352.

DATA REDUCTION

The local heat transfer coefficient is calculated from the local net heat flux, the local wall temperature, and the local bulk mean air temperature difference as

\[ h = q_{\text{net}} / (T_w - T_b) \]  

Local net heat flux \( q_{\text{net}} \) is calculated by subtracting heat loss from the power supplied to the heaters. Heat loss is estimated by running a separate test on the same test configuration but without air flow in the duct. Characteristic heat loss lines are obtained for each of the nine heater surfaces. For simplicity, all the heat flux calculations employ a smooth heater area with no rib surfaces included in the calculations. The local wall temperature \( T_w \) is obtained from the thermocouple reading of that heater section. The local bulk temperature \( T_b \) uses a linear interpolation between measured inlet and outlet bulk temperatures.

To reduce the influence of the flow Reynolds number on the heat transfer coefficient, the local Nusselt number is normalized by the Nusselt number for fully developed turbulent flow in smooth stationary circular pipe correlated by Dittus-Boelter/McAdams (Rohsenow and Choi [1961]) as

\[ Nu / Nu_0 = \left( \frac{hD}{k} \right) / (0.023 \Re^{0.8} \Pr^{0.4}) \]  

Prandtl number \( \Pr \) for air is 0.72. Air properties are based on the average of the inlet and outlet bulk mean air temperatures.

The uncertainty of the local heat transfer coefficient depends on the uncertainties in the local wall and bulk air temperature difference, and net heat input for each heater section. This uncertainty increases for decreasing both the local wall to air temperature difference and the net heat flux. Based on the method described by Kline and McClintock [1953], The typical uncertainty in the Nusselt number is estimated to be less than 8 percent for Reynolds numbers larger than 10,000. However the maximum uncertainty could be 25 percent for the lower heat transfer coefficient at the lowest Reynolds number tested (Re = 2,500).

RESULTS AND DISCUSSION

Figures 6 and 7 show the local distribution of Nusselt number ratios at leading and trailing surfaces for rotating and nonrotating situations. Each data point presented is the averaged result over a heater section of two hydraulic diameters \( D \) long. In a stationary duct both ribbed sides should have identical Nusselt number ratio profiles. Figures 6 and 7 show that there are differences in the local Nusselt number ratios between two ribbed sides especially in the first pass for both rib configurations. However, note that the profiles of the Nusselt number ratios are similar for the leading and trailing sides at nonrotating condition.

Figure 6 presents results for \( \Re = 5,000 \). The stationary duct results in Figure 6 show the peak Nusselt numbers, for both rib configurations, occur at a downstream location of the inlet rather than at the entrance region of the heated duct. Earlier study of 60° parallel rib in a stationary frame
by Han et al. [1991] also showed an increase in Nusselt number downstream of the entrance. However, the experimental set-up of Han et al. [1991] had a larger hydraulic diameter (2 inches) and the rib height to hydraulic diameter ratio was smaller than the ratio used in this work. In the first pass, the peak Nusselt numbers for both rib configurations, occur near $x/D = 5$ and is believed to be due to the generation of secondary flows induced by the rib orientations. As the flow develops through the duct, the Nusselt number ratio decreases followed by an increase near the exit of the first pass. Nusselt number ratios increase at the exit of the first pass, $x/D = 11$, because of the 180° bend. The inlet of the second pass starts from the sharp bend and the Nusselt number profiles differ from those of the first pass. The stationary duct staggered rib does better than the other rib configuration in the second pass. Nusselt number ratios initially drop after the bend and then increase followed by a decrease (may be toward an equilibrium value). The initial drop is due to the decay of turbulence produced by the sharp 180° bend. The rib
FIGURE 5  Cut-away view of the test section with 60° staggered ribs.
Figure 6 shows that the difference of Nusselt number ratios between leading and trailing sides increase with rotation, although the distribution patterns remain similar to those about the stationary duct. The drop in first-pass leading and second-pass trailing Nusselt numbers are more than in the other two-ribbed surfaces. Johnson et al. (1994) observed similar behavior with 45° skewed ribs. In an earlier study, Rohte and Johnston (1979) experimentally demonstrated that the turbulent shear layer past a backstep for low blockage ratio is more affected under stabilizing rotation (leading side of radially outward flow) than destabilizing rotation (trailing side of radially outward flow). The changes in the first-pass leading and trailing sides are similar to a smooth channel (Han et al. 1993). The leading-side Nusselt number ratio decreases and trailing-side Nusselt number ratio increases with rotation. However, the second-pass results show a reduction in Nusselt number ratio in both sides with rotation.
second-pass leading side Nusselt number ratios are higher than the second-pass trailing side Nusselt number ratios. The Coriolis force in rotation sets up secondary flow in the rotating duct that causes an increase in Nusselt number for the first-pass trailing and the second-pass leading surfaces, and a decrease in Nusselt number for the first-pass leading and second-pass trailing surfaces.

Beside the Coriolis force in heated rotating ducts, the centrifugal force induces buoyancy effects that can play a major role in the heat transfer patterns. In the first pass, the velocity and turbulence at the inlet of the test section are higher near the trailing side than the leading side due to the Coriolis effect in the unheated starting length. Air at higher velocity regions remains cooler and denser than the lower velocity region. Therefore, air near the trailing side gets accelerated by a higher centrifugal buoyancy force and air near the leading side decelerates to maintain the same mass flow rate through the duct. Therefore, both the rotational buoyancy force and Coriolis force in the first pass increase the heat transfer from the trailing side and decrease heat transfer from the leading side. Unlike the first pass, the second pass has radially inward flow. The centrifugal buoyancy in the second pass accelerates the warmer low-velocity air and decelerates the denser high-velocity air. It can be observed from the Nusselt number profiles of Figure 6, that the effects of rib orientations and

FIGURE 7 Local Nusselt number distribution for Re = 25,000.
Coriolis force in the second pass are clearly present. This is because the Nusselt number distribution profiles with rotation are similar to the nonrotating condition and the trailing-side Nusselt numbers are lower than the leading-side Nusselt numbers. However, the decrease in Nusselt number with rotation for the ribbed sides must be significantly influenced by the centrifugal buoyancy force. It can be argued that the rotational buoyancy accelerates the warmer fluid near the two smooth walls and, as a result, there is a reduction in velocity and heat transfer near and from the ribbed sides.

Figure 7 shows the results of Re = 25,000. The Nusselt number ratio distributions are similar to those for Re = 5,000 (Figure 6). The main difference observed with rotation is a significant increase in Nusselt number ratio in the second pass and in most of the first pass for staggered ribs. The effect of rotation for Re = 25,000 where the rotation number is less (Ro = 0.0352) compared to Re = 5,000 (Ro = 0.176) is expected to be smaller. A possible reason for the enhancement of Nusselt number with rotation for staggered ribs can be a result of the interaction among the different secondary flow structures. Dutta et al. [1994] numerically showed that the vortex structure formed by rotation in a square channel at different rotational situations, can vary between four and two vortex structures. Figure 3 showed the schematic vortex structures generated by rib configurations and rotation. It is possible that the rotation induces a four-vortex structure for staggered ribs at Ro = 0.0352. The four-vortex structure of the staggered rib (no-rotation) increases the core flow in the channel. The four-vortex structure of rotation opposes and reduces the vortex effect of the staggered rib. Therefore, it can be argued that as more air from the core gets redistributed toward the walls of the inter-rib region, the Nusselt number increases with rotation from a no-rotation situation.

Figure 8 shows the differences in Nusselt number ratios between the trailing and the leading sides. Smooth duct data (Han et al. [1993]) are included in this figure for comparison. The ribbed ducts show larger differences than smooth duct at high rotation number (Ro = 0.352). It should be noted that although the differences are larger in ribbed ducts, the ratios of the Nusselt numbers for the two sides (\(\frac{Nu_{\text{trailing}}}{Nu_{\text{leading}}}\)) are less for ribbed ducts than the smooth rotating channel. The difference between trailing- and leading-side Nusselt number ratios gradually approaches zero with a reduction in rotation number. The smooth duct data show the effects of vortical structure and difference of turbulence levels in leading and trailing sides from the rotational effects (Johnston et al. [1972]). Ribbed ducts also include the effects on separation process from higher and lower turbulence levels at pressure and suction sides, respectively. The influence of secondary flows generated by skewed rib configurations on flow separation are more complicated. Centrifugal buoyancy force in the first pass (radially outward flow), tries to widen the differences between trailing- and leading-side heat transfer rates. Centrifugal buoyancy in the second pass (radially inward flow) tries to equate the heat transfer from two sides. Figure 8 shows that the first-pass differences for smooth duct are always higher than second-pass differences. The other two rib configurations follow this trend.

Figures 9 and 10 show the effect of rotation number on heat transfer. Figure 9 also shows local effects under rotation, which include Coriolis effect, centrifugal buoyancy effect, effects from secondary flows generated from rib orientations, and the effect of 180° bend on Nusselt number ratio. Rather than show a uniform increase or decrease with rotation number, as observed in smooth rotating channels (Han et al. [1993]), there are variations of the Nusselt number ratios with the rotation numbers. Sides with stabilizing rotation (first-leading and second-trailing) generally show a reduction in Nusselt number ratio with an increase in rotation number. The other two ribbed sides show an opposite trend.

In Figure 10 the surface averaged Nusselt number ratios are plotted for four-rib bearing sides of the two-pass channel. Results from the nonrotating Re = 25,000 are plotted for the stationary case (Ro = 0). With an increase in rotation number the first-pass leading and second-pass trailing show mostly a decrease in Nusselt number ratio, whereas, the first-pass trailing and second-pass leading show mostly an increase in average Nusselt number ratio with rotation number. Earlier published results of 90° transverse (Parsons et al. [1994]) and 60° parallel ribs (Zhang et al. [1993b], 60° parallel ribs oriented differently in two passes) are included for comparison. The present 60° parallel rib does better in the suction sides of the two passes (first-leading and second-trailing) than the previous 60° parallel rib configuration, which shows higher Nusselt number ratios on the other two ribbed sides. The staggered rib shows a sharp increase in average Nusselt number ratios at low rotation number from the stationary duct results. As discussed earlier, this can be from a favorable secondary flow vortex structure induced by rotation. For 90° transverse and 60° parallel ribs the rate of increase of the average Nusselt number ratios in the second-pass leading is less than for the first-pass trailing. Whereas, for 60° staggered ribs the rate of increase of the average Nusselt number ratios with rotation is similar for the first-pass trailing and the second-pass leading surfaces. Compared to the skewed ribs, the performance of the 90° rib is worse on the first-pass leading and the second-pass trailing. The first-pass leading and the second-pass trail-
parallel rib, (C) Staggered rib
Smooth surface (Han et al. [1993])

\[
\frac{(\text{Nu}_D - \text{Nu}_{D0})}{\text{Nu}_0}
\]

**Figure 8** Differences between local Nusselt number ratios at different rotation numbers.

The trailing-side Nusselt number ratio in the first pass \((x/D = 5 \text{ and } 9)\) increases with \(Ro\) and the leading-side Nusselt number ratio decreases with \(Ro\) at higher rotation numbers. Though the trailing-side Nusselt number ratio at \(x/D = 9\) for the parallel rib is higher than the staggered rib at a higher rotation number, the leading-side Nusselt number ratios for the staggered rib are higher than the parallel rib results. The surface averaged Nusselt number of Figure 10 showed the leading side of the first pass has higher Nusselt number ratios with staggered ribs. The differences between the leading- and trailing-side Nusselt numbers in Figure 8 are less for the
FIGURE 9 Effect of rotation number on Nusselt number ratios for the two rib configurations at selected axial locations.
staggered rib in most situations in the first pass than the 60° parallel rib. Therefore, a staggered rib is more suitable in the first pass than a 60° parallel rib.

The second-pass ($x/D = 15$ and 19) Nusselt number ratios of Figure 9 showed that 60° parallel ribs are more advantageous than the staggered rib. The leading-side results for the staggered rib are better than the parallel rib, but the trailing-side Nusselt number ratios for the parallel rib are higher than the staggered rib results at higher rotation numbers. The surface averaged Nusselt number ratios of Figure 10 favors the parallel rib. The surface averaged Nusselt number ratios for the second-pass trailing are higher for the parallel rib than the staggered rib. The differences in Nusselt numbers between the leading and the trailing surfaces for the parallel rib is less than the smooth surface (Figure 8). Therefore, the 60° parallel rib is more suitable in the second pass than the staggered rib.

**CONCLUSION**

The present work focuses on an experimental study of heat transfer from rotating ribbed surfaces. Two types of
ribs (60° parallel and 60° staggered) are used, and the results are compared with previous results of 90° transverse ribs, 60° parallel ribs, and smooth wall results. The experimental work is carried out on a two-pass square duct that has both radially outward and inward flow situations. It is concluded that, the better heat transfer results with 60° skewed ribs over standard 90° transverse ribs in a rotating duct at moderate-to-high rotation numbers are as significant as a stationary duct at the first-pass leading and the second-pass trailing surfaces, which are the most important surfaces for heat transfer. Comparative studies of the two rib configurations (60° parallel and 60° staggered) suggest that the staggered rib is more suitable to augment heat transfer with rotation for the first pass and the 60° parallel rib for the second pass.

Acknowledgments

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Nomenclature

<table>
<thead>
<tr>
<th>Symbol</th>
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<tbody>
<tr>
<td>$D$</td>
<td>hydraulic diameter, 0.0127 m</td>
</tr>
<tr>
<td>$h$</td>
<td>heat transfer coefficient, w/m² °C</td>
</tr>
<tr>
<td>$\dot{q}_{\text{net}}$</td>
<td>net heat flux, w/m²</td>
</tr>
<tr>
<td>$R$</td>
<td>Reynolds number, $VD/\nu$</td>
</tr>
<tr>
<td>$\tau_w$</td>
<td>local wall temperature, °C</td>
</tr>
<tr>
<td>$V$</td>
<td>average flow velocity, m/s</td>
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<tr>
<td>$x$</td>
<td>distance measured along the duct axis from the start of heating, m</td>
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Greek Letters

<table>
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<th>Symbol</th>
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<tr>
<td>$\nu$</td>
<td>kinematic viscosity of air, m²/s</td>
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<tr>
<td>$\Omega$</td>
<td>rotational speed of the two-pass channel, rad/s (1 rpm = 0.105 rad/s)</td>
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References


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