Research Article

Periodic Solutions and Nontrivial Periodic Solutions for a Class of Rayleigh-Type Equation with Two Deviating Arguments

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The Rayleigh equation with two deviating arguments $x''(t) + f(x'(t)) + g_1(t, x(t - \tau_1(t))) + g_2(t, x(t - \tau_2(t))) = e(t)$ is studied. By using Leray-Schauder index theorem and Leray-Schauder fixed point theorem, we obtain some new results on the existence of periodic solutions, especially for the existence of nontrivial periodic solutions to this equation. The results are illustrated with two examples, which cannot be handled using the existing results.

1. Introduction

Consider the Rayleigh equation with two deviating arguments in the form of

$$x''(t) + f(x'(t)) + g_1(t, x(t - \tau_1(t))) + g_2(t, x(t - \tau_2(t))) = e(t), \quad (1)$$

where

$$\begin{aligned} f &\in C(R, R), \quad g_i \in C(R \times R, R), \quad i = 1, 2, \\ e, \tau_i &\in C(R, R), \quad i = 1, 2, \\ g_i(t + T, x) &= g_i(t, x), \quad \tau_i(t + T) = \tau_i(t), \quad i = 1, 2, \\ e(t + T) &= e(t). \end{aligned} \quad (2)$$

The dynamic behaviors of Rayleigh equation have been widely investigated due to their applications in many fields such as physics, mechanics, and the engineering technique fields. For example, an excess voltage of ferroresonance, known as some kind of nonlinear resonance having long duration that arises from the magnetic saturation of inductance in an oscillating circuit of a power system, and a boosted excess voltage can give rise to some problems in relay protection. To probe this mechanism, a mathematical model was proposed in [1–3], which is a special case of the Rayleigh

equation with two delays. This implies that (1) can represent analog voltage transmission. In a mechanical problem, f usually represents a damping or friction term, g_i represents the restoring force, e is an externally applied force, and τ_i is the time lag of the restoring force (see [4]). Some other examples in practical problems concerning physics and engineering technique fields can be found in [5–7].

At the same time, the periodic solutions for Rayleigh equations with two deviating arguments have been studied by authors [8–10] under the assumption of

$$f(0) = 0 \quad \text{or} \quad f(t, 0) = 0. \quad (3)$$

It is not difficult to see that if $g_1(t, 0) + g_2(t, 0) \neq e(t)$, then the periodic solution obtained in [8–10] must be nontrivial. But if $g_1(t, 0) + g_2(t, 0) \equiv e(t)$, then the periodic solution obtained in [8–10] may be trivial under the assumption of $f(0) = 0$ or $f(t, 0) = 0$. And if the periodic solution is unique, then it must be trivial. Thus, it is worth discussing the existence of the nontrivial periodic solutions of Rayleigh equations with two deviating arguments in this case.

The main purpose of this paper is to establish sufficient conditions for the existence of periodic solution, especially for the existence of nontrivial periodic solutions of (1) by using the Leray-Schauder index theory. We remark that our methods are different from those used in [8–10] to some degree. In particular, two examples are also given to illustrate the effectiveness of our results.

For ease of exposition, throughout this paper, we assume that $T > 0$.

On the other hand, the following assumptions are used in this paper.

(H₁) $f \in C(R, R)$, $g_i \in C(R^2, R)$, $\tau_i \in C(R, R)$, $i = 1, 2$, $e \in C(R, R)$, $g_i(t + T, x) = g_i(t, x)$, $\tau_i(t + T) = \tau_i(t)$, $i = 1, 2$, and $e(t + T) = e(t)$.

(H₂) $f(0) = 0$, and there exists $\gamma > 0$ such that $xf(x) \geq \gamma|x|^2$, for all $x \in R$ (or $xf(x) \leq -\gamma|x|^2$, for all $x \in R$).

(H₃) g_i is differentiable with respect to t , and there exist $a_i \geq 0$, $b_i > 0$, $i = 1, 2$, such that

$$|g'_{it}(t, x)| \leq a_i + b_i |x|, \quad \forall (t, x) \in R^2, \quad i = 1, 2. \quad (4)$$

(H₄) There exist $l_i > 0$ such that $|g_i(t, x_1) - g_i(t, x_2)| \leq l_i|x_1 - x_2|$, for all $t \in R$, $x_1, x_2 \in R$, $i = 1, 2$.

(H₅) There exists $d > 0$, such that for $|x| > d$, we have

$$(g_1(t, x) + g_2(t, x) - e(t)) \operatorname{sign} x > 0, \quad \forall t \in R. \quad (5)$$

(H₆) There exist integers n_i such that $\delta_i \triangleq \max_{t \in [0, T]} |\tau_i(t) - n_i T| \leq T$, $i = 1, 2$.

The main results of the present paper are as follows.

Theorem 1. *If (H₁)–(H₆) hold, and $(b_1 + b_2)\gamma^{-1}T^2 + 2\sqrt{2}(l_1\delta_1 + l_2\delta_2)\pi^2\gamma^{-1} < 2\pi^2$, then (1) has at least one periodic solution.*

The proof of Theorem 1 will be given in Section 3.

Remark 2. If $g_1(t, 0) + g_2(t, 0) \neq e(t)$, then the periodic solution obtained by Theorem 1 must be nontrivial. If $g_1(t, 0) + g_2(t, 0) \equiv e(t)$, we could not conclude whether or not the periodic solution is nontrivial.

So we give the following conditions:

(H₇) $g'_{ix}(t, x)$, $i = 1, 2$ exist and are continuous, $g'_{1x}(t, 0) + g'_{2x}(t, 0) \equiv \text{const.} < 0$.

(H₈) $\lim_{x \rightarrow 0} (f(x)/x) = r \geq \gamma$.

Theorem 3. *If (H₁)–(H₈) hold, and $g_1(t, 0) + g_2(t, 0) \equiv e(t)$, $(b_1 + b_2)\gamma^{-1}T^2 + 2\sqrt{2}(l_1\delta_1 + l_2\delta_2)\pi^2\gamma^{-1} < 2\pi^2$, then (1) has at least one nontrivial periodic solution.*

The proof of Theorem 3 will be given in Section 4.

2. Preliminaries

In this section, to establish the periodic solutions of (1), we provide some background definitions and some well-known results, which are crucial in our arguments.

Let X be a real Banach space, and let $A : X \rightarrow X$ be a completely continuous operator.

Definition 4 (see [11, 12]). If x_0 is an isolated fixed point of A , then the fixed point index at x_0 of A is defined by

$$i(A, x_0) = \deg(I - A, \Omega, \theta), \quad (6)$$

where Ω is a neighborhood of x_0 , which satisfies that x_0 is the unique fixed point in Ω of A .

Definition 5 (see [11, 12]). If there exists $x_0 \in X$ with $x_0 \neq \theta$ such that $Ax_0 = \lambda x_0$, $\lambda \in R$, then λ is called an eigenvalue of operator A and x_0 is called the eigenfunction of operator A corresponding to λ .

Definition 6. Let $u : R \rightarrow R$ be continuous. $u(t)$ is said to be periodic on R if

$$u(t + T) = u(t), \quad \forall t \in R. \quad (7)$$

Lemma 7 (see [11, 12], index theorem of Leray-Schauder). Suppose that x_0 is a fixed point of A , A is Fréchet differentiable at x_0 , and 1 is not the eigenvalue of $A'(x_0)$. Then x_0 is an isolated fixed point of A , and

$$i(A, x_0) = i(A'(x_0), 0) = (-1)^\beta, \quad (8)$$

where β is equal to the sum of the algebraic multiplicities of all of the eigenvalues $\lambda_i > 1$ of $A'(x_0)$.

Lemma 8 (see [11, 12], fixed point theorem of Leray-Schauder). Let X be a real Banach space and $A : X \rightarrow X$ be a completely continuous operator. If

$$\{x : x \in X, x = \lambda Ax, 0 < \lambda < 1\} \quad (9)$$

is bounded, then A has a fixed point $x^* \in \Omega$, where

$$\Omega = \{x : x \in X, \|x\| \leq l\}, \quad l = \sup \{x : x \in X, x = \lambda Ax, 0 < \lambda < 1\}. \quad (10)$$

Lemma 9. Suppose $x(t) \in C^1[0, T]$ and $x(0) = x(T) = 0$. Then $\int_0^T |x(t)|^2 dt \leq (T^2/\pi^2) \int_0^T |x'(t)|^2 dt$.

Lemma 10 (see [13]). Let $0 \leq \alpha \leq T$ be constant, $s \in C(R, R)$ be periodic with period T , and $\max_{t \in [0, T]} |s(t)| \leq \alpha$. Then for any $u \in C^1(R, R)$ which is periodic with period T , we have

$$\int_0^T |u(t) - u(t - s(t))|^2 dt \leq 2\alpha^2 \int_0^T |u'(t)|^2 dt. \quad (11)$$

3. Proof of Theorem 1

In this section, we will use Lemma 8 to prove Theorem 1.

Proof. Let

$$\begin{aligned} X &= \{x : x \in C(R, R), x(t + T) = x(t)\}, \\ Y &= \{x : x \in C^1(R, R), x(t + T) = x(t)\}. \end{aligned} \quad (12)$$

Then X and Y are real Banach spaces endowed with the norm

$$\begin{aligned} \|x\|_\infty &= \max_{t \in [0, T]} |x(t)|, \\ \|x\| &= \|x\|_\infty + \|x'\|_\infty, \end{aligned} \quad (13)$$

respectively.

Choosing $m > 0$ with $m \neq (2k\pi/T)^2$ ($k = 1, 2, \dots$), then

$$x''(t) + mx(t) = 0 \quad (14)$$

has only trivial solution in Y .

In fact, it is easy to see that the general solution of $x''(t) + mx(t) = 0$ is

$$x(t) = c_1 \sin \sqrt{mt} + c_2 \cos \sqrt{mt}. \quad (15)$$

By the periodic boundary conditions we obtain that $x = 0$ is its unique solution in Y . Then for $h \in X$,

$$-x''(t) - mx(t) = h(t) \quad (16)$$

has unique solution $x \in Y$. Writing $x = Kh$, then $K : X \rightarrow Y$ is a completely continuous operator.

Define an operator $G : Y \rightarrow X$ by

$$\begin{aligned} (Gx)(t) &= f(x'(t)) + g_1(t, x(t - \tau_1(t))) \\ &\quad + g_2(t, x(t - \tau_2(t))) - mx(t) - e(t), \quad x \in Y. \end{aligned} \quad (17)$$

Then $G : Y \rightarrow X$ is continuous and bounded. Let $A = KG : Y \rightarrow Y$. Then A is also a completely continuous operator. By Lemma 8, if

$$\{x : x \in Y, x = \lambda Ax, 0 < \lambda < 1\} \quad (18)$$

is bounded in Y , then A has a fixed point in Y . Thus, (1) has periodic solution.

Now suppose that $x \in Y$, $0 < \lambda < 1$ satisfying $x = \lambda Ax$. Then $x(t)$ is a solution of

$$\begin{aligned} x''(t) + \lambda f(x'(t)) + \lambda g_1(t, x(t - \tau_1(t))) \\ + \lambda g_2(t, x(t - \tau_2(t))) + (1 - \lambda)mx(t) \\ = \lambda e(t). \end{aligned} \quad (19)$$

Let $\underline{t}, \bar{t}$ be the maximum point and minimum point of $x(t)$ on $[0, T]$, respectively. Then

$$x'(\underline{t}) = 0, \quad x''(\underline{t}) \leq 0, \quad x'(\bar{t}) = 0, \quad x''(\bar{t}) \geq 0. \quad (20)$$

Noticing $f(0) = 0$, we have

$$\begin{aligned} \lambda g_1(\underline{t}, x(\underline{t} - \tau_1(\underline{t}))) + \lambda g_2(\underline{t}, x(\underline{t} - \tau_2(\underline{t}))) \\ + (1 - \lambda)mx(\underline{t}) - \lambda e(\underline{t}) \geq 0, \\ \lambda g_1(\bar{t}, x(\bar{t} - \tau_1(\bar{t}))) + \lambda g_2(\bar{t}, x(\bar{t} - \tau_2(\bar{t}))) \\ + (1 - \lambda)mx(\bar{t}) - \lambda e(\bar{t}) \leq 0, \end{aligned} \quad (21)$$

and hence, there exists $\xi \in [\underline{t}, \bar{t}]$ (or $[\bar{t}, \underline{t}]$) such that

$$\begin{aligned} \lambda g_1(\xi, x(\xi - \tau_1(\xi))) + \lambda g_2(\xi, x(\xi - \tau_2(\xi))) \\ - \lambda e(\xi) + (1 - \lambda)mx(\xi) = 0, \end{aligned} \quad (22)$$

which implies that

$$\begin{aligned} \lambda [g_1(\xi, x(\xi - \tau_1(\xi))) + g_2(\xi, x(\xi - \tau_2(\xi))) - e(\xi)] \\ \times \text{sign } x(\xi) + (1 - \lambda)m|x(\xi)| = 0. \end{aligned} \quad (23)$$

From (H_5) we know that $|x(\xi)| \leq d$ (see [10]). So we have

$$|x(t)| = \left| x(\xi) + \int_{\xi}^t x'(s) ds \right| \leq d + \sqrt{T} \|x'\|_{L^2}, \quad (24)$$

and then,

$$\|x\|_{\infty} \leq d + \sqrt{T} \|x'\|_{L^2}, \quad (25)$$

where $\|\cdot\|_{L^2}$ is the norm of $L^2[0, T]$.

Multiplying (19) with $x'(t)$ and integrating from 0 to T , then we have

$$\begin{aligned} \lambda \int_0^T f(x'(t)) x'(t) dt &= -\lambda \int_0^T g_1(t, x(t - \tau_1(t))) x'(t) dt \\ &\quad + \lambda \int_0^T g_2(t, x(t - \tau_2(t))) x'(t) dt \\ &\quad + \lambda \int_0^T e(t) x'(t) dt. \end{aligned} \quad (26)$$

By (H_2) we know that

$$\int_0^T f(x'(t)) x'(t) dt \geq \gamma \int_0^T |x'(t)|^2 dt. \quad (27)$$

By Hölder's inequality, from (26) and (27) we have

$$\begin{aligned} \gamma \int_0^T |x'|^2 dt &\leq \left| \int_0^T g_1(t, x(t - \tau_1(t))) x'(t) dt \right| \\ &\quad + \left| \int_0^T g_2(t, x(t - \tau_2(t))) x'(t) dt \right| + \|e\|_{L^2} \|x'\|_{L^2} \\ &\leq \int_0^T |g_1(t, x(t - \tau_1(t))) - g_1(t, x(t))| |x'(t)| dt \\ &\quad + \left| \int_0^T g_1(t, x(t)) x'(t) dt \right| \\ &\quad + \int_0^T |g_2(t, x(t - \tau_2(t))) - g_2(t, x(t))| |x'(t)| dt \\ &\quad + \left| \int_0^T g_2(t, x(t)) x'(t) dt \right| + \|e\|_{L^2} \|x'\|_{L^2}. \end{aligned} \quad (28)$$

Since $\int_0^{x(t)} g_i(t, v) dv$, $i = 1, 2$ are T -periodic, differentiable, and

$$\begin{aligned} \frac{d}{dt} \int_0^{x(t)} g_i(t, v) dv &= g_i(t, x(t)) x'(t) \\ &\quad + \int_0^{x(t)} g'_{it}(t, v) dv, \quad i = 1, 2, \end{aligned} \quad (29)$$

we have

$$\int_0^T g_i(t, x(t)) x'(t) dt = - \int_0^T dt \int_0^{x(t)} g'_{it}(t, v) dv, \quad (30)$$

$i = 1, 2.$

Combining (28) and (30) with (H_3) and (H_4) we obtain

$$\begin{aligned} & \gamma \int_0^T x'^2(t) dt \\ & \leq l_1 \int_0^T |x(t) - x(t - \tau_1(t))| |x'(t)| dt \\ & \quad + l_2 \int_0^T |x(t) - x(t - \tau_2(t))| |x'(t)| dt \\ & \quad + \int_0^T dt \int_0^{|x(t)|} (a_1 + b_1 |v|) dv \\ & \quad + \int_0^T dt \int_0^{|x(t)|} (a_2 + b_2 |v|) dv + \|e\|_{L^2} \|x'\|_{L^2} \quad (31) \\ & \leq l_1 \|x'\|_{L^2} \left(\int_0^T |x(t) - x(t - \tau_1(t) - n_1 T)|^2 dt \right)^{1/2} \\ & \quad + l_2 \|x'\|_{L^2} \left(\int_0^T |x(t) - x(t - \tau_2(t) - n_2 T)|^2 dt \right)^{1/2} \\ & \quad + (a_1 + a_2) \int_0^T |x(t)| dt + \frac{b_1 + b_2}{2} \int_0^T |x(t)|^2 dt \\ & \quad + \|e\|_{L^2} \|x'\|_{L^2}. \end{aligned}$$

By Lemma 9 we have

$$\begin{aligned} \int_0^T |x(t)|^2 dt &= \int_{\xi}^{\xi+T} |x(t)|^2 dt = \int_{\xi}^{\xi+T} |x(t) - x(\xi)|^2 dt \\ &\quad + 2x(\xi) \int_{\xi}^{\xi+T} x(t) dt - |x(\xi)|^2 T \\ &\leq \frac{T^2}{\pi^2} \int_{\xi}^{\xi+T} |x'(t)|^2 dt + 2d \int_0^T |x(t)| dt \quad (32) \\ &\quad - |x(\xi)|^2 T \\ &\leq \frac{T^2}{\pi^2} \int_0^T |x'(t)|^2 dt + 2d \int_0^T |x(t)| dt. \end{aligned}$$

By (32) and Hölder's inequality we have

$$\begin{aligned} & \int_0^T |x(t)| dt \\ & \leq \sqrt{T} \left(\int_0^T |x(t)|^2 dt \right)^{1/2} \\ & \leq \sqrt{T} \left(\frac{T^2}{\pi^2} \int_0^T |x'(t)|^2 dt + 2d \int_0^T |x(t)| dt \right)^{1/2} \quad (33) \\ & \leq \frac{T^{3/2}}{\pi} \left(\int_0^T |x'(t)|^2 dt \right)^{1/2} + \sqrt{2dT} \left(\int_0^T |x(t)| dt \right)^{1/2} \\ & \leq \frac{T^{3/2}}{\pi} \left(\int_0^T |x'(t)|^2 dt \right)^{1/2} + \frac{1}{4} \int_0^T |x(t)| dt + 2dT. \end{aligned}$$

From (33) we have

$$\int_0^T |x(t)| dt \leq \frac{4T^{3/2}}{3\pi} \left(\int_0^T |x'(t)|^2 dt \right)^{1/2} + \frac{8dT}{3}. \quad (34)$$

By (H_6) and Lemma 10 we have

$$\begin{aligned} & \left(\int_0^T |x(t) - x(t - \tau_i(t) - n_i T)|^2 dt \right)^{1/2} \\ & \leq \sqrt{2} \delta_i \left(\int_0^T |x'(t)|^2 dt \right)^{1/2}, \quad i = 1, 2. \end{aligned} \quad (35)$$

Thus, it follows from (31), (32), (34), and (35) that

$$\begin{aligned} \gamma \|x'\|_{L^2}^2 &\leq \left(\sqrt{2} l_1 \delta_1 + \sqrt{2} l_2 \delta_2 + \frac{(b_1 + b_2) T^2}{2\pi^2} \right) \|x'\|_{L^2}^2 \\ &\quad + \left[\frac{4T^{3/2}}{3\pi} [(a_1 + a_2) + (b_1 + b_2) d] + \|e\|_{L^2} \right] \|x'\|_{L^2} \\ &\quad + \frac{8dT}{3} [(a_1 + a_2) + (b_1 + b_2) d]. \end{aligned} \quad (36)$$

Combining this with $(b_1 + b_2) \gamma^{-1} T^2 + 2\sqrt{2}(l_1 \delta_1 + l_2 \delta_2) \pi^2 \gamma^{-1} < 2\pi^2$, we know that there exists c_1 such that $\|x'\|_{L^2} \leq c_1$. Then

$$\|x\|_{\infty} \leq d + \sqrt{T} c_1 \triangleq M_1. \quad (37)$$

Multiplying (19) with $x''(t)$ and integrating from 0 to T , then we have

$$\begin{aligned} & \|x''\|_{L^2}^2 \\ & \leq \left| -\lambda \int_0^T g_1(t, x(t - \tau_1(t))) x''(t) dt \right. \\ & \quad - \lambda \int_0^T g_2(t, x(t - \tau_2(t))) x''(t) dt \\ & \quad - (1 - \lambda) m \int_0^T x(t) x''(t) dt \\ & \quad \left. + \lambda \int_0^T e(t) x''(t) dt \right| \\ & \leq (g_{1M_1} + g_{2M_1}) \sqrt{T} \|x''\|_{L^2} + mM_1 \sqrt{T} \|x''\|_{L^2} \\ & \quad + \|e\|_{L^2} \|x''\|_{L^2}, \end{aligned} \quad (38)$$

where

$$\begin{aligned} g_{1M_1} &= \max_{t \in [0, T], \|x\|_{\infty} \leq M_1} |g_1(t, x(t))|, \\ g_{2M_1} &= \max_{t \in [0, T], \|x\|_{\infty} \leq M_1} |g_2(t, x(t))|. \end{aligned} \quad (39)$$

Thus,

$$\begin{aligned} \|x''\|_{L^2} &\leq (g_{1M_1} + g_{2M_1}) \sqrt{T} + mM_1 \sqrt{T} \\ &\quad + \|e\|_{L^2} \triangleq M_2. \end{aligned} \quad (40)$$

Selecting $\eta \in [0, T]$ such that $x'(\eta) = 0$, then we have

$$|x'(t)| \leq \int_0^T |x''(t)| dt \leq \sqrt{T} M_2. \quad (41)$$

Thus, from (35) and (41) we know that $\|x\| \leq M_1 + \sqrt{T} M_2 \triangleq M$. It follows that

$$\{x : x \in Y, x = \lambda Ax, 0 < \lambda < 1\} \quad (42)$$

is bounded. Therefore, by Lemma 8, we obtain that A has a fixed point $x^* \in \Omega$, where $\Omega = \{x : x \in Y, \|x\| \leq M\}$, and hence, it follows that (1) has a periodic solution. $\square$

4. Proof of Theorem 3

In this section, we will use Lemma 7 to prove Theorem 3.

Proof. Let X, Y, A be defined and same as in Section 3. We prove that A is Fréchet differentiable at 0 and

$$\begin{aligned} (A'(0)x)(t) &= K \left[rx'(t) + (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) - mx(t) \right]. \end{aligned} \quad (43)$$

In fact, by $g_1(t, 0) + g_2(t, 0) \equiv e(t)$, we have $A0 = 0$. Let

$$\begin{aligned} (Bx)(t) &= K \left[rx'(t) + (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) - mx(t) \right]. \end{aligned} \quad (44)$$

We have

$$\begin{aligned} \|Ax - Bx\| &= \|K \left[f(x'(t)) + g_1(t, x(t - \tau_1(t))) \right. \\ &\quad \left. + g_2(t, x(t - \tau_2(t))) - mx(t) - e(t) \right] \\ &\quad \left. - K \left[rx'(t) + (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) - mx(t) \right] \right\| \\ &\leq \|K\| \left\| \left[\frac{f(x'(t))}{x'(t)} - r \right] x'(t) + g_1(t, x(t - \tau_1(t))) \right. \\ &\quad \left. + g_2(t, x(t - \tau_2(t))) - e(t) \right. \\ &\quad \left. - (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) \right\| \\ &\leq \|K\| \left\| \left[\frac{f(x'(t))}{x'(t)} - r \right] + g'_{1x}(t, \theta_1 x(t - \tau_1(t))) \right. \\ &\quad \left. - g'_{1x}(t, 0) + g'_{2x}(t, \theta_2 x(t - \tau_2(t))) \right. \\ &\quad \left. - g'_{2x}(t, 0) \right\| \|x\|, \end{aligned} \quad (45)$$

where $\theta_i = \theta_i(t) \in (0, 1)$, $i = 1, 2$. Thus,

$$\lim_{\|x\| \rightarrow 0} \frac{\|Ax - A0 - Bx\|}{\|x\|} = 0, \quad (46)$$

which implies that $A'(0)$ exists and $A'(0) = B$.

Now, we prove that 1 is not the eigenvalue of $A'(0)$.

In fact, if there exists $x \in Y$ satisfying $A'(0)x = x$, then x is a solution of

$$\begin{aligned} -x''(t) - mx(t) &= rx'(t) + (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) \\ &\quad - mx(t). \end{aligned} \quad (47)$$

Integrating (47) from 0 to T , we obtain

$$\int_0^T (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) dt = 0. \quad (48)$$

By (H_7) , there exists $t_0 \in [0, T]$ such that $x(t_0) = 0$.

Multiplying (47) with $x'(t)$ and integrating it from 0 to T , we have

$$\int_0^T r|x'|^2 dt = - \int_0^T (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t)x'(t) dt. \quad (49)$$

By (H_7) we have

$$\int_0^T r|x'|^2 dt = 0, \quad (50)$$

which implies that $x'(t) \equiv 0$, and hence $x(t) \equiv 0$. It follows that 1 is not the eigenvalue of $A'(0)$. By Lemma 7 we have

$$i(A, 0) = i(A'(0), 0) = (-1)^\beta. \quad (51)$$

Next we prove that $\beta = 1$. We can prove that

$$\begin{aligned} -x''(t) - mx(t) &= rx'(t) + (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) - mx(t) \end{aligned} \quad (52)$$

has only one eigenvalue in $(0, 1)$ and the algebraic multiplicity of it is 1.

In fact, it is easy to see that $\lambda \in (0, 1)$ is the eigenvalue of (52) if and only if $m + \lambda(g'_{1x}(t, 0) + g'_{2x}(t, 0) - m) = 0$, that is,

$$\lambda = \lambda_0 = \frac{m}{m - g'_{1x}(t, 0) - g'_{2x}(t, 0)} \in (0, 1). \quad (53)$$

In this case, (52) degenerates to

$$x''(t) + \frac{m}{m - g'_{1x}(t, 0) - g'_{2x}(t, 0)}rx'(t) = 0. \quad (54)$$

It is not difficult to see that all the solutions of (54) are $x(t) \equiv \text{const}$. Therefore the geometric multiplicity of $\lambda_0 = m/(m - g'_{1x}(t, 0) - g'_{2x}(t, 0))$ is 1.

If $(I - \lambda_0 A'(0))^2 x(t) = 0$, then $(I - \lambda_0 A'(0))x(t) = c = \text{const}$, and hence, $x(t)$ satisfies

$$\begin{aligned} & -x''(t) - mx(t) \\ & = \lambda_0 (rx'(t) + (g'_{1x}(t, 0) + g'_{2x}(t, 0))x(t) - mx(t)) + mc, \end{aligned} \quad (55)$$

which implies that

$$-x''(t) = \lambda_0 rx'(t) + mc. \quad (56)$$

Integrating (56) from 0 to T , we obtain $c = 0$, which implies that $x(t) \equiv \text{const}$. It follows that the algebraic multiplicity of $\lambda_0 = m/(m - g'_{1x}(t, 0) - g'_{2x}(t, 0))$ is also 1, and therefore, $\beta = 1$. So $i(A, 0) = -1$. Letting $\delta > 0$ small enough, then we have

$$\deg(I - A, B_\delta, 0) = -1. \quad (57)$$

On the other hand, by the proof of Theorem 1, we know that there exists $M > 0$ such that if $x \in Y$, $0 < \lambda < 1$ satisfying $x = \lambda Ax$, then $\|x\| \leq M$. Suppose $Ax \neq x$ for $x \in Y$, $\|x\| = M$. By the homotopy invariance of topological degree, we have

$$\deg(I - A, B_M, 0) = 1, \quad (58)$$

where $B_M = \{x \in Y : \|x\| < M\}$. It follows from (57) and (58) that

$$\deg(I - A, B_M \setminus \bar{B}_\delta, 0) = 2, \quad (59)$$

which implies that A has a fixed point x in $B_M \setminus \bar{B}_\delta$ and x is a nontrivial periodic solution of (1). $\square$

5. Two Examples

In this section, we give two examples to demonstrate the results obtained in previous sections.

Example 1. Consider the forced Rayleigh-type equation with period 2π :

$$\begin{aligned} & x''(t) + f(x'(t)) + g_1(t, x(t - \tau_1(t))) \\ & + g_2(t, x(t - \tau_2(t))) = e(t), \end{aligned} \quad (60)$$

where

$$f(x) = \begin{cases} e^x - 1, & x \geq 0, \\ 1 - e^{-x}, & x \leq 0, \end{cases} \quad (61)$$

$$\begin{aligned} g_1(t, x) &= \frac{1}{9} \sin^2 t x(t - \theta \cos t) + \cos t, \\ g_2(t, x) &= \frac{1}{9} \cos^2 t x(t - \theta \sin t) + \sin t, \end{aligned} \quad (62)$$

$$\begin{aligned} e(t) &= \sin t, & \tau_1(t) &= \theta \cos t, \\ \tau_2(t) &= \theta \sin t, & \theta &\in (0, 1). \end{aligned}$$

Conclusion. Equation (60) has at least one periodic solution with period 2π .

Proof. By (61) and (62), it is not difficult to see that the condition (H_1) holds, $T = 2\pi$, and

$$\begin{aligned} & |f(0)| = 0, \\ & |g'_{1t}(t, x)| = \left| \frac{1}{9} x \sin 2t - \sin t \right| \\ & \leq \frac{1}{9} |x| + 1, \quad \forall (t, x) \in \mathbb{R}^2, \\ & |g'_{2t}(t, x)| = \left| -\frac{1}{9} x \sin 2t - \cos t \right| \\ & \leq \frac{1}{9} |x| + 1, \quad \forall (t, x) \in \mathbb{R}^2. \end{aligned} \quad (63)$$

On the other hand, let $d = 9$, $\gamma = 1$, $\delta_i = \theta$, $l_i = 1/9$, $b_i = 1/9$, and $i = 1, 2$. If $\theta \in (0, \sqrt{2}/2)$, then

$$\begin{aligned} & xf(x) \geq |x|^2, \quad \forall x \in \mathbb{R}, \\ & (b_1 + b_2) \gamma^{-1} T^2 + 2\sqrt{2} (l_1 \delta_1 + l_2 \delta_2) \pi^2 \gamma^{-1} < 2\pi^2. \end{aligned} \quad (64)$$

Hence, (H_1) – (H_6) are satisfied. Thus, by Theorem 1, (60) has at least one periodic solution with period 2π . $\square$

Example 2. If we replace $g_1(t, x)$, $g_2(t, x)$, and $e(t)$ in Example 1 by

$$\begin{aligned} g_1(t, x(t - \tau_1(t))) &= \frac{1}{9} \sin^2 t \cos \left(x(t - \tau_1(t)) + \frac{1}{2} \right) \\ &+ \frac{1}{10} x(t - \tau_1(t)) \sin \frac{1}{2}, \quad \forall (t, x) \in \mathbb{R}^2, \\ g_2(t, x(t - \tau_2(t))) &= \frac{1}{9} \cos^2 t \cos \left(x(t - \tau_2(t)) + \frac{1}{2} \right) \\ &+ \sin t, \quad \forall (t, x) \in \mathbb{R}^2, \end{aligned} \quad (65)$$

$e(t) = \sin t + (1/9) \cos(1/2)$, then we can obtain that $g_1(t, 0) + g_2(t, 0) \equiv e(t)$, for all $t \in \mathbb{R}$, and

$$\begin{aligned} g'_{1t}(t, x) &= \frac{1}{9} \sin 2t \cos \left(x + \frac{1}{2} \right), \\ g'_{2t}(t, x) &= -\frac{1}{9} \sin 2t \cos \left(x + \frac{1}{2} \right) + \cos t, \\ |g'_{1t}(t, x)| &\leq \frac{1}{9} |x| + \frac{1}{18}, \\ |g'_{2t}(t, x)| &\leq \frac{1}{9} |x| + \frac{37}{18}, \quad \forall (t, x) \in \mathbb{R}^2, \end{aligned} \quad (66)$$

$$g'_{1x}(t, x) = -\frac{1}{9} \sin^2 t \sin \left(x + \frac{1}{2} \right) + \frac{1}{10} \sin \frac{1}{2},$$

$$g'_{2x}(t, x) = -\frac{1}{9} \cos^2 t \sin \left(x + \frac{1}{2} \right), \quad \forall t \in \mathbb{R},$$

$$\begin{aligned} g'_{1x}(t, 0) + g'_{2x}(t, 0) &= -\frac{1}{9} \sin \frac{1}{2} + \frac{1}{10} \sin \frac{1}{2} < 0, \\ &\forall t \in \mathbb{R}. \end{aligned}$$

Let $d = (10/(9 \sin(1/2))) + ((10 \sin(1/2))/(9 \cos(1/2)))$, $\gamma = 1$, $\delta_i = \theta$, $i = 1, 2$, $l_1 = 19/90$, $l_2 = 1/9$, $b_1 = 1/9$, and $b_2 = 1/9$. If $\theta \in (0, 1)$; then

$$xf(x) \geq |x|^2, \quad \forall x \in R, \quad (67)$$

$$(b_1 + b_2) \gamma^{-1} T^2 + 2\sqrt{2} (l_1 \delta_1 + l_2 \delta_2) \pi^2 \gamma^{-1} < 2\pi^2.$$

Hence, $(H_1)-(H_8)$ are satisfied. Thus, by Theorem 3, (60) has at least one nontrivial periodic solution with period 2π .

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