Research Article

Energy Scattering for Schrödinger Equation with Exponential Nonlinearity in Two Dimensions

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When the spatial dimensions $n = 2$, the initial data $u_0 \in H^1$, and the Hamiltonian $H(u_0) \leq 1$, we prove that the scattering operator is well defined in the whole energy space $H^1(\mathbb{R}^2)$ for nonlinear Schrödinger equation with exponential nonlinearity $(e^{\lambda|u|^2} - 1)u$, where $0 < \lambda < 4\pi$.

1. Introduction

We consider the Cauchy problem for the following nonlinear Schrödinger equation:

$$iu_t + \Delta u = f(u), \quad (1)$$

$$f(u) := (e^{\lambda|u|^2} - 1)u, \quad (2)$$

in two spatial dimensions with initial data $u_0 \in H^1$ and $0 < \lambda < 4\pi$. Solutions of the above problem satisfy the conservation of mass and Hamiltonian:

$$M(u; t) := \int_{\mathbb{R}^2} |u|^2 dx = M(u_0), \quad (3)$$

$$H(u; t) := \int_{\mathbb{R}^2} (|\nabla u|^2 + F(u)) dx = H(u_0),$$

where

$$F(u) = \frac{1}{\lambda} (e^{\lambda|u|^2} - \lambda|u|^2 - 1). \quad (4)$$

Nakamura and Ozawa [1] showed the existence and uniqueness of the scattering operator of (1) with (2). Then, Wang [2] proved the smoothness of this scattering operator. However, both of these results are based on the assumption of small initial data u_0 . In this paper, we remove this assumption and show that for arbitrary initial data $u_0 \in H^1(\mathbb{R}^2)$ and $H(u_0) \leq 1$, the scattering operator is always well defined.

Wang et al. [3] proved the energy scattering theory of (1) with $f(u) = (e^{\lambda|u|^2} - 1 - \lambda|u|^2 - (\lambda^2/2)|u|^4)u$, where $\lambda \in \mathbb{R}$ and the spatial dimension $n = 1$. Ibrahim et al. [4] showed the existence and asymptotic completeness of the wave operators for (1) with $f(u) = (e^{\lambda|u|^2} - 1 - \lambda|u|^2)u$ when the spatial dimensions $n = 2$, $\lambda = 4\pi$, and $H(u_0) \leq 1$. Under the same assumptions as [4], Colliander et al. [5] proved the global well-posedness of (1) with (2).

Theorem 1. Assume that $u_0 \in H^1(\mathbb{R}^2)$, $H(u_0) \leq 1$, and $\lambda = 4\pi$. Then problem (1) with (2) has a unique global solution u in the class $C(\mathbb{R}, H^1(\mathbb{R}^2))$.

Remark 2. In fact, by the proof in [5], the global well-posedness of (1) with (2) is also true for $0 < \lambda \leq 4\pi$.

In this paper, we further study the scattering of this problem. Note that $f(u) = (e^{\lambda|u|^2} - 1)u = \sum_{k=1}^{\infty} (\lambda^k/k!)|u|^{2k}u$. Nakanishi [6] proved the existence of the scattering operators in the whole energy space $H^1(\mathbb{R}^2)$ for (1) with $f(u) = |u|^p u$ when $p > 2$. Then, Killip et al. [7] and Dodson [8] proved the existence of the scattering operators in $L^2(\mathbb{R}^2)$ for (1) with $f(u) = |u|^2 u$. Inspired by these two works, we use the concentration compactness method, which was introduced by Kenig and Merle in [9], to prove the existence of the scattering operators for (1) with (2).

For convenience, we write (1) and (2) together; that is,

$$iu_t + \Delta u = f(u) := (e^{\lambda|u|^2} - 1)u, \quad u(0, x) = u_0, \quad (5)$$

where $u_0 \in H^1(\mathbb{R}^2)$ and $0 < \lambda < 4\pi$. Our main result is as follows.

Theorem 3. *Assume that the initial data $u_0 \in H^1(\mathbb{R}^2)$, $H(u_0) \leq 1$, and $0 < \lambda < 4\pi$. Let u be a global solution of (5). Then*

$$\|u\|_{L^4_{t,x}(\mathbb{R} \times \mathbb{R}^2)} < \infty. \quad (6)$$

In Section 2, Lemma 9 will show us that Theorem 3 implies the following scattering result.

Theorem 4. *Assume that the initial data $u_0 \in H^1(\mathbb{R}^2)$, $H(u_0) \leq 1$, and $0 < \lambda < 4\pi$. Then the solution of (5) is scattering in the energy space $H^1(\mathbb{R}^2)$.*

We will prove Theorem 3 by contradiction in Section 5. In Section 2, we give some nonlinear estimates. In Section 3, we prove the stability of solutions. In Section 4, we give a new profile decomposition for H^1 sequence which will be used to prove concentration compactness.

Now, we introduce some notations:

$$\begin{aligned} G(u) &:= \bar{u}f(u) - F(u) = e^{\lambda|u|^2}|u|^2 - \frac{1}{\lambda}(e^{\lambda|u|^2} - 1) \\ &= \sum_{k=1}^{\infty} \frac{k\lambda^k |u|^{2k+2}}{(k+1)!}, \end{aligned} \quad (7)$$

$$E = E(u; t) := M(u; t) + H(u; t).$$

We define

$$\begin{aligned} \|u\|_{H^s_q(\mathbb{R}^2)} &:= \|(I - \Delta)^{s/2} u\|_{L^q(\mathbb{R}^2)}, \\ \|u\|_{\dot{H}^s_q(\mathbb{R}^2)} &:= \|(-\Delta)^{s/2} u\|_{L^q(\mathbb{R}^2)}. \end{aligned} \quad (8)$$

For Banach space $X = H^s_q(\mathbb{R}^2)$, $\dot{H}^s_q(\mathbb{R}^2)$, or $L^q(\mathbb{R}^2)$, we denote

$$\|u\|_{L^p(\mathbb{R}; X)} := \left(\int_{\mathbb{R}} \|u(t)\|_X^p dt \right)^{1/p}. \quad (9)$$

When $q = r$, $L^q_t L^r_x$ is abbreviated to $L^q_{t,x}$. When q or r is infinity or when the domain $\mathbb{R} \times \mathbb{R}^2$ is replaced by $I \times \mathbb{R}^2$, we make the usual modifications. Specially, we denote

$$S(u) := \|u\|_{L^4_{t,x}(\mathbb{R} \times \mathbb{R}^2)}^4. \quad (10)$$

For $t_0 \in \mathbb{R}$, we split $S(u) = S_{\leq t_0}(u) + S_{\geq t_0}(u)$, where

$$\begin{aligned} S_{\leq t_0}(u) &:= \int_{-\infty}^{t_0} \int_{\mathbb{R}^2} |u(t, x)|^4 dx dt, \\ S_{\geq t_0}(u) &:= \int_{t_0}^{+\infty} \int_{\mathbb{R}^2} |u(t, x)|^4 dx dt. \end{aligned} \quad (11)$$

For any two Banach spaces X and Y , $\|\cdot\|_{X \cap Y} := \max\{\|\cdot\|_X, \|\cdot\|_Y\}$. C denotes positive constant. If C depends upon some parameters, such as λ , we will indicate this with $C(\lambda)$.

Remark 5. Note that $0 < \lambda < 4\pi$ in Theorem 3; we only need to prove the result for $0 < \lambda < 4(1-4\varepsilon)^2\pi$, $\varepsilon \in (0, 1/8)$. Hence, we always suppose that $0 < \lambda < 4(1-4\varepsilon)^2\pi$ in the context.

Moreover, we always suppose that the initial data u_0 of (5) satisfies $u_0 \in H^1(\mathbb{R}^2)$ and $H(u_0) \leq 1$.

2. Nonlinear Estimates

In order to estimate (2), we need the following Trudinger-type inequality.

Lemma 6 (see [10]). *Let $\lambda \in [0, 4\pi)$. Then for all $u \in H^1(\mathbb{R}^2)$ satisfying $\|\nabla u\|_{L^2(\mathbb{R}^2)} \leq 1$, one has*

$$\int_{\mathbb{R}^2} (e^{\lambda|u|^2} - 1) dx \leq C(\lambda) \|u\|_{L^2(\mathbb{R}^2)}^2. \quad (12)$$

Note that for all $\alpha \geq 1$,

$$(e^{\lambda|u|^2} - 1)^\alpha \leq e^{\lambda\alpha|u|^2} - 1. \quad (13)$$

By Lemma 6 and Hölder inequality, for $\lambda \in (0, 4\pi)$ and for all $\beta \geq 0$, we have

$$\begin{aligned} \int_{\mathbb{R}^2} (e^{\lambda|u|^2} - 1) |u|^\beta dx &\leq \|e^{\lambda|u|^2} - 1\|_{L^{1/(1-\varepsilon)}_x(\mathbb{R}^2)} \|u\|_{L^{\beta/\varepsilon}_x(\mathbb{R}^2)}^\beta \\ &\leq \|u\|_{L^2(\mathbb{R}^2)}^2 \|u\|_{H^1(\mathbb{R}^2)}^\beta \\ &\leq C(\lambda, \beta) \|u\|_{L^2(\mathbb{R}^2)}^2, \end{aligned} \quad (14)$$

and thus

$$\begin{aligned} \int_{\mathbb{R}^2} (e^{\lambda|u|^2} - \lambda|u|^2 - 1) dx &\leq \lambda \int_{\mathbb{R}^2} (e^{\lambda|u|^2} - 1) |u|^2 dx \\ &\leq C(\lambda) \|u\|_{L^2(\mathbb{R}^2)}^2. \end{aligned} \quad (15)$$

Lemma 7 (Strichartz estimates). *For $s = 0$ or 1 ,*

$$2 \leq r, \quad p < \infty, \quad \frac{1}{\gamma(p)} + \frac{1}{p} = \frac{1}{2}, \quad \frac{1}{p} + \frac{1}{p'} = 1, \quad (16)$$

(the pairs $(\gamma(p), p)$ were called admissible pairs) we have

$$\begin{aligned} \|e^{i(t-t_0)\Delta} u(t_0)\|_{L^{\gamma(p)}(\mathbb{R}; H^s_p)} &\leq C \|u(t_0)\|_{H^s(\mathbb{R}^2)}, \\ \left\| \int_{t_0}^t e^{i(t-\tau)\Delta} f(u(\tau)) d\tau \right\|_{L^{\gamma(p)}(\mathbb{R}; H^s_p)} &\leq C \|f(u)\|_{L^{\gamma(r)'}(\mathbb{R}; H^s_{r'})}. \end{aligned} \quad (17)$$

Lemma 8 (see [3, Proposition 2.3]). *Let $1 < r < p < \infty$ be fixed indices. Then for any $q \in [p, \infty)$,*

$$\|u\|_{L^q(\mathbb{R}^2)} \leq C(p, r) q^{1/r' + p/rq} \|u\|_{L^p(\mathbb{R}^2)}^{p/q} \|u\|_{\dot{H}^{2/r'}_r(\mathbb{R}^2)}^{1-p/q}. \quad (18)$$

As shown in [6, 11], to obtain the scattering result, it suffices to show that any finite energy solution has a finite global space-time norm. In fact, if Theorem 3 is true, we have the following theorem.

Lemma 9 (Theorem 3 implies Theorem 4). *Let u be a global solution of (5), $H(u) \leq 1$, and $\|u\|_{L^4_{t,x}(\mathbb{R} \times \mathbb{R}^2)} < \infty$. Then, for all admissible pairs, we have*

$$\|u\|_{L^{p(p)}(\mathbb{R}; H^1_p)} < \infty. \quad (19)$$

Moreover, there exist $u_{\pm} \in H^1$ such that

$$\lim_{t \rightarrow \pm\infty} \|u - e^{it\Delta} u_{\pm}\|_{H^1(\mathbb{R}^2)} = 0. \quad (20)$$

Proof. Defining $X(I) = L^{2/(1-2\epsilon)}(I; H^1_{1/\epsilon})$, $Y(I) = L^4(I; H^1_4)$, by Strichartz estimates, (14) and (15),

$$\begin{aligned} & \|u\|_{X(I) \cap Y(I)} \\ & \leq C \|u(S)\|_{H^1(\mathbb{R}^2)} + C \left\| \lambda |u|^2 u \right\|_{L^{4/3}(I; H^1_{4/3})} \\ & \quad + C \left\| \left(e^{\lambda |u|^2} - \lambda |u|^2 - 1 \right) u \right\|_{L^{2/(1+2\epsilon)}(I; H^1_{1/(1-\epsilon)})} \\ & \leq C(E) + C \|u\|_{L^4_{t,x}(I \times \mathbb{R}^2)}^2 \|u\|_{Y(I)} \\ & \quad + C \|u\|_{L^4_{t,x}(I \times \mathbb{R}^2)}^{8\epsilon} \|u\|_{X(I)} \\ & \quad \times \left\| \left(e^{\lambda |u|^2} - \lambda |u|^2 - 1 \right) u^{-8\epsilon} \right\|_{L^\infty_{t,x} L^{1/(1-4\epsilon)}(I \times \mathbb{R}^2)} \\ & \quad + C \|u\|_{L^4_{t,x}(I \times \mathbb{R}^2)}^{8\epsilon} \|u\|_{X(I)} \\ & \quad \times \left\| \left(e^{\lambda |u|^2} - 1 \right) |u|^{2-8\epsilon} \right\|_{L^\infty_{t,x} L^{1/(1-4\epsilon)}(I \times \mathbb{R}^2)} \\ & \leq C(E) + C(E) \left(\|u\|_{L^4_{t,x}(I \times \mathbb{R}^2)}^2 + \|u\|_{L^4_{t,x}(I \times \mathbb{R}^2)}^{8\epsilon} \right) \|u\|_{X(I) \cap Y(I)}. \end{aligned} \quad (21)$$

Using the same way as in Bourgain [12], one can split $\mathbb{R}$ into finitely many pairwise disjoint intervals:

$$\mathbb{R} = \bigcup_{j=1}^J I_j, \quad \|u\|_{L^4_{t,x}(I_j \times \mathbb{R}^2)} \leq \eta, \quad C(E) \left(\eta^2 + \eta^{8\epsilon} \right) \leq \frac{1}{2}. \quad (22)$$

By (21),

$$\|u\|_{L^{2/(1-2\epsilon)}(I_j; H^1_{1/\epsilon}) \cap L^4(I_j; H^1_4)} \leq C(E). \quad (23)$$

Since $\|u\|_{L^\infty(I_j; H^1)} \leq C(E)$ and $\epsilon \in (0, 1/8)$ can be chosen small arbitrarily, by interpolation,

$$\|u\|_{L^{p(p)}(I_j; H^1_p)} \leq C(E), \quad (24)$$

for all admissible pairs and $j = 1, 2, \dots, J$. The desired result (19) follows.

Let

$$u_{\pm} = u_0 - i \int_0^{\pm\infty} e^{-i\tau\Delta} f(u(\tau)) d\tau. \quad (25)$$

By (19) and (21),

$$\|u_{\pm}\|_{H^1(\mathbb{R}^2)} \lesssim 1 + \|u\|_{Y(0, \pm\infty)}^3 + \|u\|_{X(0, \pm\infty) \cap Y(0, \pm\infty)}^{1+8\epsilon} \lesssim 1. \quad (26)$$

Thus, $u_{\pm}$ were well defined and belong to H^1 . Since

$$\begin{aligned} \|u - e^{it\Delta} u_{\pm}\|_{H^1(\mathbb{R}^2)} &= \left\| \int_t^{\pm\infty} e^{i(t-\tau)\Delta} f(u(\tau)) d\tau \right\|_{H^1(\mathbb{R}^2)} \\ &\lesssim \|u\|_{Y(t, \pm\infty)}^3 + \|u\|_{X(t, \pm\infty) \cap Y(t, \pm\infty)}^{1+8\epsilon}, \end{aligned} \quad (27)$$

$$\|u\|_{X(0, \pm\infty) \cap Y(0, \pm\infty)} \lesssim 1,$$

we must have

$$\begin{aligned} & \lim_{t \rightarrow \pm\infty} \|u - e^{it\Delta} u_{\pm}\|_{H^1(\mathbb{R}^2)} \\ & \leq \lim_{t \rightarrow \pm\infty} \left(\|u\|_{Y(t, \pm\infty)}^3 + \|u\|_{X(t, \pm\infty) \cap Y(t, \pm\infty)}^{1+8\epsilon} \right) = 0. \end{aligned} \quad (28)$$

(20) was proved. $\square$

3. Stability

Lemma 10 (stability). *For any $A > 0$ and $\sigma > 0$, there exists $\delta > 0$ with the following property: suppose that $u : [0, \infty) \times \mathbb{R}^2 \rightarrow \mathbb{C}$ satisfies $\|\nabla u\|_{L^2_x(\mathbb{R}^2)} \leq 1$ for all $t \in [0, \infty)$, $\|u\|_{L^4_{t,x}([0, \infty) \times \mathbb{R}^2)} \leq A$ and approximately solves (5) in the sense that*

$$\begin{aligned} & \left\| \int_0^t e^{i(t-\tau)\Delta} (iu_t + \Delta u - f(u))(\tau) d\tau \right\|_{L^4_{t,x} \cap L^{2/(1-2\epsilon)}_{t,x} L^{1/\epsilon}_x([0, \infty) \times \mathbb{R}^2)} \\ & \leq \delta. \end{aligned} \quad (29)$$

Then for any initial data $v(0) \in H^1(\mathbb{R}^2)$ satisfying $H(v(0)) \leq 1$ and $\|u(0) - v(0)\|_{L^2_x(\mathbb{R}^2)} \leq \delta$, there is a unique global solution v to (5) satisfying $\|u - v\|_{L^4_{t,x}([0, \infty) \times \mathbb{R}^2)} \leq \sigma$.

Proof. Denote $v = u + w$, then

$$iw_t + \Delta w = (f(u + w) - f(u)) - (iu_t + \Delta u - f(u)) \quad (30)$$

and $\|w(0)\|_{L^2(\mathbb{R}^2)} \leq \delta$. Let $X = L^\infty_t L^2_x \cap L^4_{t,x} \cap L^{2/(1-2\epsilon)}_t L^{1/\epsilon}_x$. By the similar estimates as (21), we have

$$\begin{aligned} \|w\|_X &\leq \delta + \left\| \int_0^t e^{i(t-\tau)\Delta} (f(u + w) - f(u))(\tau) d\tau \right\|_X + \delta \\ &\leq 2\delta + C \| |u + w|^2 (u + w) - |u|^2 u \|_{L^{4/3}_{t,x}} \\ &\quad + C \left\| \left(e^{\lambda |u+w|^2} - \lambda |u + w|^2 - 1 \right) (u + w) \right. \\ &\quad \left. - \left(e^{\lambda |u|^2} - \lambda |u|^2 - 1 \right) u \right\|_{L^{2/(1+2\epsilon)}_t L^{1/(1-\epsilon)}_x} \end{aligned}$$

$$\begin{aligned}
&\leq 2\delta + C \| |u+w|^2 |w| + |u|^2 |w| \|_{L_{t,x}^{4/3}} \\
&\quad + C \left\| \left(e^{\lambda|u+w|^2} - \lambda|u+w|^2 - 1 \right) |w| \right. \\
&\quad \left. + \left(e^{\lambda|u|^2} - \lambda|u|^2 - 1 \right) |w| \right\|_{L_t^{2/(1+2\epsilon)} L_x^{1/(1-\epsilon)}} \\
&\quad + C \left\| \left(e^{\lambda|u+w|^2} - 1 \right) |u+w|^2 |w| \right. \\
&\quad \left. + \left(e^{\lambda|u|^2} - 1 \right) |u|^2 |w| \right\|_{L_t^{2/(1+2\epsilon)} L_x^{1/(1-\epsilon)}} \\
&\leq 2\delta + C \left(\|u\|_{L_{t,x}^4}^2 + \|u+w\|_{L_{t,x}^4}^2 \right) \|w\|_{L_{t,x}^4} \\
&\quad + C \left(\|u\|_{L_{t,x}^4}^{8\epsilon} + \|u+w\|_{L_{t,x}^4}^{8\epsilon} \right) \|w\|_{L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}} \\
&\leq 2\delta + C \left(\|u\|_{L_{t,x}^4}^2 + \|u\|_{L_{t,x}^4}^{8\epsilon} \right) \|w\|_X + \|w\|_X^{1+8\epsilon} + \|w\|_X^3.
\end{aligned} \tag{31}$$

Then we subdivide the time interval $[0, \infty)$ into finite subintervals $I_j = [t_j, t_{j+1}]$, $j = 1, \dots, J$, $J = J(A)$ such that

$$C \left(\|u\|_{L_{t,x}^4(I_j \times \mathbb{R}^2)}^2 + \|u\|_{L_{t,x}^4(I_j \times \mathbb{R}^2)}^{8\epsilon} \right) < \frac{1}{4} \tag{32}$$

for each j . Let $\delta = \delta(A, J)$ be small such that

$$C8^{2J}\delta \ll 1, \quad 8^{2J}\delta \ll \sigma. \tag{33}$$

Then by (31) on $I_1 \times \mathbb{R}^2$, we have $\|w\|_{X(I_1 \times \mathbb{R}^2)} \leq 8\delta$ and

$$\begin{aligned}
&\|w(t_2)\|_{L_x^2(\mathbb{R}^2)} \\
&\leq \|e^{i(t_2-t_1)\Delta} w(t_1)\|_{L_x^2(\mathbb{R}^2)} \\
&\quad + \left\| \int_{t_1}^{t_2} e^{i(t_2-\tau)\Delta} (f(u+w) - f(u))(\tau) d\tau \right\|_{L_x^2(\mathbb{R}^2)} + \delta \\
&\leq 2\delta + \delta + C \left(\|u\|_{L_{t,x}^4(I_1 \times \mathbb{R}^2)}^2 + \|u\|_{L_{t,x}^4(I_1 \times \mathbb{R}^2)}^{8\epsilon} \right) \|w\|_{X(I_1 \times \mathbb{R}^2)} \\
&\quad + \|w\|_{X(I_1 \times \mathbb{R}^2)}^{1+8\epsilon} + \|w\|_{X(I_1 \times \mathbb{R}^2)}^3 \\
&\leq 2\delta + 4 \cdot 8\delta \leq 8^2\delta.
\end{aligned} \tag{34}$$

Using the same analysis as above, we can get $\|w\|_{X(I_2 \times \mathbb{R}^2)} \leq 8^3\delta$. Iterating this for $I_2, I_3, \dots, I_J$, we obtain $\|w\|_{X([0, \infty) \times \mathbb{R}^2)} \leq 8^{2J}\delta \leq \sigma$; the desired result was obtained. $\square$

4. Linear Profile Decomposition

In this section, we will give the linear profile decomposition for Schrödinger equation in $H^1(\mathbb{R}^2)$. First, we give some definitions and lemmas.

Definition 11 (symmetry group, [13]). For any phase $\theta \in \mathbb{R}/2\pi\mathbb{Z}$, position $x_0 \in \mathbb{R}^2$, frequency $\xi_0 \in \mathbb{R}^2$, and scaling

parameter $\lambda > 0$, we define the unitary transformation $\mathcal{G}_{\theta, \xi_0, x_0, \lambda} : L_x^2(\mathbb{R}^2) \rightarrow L_x^2(\mathbb{R}^2)$ by the formula

$$\mathcal{G}_{\theta, \xi_0, x_0, \lambda} f(x) := \frac{1}{\lambda} e^{i\theta} e^{ix \cdot \xi_0} f\left(\frac{x - x_0}{\lambda}\right). \tag{35}$$

We let G be the collection of such transformations; this is a group with identity $\mathcal{G}_{0,0,0,1}$, inverse $\mathcal{G}_{\theta, \xi_0, x_0, \lambda}^{-1} = \mathcal{G}_{-\theta - x_0 \xi_0, -\lambda \xi_0, -x_0/\lambda, 1/\lambda}$, and group law

$$\mathcal{G}_{\theta, \xi_0, x_0, \lambda} \mathcal{G}_{\theta', \xi'_0, x'_0, \lambda'} = \mathcal{G}_{\theta+\theta' - x_0 \xi'_0/\lambda, \xi_0 + \xi'_0/\lambda, x_0 + \lambda x'_0, \lambda \lambda'}. \tag{36}$$

If $u : I \times \mathbb{R}^2 \rightarrow \mathbb{C}$ is a function, we define $T_{\mathcal{G}_{\theta, \xi_0, x_0, \lambda}} u : \lambda^2 I \times \mathbb{R}^2 \rightarrow \mathbb{C}$, where $\lambda^2 I := \{\lambda^2 t : t \in I\}$ by the formula

$$(T_{\mathcal{G}_{\theta, \xi_0, x_0, \lambda}} u)(t, x) := \frac{1}{\lambda} e^{i\theta} e^{ix \cdot \xi_0} e^{-it|\xi_0|^2} u\left(\frac{t}{\lambda^2}, \frac{x - x_0 - 2\xi_0 t}{\lambda}\right), \tag{37}$$

or equivalently

$$(T_{\mathcal{G}_{\theta, \xi_0, x_0, \lambda}} u)(t) = \mathcal{G}_{\theta - t|\xi_0|^2, \xi_0, x_0 + 2\xi_0 t, \lambda} \left(u\left(\frac{t}{\lambda^2}\right)\right). \tag{38}$$

If $g \in G$, we can easily prove that $M(T_g u) = M(u)$ and $S(T_g u) = S(u)$.

Definition 12 (enlarged group, [13]). For any phase $\theta \in \mathbb{R}/2\pi\mathbb{Z}$, position $x_0 \in \mathbb{R}^2$, frequency $\xi_0 \in \mathbb{R}^2$, scaling parameter $\lambda > 0$, and time t_0 , we define the unitary transformation $\mathcal{G}_{\theta, \xi_0, x_0, \lambda, t_0} : L_x^2(\mathbb{R}^2) \rightarrow L_x^2(\mathbb{R}^2)$ by the formula

$$\mathcal{G}_{\theta, \xi_0, x_0, \lambda, t_0} = \mathcal{G}_{\theta, \xi_0, x_0, \lambda} e^{it_0 \Delta}, \tag{39}$$

or in other words

$$\mathcal{G}_{\theta, \xi_0, x_0, \lambda, t_0} f(x) := \frac{1}{\lambda} e^{i\theta} e^{ix \cdot \xi_0} (e^{it_0 \Delta} f)\left(\frac{x - x_0}{\lambda}\right). \tag{40}$$

Let G' be the collection of such transformations. We also let G' act on global space-time function $u : \mathbb{R} \times \mathbb{R}^2 \rightarrow \mathbb{C}$ by defining

$$\begin{aligned}
&T_{\mathcal{G}_{\theta, \xi_0, x_0, \lambda, t_0}} u(t, x) \\
&:= \frac{1}{\lambda} e^{i\theta} e^{ix \cdot \xi_0} e^{-it|\xi_0|^2} (e^{it_0 \Delta} u)\left(\frac{t}{\lambda^2}, \frac{x - x_0 - 2\xi_0 t}{\lambda}\right),
\end{aligned} \tag{41}$$

or equivalently

$$(T_{\mathcal{G}_{\theta, \xi_0, x_0, \lambda, t_0}} u)(t) = \mathcal{G}_{\theta - t|\xi_0|^2, \xi_0, x_0 + 2\xi_0 t, \lambda, t_0} \left(u\left(\frac{t}{\lambda^2}\right)\right). \tag{42}$$

Lemma 13 (linear profiles for L^2 sequence, [14]). Let u_n be a bounded sequence in $L_x^2(\mathbb{R}^2)$. Then (after passing to a subsequence if necessary) there exists a family $\phi^{(j)}$, $j = 1, 2, \dots$ of functions in $L_x^2(\mathbb{R}^2)$ and group elements $g_n^{(j)} \in G'$ for $j, n = 1, 2, \dots$ such that one has the decomposition

$$u_n = \sum_{j=1}^l g_n^{(j)} \phi^{(j)} + w_n^{(l)} \tag{43}$$

for all $l = 1, 2, \dots$; here, $w_n^{(l)} \in L_x^2(\mathbb{R}^2)$ is such that its linear evolution has asymptotically vanishing scattering size:

$$\lim_{l \rightarrow \infty} \limsup_{n \rightarrow \infty} S(e^{it\Delta} w_n^{(l)}) = 0. \quad (44)$$

Moreover, for any $j \neq j'$,

$$\begin{aligned} \frac{\lambda_n^{(j')}}{\lambda_n^{(j)}} + \frac{\lambda_n^{(j)}}{\lambda_n^{(j')}} + \lambda_n^{(j)} \lambda_n^{(j')} \left| \xi_n^{(j)} - \xi_n^{(j')} \right|^2 + \frac{\left| x_n^{(j)} - x_n^{(j')} \right|^2}{\lambda_n^{(j)} \lambda_n^{(j')}} \\ + \frac{\left| t_n^{(j)} (\lambda_n^{(j)})^2 - t_n^{(j')} (\lambda_n^{(j')})^2 \right|}{\lambda_n^{(j)} \lambda_n^{(j')}} \rightarrow \infty. \end{aligned} \quad (45)$$

Furthermore, for any $l \geq 1$, one has the mass decoupling property

$$\lim_{n \rightarrow \infty} \left[M(u_n) - \sum_{j=1}^l M(\phi^{(j)}) - M(w_n^{(l)}) \right] = 0. \quad (46)$$

For any $j \leq l$, we have

$$(g_n^{(j)})^{-1} w_n^{(l)} \rightarrow 0 \quad \text{weakly in } L_x^2(\mathbb{R}^2). \quad (47)$$

Remark 14. If the orthogonal condition (45) holds, then (see [14])

$$\begin{aligned} \lim_{n \rightarrow \infty} \langle g_n^{(j)} \phi^{(j)}, g_n^{(j')} \phi^{(j')} \rangle_{L^2(\mathbb{R}^2)} &= 0, \quad j \neq j', \\ \lim_{n \rightarrow \infty} \langle g_n^{(j)} \phi^{(j)}, w_n^{(l)} \rangle_{L^2(\mathbb{R}^2)} &= 0. \end{aligned} \quad (48)$$

Moreover, if $v^{(j)}, v^{(j')} \in L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)$, then (see [14, 15]), for any $0 < \theta < 1$,

$$\lim_{n \rightarrow \infty} \left\| \left| T_{g_n^{(j)}} v^{(j)} \right|^{1-\theta} \left| T_{g_n^{(j')}} v^{(j')} \right|^\theta \right\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} = 0. \quad (49)$$

If $v^{(1)}, \dots, v^{(l)} \in L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)$, then (see [16, Lemma 5.5])

$$\lim_{n \rightarrow \infty} S \left(\sum_{j=1}^l g_n^{(j)} v^{(j)} \right) \leq \sum_{j=1}^l S(v^{(j)}). \quad (50)$$

Remark 15. As each linear profile $\phi^{(j)}$ in Lemma 13 is constructed in the sense that

$$e^{it_n^{(j)} \Delta} \left\{ e^{i\xi_n^{(j)} \cdot x} \lambda_n^{(j)} u_n(\lambda_n^{(j)} x) \right\} (x + x_n^{(j)}) \rightarrow \phi^{(j)} \quad (51)$$

weakly in $L_x^2(\mathbb{R}^2)$ (see [14]), after passing to a subsequence in n , rearrangement, translation, and refining $\phi^{(j)}$ accordingly, we may assume that the parameters satisfy the following properties:

- (i) $t_n^{(j)} \rightarrow \pm\infty$ as $n \rightarrow \infty$, or $t_n^{(j)} \equiv 0$ for all n, j ;
- (ii) $\lambda_n^{(j)} \rightarrow 0$ or ∞ as $n \rightarrow \infty$, or $\lambda_n^{(j)} \equiv 1$ for all n, j ;

- (iii) $|\xi_n^{(j)}| \rightarrow \infty$ as $n \rightarrow \infty$, or $\xi_n^{(j)} \equiv \xi^{(j)}$ with $|\xi^{(j)}| < \infty$;
- (iv) when $\lambda_n^{(j)} \equiv 1$, $\xi_n^{(j)} \equiv \xi^{(j)}$ and $|\xi^{(j)}| < \infty$, we can let $\xi^{(j)} \equiv 0$.

Our main result in this section is the following lemma.

Lemma 16 (linear profiles for H^1 sequence). *Let u_n be a bounded sequence in $H^1(\mathbb{R}^2)$. Then up to a subsequence, for any $J \geq 1$, there exists a sequence ϕ_α^J in $H^1(\mathbb{R}^2)$ and a sequence of group elements $g_{n\alpha} = g_{\theta_{n\alpha}, \xi_{n\alpha}, x_{n\alpha}, \lambda_{n\alpha}, t_{n\alpha}} \in G'$ such that*

$$u_n = \sum_{\alpha=1}^J g_{n\alpha} \phi_\alpha^J + R(n, J). \quad (52)$$

Here, for each α , $\lambda_{n\alpha}$ and $\xi_{n\alpha}$ must satisfy

$$“\lambda_{n\alpha} \equiv 1 \text{ and } \xi_{n\alpha} \equiv 0,” \quad \text{or} \quad “\lambda_{n\alpha} \rightarrow \infty.” \quad (53)$$

$R(n, J) \in H^1(\mathbb{R}^2)$ is such that

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} S(e^{it\Delta} R(n, J)) = 0. \quad (54)$$

Moreover, for any $\alpha \neq \alpha'$, one has the same orthogonal conditions as (45). For any $J \geq 1$, one has the following decoupling properties:

$$\lim_{n \rightarrow \infty} \left\{ \|u_n\|_{L^2(\mathbb{R}^2)}^2 - \sum_{\alpha=1}^J \|\phi_\alpha^J\|_{L^2(\mathbb{R}^2)}^2 - \|R(n, J)\|_{L^2(\mathbb{R}^2)}^2 \right\} = 0, \quad (55)$$

$$\begin{aligned} \lim_{n \rightarrow \infty} \left\{ \|u_n\|_{H^1(\mathbb{R}^2)}^2 - \sum_{\alpha=1}^J \|g_{n\alpha} \phi_\alpha^J\|_{H^1(\mathbb{R}^2)}^2 - \|R(n, J)\|_{H^1(\mathbb{R}^2)}^2 \right\} \\ = 0, \end{aligned} \quad (56)$$

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \left\{ H(u_n) - \sum_{\alpha=1}^J H(g_{n\alpha} \phi_\alpha^J) - H(R(n, J)) \right\} = 0. \quad (57)$$

Proof. Let

$$\square_k = \mathcal{F}^{-1} \chi_k \mathcal{F}, \quad \chi_k = \begin{cases} 1 & 2^{k-1} < |\xi| \leq 2^k, \\ 0 & \text{else.} \end{cases} \quad (58)$$

Then, we have

$$\begin{aligned} u_n &= \sum_{k=-\infty}^{+\infty} \square_k u_n := \sum_{|k| \leq N} \square_k u_n + R_N, \\ \|u_n\|_{L^2(\mathbb{R}^2)}^2 &= \sum_{|k| \leq N} \|\square_k u_n\|_{L^2(\mathbb{R}^2)}^2 + \|R_N\|_{L^2(\mathbb{R}^2)}^2, \\ \|u_n\|_{H^1(\mathbb{R}^2)}^2 &= \sum_{|k| \leq N} \|\square_k u_n\|_{H^1(\mathbb{R}^2)}^2 + \|R_N\|_{H^1(\mathbb{R}^2)}^2, \end{aligned} \quad (59)$$

$$\lim_{N \rightarrow \infty} \limsup_{n \rightarrow \infty} \|R_N\|_{L^2(\mathbb{R}^2)} = 0.$$

By Lemma 13, after passing to a subsequence if necessary, we can obtain

$$\square_k u_n = \sum_{j=1}^{l_k} g_{nk}^{(j)} \phi_k^{(j)} + w_{nk}^{(l_k)} \quad (60)$$

with the stated properties (i)–(iv) in Remark 15 and (43)–(47). Denote

$$\begin{aligned} \Lambda_{1,0} &= \{(k, j) \mid |k| \leq N, 1 \leq j \leq l_k, \lambda_{nk}^{(j)} \equiv 1, \xi_{nk}^{(j)} \equiv 0\}, \\ \Lambda_{1,\infty} &= \{(k, j) \mid |k| \leq N, 1 \leq j \leq l_k, \lambda_{nk}^{(j)} \equiv 1, |\xi_{nk}^{(j)}| \rightarrow \infty\}, \\ \Lambda_0 &= \{(k, j) \mid |k| \leq N, 1 \leq j \leq l_k, \lambda_{nk}^{(j)} \rightarrow 0\}, \\ \Lambda_{\infty,0} &= \{(k, j) \mid |k| \leq N, 1 \leq j \leq l_k, \lambda_{nk}^{(j)} \rightarrow \infty, \xi_{nk}^{(j)} \equiv \xi_k^{(j)}, \\ &\quad |\xi_k^{(j)}| < 2^{k-1}\}, \\ \Lambda_{\infty,1} &= \{(k, j) \mid |k| \leq N, 1 \leq j \leq l_k, \lambda_{nk}^{(j)} \rightarrow \infty, \xi_{nk}^{(j)} \equiv \xi_k^{(j)}, \\ &\quad |\xi_k^{(j)}| \in [2^{k-1}, 2^k]\}, \\ \Lambda_{\infty,\infty} &= \{(k, j) \mid |k| \leq N, 1 \leq j \leq l_k, \lambda_{nk}^{(j)} \rightarrow \infty, |\xi_{nk}^{(j)}| \rightarrow \infty \\ &\quad \text{or } \xi_{nk}^{(j)} \equiv \xi_k^{(j)}, |\xi_k^{(j)}| > 2^k\}. \end{aligned} \quad (61)$$

Step 1. We prove that

$$u_n = \sum_{(k,j) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}} g_{nk}^{(j)} \phi_k^{(j)} + R \quad (62)$$

with $\phi_k^{(j)} \in H^1$ and for each fixed N ,

$$\lim_{n \rightarrow \infty} \left\{ \|u_n\|_{L^2(\mathbb{R}^2)}^2 - \sum_{(k,j) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}} \|\phi_k^{(j)}\|_{L^2(\mathbb{R}^2)}^2 - \|R\|_{L^2(\mathbb{R}^2)}^2 \right\} = 0, \quad (63)$$

$$\lim_{n \rightarrow \infty} \left\{ \|u_n\|_{\dot{H}^1(\mathbb{R}^2)}^2 - \sum_{(k,j) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}} \|g_{nk}^{(j)} \phi_k^{(j)}\|_{\dot{H}^1(\mathbb{R}^2)}^2 - \|R\|_{\dot{H}^1(\mathbb{R}^2)}^2 \right\} = 0,$$

$$\lim_{N \rightarrow \infty} \lim_{l_k \rightarrow \infty} \limsup_{n \rightarrow \infty} S(e^{it\Delta} R) = 0, \quad (64)$$

where

$$R = R_N + R_w, \quad R_w = \sum_{|k| \leq N} w_{nk}^{(l_k)}. \quad (65)$$

By (44) and $\lim_{N \rightarrow \infty} \limsup_{n \rightarrow \infty} \|R_N\|_{L^2(\mathbb{R}^2)} = 0$, (64) holds obviously. For (62), we prove it by induction. For every k , suppose that

$$\square_k u_n = g_{nk}^{(1)} \phi_k^{(1)} + w_{nk}^{(1)}. \quad (66)$$

Case 1. If $(k, 1) \in \Lambda_{1,\infty} \cup \Lambda_0 \cup \Lambda_{\infty,0} \cup \Lambda_{\infty,\infty}$, we have $\phi_k^{(1)} = 0$. In fact, by (66),

$$\phi_k^{(1)} = (g_{nk}^{(1)})^{-1} \square_k u_n - (g_{nk}^{(1)})^{-1} w_{nk}^{(1)}. \quad (67)$$

Thus,

$$\begin{aligned} \|\phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 &= \left\langle (g_{nk}^{(1)})^{-1} \square_k u_n - (g_{nk}^{(1)})^{-1} w_{nk}^{(1)}, \phi_k^{(1)} \right\rangle_{L^2(\mathbb{R}^2)} \\ &= \langle u_n, \square_k g_{nk}^{(1)} \phi_k^{(1)} \rangle_{L^2(\mathbb{R}^2)} \\ &\quad - \left\langle (g_{nk}^{(1)})^{-1} w_{nk}^{(1)}, \phi_k^{(1)} \right\rangle_{L^2(\mathbb{R}^2)}. \end{aligned} \quad (68)$$

Using (47),

$$\left\langle (g_{nk}^{(1)})^{-1} w_{nk}^{(1)}, \phi_k^{(1)} \right\rangle_{L^2(\mathbb{R}^2)} \rightarrow 0 \quad \text{as } n \rightarrow \infty. \quad (69)$$

By direct calculation,

$$\begin{aligned} \square_k g_{nk}^{(1)} \phi_k^{(1)} &= \mathcal{F}^{-1} \chi_k(\xi) \lambda_{nk}^{(1)} e^{i\theta_{nk}^{(1)}} e^{-ix_{nk}^{(1)}(\xi + \xi_{nk}^{(1)})} e^{-it_{nk}^{(1)}(\lambda_{nk}^{(1)})^2 |\xi + \xi_{nk}^{(1)}|^2} \\ &\quad \times \widehat{\phi_k^{(1)}}(\lambda_{nk}^{(1)}(\xi + \xi_{nk}^{(1)})) \\ &= \frac{1}{\lambda_{nk}^{(1)}} e^{i\theta_{nk}^{(1)}} e^{ix_{nk}^{(1)}} \left(\mathcal{F}^{-1} \chi_k \left(\frac{\xi}{\lambda_{nk}^{(1)}} + \xi_{nk}^{(1)} \right) e^{-it_{nk}^{(1)} |\xi|^2} \widehat{\phi_k^{(1)}} \right) \\ &\quad \times \left(\frac{x - x_{nk}^{(1)}}{\lambda_{nk}^{(1)}} \right). \end{aligned} \quad (70)$$

Let $n \rightarrow \infty$. When $(k, 1) \in \Lambda_{1,\infty}$,

$$\|\square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \leq \int_{2^{k-1} \leq |\xi + \xi_{nk}^{(1)}| \leq 2^k} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \rightarrow 0. \quad (71)$$

When $(k, 1) \in \Lambda_0$,

$$\|\square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \leq \int_{\lambda_{nk}^{(1)} 2^{k-1} \leq |\xi + \lambda_{nk}^{(1)} \xi_{nk}^{(1)}| \leq \lambda_{nk}^{(1)} 2^k} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \rightarrow 0. \quad (72)$$

When $(k, 1) \in \Lambda_{\infty,0}$,

$$\|\square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \leq \int_{|\xi| \geq \lambda_{nk}^{(1)} (2^{k-1} - |\xi_{nk}^{(1)}|)} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \rightarrow 0. \quad (73)$$

When $(k, 1) \in \Lambda_{\infty, \infty}$,

$$\|\square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \leq \int_{|\xi| \geq \lambda_{nk}^{(1)} (|\xi_k^{(1)}| - 2^k)} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \longrightarrow 0. \quad (74)$$

By (68)–(74), $\|\phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 = 0$ and thus $\phi_k^{(1)} = 0$.

Case 2. If $(k, 1) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}$, we can prove

$$\|g_{nk}^{(1)} \phi_k^{(1)} - \square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)} \longrightarrow 0 \quad \text{as } n \longrightarrow \infty. \quad (75)$$

By absorbing the error into $w_{nk}^{(1)}$, we can suppose $g_{nk}^{(1)} \phi_k^{(1)} = \square_k g_{nk}^{(1)} \phi_k^{(1)}$. Since $\square_k g_{nk}^{(1)} \phi_k^{(1)} \in H^1$ for each fixed n , we must have $\phi_k^{(1)} \in H^1$.

Now, we begin to prove (75). Let $\chi_{A_k^{(1)}}$ be the characteristic function of the set $A_k^{(1)}$ and $P_{A_k^{(1)}} = \mathcal{F}^{-1} \chi_{A_k^{(1)}} \mathcal{F}$, and then

$$\begin{aligned} g_{nk}^{(1)} (P_{A_k^{(1)}} \phi_k^{(1)} + P_{A_k^{(1)}} (g_{nk}^{(1)})^{-1} w_{nk}^{(1)}) &= P_{A_k^{(1)}} (g_{nk}^{(1)} \phi_k^{(1)} + w_{nk}^{(1)}) \\ &= P_{A_k^{(1)}} \square_k u_n, \end{aligned} \quad (76)$$

where

$$P_{A_k^{(1)}} = \mathcal{F}^{-1} \chi_{A_k^{(1)}} (\lambda_{nk}^{(1)} (\xi - \xi_k^{(1)})) \mathcal{F}. \quad (77)$$

Note that

$$\begin{aligned} \langle P_{A_k^{(1)}} \phi_k^{(1)}, P_{A_k^{(1)}} (g_{nk}^{(1)})^{-1} w_{nk}^{(1)} \rangle_{L^2(\mathbb{R}^2)} \\ = \langle P_{A_k^{(1)}} \phi_k^{(1)}, (g_{nk}^{(1)})^{-1} w_{nk}^{(1)} \rangle_{L^2(\mathbb{R}^2)} \longrightarrow 0. \end{aligned} \quad (78)$$

We have

$$\begin{aligned} \lim_{n \rightarrow \infty} \left(\|P_{A_k^{(1)}} \square_k u_n\|_{L^2(\mathbb{R}^2)}^2 - \|P_{A_k^{(1)}} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \right. \\ \left. - \|P_{A_k^{(1)}} (g_{nk}^{(1)})^{-1} w_{nk}^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \right) = 0. \end{aligned} \quad (79)$$

When $(k, 1) \in \Lambda_{1,0}$, we have $P_{A_k^{(1)}} = P_{A_k^{(1)}}$. Choosing $A_k^{(1)} = \{\xi \mid |\xi| \leq 2^{k-1} \text{ or } |\xi| > 2^k\}$, then by (79), $P_{A_k^{(1)}} \phi_k^{(1)} = 0$, the desired result follows.

When $(k, 1) \in \Lambda_{\infty,1}$ and $|\xi_k^{(1)}| \in (2^{k-1}, 2^k)$, we have

$$\begin{aligned} \|g_{nk}^{(1)} \phi_k^{(1)} - \square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \\ \leq \int_{|\xi + \lambda_{nk}^{(1)} \xi_k^{(1)}| \leq \lambda_{nk}^{(1)} 2^{k-1}} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \\ + \int_{|\xi + \lambda_{nk}^{(1)} \xi_k^{(1)}| \geq \lambda_{nk}^{(1)} 2^k} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \\ \leq \int_{|\xi| \geq \lambda_{nk}^{(1)} (|\xi_k^{(1)}| - 2^{k-1})} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \\ + \int_{|\xi| \geq \lambda_{nk}^{(1)} (2^k - |\xi_k^{(1)}|)} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \longrightarrow 0 \\ \text{as } n \longrightarrow \infty. \end{aligned} \quad (80)$$

When $(k, 1) \in \Lambda_{\infty,1}$ and $|\xi_k^{(1)}| = 2^k$, we denote $\xi = (\xi_1, \xi_2)$ and $\xi_k^{(1)} = (\xi_{k1}^{(1)}, \xi_{k2}^{(1)})$. The line $\xi_2 = -(\xi_{k1}^{(1)}/\xi_{k2}^{(1)})\xi_1$ (when $\xi_{k2}^{(1)} = 0$, we use the line $\xi_1 = 0$ instead) separates the frequency space $\mathbb{R}^2$ into two half-planes. We let $A_k^{(1)}$ to be the half-plane which contains the point $\xi_k^{(1)}$, and then

$$P_{A_k^{(1)}} \square_k u_n = \mathcal{F}^{-1} \chi_{A_k^{(1)}} (\lambda_{nk}^{(1)} (\xi - \xi_k^{(1)})) \chi_k \mathcal{F} u_n = 0. \quad (81)$$

By (79), we have $P_{A_k^{(1)}} \phi_k^{(1)} = 0$. Note that

$$\begin{aligned} \lim_{n \rightarrow \infty} \|g_{nk}^{(1)} (I - P_{A_k^{(1)}}) \phi_k^{(1)} - \square_k g_{nk}^{(1)} \phi_k^{(1)}\|_{L^2(\mathbb{R}^2)}^2 \\ = \lim_{n \rightarrow \infty} \left\| \left[(1 - \chi_{A_k^{(1)}}) (\cdot) - \chi_k \left(\frac{\cdot}{\lambda_{nk}^{(1)}} + \xi_k^{(1)} \right) \right] \mathcal{F} \phi_k^{(1)} \right\|_{L^2(\mathbb{R}^2)}^2 \\ \leq \lim_{n \rightarrow \infty} \int_{\xi \in \mathbb{R}^2 \setminus (A_k^{(1)} \cup \{2^{k-1} \leq |\xi/\lambda_{nk}^{(1)} + \xi_k^{(1)}| \leq 2^k\})} |\mathcal{F} \phi_k^{(1)}|^2 d\xi \\ = 0. \end{aligned} \quad (82)$$

(75) holds.

When $(k, 1) \in \Lambda_{\infty,1}$ and $|\xi_k^{(1)}| = 2^{k-1}$, let $A_k^{(1)}$ be the half-plane which does NOT contain the point $\xi_k^{(1)}$; we can prove (75) similarly as above.

By the proof above, we get

$$\square_k u_n = g_{nk}^{(1)} \phi_k^{(1)} + w_{nk}^{(1)}, \quad g_{nk}^{(1)} \phi_k^{(1)} = \square_k g_{nk}^{(1)} \phi_k^{(1)}, \quad \phi_k^{(1)} \in H^1 \quad (83)$$

and $(k, 1) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}$. Denote $u_n^{(1)} = u_n - g_{nk}^{(1)} \phi_k^{(1)}$ and suppose

$$\square_k u_n^{(1)} = g_{nk}^{(2)} \phi_k^{(2)} + w_{nk}^{(2)}. \quad (84)$$

Repeating the proof above, we can get $g_{nk}^{(2)} \phi_k^{(2)} = \square_k g_{nk}^{(2)} \phi_k^{(2)}$, $\phi_k^{(2)} \in H^1$, and $(k, 2) \in \Lambda_{1,0} \cup \Lambda_{\infty,1} \dots$; by induction, we obtain (62).

By the orthogonal condition (45), following the proof in [14], we can obtain that for fix k and for all $j \neq j'$,

$$\begin{aligned} \lim_{n \rightarrow \infty} \langle g_{nk}^{(j)} \phi_k^{(j)}, g_{nk}^{(j')} \phi_k^{(j')} \rangle_{\dot{H}^1(\mathbb{R}^2)} \\ = \lim_{n \rightarrow \infty} \langle \square_k g_{nk}^{(j)} \phi_k^{(j)}, \square_k g_{nk}^{(j')} \phi_k^{(j')} \rangle_{\dot{H}^1(\mathbb{R}^2)} = 0, \\ \lim_{n \rightarrow \infty} \langle g_{nk}^{(j)} \phi_k^{(j)}, w_{nk}^{(l_k)} \rangle_{\dot{H}^1(\mathbb{R}^2)} \\ = \lim_{n \rightarrow \infty} \langle \square_k g_{nk}^{(j)} \phi_k^{(j)}, \square_k w_{nk}^{(l_k)} \rangle_{\dot{H}^1(\mathbb{R}^2)} = 0, \end{aligned} \quad (85)$$

(63) were proved.

Step 2. For arbitrary $(k_1, j_1), (k_2, j_2) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}$, we define $(k_1, j_1) \sim (k_2, j_2)$ if the orthogonal condition (45) is NOT true for any subsequence; that is,

$$\limsup_{n \rightarrow \infty} \left(\frac{\lambda_{nk_1}^{(j_1)}}{\lambda_{nk_2}^{(j_2)}} + \frac{\lambda_{nk_2}^{(j_2)}}{\lambda_{nk_1}^{(j_1)}} + \lambda_{nk_2}^{(j_2)} \lambda_{nk_1}^{(j_1)} \left| \xi_{nk_2}^{(j_2)} - \xi_{nk_1}^{(j_1)} \right|^2 \right. \\ \left. + \frac{\left| x_{nk_2}^{(j_2)} - x_{nk_1}^{(j_1)} \right|^2}{\lambda_{nk_2}^{(j_2)} \lambda_{nk_1}^{(j_1)}} + \frac{\left| t_{nk_2}^{(j_2)} (\lambda_{nk_2}^{(j_2)})^2 - t_{nk_1}^{(j_1)} (\lambda_{nk_1}^{(j_1)})^2 \right|}{\lambda_{nk_2}^{(j_2)} \lambda_{nk_1}^{(j_1)}} \right) \\ < \infty. \quad (86)$$

By the definition above, if $(k_1, j_1) \sim (k_2, j_2)$, we have

$$\lambda_{nk_1}^{(j_1)} \sim \lambda_{nk_2}^{(j_2)}, \quad \lambda_{nk_1}^{(j_1)} \xi_{nk_1}^{(j_1)} \sim \lambda_{nk_2}^{(j_2)} \xi_{nk_2}^{(j_2)}, \\ \frac{x_{nk_1}^{(j_1)}}{\lambda_{nk_1}^{(j_1)}} \sim \frac{x_{nk_2}^{(j_2)}}{\lambda_{nk_2}^{(j_2)}}, \quad t_{nk_1}^{(j_1)} \sim t_{nk_2}^{(j_2)}. \quad (87)$$

Note that

$$g_{\theta, \xi_0, x_0, \lambda, t_0} f(x) := \frac{1}{\lambda} e^{i\theta} e^{i(x/\lambda) \cdot \lambda \xi_0} (e^{it_0 \Delta} f) \left(\frac{x}{\lambda} - \frac{x_0}{\lambda} \right). \quad (88)$$

By Remark 15, we can put these two profiles together as one profile. Then, by denoting $(\Lambda_{1,0} \cup \Lambda_{\infty,1})/\sim$ as $\{1, 2, \dots, J\}$, we can obtain the sequence ϕ_α^J , $\alpha = 1, 2, \dots, J$; and (52)–(56) were proved.

Specially, since $\phi_k^{(j)} \in H^1$ for each $(k, j) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}$, we have

$$\sum_{(k,j) \in \Lambda_{1,0} \cup \Lambda_{\infty,1}} \phi_k^{(j)} \in H^1 \quad (89)$$

for fixed N and l_k , and hence $\phi_\alpha^J \in H^1(\mathbb{R}^2)$ for any fixed J and $1 \leq \alpha \leq J$.

Step 3. We prove (57) now. By (56), we only need to prove that for all $m \in \mathbb{N}$, $m \geq 2$,

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \left\{ \|u_n\|_{L^{2m}(\mathbb{R}^2)}^{2m} - \sum_{\alpha=1}^J \|g_{n\alpha} \phi_\alpha^J\|_{L^{2m}(\mathbb{R}^2)}^{2m} \right. \\ \left. - \|R(n, J)\|_{L^{2m}(\mathbb{R}^2)}^{2m} \right\} = 0. \quad (90)$$

As

$$\|R(n, J)\|_{L^{2m}(\mathbb{R}^2)} \leq \|R(n, J)\|_{L^4(\mathbb{R}^2)}^{1/(m-1)} \|R(n, J)\|_{H^1(\mathbb{R}^2)}^{(m-2)/(m-1)} \quad (91)$$

and for $1/4 < \theta < 1/2$,

$$\|R(n, J)\|_{L^4(\mathbb{R}^2)} \\ \leq \|e^{it\Delta} R(n, J)\|_{L_t^\infty L_x^4(\mathbb{R} \times \mathbb{R}^2)} \\ \leq \|e^{it\Delta} R(n, J)\|_{H_4^\theta L_x^4(\mathbb{R} \times \mathbb{R}^2)} \\ = \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1-2\theta} \|e^{it\Delta} R(n, J)\|_{H_4^{1/2} L_x^4(\mathbb{R} \times \mathbb{R}^2)}^{2\theta} \\ = \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1-2\theta} \|e^{it\Delta} R(n, J)\|_{L_t^4 H_4^1(\mathbb{R} \times \mathbb{R}^2)}^{2\theta} \\ \leq \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1-2\theta}. \quad (92)$$

By (54), we have

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|R(n, J)\|_{L^{2m}(\mathbb{R}^2)}^{2m} = 0. \quad (93)$$

We separate the set $1 \leq \alpha \leq J$ into two subsets:

$$\Lambda_1 = \{\alpha \mid 1 \leq \alpha \leq J, \lambda_{n\alpha} \equiv 1, \xi_{n\alpha} \equiv 0\}, \\ \Lambda_\infty = \{\alpha \mid 1 \leq \alpha \leq J, \lambda_{n\alpha} \rightarrow \infty\}. \quad (94)$$

When $\alpha \in \Lambda_\infty$,

$$\lim_{n \rightarrow \infty} \|g_{n\alpha} \phi_\alpha^J\|_{L^{2m}(\mathbb{R}^2)} = \lim_{n \rightarrow \infty} (\lambda_{n\alpha})^{-1+1/m} \|e^{it_{n\alpha} \Delta} \phi_\alpha^J\|_{L^{2m}(\mathbb{R}^2)} \\ \leq \lim_{n \rightarrow \infty} (\lambda_{n\alpha})^{-1+1/m} \|\phi_\alpha^J\|_{H^1(\mathbb{R}^2)} = 0. \quad (95)$$

Hence, in order to prove (90), one only needs to prove

$$\lim_{n \rightarrow \infty} \left\{ \left\| \sum_{\alpha \in \Lambda_1} g_{n\alpha} \phi_\alpha^J \right\|_{L^{2m}(\mathbb{R}^2)}^{2m} - \sum_{\alpha \in \Lambda_1} \|g_{n\alpha} \phi_\alpha^J\|_{L^{2m}(\mathbb{R}^2)}^{2m} \right\} = 0. \quad (96)$$

If $\alpha \in \Lambda_1$ and $t_{n\alpha} \rightarrow \infty$, for a function $\tilde{\phi}_\alpha^J \in \dot{H}^{1/2} \cap L^{4/3}$, we have

$$\|g_{n\alpha} \phi_\alpha^J\|_{L^4(\mathbb{R}^2)} \leq \|g_{n\alpha} \phi_\alpha^J - g_{n\alpha} \tilde{\phi}_\alpha^J\|_{L^4(\mathbb{R}^2)} + \|g_{n\alpha} \tilde{\phi}_\alpha^J\|_{L^4(\mathbb{R}^2)} \\ \leq \|\phi_\alpha^J - \tilde{\phi}_\alpha^J\|_{\dot{H}^{1/2}(\mathbb{R}^2)} + |t_{n\alpha}|^{-1/2} \|\tilde{\phi}_\alpha^J\|_{L^{4/3}(\mathbb{R}^2)}. \quad (97)$$

By approximating ϕ_α^J by $\tilde{\phi}_\alpha^J \in C_c^\infty$ in $\dot{H}^{1/2}$ and sending $n \rightarrow \infty$, we have $\|g_{n\alpha} \phi_\alpha^J\|_{L^4(\mathbb{R}^2)} \rightarrow 0$. Note that $g_{n\alpha} \phi_\alpha^J \in H^1$; we obtain $\|g_{n\alpha} \phi_\alpha^J\|_{L^{2m}(\mathbb{R}^2)} \rightarrow 0$ for all $m \geq 2$.

If $\alpha \in \Lambda_1$ and $t_{n\alpha} \equiv 0$, we have orthogonal condition $|x_{n\alpha} - x_{n\alpha'}| \rightarrow \infty$ for any $\alpha \neq \alpha'$. Thus,

$$\begin{aligned} & \lim_{n \rightarrow \infty} \left\{ \left\| \sum_{\alpha \in \Lambda_1, t_{n\alpha} \equiv 0} g_{n\alpha} \phi_\alpha^J \right\|_{L^{2m}(\mathbb{R}^2)}^{2m} - \sum_{\alpha \in \Lambda_1, t_{n\alpha} \equiv 0} \left\| g_{n\alpha} \phi_\alpha^J \right\|_{L^{2m}(\mathbb{R}^2)}^{2m} \right\} \\ &= \lim_{n \rightarrow \infty} \left\{ \left\| \sum_{\alpha \in \Lambda_1, t_{n\alpha} \equiv 0} e^{i\theta_{n\alpha}} \phi_\alpha^J(\cdot - x_{n\alpha}) \right\|_{L^{2m}(\mathbb{R}^2)}^{2m} \right. \\ & \quad \left. - \sum_{\alpha \in \Lambda_1, t_{n\alpha} \equiv 0} \left\| \phi_\alpha^J(\cdot - x_{n\alpha}) \right\|_{L^{2m}(\mathbb{R}^2)}^{2m} \right\} = 0. \end{aligned} \quad (98)$$

(96) holds and then (57) was proved. $\square$

5. The Proof of Theorem 3

Let u be a solution of (5), $H(u_0) \leq 1$; by Strichartz estimate and (21),

$$\begin{aligned} & \|u\|_{L_{t,x}^4 \cap L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}(\mathbb{R} \times \mathbb{R}^2)} \\ & \leq C \|u_0\|_{L^2(\mathbb{R}^2)} + C \|u\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^3 \\ & \quad + C(E) \|u\|_{L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}(\mathbb{R} \times \mathbb{R}^2)} \|u\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{8\epsilon} \quad (99) \\ & \leq C \|u_0\|_{L^2(\mathbb{R}^2)} + C \|u\|_{L_{t,x}^4 \cap L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}(\mathbb{R} \times \mathbb{R}^2)}^3 \\ & \quad + C(E) \|u\|_{L_{t,x}^4 \cap L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}(\mathbb{R} \times \mathbb{R}^2)}^{1+8\epsilon}. \end{aligned}$$

When $\|u_0\|_{L^2(\mathbb{R}^2)} \ll 1$, by standard continuity argument, we have

$$\|u\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} \leq C \|u_0\|_{L^2(\mathbb{R}^2)} < \infty. \quad (100)$$

Hence, if $M(u) \ll 1$, then $\|u\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} < \infty$. In particular, we have scattering in both directions.

For any mass $m \geq 0$, we define

$$A(m) := \sup \{S(u) : u \text{ is the global solution of (5)},$$

$$M(u) \leq m, H(u) \leq 1\}. \quad (101)$$

Then $A : [0, +\infty) \rightarrow [0, +\infty]$ is a monotone increasing function of m . As A is left-continuous and finite for small m , there must exist a unique critical mass $m_0 \in (0, +\infty]$ such that $A(m)$ is finite for all $m < m_0$ but infinite for all $m \geq m_0$.

To prove Theorem 3, one only needs to prove that the critical mass m_0 is infinite. We will prove that by contradiction.

Proposition 17. Suppose that the critical mass m_0 is finite. Let $u_n : \mathbb{R} \times \mathbb{R}^2 \rightarrow \mathbb{C}$ for $n = 1, 2, \dots$ be a sequence of solutions and let $t_n \in \mathbb{R}$ be a sequence of times such that $H(u_n) \leq 1$, $\limsup_{n \rightarrow \infty} M(u_n) = m_0$, and

$$\lim_{n \rightarrow \infty} S_{\geq t_n}(u_n) = \lim_{n \rightarrow \infty} S_{\leq t_n}(u_n) = +\infty. \quad (102)$$

Then there exists a sequence of $x_n = x_n(t_n) \in \mathbb{R}^2$ such that $u_n(t_n, x + x_n)$ has a subsequence which converges strongly into $L_x^2(\mathbb{R}^2)$. Especially, the Hamiltonian of the limiting function is not greater than 1.

Proof. We can take $t_n = 0$ for all n by translating u_n in time. Thus,

$$\lim_{n \rightarrow \infty} S_{\geq 0}(u_n) = \lim_{n \rightarrow \infty} S_{\leq 0}(u_n) = +\infty. \quad (103)$$

By Lemma 16, up to a subsequence if necessary, we have

$$u_n(0) = \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} g_{n\alpha} \phi_\alpha^J + R(n, J), \quad (104)$$

where Λ_1 and Λ_∞ were defined by (94). Suppose that

$$g_{n\alpha} = h_{n\alpha} e^{it_{n\alpha}\Delta}, \quad (105)$$

where $t_{n\alpha} \in \mathbb{R}$ and $h_{n\alpha} \in G$. By (55),

$$\sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} M(\phi_\alpha^J) \leq \limsup_{n \rightarrow \infty} M(u_n(0)) \leq m_0. \quad (106)$$

Hence,

$$\sup_{\alpha \in \Lambda_1 \cup \Lambda_\infty} M(\phi_\alpha^J) \leq m_0. \quad (107)$$

Suppose

$$\sup_{\alpha \in \Lambda_1 \cup \Lambda_\infty} M(\phi_\alpha^J) \leq m_0 - \sigma. \quad (108)$$

For some $\sigma > 0$, we will prove that this leads to a contradiction: defining the nonlinear profile $v_\alpha^J : \mathbb{R} \times \mathbb{R}^2 \rightarrow \mathbb{C}$ as follows.

(i) When $\alpha \in \Lambda_1$,

- (1) if $t_{n\alpha} \equiv 0$, we define v_α^J to be the global solution of (5) with initial data $v_\alpha^J(0) = \phi_\alpha^J$;
- (2) if $t_{n\alpha} \rightarrow +\infty$, we define v_α^J to be the global solution of (5) which scatters to $e^{it\Delta}\phi_\alpha^J$ in H^1 when $t \rightarrow +\infty$;
- (3) if $t_{n\alpha} \rightarrow -\infty$, we define v_α^J to be the global solution of (5) which scatters to $e^{it\Delta}\phi_\alpha^J$ in H^1 when $t \rightarrow -\infty$.

(ii) When $\alpha \in \Lambda_\infty$,

- (1) if $t_{n\alpha} \equiv 0$, we define v_α^J to be the global solution of $iu_t + \Delta u = |u|^2 u$ with initial data $v_\alpha^J(0) = \phi_\alpha^J$;
- (2) if $t_{n\alpha} \rightarrow +\infty$, we define v_α^J to be the global solution of $iu_t + \Delta u = |u|^2 u$ which scatters to $e^{it\Delta}\phi_\alpha^J$ in H^1 when $t \rightarrow +\infty$;
- (3) if $t_{n\alpha} \rightarrow -\infty$, we define v_α^J to be the global solution of $iu_t + \Delta u = |u|^2 u$ which scatters to $e^{it\Delta}\phi_\alpha^J$ in H^1 when $t \rightarrow -\infty$.

For each $\alpha \in \Lambda_1$, v_α^J is well defined by the definition of $A(m)$, (100), and (108). For each $\alpha \in \Lambda_\infty$, v_α^J is well defined by the scattering of cubic Schrödinger equation in $L^2(\mathbb{R}^2)$ (see [7, 8]) and the same analysis as in Lemma 9.

Now, we define

$$\tilde{u}_n = \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} T_{h_{n\alpha}} [v_\alpha^J(\cdot + t_{n\alpha})] + e^{it\Delta} R(n, J) \quad (109)$$

for $n, J = 1, 2, \dots$, and then we have the following two lemmas.

Lemma 18

$$\begin{aligned} \lim_{n \rightarrow \infty} M(\tilde{u}_n(0) - u_n(0)) &= 0, \\ \lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|\tilde{u}_n\|_{\dot{H}^1(\mathbb{R}^2)}^2 &\leq 1. \end{aligned} \quad (110)$$

Proof. Since

$$\begin{aligned} &\lim_{n \rightarrow \infty} M(\tilde{u}_n(0) - u_n(0)) \\ &\leq \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} \lim_{n \rightarrow \infty} M(T_{h_{n\alpha}} [v_\alpha^J(\cdot + t_{n\alpha})](0) - g_{n\alpha} \phi_\alpha^J) \\ &= \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} \lim_{n \rightarrow \infty} M(h_{n\alpha} [v_\alpha^J(t_{n\alpha})] - h_{n\alpha} e^{it_{n\alpha}\Delta} \phi_\alpha^J) \\ &= \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} \lim_{n \rightarrow \infty} M(v_\alpha^J(t_{n\alpha}) - e^{it_{n\alpha}\Delta} \phi_\alpha^J) = 0, \end{aligned} \quad (111)$$

we get the first equality.

Note that $T_{h_{n\alpha}}$ maps the solutions of $iu_t + \Delta u = |u|^2 u$ from one to another; using (57) and energy conservation, we have

$$\begin{aligned} &\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|\tilde{u}_n\|_{\dot{H}^1(\mathbb{R}^2)}^2 \\ &\leq \lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \left\{ \sum_{\alpha=1}^J \|T_{h_{n\alpha}} [v_\alpha^J(\cdot + t_{n\alpha})]\|_{\dot{H}^1(\mathbb{R}^2)}^2 \right. \\ &\quad \left. + \|R(n, J)\|_{\dot{H}^1(\mathbb{R}^2)}^2 \right\} \\ &\leq \lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \left\{ \sum_{\alpha=1}^J H(g_{n\alpha} \phi_\alpha^J) + H(R(n, J)) \right\} \\ &\leq \lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} H(u_n) \leq 1. \end{aligned} \quad (112)$$

Lemma 19. *If*

$$\begin{aligned} &\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|\tilde{u}_n\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} < \infty, \\ &\|v_\alpha^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} < \infty \quad (\forall \alpha), \end{aligned} \quad (113)$$

then

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \left\| \int_0^t e^{i(t-\tau)\Delta} ((i\partial_t + \Delta) \tilde{u}_n - f(\tilde{u}_n))(\tau) d\tau \right\|_X = 0, \quad (114)$$

where $X = L_{t,x}^4 \cap L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}(\mathbb{R} \times \mathbb{R}^2)$.

Proof. Denote

$$v_{n\alpha}^J = T_{h_{n\alpha}} [v_\alpha^J(\cdot + t_{n\alpha})]. \quad (115)$$

By the definition of $\tilde{u}_n$, we have

$$\begin{aligned} \tilde{u}_n &= \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} v_{n\alpha}^J + e^{it\Delta} R(n, J), \\ (i\partial_t + \Delta) \tilde{u}_n &= \sum_{\alpha \in \Lambda_1} f(v_{n\alpha}^J) + \sum_{\alpha \in \Lambda_\infty} |v_{n\alpha}^J|^2 v_{n\alpha}^J. \end{aligned} \quad (116)$$

Thus, by triangle inequality, it suffices to show that

$$\begin{aligned} &\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \left\| \int_0^t e^{i(t-\tau)\Delta} (f(\tilde{u}_n - e^{it\Delta} R(n, J)) \right. \\ &\quad \left. - f(\tilde{u}_n))(\tau) d\tau \right\|_X = 0, \end{aligned} \quad (117)$$

$$\begin{aligned} &\lim_{n \rightarrow \infty} \left\| \int_0^t e^{i(t-\tau)\Delta} \left(f \left(\sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} v_{n\alpha}^J \right) \right. \right. \\ &\quad \left. \left. - \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} f(v_{n\alpha}^J) \right) (\tau) d\tau \right\|_X = 0, \end{aligned} \quad (118)$$

$$\begin{aligned} &\lim_{n \rightarrow \infty} \left\| \int_0^t e^{i(t-\tau)\Delta} \sum_{\alpha \in \Lambda_\infty} (f(v_{n\alpha}^J) - |v_{n\alpha}^J|^2 v_{n\alpha}^J) (\tau) d\tau \right\|_X \\ &= 0. \end{aligned} \quad (119)$$

By Lemma 18, let J and n be sufficiently large; we have

$$\|\tilde{u}_n\|_{\dot{H}^1(\mathbb{R}^2)}^2 \leq \frac{1}{1-4\epsilon}, \quad \|\tilde{u}_n - e^{it\Delta} R(n, J)\|_{\dot{H}^1(\mathbb{R}^2)}^2 \leq \frac{1}{1-4\epsilon}. \quad (120)$$

Since we have supposed $0 < \lambda < 4(1-4\epsilon)^2\pi$, Lemma 6 can also be used here. Then by the same estimates as (31),

$$\begin{aligned} &\left\| \int_0^t e^{i(t-\tau)\Delta} (f(\tilde{u}_n - e^{it\Delta} R(n, J)) - f(\tilde{u}_n))(\tau) d\tau \right\|_X \\ &\leq \left(\|\tilde{u}_n\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^2 + \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^2 \right) \\ &\quad \times \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} \\ &\quad + \left(\|\tilde{u}_n\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{8\epsilon} + \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{8\epsilon} \right) \\ &\quad \times \|e^{it\Delta} R(n, J)\|_{L_t^{2/(1-2\epsilon)} L_x^{1/\epsilon}(\mathbb{R} \times \mathbb{R}^2)}. \end{aligned} \quad (121)$$

By (54),

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} = 0. \quad (122)$$

As

$$\begin{aligned} & \|e^{it\Delta} R(n, J)\|_{L_t^{2/(1-2\varepsilon)} L_x^{1/\varepsilon}(\mathbb{R} \times \mathbb{R}^2)} \\ & \leq \|e^{it\Delta} R(n, J)\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{2\varepsilon/(1-2\varepsilon)} \|e^{it\Delta} R(n, J)\|_{L_t^{2/(1-\varepsilon)} L_x^{2/\varepsilon}(\mathbb{R} \times \mathbb{R}^2)}^{(1-4\varepsilon)/(1-2\varepsilon)}, \\ & \|e^{it\Delta} R(n, J)\|_{L_t^{2/(1-\varepsilon)} L_x^{2/\varepsilon}(\mathbb{R} \times \mathbb{R}^2)} \leq \|R(n, J)\|_{L_x^2(\mathbb{R}^2)} < \infty, \end{aligned} \quad (123)$$

we have

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|e^{it\Delta} R(n, J)\|_{L_t^{2/(1-2\varepsilon)} L_x^{1/\varepsilon}(\mathbb{R} \times \mathbb{R}^2)} = 0. \quad (124)$$

By $\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|\tilde{u}_n\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} < \infty$, (117) was obtained.

Using Strichartz estimate,

$$\begin{aligned} & \left\| \int_0^t e^{i(t-\tau)\Delta} \left(f(v_{n\alpha}^J) - |v_{n\alpha}^J|^2 v_{n\alpha}^J \right) (\tau) d\tau \right\|_X \\ & \leq \left\| \sum_{m=2}^{\infty} \frac{|v_{n\alpha}^J|^{2m} v_{n\alpha}^J}{m!} \right\|_{L_{t,x}^{4/3}(\mathbb{R} \times \mathbb{R}^2)} \\ & \leq \sum_{m=2}^{\infty} \frac{1}{m!} \|v_{n\alpha}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} \|v_{n\alpha}^J\|_{L_{t,x}^{4m}(\mathbb{R} \times \mathbb{R}^2)}^{2m} \\ & \leq \sum_{m=2}^{\infty} \frac{1}{m!} (\lambda_{n\alpha})^{-2(m-1)} \|v_{n\alpha}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} \|v_{n\alpha}^J\|_{L_{t,x}^{4m}(\mathbb{R} \times \mathbb{R}^2)}^{2m}. \end{aligned} \quad (125)$$

By Lemma 8,

$$\|v_{n\alpha}^J\|_{L_{t,x}^{4m}(\mathbb{R} \times \mathbb{R}^2)} \leq C m^{1/2+1/2m} \|v_{n\alpha}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1/m} \|v_{n\alpha}^J\|_{L^\infty(\mathbb{R}, \dot{H}^1)}^{1-1/m}. \quad (126)$$

Note that $\|v_{n\alpha}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} < \infty$. Equation (119) was obtained.

To prove Lemma 19, it is only left to prove (118). Note that

$$\left| f\left(\sum_{\alpha=1}^J z_\alpha\right) - \sum_{\alpha=1}^J f(z_\alpha) \right| \leq \sum_{\alpha \neq \alpha'} |z_\alpha| \left| e^{\lambda|z_{\alpha'}|^2} - 1 \right|. \quad (127)$$

By (57), we have

$$\left\| \int_0^t e^{i(t-\tau)\Delta} \left(f\left(\sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} v_{n\alpha}^J\right) - \sum_{\alpha \in \Lambda_1 \cup \Lambda_\infty} f(v_{n\alpha}^J) \right) (\tau) d\tau \right\|_X$$

$$\begin{aligned} & \lesssim \sum_{\alpha \neq \alpha'} \left(\|v_{n\alpha}^J v_{n\alpha'}^J\|_{L_{t,x}^{4/3}(\mathbb{R} \times \mathbb{R}^2)}^2 \right. \\ & \quad \left. + \|v_{n\alpha}^J \left(e^{\lambda|v_{n\alpha'}^J|^2} - \lambda|v_{n\alpha'}^J|^2 - 1 \right)\|_{L_t^{2/(1+2\varepsilon)} L_x^{1/(1-\varepsilon)}(\mathbb{R} \times \mathbb{R}^2)} \right) \\ & \lesssim \sum_{\alpha \neq \alpha'} \|v_{n\alpha}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1/2} \|v_{n\alpha'}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1/2} \|v_{n\alpha'}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)} \\ & \quad + \sum_{\alpha \neq \alpha'} \|v_{n\alpha}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{\varepsilon/(1+\varepsilon)} \|v_{n\alpha'}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{1/(1+\varepsilon)} \|v_{n\alpha'}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{4\varepsilon(1+\varepsilon)/(1-2\varepsilon)} \\ & \quad \times \|v_{n\alpha}^J\|_{L_t^{2/(1-2\varepsilon-4\varepsilon^2)/(1-2\varepsilon-6\varepsilon^2)} L_x^{(1-2\varepsilon-4\varepsilon^2)/\varepsilon^2}(\mathbb{R} \times \mathbb{R}^2)}^{(1-2\varepsilon-4\varepsilon^2)/(1-2\varepsilon)} \\ & \quad \times \|v_{n\alpha'}^J\|_{L_{t,x}^4(\mathbb{R} \times \mathbb{R}^2)}^{-4\varepsilon/(1-2\varepsilon)} \\ & \quad \times \left(e^{\lambda|v_{n\alpha'}^J|^2} - \lambda|v_{n\alpha'}^J|^2 - 1 \right) \Big\|_{L_t^\infty L_x^{(1-2\varepsilon)/(1-4\varepsilon)}(\mathbb{R} \times \mathbb{R}^2)}. \end{aligned} \quad (128)$$

By (49), we immediately obtain (118).

If

$$\sup_{\alpha \in \Lambda_1 \cup \Lambda_\infty} M(\phi_\alpha) \leq m_0 - \sigma \quad (129)$$

for some $\sigma > 0$, by the definition of $A(m)$ and (100), we have

$$A(m) \leq Bm \quad \forall m \in [0, m_0 - \sigma], \quad (130)$$

where $B = B(\sigma) \in (0, +\infty)$. Then v_α^J satisfies

$$M(v_\alpha^J) = M(\phi_\alpha^J) \leq m_0 - \sigma, \quad (131)$$

$$S(v_\alpha^J) \leq A(M(\phi_\alpha^J)) \leq BM(\phi_\alpha^J). \quad (132)$$

By (50), (106), and (132), we have

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} S(\tilde{u}_n) \leq Bm_0. \quad (133)$$

Using Lemmas 18 and 19, we have

$$M(\tilde{u}_n(0) - u_n(0)) \leq \delta, \quad S(\tilde{u}_n) \leq 2Bm_0,$$

$$\begin{aligned} & \left\| \int_0^t e^{i(t-\tau)\Delta} ((i\partial_t + \Delta) \tilde{u}_n - f(\tilde{u}_n)) (\tau) d\tau \right\|_{L_{t,x}^4 \cap L_t^{2/(1-2\varepsilon)} L_x^{1/\varepsilon}(\mathbb{R} \times \mathbb{R}^2)} \\ & \leq \delta, \end{aligned} \quad (134)$$

for $\delta > 0$ sufficiently small, $J = J(\delta)$, and $n = n(J, \delta)$ sufficiently large. By Lemma 10, we obtain that $S(u_n) \leq 3Bm_0$ which contradicts (103). Thus, (129) fails for all $\sigma > 0$, and then

$$\sup_{\alpha \in \Lambda_1 \cup \Lambda_\infty} M(\phi_\alpha) = m_0. \quad (135)$$

Comparing this with (106), we have

$$u_n(0) = h_n e^{it_n \Delta} \phi + R_n \quad (136)$$

with t_n converging to $\pm\infty$ or $t_n \equiv 0$, $h_n \in G$, $M(\phi) = m_0$, and

$$M(R_n) \rightarrow 0, \quad S(e^{it\Delta} R_n) \rightarrow 0 \quad \text{as } n \rightarrow \infty. \quad (137)$$

Specially, the parameters λ_n, ξ_n of h_n must satisfy

$$“\lambda_n \equiv 1 \text{ and } \xi_n \equiv 0,” \quad \text{or} \quad “\lambda_n \rightarrow \infty.” \quad (138)$$

If $\lambda_n \rightarrow \infty$, similar to the former case, we can define the approximate solution

$$\tilde{u}_n = T_{h_n}[\nu(\cdot + t_n)] + e^{it\Delta} R_n. \quad (139)$$

By the scattering of cubic Schrödinger equation, we have $S(\nu) < \infty$ and $\lim_{n \rightarrow \infty} S(\tilde{u}_n) < \infty$. By Lemmas 18, 19, and 10, for n sufficiently large, we obtain $S(u_n) < \infty$ which contradicts (103).

If $\lambda_n \equiv 1$, $\xi_n \equiv 0$, and $t_n \rightarrow +\infty$, by Strichartz estimate and monotone convergence, we have

$$\lim_{n \rightarrow \infty} S_{\geq 0}(e^{it\Delta} e^{it_n\Delta} \phi) = 0. \quad (140)$$

Thus,

$$\lim_{n \rightarrow \infty} S_{\geq 0}(e^{it\Delta} h_n e^{it_n\Delta} \phi) = \lim_{n \rightarrow \infty} S_{\geq 0}(e^{it\Delta} e^{it_n\Delta} \phi) = 0. \quad (141)$$

Since $\lim_{n \rightarrow \infty} S(e^{it\Delta} R_n) = 0$, we can see from (136) that

$$\lim_{n \rightarrow \infty} S_{\geq 0}(e^{it\Delta} u_n(0)) = 0. \quad (142)$$

Hence,

$$\begin{aligned} \lim_{n \rightarrow \infty} \left\| \int_0^t e^{i(t-\tau)\Delta} f(e^{it\Delta} u_n(0))(\tau) d\tau \right\|_{L^4_{t,x} \cap L^{2/(1-2\epsilon)}_t L^{1/\epsilon}_x([0, \infty) \times \mathbb{R}^2)} \\ = 0. \end{aligned} \quad (143)$$

By Lemma 10 (with $e^{it\Delta} u_n(0)$ as the approximate solution and $u_n(0)$ as the initial data), we have

$$\lim_{n \rightarrow \infty} S_{\geq 0}(u_n) = 0, \quad (144)$$

which contradicts one of the estimates in (103).

If $\lambda_n \equiv 1$, $\xi_n \equiv 0$, and $t_n \rightarrow -\infty$, the argument is the same and we can obtain a contradiction by using the other half of (103).

Now, the only case left is $\lambda_n \equiv 1$, $\xi_n \equiv 0$, and $t_n \equiv 0$. In this case, we have

$$M(u_n(0) - h_n\phi) = M(R_n) \rightarrow 0 \quad \text{as } n \rightarrow \infty. \quad (145)$$

Thus, $(h_n)^{-1} u_n(0) = e^{i\theta_n} u_n(0, x + x_n)$ converges to ϕ in $L^2_x(\mathbb{R}^2)$. Since $H(\phi) = H(g_n\phi) \leq 1$, after passing to a subsequence if necessary and rotating ϕ , the desired result follows. $\square$

Let $\{u_n\}$ be the sequence given in Proposition 17 and satisfy $M(u_n) \leq m_0$ and suppose that $u_n(0, x + x_n)$ converges to u_0 strongly in $L^2_x(\mathbb{R}^2)$; then $M(u_0) \leq m_0$ and $H(u_0) \leq 1$.

Let u be the global solution with initial data $u(0) = u_0$; by Lemma 10, we must have

$$S_{\geq 0}(u) = S_{\leq 0}(u) = +\infty. \quad (146)$$

By the definition of m_0 , $M(u_0) \geq m_0$ and hence $M(u_0) = m_0$. Since u is locally in $L^4_{t,x}$, for all $t_n \in \mathbb{R}$, we have

$$S_{\geq t_n}(u) = S_{\leq t_n}(u) = +\infty. \quad (147)$$

Using Proposition 17 for $\{u(t_n)\}$, we have that $u(t_n, x + x(t_n))$ converges into $L^2_x(\mathbb{R}^2)$. By Ascoli-Arzelà Theorem, we have the following.

Proposition 18. *Suppose that the critical mass m_0 is finite. Then there exists a global solution u with mass m_0 , and for every $\eta > 0$ there exists $0 < C(\eta) < \infty$ such that*

$$\int_{|x-x(t)| \geq C(\eta)} |u(t, x)|^2 dx + \int_{|\xi-\xi(t)| \geq C(\eta)} |\hat{u}(t, \xi)|^2 d\xi \leq \eta \quad (148)$$

for all $t \in \mathbb{R}$, where the functions $x, \xi : \mathbb{R} \rightarrow \mathbb{R}^2$.

Proposition 19. *The solution described in Proposition 18 does not exist.*

Once we proved Proposition 19, we can say that $m_0 = \infty$ and thus Theorem 3 is true. In order to prove Proposition 19, we need the following two lemmas.

Lemma 20 (see [6, Lemma 5.2]). *Let u be a global solution of (5). Then one has*

$$\iint_{\mathbb{R} \times \mathbb{R}^2} \frac{\langle t \rangle^2 G(u)}{\langle t \rangle^3 + |x|^3} dx dt \leq C(E), \quad (149)$$

where $\langle t \rangle = \sqrt{1 + |t|^2}$.

Lemma 21 (see [6, Lemma 6.2]). *Let u be a global solution of (5). Let B be a compact subset of $\mathbb{R}^2$. Then for any $R > 0$ and $T > 0$, one has*

$$\int_{B(R)} |u(T, x)|^2 dx \geq \int_B |u(0, x)|^2 dx - \frac{C(E)T}{R}, \quad (150)$$

where $B(R) := \{x \in \mathbb{R}^2 \mid \exists y \in B \text{ s.t. } |x - y| < R\}$.

Proof of Proposition 19. By Lemma 21, choosing η sufficiently small,

$$\begin{aligned} & \int_{|x-x(0)| \leq C(\eta)+R|t|} |u(t, x)|^2 dx \\ & \geq \int_{|x-x(0)| \leq C(\eta)} |u(0, x)|^2 dx - \frac{C(E)}{R} \\ & \geq m_0 - \eta - \frac{C(E)}{R}. \end{aligned} \quad (151)$$

By Proposition 18,

$$\int_{|x-x(t)| \leq C(\eta)} |u(t, x)|^2 dx \geq m_0 - \eta. \quad (152)$$

For a fixed large number R , we must have $|x(t) - x(0)| \leq 2C(\eta) + R|t|$. By Lemma 20 and Hölder inequality,

$$\begin{aligned} \infty &> \iint_{\mathbb{R} \times \mathbb{R}^2} \frac{\langle t \rangle^2 |u|^4}{\langle t \rangle^3 + |x|^3} dx dt \\ &\geq \int_{\mathbb{R}} \int_{|x-x(t)| \leq C(\eta)} \frac{\langle t \rangle^2 |u|^4}{\langle t \rangle^3 + |x|^3} dx dt \\ &\geq \int_{\mathbb{R}} \int_{|x-x(t)| \leq C(\eta)} \frac{|u|^4}{\langle t \rangle + 1} dx dt \\ &\geq \int_{\mathbb{R}} \int_{|x-x(t)| \leq C(\eta)} \frac{|u|^2}{\langle t \rangle + 1} dx dt \\ &\geq (m_0 - \eta) \int_{\mathbb{R}} \frac{1}{\langle t \rangle + 1} dx dt = \infty. \end{aligned} \quad (153)$$

This is a contradiction. Proposition 19 was obtained. $\square$

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