

Research Article

The Partial Second Boundary Value Problem of an Anisotropic Parabolic Equation

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Consider an anisotropic parabolic equation with the variable exponents $v_t = \sum_{i=1}^n (b_i(x, t) |v_{x_i}|^{p_i(x)-2} v_{x_i})_{x_i} + f(v, x, t)$, where $b_i(x, t) \in C^1(\overline{Q_T})$, $p_i(x) \in C^1(\overline{\Omega})$, $p_i(x) > 1$, $b_i(x, t) \geq 0$, $f(v, x, t) \geq 0$. If $\{b_i(x, t)\}$ is degenerate on $\Gamma_2 \subset \partial\Omega$, then the second boundary value condition is imposed on the remaining part $\partial\Omega \setminus \Gamma_2$. The uniqueness of weak solution can be proved without the boundary value condition on Γ_2 .

1. Introduction

The evolutionary $p(x)$ -Laplacian equation,

$$u_t = \operatorname{div}(|\nabla u|^{p(x)-2} \nabla u), \quad (x, t) \in Q_T = \Omega \times (0, T), \quad (1)$$

possesses some interesting mechanical properties in the presence of an electromagnetic field [1, 2], the well-posedness of weak solutions to the first initial-boundary value problem of (1) had been researched in [3–7], etc. Here, $\Omega \subset \mathbb{R}^N$ is a bounded domain with the smooth boundary $\partial\Omega$, and $p(x) > 1$ is a $C^1(\overline{\Omega})$ function. Sobolev spaces play an important role in the theory of evolutionary $p(x)$ -Laplacian equation. In recent years, the generalized Orlicz-Lebesgue spaces $L^p(x)$ and the corresponding generalized Orlicz-Sobolev spaces $W^{1,p(x)}$ have attracted more and more attention. The spaces $L^{p(\cdot)}$ are special cases of the generalized Orlicz-Spaces originated by Nakano and developed by Musielak and Orlicz (see [8–10]). We refer to [11–13] for properties of the spaces $L^{p(\cdot)}$ and $W^{k,p(\cdot)}$ such as reflexivity, denseness of smooth functions, and Sobolev type embeddings. The study of these spaces has been stimulated by problems of elasticity, fluid dynamics, calculus of variations, and differential equations with $p(x)$ -growth conditions. Roughly speaking, the interest in variable exponent spaces comes not only from their mathematical curiosity but also from their relevance in many applications such as fluid dynamics, elasticity theory,

differential equations with nonstandard growth conditions, and image restoration. In addition, the study of the weak solution in other spaces such as Orlicz-Morrey space and $\dot{B}_{\infty, \infty}^{-1}$ space is a research problem (see [14–18]).

The so-called anisotropic evolutionary $\vec{p}(x)$ -Laplacian equation,

$$u_t = \sum_{i=1}^N \frac{\partial}{\partial x_i} \left(|u_{x_i}|^{p_i(x)-2} u_{x_i} \right), \quad (2)$$

comes closer to the truth than equation (1) [19–21], where $\vec{p}(x) = (p_1(x), p_2(x), \dots, p_N(x))$. Recently, Zhan et al. [22–24] considered the first initial-boundary value problem to the equation:

$$u_t = \sum_{i=1}^N \frac{\partial}{\partial x_i} \left(b_i(x) |u_{x_i}|^{p_i(x)-2} u_{x_i} \right), \quad (3)$$

where $b_i(x) \in C^1(\overline{\Omega})$ ($i = 1, 2, \dots, N$) satisfies

$$b_i(x) > 0 \quad \text{if } x \in \Omega \quad (4)$$

$$\text{and } b_i(x) = 0 \quad \text{if } x \in \partial\Omega.$$

We have shown that this condition may act as the role of the Dirichlet boundary condition

$$u(x, t) = 0, \quad (x, t) \in \partial\Omega \times (0, T) \quad (5)$$

to assure the stability of weak solutions to (3).

In this paper, we will consider the anisotropic parabolic equation

$$v_t = \sum_{i=1}^n \left(b_i(x, t) |v_{x_i}|^{p_i(x)-2} v_{x_i} \right)_{x_i} + f(v, x, t), \quad (6)$$

$$(x, t) \in Q_T,$$

with the initial value

$$v(x, 0) = v_0(x), \quad x \in \Omega, \quad (7)$$

and with a partial second boundary value condition

$$\frac{\partial v}{\partial n} = 0, \quad (x, t) \in \Gamma_1 \times (0, T), \quad (8)$$

where $b_i(x, t) \in C^1(\overline{Q_T})$, $p_i(x) \in C^1(\overline{\Omega})$, $p_i(x) > 1$, $b_i(x, t) \geq 0$, $f(v(x, t), x, t) \geq 0$ and at least there is a point $x_0 \in \Omega$ such that $f(v(x_0, t), x_0, t) > 0$, Ω is a bounded domain with a smooth boundary $\partial\Omega$. We first assume that

$$\begin{aligned} \partial\Omega &= \Gamma_1 \cup \Gamma_2, \\ \Gamma_1^0 \cap \Gamma_2^0 &= \emptyset, \end{aligned} \quad (9)$$

where Γ_1^0 and Γ_2^0 are the interior of Γ_1 and Γ_2 , which are relatively open subset of $\partial\Omega$. We mainly assume that

$$b_i(x, t) = 0, \quad (x, t) \in \Gamma_2 \times [0, T], \quad (10)$$

denote that $p^+ = \max_{x \in \overline{\Omega}} p(x)$, $p^- = \min_{x \in \overline{\Omega}} p(x)$, for any $p(x) \in C^1(\overline{\Omega})$, and let

$$\begin{aligned} p_0 &= \min_{x \in \overline{\Omega}} \{p_1(x), p_2(x), \dots, p_{n-1}(x), p_n(x)\}, \\ p^0 &= \max_{x \in \overline{\Omega}} \{p_1(x), p_2(x), \dots, p_{n-1}(x), p_n(x)\}, \end{aligned} \quad (11)$$

for any $\vec{p}(x)$. As for the anisotropic function spaces and their applications to anisotropic equations, one can refer to [25, 26] and the references therein.

Definition 1. If a function $v(x, t)$ satisfies that

$$\begin{aligned} v &\in L^\infty(Q_T), \\ \frac{\partial v}{\partial t} &\in L^2(Q_T), \\ b_i(x, t) |v_{x_i}|^{p_i(x)} &\in L^2(0, T; L^1(\Omega)), \end{aligned} \quad (12)$$

for any function $\varphi \in C^1(\overline{Q_T})$, $\varphi_{x_i} \in L^2(0, T; L_{loc}^{p_i(x)}(\Omega))$, $\varphi|_{(x,t) \in \Gamma_2 \times [0, T]} = 0$,

$$\begin{aligned} &\iint_{Q_T} \left[\frac{\partial v}{\partial t} \varphi + \sum_{i=1}^n b_i(x, t) |v_{x_i}|^{p_i(x)-2} v_{x_i} \varphi_{x_i} \right] dx dt \\ &= \iint_{Q_T} f(v, x, t) \varphi dx dt, \end{aligned} \quad (13)$$

and, for any $\phi(x) \in C_0^\infty(\Omega)$,

$$\lim_{t \rightarrow 0} \int_{\Omega} u(x, t) \phi(x) dx = \int_{\Omega} u_0(x) \phi(x) dx, \quad (14)$$

then we say $v(x, t)$ is a weak solution of equation (6) with the initial value (7) and with the partial second value condition (8).

The main results are the following theorems.

Theorem 2. Suppose that $p_0 \geq 2$, $b_i(x, t)$ satisfies (10) and f, g are two C^1 functions satisfying

$$f(v, s, t) \geq g(v). \quad (15)$$

Let

$$v_0(x) \in L^\infty(\Omega), \quad (16)$$

$$b_i(x, 0) v_{0x_i}(x) \in L^2(0, T; L^{p_i(x)}(\Omega)).$$

Then we have a positive constant T_1 such that there exists a weak solution of (6) $v(x, t) \in L^\infty(Q_{T_1})$ with the initial condition (7) and with the partial second boundary value condition (8).

We would like to suggest that, since $f(v(x, t), x, t) \geq 0$ and at least there is a point $x_0 \in \Omega$ such that $f(v(x_0, t), x_0, t) > 0$, then the weak solution $v(x, t)$ generally blows up in a finite time [27]. However, the uniqueness of weak solution still may be true.

Theorem 3. If $b_i(x, t)$ satisfies (10) and f is a C^1 function, and when x is near to Γ_2 ,

$$|b_i(x, t)| \leq cd(x)^{p_i^+ - 1} \quad (17)$$

$v(x, t)$ and $u(x, t)$ are two solutions of (6) with the same partial second boundary value condition

$$\frac{\partial v}{\partial n}(x, t) = \frac{\partial u}{\partial n}(x, t) = 0, \quad (x, t) \in \Gamma_1 \times (0, T), \quad (18)$$

and with the same initial values

$$v_0(x) = u_0(x), \quad x \in \Omega; \quad (19)$$

then

$$v(x, t) = u(x, t), \quad (x, t) \in \Omega \times (0, T). \quad (20)$$

Here $T > T_1$ is the blow-up time of the weak solutions.

At the end of the introduction, we would like to suggest that it is an interesting research problem to study the anisotropic parabolic equation (6) for either (p, q) -Laplacian or p -biharmonic (see [28–30]).

2. The Existence of Solutions

Consider the following asymptotic problem:

$$\begin{aligned} v_t &= \sum_{i=1}^n \left((b_i(x, t) + \varepsilon) (|v_{x_i}|^2 + \varepsilon)^{(p_i(x)-2)/2} v_{x_i} \right)_{x_i} \\ &+ f(v, x, t), \quad (x, t) \in Q_T, \end{aligned} \quad (21)$$

with the initial value

$$v(x, 0) = v_{0\varepsilon}(x) + \varepsilon, \quad x \in \Omega, \tag{22}$$

$$u(x, t) = \varepsilon, \quad (x, t) \in \Gamma_2 \times (0, T), \tag{23}$$

and a partial second boundary value condition

$$\frac{\partial v}{\partial n} = 0, \quad (x, t) \in \Gamma_1 \times (0, T), \tag{24}$$

where $0 < \varepsilon < 1$, $v_{0\varepsilon}(x) \in C_0^\infty(\Omega)$ such that

$$\begin{aligned} \|v_{0\varepsilon}\|_{L^\infty(\Omega)} &\leq \|v_0\|_{L^\infty(\Omega)}, \\ \|v_{0\varepsilon x_i}\|_{L^{p_i(x)}(\Omega)} &\leq \|v_{0x_i}\|_{L^{p_i(x)}(\Omega)}, \end{aligned} \tag{25}$$

and $v_{0\varepsilon}(x) + \varepsilon \rightarrow v_0(x)$ in $W^{1, \vec{p}(\cdot)}(\Omega)$. This is possible; only we assume that $\vec{p}(x) = \{p_i(x)\}$ and every $p_i(x)$ has the logarithmic Hölder continuity [11–13]. Then similar as the usual p -Laplacian equation, one can show that problem (21)–(24) has a classical solution u_ε by the classical theory for parabolic equations, and $v_\varepsilon \geq \varepsilon > 0$ provided that $f(v(x, t), x, t) \geq 0$ and at least there is a point $x_0 \in \Omega$ such that $f(v(x_0, t), x_0, t) > 0$ ([31], Theorem 4.1). We first give a lemma in a similar way as Lemma 2.1 in [32].

Lemma 4. *If $f(v, x, t) \leq g(v)$ and $g(v) \in C^1(0, \infty)$, then there exists a $T_1 < T$ such that*

$$\|v_\varepsilon\|_{L^\infty(Q_{T_1})} \leq c(T_1), \tag{26}$$

where $c = c(T_1)$ represents c dependent on T_1 .

Proof. Let $w(t)$ be the solution of the ordinary differential equation

$$\frac{dw}{dt} = g(w), \tag{27}$$

$$w(0) = \|v_0(x) + 1\|_{L^\infty(\Omega)}. \tag{28}$$

It is well known that there is a local solution $w(t)$, $t \in [0, T_0] \subset [0, T]$, where $T_0 = T_0(\|v_0(x) + 1\|_{L^\infty(\Omega)})$ ([33], Chapter 5). Let $u = v_\varepsilon - w$. One has

$$\begin{aligned} u_t - \sum_{i=1}^n \left((b_i(x, t) + \varepsilon) \left(|(u+w)_{x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} \right. \\ \left. \cdot (u+w)_{x_i} \right)_{x_i} = f(v, x, t) - g(w). \end{aligned} \tag{29}$$

Since $f(u, x, t) \geq g(u)$, one has

$$\begin{aligned} f(v_\varepsilon, x, t) - g(w) &\leq g(v_\varepsilon) - g(w) \\ &= (v_\varepsilon - w) \int_0^1 g'(\theta v_\varepsilon + (1-\theta)w) d\theta \\ &= c_\varepsilon(x, t)(v_\varepsilon - w). \end{aligned} \tag{30}$$

From (29),

$$\begin{aligned} u_t - \sum_{i=1}^n \left((b_i(x, t) + \varepsilon) \left(|(u+w)_{x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} \right. \\ \left. \cdot (u+w)_{x_i} \right)_{x_i} - c_\varepsilon(x, t)u \leq 0, \end{aligned} \tag{31}$$

with the mixed boundary condition

$$\begin{aligned} u = v_\varepsilon - w \leq \varepsilon - \|u_0(x) + 1\|_{L^\infty(\Omega)} \leq 0, \\ (x, t) \in \Gamma_1 \times [0, T_0], \end{aligned} \tag{32}$$

$$\frac{\partial u}{\partial n} = 0, \quad (x, t) \in \Gamma_2 \times [0, T_0], \tag{33}$$

and

$$\begin{aligned} u_0(x) = [(v_{0\varepsilon}(x) + \varepsilon) - \|u_0(x) + 1\|_{L^\infty(\Omega)}] \leq 0, \\ x \in \overline{\Omega}. \end{aligned} \tag{34}$$

By the Hopf maximum principle, one has

$$u(x, t) \leq 0, \quad (x, t) \in \Omega \times [0, T_0]. \tag{35}$$

Hence, for any given $T_1 \in (0, T_0)$, one has

$$\|v_\varepsilon\|_{L^\infty(Q_{T_1})} \leq c(T_1). \tag{36}$$

□

By multiplying (21) with v_ε and integrating over Ω , one has

$$\|(b_i(x, t) + \varepsilon) v_{\varepsilon x_i}\|_{L^{p_i(x)}(\Omega)} \leq c(T_1). \tag{37}$$

Lemma 5. *If $f(v, x, t) \geq g(v)$ and $g(v) \in C^1(0, \infty)$, then there exists a $T_1 < T$ such that*

$$\left\| \frac{\partial v_\varepsilon}{\partial t} \right\|_{L^2(Q_{T_1})} \leq c(T_1). \tag{38}$$

Proof. By multiplying (21) with $\partial u_\varepsilon / \partial t$ and integrating over Q_{T_1} , one has

$$\begin{aligned} v_{\varepsilon t} = - \sum_{i=1}^n \iint_{Q_{T_1}} \left((b_i(x, t) + \varepsilon) \left(|v_{x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} \right. \\ \left. \cdot v_{x_i} \right)_{x_i} \frac{\partial v_{\varepsilon x_i}}{\partial t} dx dt \\ + \iint_{Q_{T_1}} f(v_\varepsilon, x, t) \frac{\partial v_\varepsilon}{\partial t} dx dt. \end{aligned} \tag{39}$$

Since

$$\begin{aligned} \left((b_i(x, t) + \varepsilon) \left(|v_{x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} v_{x_i} \right)_{x_i} \frac{\partial v_{\varepsilon x_i}}{\partial t} \\ = \frac{1}{2} (b_i(x, t) + \varepsilon) \frac{\partial}{\partial t} \int_0^{|v_{\varepsilon x_i}|^2} (s + \varepsilon)^{(p_i(x)-2)/2} ds, \end{aligned} \tag{40}$$

accordingly

$$\begin{aligned}
& \iint_{Q_{T_1}} \left((b_i(x, t) + \varepsilon) \left(|v_{x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} v_{x_i} \right)_{x_i} \frac{\partial v_{\varepsilon x_i}}{\partial t} dx dt \\
&= \frac{1}{2} \iint_{Q_{T_1}} (b_i(x, t) + \varepsilon) \frac{\partial}{\partial t} \int_0^{|v_{x_i}|^2} (s + \varepsilon)^{(p_i(x)-2)/2} ds dx dt \\
&= \frac{1}{2} \\
&\cdot \iint_{Q_{T_1}} \frac{\partial}{\partial t} \left(\int_0^{|v_{x_i}|^2} (b_i(x, t) + \varepsilon) (s + \varepsilon)^{(p_i(x)-2)/2} ds \right) dx dt \\
&- \frac{1}{2} \iint_{Q_{T_1}} b_{it}(x, t) \int_0^{|v_{x_i}|^2} (s + \varepsilon)^{(p_i(x)-2)/2} ds dx dt = \frac{1}{2} \\
&\cdot \int_{\Omega} \int_0^{|v_{\varepsilon x_i}(x, T_1)|^2} (b_i(x, T_1) + \varepsilon) (s + \varepsilon)^{(p_i(x)-2)/2} ds dx - \frac{1}{2} \\
&\cdot \int_{\Omega} \int_0^{|v_{\varepsilon x_i}|^2} (b_i(x, 0) + \varepsilon) (s + \varepsilon)^{(p_i(x)-2)/2} ds dx - \frac{1}{2} \\
&\cdot \iint_{Q_{T_1}} b_{it}(x, t) \int_0^{|v_{x_i}|^2} (s + \varepsilon)^{(p_i(x)-2)/2} ds dx dt.
\end{aligned} \tag{41}$$

Thus by (37), using Young's inequality, one has

$$\begin{aligned}
& \iint_{Q_{T_1}} \left| \frac{\partial v_{\varepsilon}}{\partial t} \right|^2 dx dt = \sum_{i=1}^N \int_{\Omega} \frac{1}{p_i(x)} (b_i(x, T_1) + \varepsilon) \\
&\cdot \left(|v_{x_i}(x, T_1)|^2 + \varepsilon \right)^{p_i(x)/2} dx \\
&- \sum_{i=1}^N \int_{\Omega} \frac{1}{p_i(x)} (b_i(x, 0) + \varepsilon) \left(|v_{0\varepsilon x_i}|^2 + \varepsilon \right)^{p_i(x)/2} dx \\
&- \sum_{i=1}^N \int_{\Omega} \frac{1}{p_i(x)} b_{it}(x, t) \left(|v_{\varepsilon x_i}|^2 + \varepsilon \right)^{p_i(x)/2} dx \\
&+ \iint_{Q_{T_1}} f(v_{\varepsilon}, x, t) \frac{\partial v_{\varepsilon}}{\partial t} dx dt \leq c.
\end{aligned} \tag{42}$$

□

Proof of Theorem 2. By multiplying (21) with v_{ε} and integrating over Q_{T_1} , one has

$$\iint_{Q_{T_1}} \varepsilon^{(p_i(x)-2)/2} (b_i(x, t) + \varepsilon) |v_{x_i}|^2 dx dt \leq c(T_1), \tag{43}$$

$$\begin{aligned}
& \iint_{Q_{T_1}} (b_i(x, t) + \varepsilon) \left(|v_{x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} dx dt \\
&\leq c(T_1).
\end{aligned} \tag{44}$$

By (37), (39), (43), and (44), there is a function v that satisfies

$$\begin{aligned}
v_{\varepsilon} &\longrightarrow v, \quad \text{a.e. } Q_{T_1}, \\
f(v_{\varepsilon}, x, t) &\longrightarrow f(v, x, t), \quad \text{a.e. } Q_{T_1},
\end{aligned}$$

$$\begin{aligned}
b_i(x, t)^{1/p_i(x)} v_{\varepsilon x_i} &\rightharpoonup b_i(x, t)^{1/p_i(x)} v_{x_i}, \\
&\text{in } L^1(0, T; L^{p_i(x)}(\Omega))
\end{aligned}$$

$$\frac{\partial v_{\varepsilon}}{\partial t} \rightharpoonup \frac{\partial v}{\partial t}, \quad \text{in } L^2(Q_{T_1}). \tag{45}$$

In addition, if one notices that

$$\begin{aligned}
& \left| (b_i(x, t) + \varepsilon)^{(p_i(x)-1)/p_i(x)} \left(|v_{\varepsilon x_i}|^2 + \varepsilon \right)^{p_i(x)/(p_i(x)-1)} \right. \\
&\leq (b_i(x, t) + \varepsilon) \left(|v_{\varepsilon x_i}|^2 + \varepsilon \right)^{p_i(x)(p_i(x)-2)/2(p_i(x)-1)} \\
&\cdot |v_{\varepsilon x_i}|^{p_i(x)/(p_i(x)-1)} \leq c \left((b_i(x, t) + \varepsilon) |v_{\varepsilon x_i}|^{p_i(x)} \right. \\
&+ \varepsilon^{p_i(x)(p_i(x)-2)/2(p_i(x)-1)} (b_i(x, t) + \varepsilon) \\
&\cdot |v_{\varepsilon x_i}|^{p_i(x)/(p_i(x)-1)} \left. \right);
\end{aligned} \tag{46}$$

by (43) and (44) and by that $p_i(x)/(p_i(x)-1) \leq 2$, one has

$$\begin{aligned}
& \iint_{Q_{T_1}} \left| (b_i(x, t) + \varepsilon)^{(p_i(x)-1)/p_i(x)} \right. \\
&\cdot \left. \left(|v_{\varepsilon x_i}|^2 + \varepsilon \right)^{p_i(x)/(p_i(x)-1)} dx dt \right. \\
&\leq c \left(\iint_{Q_{T_1}} (b_i(x, t) + \varepsilon) |v_{\varepsilon x_i}|^{p_i(x)} dx dt \right. \\
&+ \iint_{Q_{T_1}} \left. + \varepsilon^{p_i(x)(p_i(x)-2)/2(p_i(x)-1)} (b_i(x, t) + \varepsilon) \right. \\
&\cdot \left. |v_{\varepsilon x_i}|^{p_i(x)/(p_i(x)-1)} dx dt \right) \\
&\leq c \left(\iint_{Q_{T_1}} (b_i(x, t) + \varepsilon) |v_{\varepsilon x_i}|^{p_i(x)} dx dt \right. \\
&+ c \iint_{Q_{T_1}} \left. + \varepsilon^{p_i(x)(p_i(x)-2)/2(p_i(x)-1)} (b_i(x, t) + \varepsilon) \right. \\
&\cdot \left. \left(|v_{\varepsilon x_i}|^2 + 1 \right) dx dt \right) \leq c \iint_{Q_{T_1}} (b_i(x, t) + \varepsilon) \\
&\cdot |v_{\varepsilon x_i}|^{p_i(x)} dx dt + c \leq c(T_1).
\end{aligned} \tag{47}$$

Then, there exists $w_i \in L^1(0, T_1; L^{p_i(x)/(p_i(x)-1)}(\Omega))$, $i = 1, 2, \dots, n$ such that

$$\begin{aligned}
(b_i(x, t) + \varepsilon)^{(p_i(x)-1)/p_i(x)} \left(|v_{\varepsilon x_i}|^2 + \varepsilon \right)^{(p_i(x)-2)/2} &\rightharpoonup w_i, \\
&\text{in } L^1(0, T_1; L^{p_i(x)/(p_i(x)-1)}(\Omega)).
\end{aligned} \tag{48}$$

Using the similar method as that of the evolutionary p -Laplacian equation [34], we can deduce that

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \iint_{Q_T} \sum_{i=1}^N |v_{\varepsilon x_i}|^{p_i(x)-2} v_{\varepsilon x_i} \xi_{x_i} dx dt \\ &= \iint_{Q_T} \sum_{i=1}^N w_i \xi_{x_i} dx dt \\ &= \iint_{Q_T} \sum_{i=1}^N |v_{x_i}|^{p_i(x)-2} v_{x_i} \xi_{x_i} dx dt, \end{aligned} \tag{49}$$

for any $\xi(x, t) \in C_0^1(Q_T)$.

At last, the initial value in the sense of (14) can be found in [3]. Consequently, v_ε is the solution of (6). \square

3. The Stability

One can refer to [11–13] for the definitions of the exponent variable spaces, $(L^{q(x)}(\Omega), |\cdot|_{L^{q(x)}(\Omega)})$, $(W^{1,q(x)}(\Omega), |\cdot|_{W^{1,q(x)}(\Omega)})$, and $W_0^{1,q(x)}(\Omega)$. Also, one can find other details and recent applications to partial differential equations in [35–39].

Lemma 6 (see [11–13]). *If $p(x)$ and $q(x)$ are real functions with $1/p(x) + 1/q(x) = 1$ and $q(x) > 1$, then, for any $v \in L^{p(x)}(\Omega)$ and $u \in L^{q(x)}(\Omega)$, we have*

$$\left| \int_{\Omega} v u dx \right| \leq 2 |v|_{L^{p(x)}(\Omega)} |u|_{L^{q(x)}(\Omega)}. \tag{50}$$

Moreover,

if $|v|_{L^{q(x)}(\Omega)} = 1$, then $\int_{\Omega} |v|^{q(x)} dx = 1$,

if $|v|_{L^{q(x)}(\Omega)} > 1$,

then $|v|_{L^{q(x)}(\Omega)}^{q^-} \leq \int_{\Omega} |v|^{q(x)} dx \leq |v|_{L^{q(x)}(\Omega)}^{q^+}$, (51)

if $|v|_{L^{q(x)}(\Omega)} < 1$,

then $|v|_{L^{q(x)}(\Omega)}^{q^+} \leq \int_{\Omega} |v|^{q(x)} dx \leq |v|_{L^{q(x)}(\Omega)}^{q^-}$.

We let $g_m(s)$ be an odd function, and

$$g_m(s) = \begin{cases} 1, & s > \frac{1}{m}, \\ m^2 s^2 e^{1-m^2 s^2}, & 0 \leq s \leq \frac{1}{m}. \end{cases} \tag{52}$$

Then,

$$\lim_{m \rightarrow \infty} g_m(s) = \operatorname{sgn}(s), \quad s \in (-\infty, +\infty). \tag{53}$$

Let $\varphi(x)$ be a $C^1(\bar{\Omega})$ function satisfying

$$\begin{aligned} \varphi(x)|_{x \in \Gamma_2} &= 0, \\ \varphi(x)|_{x \in \bar{\Omega} \setminus \Gamma_2} &> 0, \end{aligned} \tag{54}$$

and

$$\Omega_m = \left\{ x \in \Omega : \varphi(x) > \frac{1}{m} \right\}. \tag{55}$$

Define

$$\varphi_m(x) = \begin{cases} 1, & \text{if } x \in \Omega_m, \\ m\varphi(x), & \text{if } x \in \Omega \setminus \Omega_m. \end{cases} \tag{56}$$

Then $\varphi_m(x)|_{x \in \Gamma_2} = 0$ and

$$\varphi_{m x_i}(x) = \begin{cases} 0, & \text{if } x \in \Omega_m, \\ m\varphi_{x_i}(x), & \text{if } x \in \Omega \setminus \Omega_m. \end{cases} \tag{57}$$

Theorem 7. *Let $v(x, t) \in L^\infty(Q_T)$ and $u(x, t) \in L^\infty(Q_T)$ be two solutions of (6) with the same partial second boundary value condition (8) and with the initial values $v_0(x)$ and $u_0(x)$. If $\partial\Omega = \Gamma_1 \cup \Gamma_2$ satisfies (9), b_i satisfies (10),*

$$|f(v, x, t) - f(u, x, t)| \leq c |v - u|, \tag{58}$$

and, for large enough m ,

$$\operatorname{ess\,sup}_{t \in [0, T]} \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |\varphi(x)_{x_i}|^{p_i(x)} dx \right)^{1/p_i^+} \leq c; \tag{59}$$

then

$$\int_{\Omega} |v(x, t) - u(x, t)| dx \leq \int_{\Omega} |v(x, 0) - u(x, 0)| dx, \tag{60}$$

a.e. $t \in [0, T]$.

Proof. Let $\chi_{[\tau, s]}$ be the characteristic function of $[\tau, s] \subseteq [0, T)$. By a process of limit, we can choose the test function as $\chi_{[\tau, s]} \varphi_m g_m(v - u)$. Then

$$\begin{aligned} & \int_{\tau}^s \int_{\Omega} \varphi_m g_m(v - u) \frac{\partial(v - u)}{\partial t} dx dt + \sum_{i=1}^n \int_{\tau}^s \int_{\Omega} b_i(x) \\ & \cdot \left(|v_{x_i}|^{p_i(x)-2} v_{x_i} - |u_{x_i}|^{p_i(x)-2} u_{x_i} \right) (v_{x_i} - u_{x_i}) \\ & \cdot g'_m(v - u) \varphi_m(x) dx dt + \sum_{i=1}^n \int_{\tau}^s \int_{\Omega} b_i(x) \\ & \cdot \left(|v_{x_i}|^{p_i(x)-2} v_{x_i} - |u_{x_i}|^{p_i(x)-2} u_{x_i} \right) (v_{x_i} - u_{x_i}) \\ & \cdot g_m(v - u) \varphi_{m x_i} dx dt = \int_{\tau}^s \int_{\Omega} [f(v, x, t) \\ & - f(u, x, t)] \varphi_m g_m(v - u) dx dt. \end{aligned} \tag{61}$$

First of all,

$$\begin{aligned} & \int_{\tau}^s \int_{\Omega} b_i(x) \left(|v_{x_i}|^{p_i(x)-2} v_{x_i} - |u_{x_i}|^{p_i(x)-2} u_{x_i} \right) \\ & \cdot (v_{x_i} - u_{x_i}) g'_m(v - u) \varphi_m(x) dx dt \geq 0. \end{aligned} \tag{62}$$

Secondly, since $v_t \in L^2(Q_T)$, $u_t \in L^2(Q_T)$, by the Lebesgue dominated convergence theorem, we have

$$\begin{aligned} & \lim_{m \rightarrow \infty} \int_{\tau}^s \int_{\Omega} \varphi_m(x) g_m(v-u) \frac{\partial(v-u)}{\partial t} dx dt \\ &= \int_{\Omega} |v-u|(x, s) dx - \int_{\Omega} |v-u|(x, \tau) dx. \end{aligned} \quad (63)$$

By (59), we have

$$\begin{aligned} & \left| \int_{\tau}^s \int_{\Omega} b_i(x, t) \left(|v_{x_i}|^{p_i(x)-2} v_{x_i} - |u_{x_i}|^{p_i(x)-2} u_{x_i} \right) \right. \\ & \quad \cdot \varphi_{m x_i} g_m(v-u) dx dt \Big| = \left| \int_{\tau}^s \int_{\Omega \setminus \Omega_m} b_i(x, t) \right. \\ & \quad \cdot \left(|v_{x_i}|^{p_i(x)-2} v_{x_i} - |u_{x_i}|^{p_i(x)-2} u_{x_i} \right) \\ & \quad \cdot \varphi_{m x_i} g_m(v-u) dx dt \Big| \leq m \int_{\tau}^s \int_{\Omega \setminus \Omega_m} b_i(x, t) \\ & \quad \cdot \left(|v_{x_i}|^{p_i(x)-1} + |u_{x_i}|^{p_i(x)-1} \right) \\ & \quad \cdot |\varphi_{x_i} g_m(v-u)| dx dt \\ & \leq cm \int_{\tau}^s \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) \right. \\ & \quad \cdot \left(|v_{x_i}|^{p_i(x)} + |u_{x_i}|^{p_i(x)} \right) dx \Big)^{1/q_i^+} \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) \right. \\ & \quad \cdot |\varphi_{x_i}|^{p_i(x)} dx \Big)^{1/p_i^+} dt \\ & \leq c \int_{\tau}^s \left[\left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |v_{x_i}|^{p_i(x)} dx \right)^{1/q_i^+} \right. \\ & \quad \left. + \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |u_{x_i}|^{p_i(x)} dx \right)^{1/q_i^+} \right] \\ & \quad \cdot m \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |\varphi_{x_i}|^{p_i(x)} dx \right)^{1/p_i^+} dt \\ & \leq c \int_{\tau}^s \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |v_{x_i}|^{p_i(x)} dx \right)^{1/q_i^+} dt \\ & \quad + c \int_{\tau}^s \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |u_{x_i}|^{p_i(x)} dx \right)^{1/q_i^+} dt, \end{aligned} \quad (64)$$

where $q_i(x) = p_i(x)/(p_i(x) - 1)$, $q_i^+ = \max_{x \in \bar{\Omega}} q_i(x)$. Then we have

$$\begin{aligned} & \lim_{m \rightarrow \infty} \left| \int_{\tau}^s \int_{\Omega} b_i(x, t) \left(|v_{x_i}|^{p_i(x)-2} v_{x_i} - |u_{x_i}|^{p_i(x)-2} u_{x_i} \right) \right. \\ & \quad \cdot \varphi_{m x_i} g_m(v-u) dx dt \Big| \end{aligned}$$

$$\begin{aligned} & \leq c \lim_{m \rightarrow \infty} \left[\left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |v_{x_i}|^{p_i(x)} dx \right)^{1/q_i^+} \right. \\ & \quad \left. + \left(\int_{\Omega \setminus \Omega_m} b_i(x, t) |u_{x_i}|^{p_i(x)} dx \right)^{1/q_i^+} \right] = 0. \end{aligned} \quad (65)$$

In addition, by (58)

$$\begin{aligned} & \lim_{m \rightarrow \infty} \left| \int_{\tau}^s \int_{\Omega} [f(v, x, t) - f(u, x, t)] \right. \\ & \quad \cdot \varphi_m g_m(v-u) dx dt \Big| \leq \int_{\tau}^s \int_{\Omega} |v(x, t) \\ & \quad - u(x, t)| dx dt. \end{aligned} \quad (66)$$

At last, let $m \rightarrow \infty$ in (61). Then

$$\int_{\Omega} |v(x, s) - u(x, s)| dx \leq \int_{\Omega} |v(x, \tau) - u(x, \tau)| dx. \quad (67)$$

By the arbitrary of τ , we have the conclusion. \square

Proof of Theorem 3. We only need to choose

$$\varphi(x) = d(x); \quad (68)$$

in Theorem 7, the conclusion is clear. \square

4. Conclusion

Comparing with the isotropic case, the anisotropic parabolic equations are closer to the truth about the applications. They reflect in the mathematical modeling of physical and mechanical processes in anisotropic continuous medium. Different from the elliptic anisotropic equations with the variable exponent that has attracted much attentions recently, more or less beyond my imagination, only a few references related to the anisotropic parabolic equations can be found. Even there are not any papers to discuss the fundamental solution of the anisotropic equation. This paper has taken this one step further. We think the contributions are mainly in two aspects. The first one lies in the fact that if the diffusion coefficient $b_i(x, t)$ is degenerate on a part of the boundary $\Gamma_1 \subset \partial\Omega$, then, only under the partial second boundary value condition, one can study the well-posedness of solutions to the anisotropic parabolic equations. The second one lies in the fact that even if the weak solutions may be blow-up in finite time provided that the source $f(x, t, u) \geq 0$, the uniqueness of weak solution still may be true.

Data Availability

There is not any data in the paper.

Conflicts of Interest

The author declares that he has no conflicts of interest.

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