The optical properties of titanate nanotubes (TiNts) intercalated with rare earths (RE) ions were intensively investigated in this study. To prepare the nanomaterials, sodium titanate nanotubes (Na-TiNts) were submitted to ion exchange reactions with different rare earth elements (RE: \(\text{Pr}^{3+}\), \(\text{Er}^{3+}\), \(\text{Nd}^{3+}\), and \(\text{Yb}^{3+}\)). To the best of our knowledge, it is the first time that these RE-TiNts were synthesized. All samples were characterized by Raman spectroscopy, X-ray powder diffraction (XRD), transmission electron microscopy (TEM), and energy dispersive X-ray spectroscopy (EDS). Furthermore, the optical properties were examined using photoluminescence spectroscopy (PL) and UV-Vis-NIR absorption spectroscopy. The PL intensity (visible range) of the RE-TiNt samples showed a strong dependence when the temperature was decreased down to 20 K. This PL enhancement probably was promoted by electronic modifications in titanate layer band gap and/or interface charge transfers due to RE ions intercalation.

1. Introduction

Oxide based nanostructures exhibit some properties that differ from their bulk counterparts [1, 2]. In particular, alkali metal titanates, usually possessing a layered or tunnel-type structure, have been subject of extensive investigations due to their excellent ion exchange ability added to the exceptional photocatalytic activities [3–6]. Kasuga et al. [7] demonstrated the possibility of producing alkali metal titanate nanoparticles with tubular morphology through hydrothermal treatment of \(\text{TiO}_2\) in the high alkaline medium. Ever since, titanate nanotubes (TiNts) have attracted both scientific and technological interest due to its unique combination of morphology, size, and functionality [8]. Furthermore, TiNts have large surface area (~180 m\(^2\) g\(^{-1}\)) and porosity that leads to a wide variety of applications involving their surfaces [9–11].

Optical properties of nanomaterials are widely investigated due to their potential applications in small devices [1–3, 8]. The optical properties of TiNts have been tuned in according to the desired application using different methods to dope and/or inserting ions into the structure of the nanotubes [3, 6, 7, 12, 13]. For example, Li and Zhong [14] have...
inserted Sn$^{2+}$ and Pb$^{2+}$ in titanate nanotubes via ion exchange reaction and these nanomaterials were applied in photocatalysis using visible light for dyes degradation. Haro-González et al. [15] have prepared Eu$^{3+}$ doped titania nanotubes via hydrothermal treatment of Eu$^{3+}$ doped TiO$_2$ nanopowders and they have studied the luminescent properties of these nanomaterials.

Photoelectronic properties of titanate nanotubes have been investigated [16]. It has been suggested that the photoluminescence emission process is from the exciton states trapped in titanate octahedral ([TiO$_6$]) units. The excitonic properties of titanate nanomaterials depend on the particular environment of the interlayer space. Therefore, the position and oxidation states of the cations intercalated can modify the [TiO$_6$] octahedron promoting or not excited states inside of its band gap with lifetime sufficient to have photoluminescence (PL) [16]. Then, there is a critical role of the layer composition and structure on charge separation of the photogenerated charge carriers. The intercalated ions can modify and control these properties of the layered nanostructures [16–18].

Rare earths (RE) doping and/or intercalating nanomaterials are widely investigated due to their excellent luminescent characteristics arising from the transitions of electronic levels which have optical applications such as lasers, optical amplifiers, phosphors, and other optical devices [19]. The choice of host materials has a great importance in the design of the rare earth-based luminescent properties for efficient applications [8, 20, 21]. In this work, in order to obtain different photoluminescence (PL) properties in titanate nanotubes, they were used as a host. Na$^+$-intercalated titanate nanotubes (Na-TiNts) were prepared by hydrothermal method before being submitted to ion exchange reactions with different rare earth ions (Pr$^{3+}$, Er$^{3+}$, Nd$^{3+}$, and Yb$^{3+}$). To the best of our knowledge, it is the first time that these nanotubes (Pr-, Er-, Nd-, and Yb-TiNts) were prepared and their optical properties studied. Thus, this work proposes a very simple route to modify the optical properties of titanate nanotubes and obtain efficient PL emitters at room and low temperatures.

2. Experimental Procedures

All chemicals (reagent grade, Sigma-Aldrich, and Vetec) were used without further purification processes. Deionized water was used to prepare all solutions.

2.1. Titanate Nanotubes Preparation. Na-TiNts were prepared using commercial TiO$_2$ (anatase) powder (purity of 99.8%). Typically, 2.0 g of the TiO$_2$ was mixed in 60 mL of 10 mol L$^{-1}$ NaOH solution and then transferred into a Teflon-lined autoclave and heated at 150°C for 72 h [22].

2.2. Ion Exchange Reactions of the Titanate Nanotubes. The as-prepared Na-TiNts are suitable for rare earths ion exchange reactions without additional process of protonation. These reactions were carried out by suspending 100 mg of Na-TiNts in 100 mL of 0.05 mol L$^{-1}$ Pr(NO$_3$)$_3$·6H$_2$O, Nd(NO$_3$)$_3$·6H$_2$O, Er(NO$_3$)$_3$·5H$_2$O, and Yb(NO$_3$)$_3$·5H$_2$O (purity of 99.9%) aqueous solutions. After this, the suspension was left under magnetic stirring for 24 h at room temperature. The solid products were isolated by centrifugation at 3500 rpm. Then, the products were washed several times with deionized water aiming to remove residual soluble ions from the precursors. Finally, the samples were dried under vacuum conditions at room temperature during 6 h to obtain the RE intercalated titanate nanotubes (RE-TiNts).

2.3. Characterization. After the Na-TiNt and Re-TiNT preparations, the Raman spectroscopy experiments were performed using a Bruker Senterra spectrometer equipped with an Olympus BX5 microscope in all samples. A solid-state laser line at 532 nm was used for measuring the Raman spectra. Power density of laser was below 10$^7$ W/m$^2$ to avoid sample overheating. A spectral resolution of 3 cm$^{-1}$ was used and measurements were carried out using a backscattering geometry. X-ray powder diffraction (XRD) patterns were obtained using a Shimadzu XRD6000 diffractometer, using CuK$\alpha$ ($\lambda = 1.5406$ Å) radiation operating at 30 mA and 40 kV. A scan rate of 1° min$^{-1}$ was used. Transmission electron microscopy (TEM) images were obtained using a JEOL-JEM-2010F operating with 200 kV. The TEM samples were prepared by dropping aqueous suspension containing powder samples onto a holey carbon-coated copper grid and letting the water to evaporate at room temperature. The elemental analysis by energy dispersive X-ray spectroscopy (EDS) was performed using an Oxford X-ray detector attached to a Tescan FEG-Mira3 scanning electron microscope. The wavelength dependence of the optical diffuse reflectivity of the powder samples was measured using Perkin-Elmer Lambda 950 UV-Vis-NIR Spectrophotometer. Diffuse reflectance spectra were acquired using the Labsphere DRA-CA-30I accessory. Photoluminescence measurements were carried out upon excitation using an Ar$^+$ laser line at 488 nm on the powder pellet samples using a microscope with 50x objective lenses to focalize the laser beam. The emission signal was collected using the objective lenses and analyzed by a Renishaw inVia spectrometer equipped with an 1800 lines/mm grating and detected by a CCD. Low-temperature measurements were performed in a helium flux cryostat (Advanced Research Systems). A digital temperature controller (Lakeshore) with 0.1 K of stability was used for monitoring the temperature in the range of 20–300 K. The spectra were measured after an appropriate stabilization time for each thermal step.

3. Results and Discussion

In Figure 1(a) the Raman spectra of intercalated TiNts are shown. Firstly, the peaks around 158 and 190 cm$^{-1}$ are assigned to lattice modes (Na$^+$-O-Ti). The peaks at 275, 448, 662, and 700 cm$^{-1}$ are assigned to Ti-O stretching in the edge-shared TiO$_6$ and the peaks at 906 cm$^{-1}$ represent the stretching vibration of the nonshared Ti-O bonds in distorted TiO$_6$ unit [6, 23, 24]. It was observed that the highest (906 cm$^{-1}$) and the two lowest (158 and 190 cm$^{-1}$) energy
modes have their wavenumbers and intensities extensively affected by the ion exchange with different RE ion species. Additionality, some of the medium energy modes (between 200 and 800 cm\(^{-1}\)) are also changed. The most affected mode is that related to Ti-O bonds whose oxygen atoms are not shared among the TiO\(_6\) units (at the corner of the TiO\(_6\) octahedral slab). Thus, the different RE ions inserted into the TiN\(_t\)s structure promote the reordering of the interlayer regions leading to these changes in Raman modes [25]. The changes can be explained as follows: the Na\(^+\) ions were substituted by hydrated RE ions that may occupy different sites in the interlayer spaces. Due to this reordering the RE-O-Ti interlayer bond could deviate from a linear connection with the different intercalated ions and leads to shifting of the band at 906 cm\(^{-1}\) toward lower frequencies [25]. This explanation is supported by observing the difference between the electronegativity as well as the covalent radii of the Na and RE atoms. In general, the more electronegative ion will attract more easily the oxygen atoms, stretching the Ti-O bonds. Also, another important aspect to be pointed out is the considerable changes of the bands below 200 cm\(^{-1}\), which are related to the interlayer structure modifications. Thus, the coordination environments for Na\(^+\) ion and other intercalated RE ions are different in the structural analysis, as also mentioned previously for other interlayer cations [23, 25].

To better understand the relationship between the crystalline structure and vibrational properties, X-ray diffraction experiments were performed because they provide information about the interlayer region and crystallinity of the layers. The X-ray diffraction data (Figure 1(b)) suggested that both structure and morphology of TiN\(_t\)s were preserved after the RE ions intercalation. The peaks around 10° (2\(\theta\)) in Na-TiN\(_t\)s, attributed to (200) plane which is related to interlayer distance [25, 26], were slightly shifted to lower 2\(\theta\) position when RE ions were intercalated indicating an increase of
interlayer distance (see Table S1 in Supplementary Material available online at https://doi.org/10.1155/2017/3809807). This increase of the interlayer distance is expected since the ionic radius of Na\(^+\) is smaller than that of RE hydrated ions [25, 27, 28]. Furthermore, it is possible to observe changes in the diagonal planes (110) and (211) which can be related to the rearrange of the RE ions into the layers of titanate to different coordination positions when it is compared with Na-TiNt [13, 22].

Figure 1(c) shows the UV-Vis diffuse reflectance spectra for pristine nanotubes and RE-TiNt samples. It can be observed that Na-TiNt have only a strong band edge absorption in the ultraviolet region (above 3.0 eV). This absorption band in the ultraviolet range can be associated with the excitation of the O2p electron to the Ti3d level [29]. All of the RE-TiNt spectra have shown an edge band corresponding to the absorption edge of the titanate nanotubes at around 3.0 eV. However, for Pr-TiNt UV-Vis spectrum, it is possible to see three bands at 2.56, 2.63, and 2.78 eV and a broad band at 2.08 eV which indicates Pr\(^{3+}\) ions absorption peaks and that this ion was inserted into the TiNt [30]. The Nd-TiNt spectrum shows four broad and intense bands at 1.67, 2.11, 2.34, and 2.42 eV attributed to Nd\(^{3+}\) ions absorptions [31]. In the spectrum of Er-TiNt bands can be seen at 1.89, 2.26, 2.37, 2.54, and 3.28 eV associated with trivalent erbium (Er\(^{3+}\)) absorption [32]. In the case of Yb-TiNt sample the absorption bands due to the ions absorption are not visible because it is out of the spectral range studied here; the absorption peak of Yb\(^{3+}\) ion is around 1.26 eV [21, 33]. On the other hand, it is possible to see the Yb-TiNt spectrum slightly broad in the edges around 2.25 and 2.95 eV which could be associated with defects promoted in titanate electronic structure by ion exchanging [13]. The changes occurred in the TiNt spectrum profile when RE ions were inserted, like a slightly flat absorption on visible range, and a shifting of the absorption edge band can be attributed to charge transfer between TiNt host and RE ions, promoting modifications into the electronic structure of the titanate layers [34–36].

All analyses aforesaid indicate that pristine nanotubes after ion exchange reactions were modified due to substitution of Na by RE ions. To evaluate these substitutions directly, the chemical compositions of RE-TiNt samples were investigated by energy dispersive X-ray spectroscopy (EDS). The Na/Ti and RE/Ti atomic ratios are shown in Table 1. The amount of Na\(^{+}\) ions in pristine nanotubes is in agreement with the Na\(_{x}\)H\(_{3}\)Ti\(_{3}\)O\(_{7}\)-nH\(_{2}\)O \((0 \leq x \leq 2)\) chemical formula [22]. After ion exchange reactions, the sodium content drastically decreased in nanotubes. On the other hand, it was possible to detect the RE ions, suggesting that ion exchange reactions were very efficient to substitute Na ions by RE. These findings were demonstrated by a very low Na/Ti atomic ratio in all RE ion exchanged samples and RE/Ti atomic ratio between 0.14 and 0.19. Also, the amount of oxygen and titanium in the samples was almost unchanged after the reactions. Then, \([\text{RE}_{x}\text{Na}_{1-x}\text{H}_{3}\text{Ti}_{3}\text{O}_{7}]^{\text{3-}}\cdot\text{nH}_{2}\text{O}\) chemical formula is suitable for RE-TiNt samples, disregarding the small amount of remaining Na\(^{+}\). Based on EDS analyses the following chemical formulas were proposed for RE-TiNt: \(\text{Pr}_{1-x}\text{Na}_{x}\text{H}_{3}\text{Ti}_{3}\text{O}_{7}\cdot\text{nH}_{2}\text{O}\); \(\text{Nd}_{1-x}\text{Na}_{x}\text{H}_{3}\text{Ti}_{3}\text{O}_{7}\cdot\text{nH}_{2}\text{O}\); \(\text{Er}_{1-x}\text{Na}_{x}\text{H}_{3}\text{Ti}_{3}\text{O}_{7}\cdot\text{nH}_{2}\text{O}\); and \(\text{Yb}_{1-x}\text{Na}_{x}\text{H}_{3}\text{Ti}_{3}\text{O}_{7}\cdot\text{nH}_{2}\text{O}\) (H\(^{+}\) was added in the chemical formulas because the aqueous solutions of rare earth salts are acids).

TEM images of the RE-TiNts (Figure 2) show that tubular morphology was preserved after the ion exchange reaction, corroborating with Raman and XRD results. Furthermore, in these images, it is not possible to observe nanoparticles decorating the nanotube on external surface as demonstrated in previous work about ion exchange reactions involving Ce\(^{3+}\) [6]. The RE-TiNt interlayer distances measured by TEM images (0.72, 0.83, 0.86, and 0.87 nm for Pr-, Er-, Nd-, and Yb-TiNt, resp.) are larger than Na-TiNt one (0.70 nm, see Supplementary Material, Figure S1). Also, it is important to notice that, for all samples, the interlayer distances measured using TEM images are lower than those obtained by XRD patterns. These contractions would be due to partial release of water molecules present in the interlayer region [26, 37], since for TEM experiments the samples were submitted to high vacuum conditions (below 10\(^{-6}\) Pa).

In Figure 3(a) the photoluminescence spectra at room temperature and atmospheric pressure of the Na-TiNts and RE-TiNts are shown. It is possible to see the PL enhancement of RE-TiNt as compared to pristine sample. Pr- and Er-TiNts had the highest luminescence intensity, being the emission band centered around 1.81 and 1.86 eV, respectively. Yb-TiNt had an emission around 2.04 eV and the intensity is lower than Er-TiNt and Pr-TiNt. Finally, Nd-TiNt had the lowest PL intensity among the RE-TiNt samples with the broad emission band centered around 1.99 eV. Additionally, Er-TiNt and Nd-TiNt spectra show that there are one (1.89 eV) and three (1.67, 2.11, and 2.14 eV) absorptions related to Er\(^{3+}\) [32] and Nd\(^{3+}\) [21].

The PL profiles are broad and typical of multiphonon process and/or interface emissions [35, 36, 38, 39]. PL spectra show a strong overlap of the peaks from the RE ions direct recombination as well, and the emission energies are lower than the direct recombination of an electron in Ti3d band with a hole in an O2p band. Thus, electronic (defect) energy levels within the band gap and/or interface effects are promoted into TiNt by RE ions modifying the electronic structure of the host [15, 21]. The electronic levels or interface densities of state were generated by the insertion of RE ions leading to the PL emission through of the TiNt lattice (host) (see Supplementary Information, Figure S2) [40].

Figure 3(b) shows the effect of the vacuum (\(\sim10^{-5}\) Pa) on the PL spectra of Na-TiNts and RE-TiNts. It is possible

### Table 1: RE/Ti atomic ratios obtained by EDS from intercalated samples. The amount of Na ions in pristine nanotubes is in agreement with the composition Na\(_{x}\)H\(_{3}\)Ti\(_{3}\)O\(_{7}\)-nH\(_{2}\)O.

<table>
<thead>
<tr>
<th>Samples</th>
<th>Na-TiNt</th>
<th>Pr-TiNt</th>
<th>Nd-TiNt</th>
<th>Er-TiNt</th>
<th>Yb-TiNt</th>
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<tbody>
<tr>
<td>Na/Ti</td>
<td>0.33</td>
<td>0.02</td>
<td>0.01</td>
<td>0.01</td>
<td>0.01</td>
</tr>
<tr>
<td>RE/Ti</td>
<td>—</td>
<td>0.19</td>
<td>0.14</td>
<td>0.18</td>
<td>0.17</td>
</tr>
</tbody>
</table>
Figure 2: TEM (a) and HRTEM (b) images of RE-TiNts. The interlayer average distances of RE-TiNts samples were shown (white lines in HRTEM images).

Figure 3: Room temperature PL spectra of RE-TiNts and Na-TiNts samples at (a) atmospheric pressure and (b) under vacuum (∼10⁻⁵ Pa). Some absorptions peaks due to RE ions inserted into the nanotubes can be seen.
to see an enhancement of the PL intensity (~2, ~5, ~14, ~24, and ~30 times more intense for Pr, Er, Na, Yb, and Nd, resp.) probably, due to water release promoted by the vacuum conditions, suggesting that the interlayer space of the nanotubes was reordered and energy of ion-layer bonds and electronic levels were modified. Now, Yb-TiNts show the highest luminescence intensity with a more defined emission peak centered at 2.05 eV. For Er- and Pr-TiNts, emission peaks at 2.05 eV and 1.95 eV were observed, respectively. The spectrum of Nd-TiNts showed a peak around 1.95 eV. On the other hand, under vacuum Na-TiNts showed a weak broad PL peak around 2.05 eV. Some extra absorption peaks were observed on the spectra of Yb-, Nd-, and Er-TiNts, as it was observed in the PL spectra at atmospheric pressure. Thus, the water release and reordering of the interlayer space allowed another enhancement of the electronic recombination (PL) promoted by the excitation through inserted levels and/or interface densities of state into the electronic structure of TiNts intercalated with RE [17,18,29].

Analyzing the PL of the RE-TiNts at atmospheric pressure (with structural water, Figure 3(a)), the intensity of PL peak could be related to the electronic effects promoted by the concentration of RE into the nanotube structures. The two most intense PL peaks are from Pr- and Er-TiNts, which have the highest concentration according to EDS quantifications (see Table 1). After the dehydration process, the electronic states promoted by vacancies compete with RE concentration changing the PL emission characteristics. Thus, it is very difficult to say, using that laser line (488 nm), which phenomena (electronic states inserted into the band gap of TiNts or interface energy states with charge transfer between RE and TiNts layers) can contribute the most to the PL emission efficiency. To understand better the PL mechanism, it is necessary to perform experiments with lower wavenumber laser lines to reach the energy of the band gap of pristine nanotubes.

The PL spectra at low temperatures were performed and they are shown in Figure 4. All of these PL peaks showed shifts with the decrease of temperature, and it may be due to some variations into the electronic structures. Furthermore, the intensity of PL peaks increase along with decreasing temperature which suggests that PL peaks could be related to shallow energy levels inserted by the ions and/or interface effects that are not thermally stable for electrons at that level [29]. In Na-TiNts spectrum (Figure 4(a)), the PL intensity increases almost linearly with the decrease of temperature and the peak position shifts to higher energies. The Nd-TiNts spectrum (Figure 4(b)) has shown the same behavior in the PL intensity and peak position, but there was stabilization at 45 K and the absorptions of Nd$^{3+}$ ions are more evident at low temperatures. For Pr-TiNts (Figure 4(c)), during the decrease of temperature, the peak position shifts toward higher energies and PL intensity increases. Below 220 K, an absorption peak appears around 2.08 eV that can be related to Pr$^{3+}$ direct absorption. For Er-TiNts (Figure 4(d)), the increase of intensity is almost linear as the temperature decreases and the peak position shows a shift to higher energy. Similar to Nd-TiNts, the Er$^{3+}$ absorptions (at 1.89, 2.11, 2.26, and 2.34 eV) are more evident at low temperatures. Yb-TiNts spectrum (Figure 4(e)) has shown the same shift to higher energy at low temperatures. However, there was a decrease of the intensity from 35 K to 9 K. This observation corroborates with an absorption plateau at UV-Vis measurements suggesting that Yb$^{3+}$ ions modify the electronic structure of titanate nanotubes. As the emission process comes from electronic structure of the titanate nanotube walls which were affected by the RE insertion, this emission might be absorbed by the lanthanide ion into the interlayer promoting the absorption peaks in the PL spectra. The result of PL enhancement during the increase of temperature from 9 K to 35 K for Yb-TiNts is probably due to thermal rectification of surface trap states or densities of states. A similar effect was observed for photoinduced PL efficiency enhancement of QDs during surface passivation [35,41,42]. At higher temperatures, the main part of nonradiative process that limits the quantum efficiency is the thermal escape, induced by multiple phonon absorption and promoting the PL quenching. At lower temperatures, the most of the excited electrons return to the shallow level via nonradiative transitions, which can lead to the PL enhancement [29].

4. Conclusions

In summary, we have reported the synthesis of titanate nanotubes intercalated with rare earths (RE) ions (RE: Pr$^{3+}$, Er$^{3+}$, Nd$^{3+}$, and Yb$^{3+}$). The morphology, structure, and optical properties of these novel titanate nanotubes were studied. The RE-TiNts samples have shown a higher PL emission than pristine Na-TiNts. The intensity and energy of the PL peaks are dependent on the temperature. The luminescence mechanism was briefly discussed and it was proposed that possible origin of the PL broad peaks would be due to the electronic levels inserted into the TiNts band gap and/or interface electronic densities introduced by the RE intercalations. Based on the results showed here, these novel nanomaterials could be used in applications involving optical devices like lasers, sensors, and phosphors.

Competing Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

Acknowledgments

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Figure 4: Temperature-dependent PL spectra: (a) Na-TiNts; (b) Nd-TiNts; (c) Pr-TiNts; (d) Er-TiNts; and (e) Yb-TiNts.
References


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