Research Article

Several Dynamical Properties for a Nonlinear Shallow Water Equation

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A nonlinear third order dispersive shallow water equation including the Degasperis-Procesi model is investigated. The existence of weak solutions for the equation is proved in the space $L^1(R) \cap BV(R)$ under certain assumptions. The Oleinik type estimate and $L^{2N}(R)$ (N is a natural number) estimate for the solution are obtained.

1. Introduction

Constantin and Lannes [1] derived the shallow water wave equation

$$\begin{aligned} u_t + u_x + \frac{3}{2}\rho uu_x + \mu(\alpha u_{xxx} + \beta u_{xxt}) \\ = \rho\mu(\gamma u_x u_{xx} + \delta uu_{xxx}), \end{aligned} \quad (1)$$

where the constants α , β , γ , δ , ρ , and μ satisfy certain conditions. As stated in [1], using suitable mathematical transformations, one can turn (1) into the form

$$u_t - u_{xxt} + 2ku_x + muu_x = au_x u_{xx} + buu_{xxx}, \quad (2)$$

where a , b , k , and m are constants. Clearly, (2) contains both the Camassa-Holm and Degasperis-Procesi models.

The aim of this paper is to investigate the existence of weak solutions for the special case of (2). Namely, we study the shallow water equation

$$\begin{aligned} \partial_t u - \partial_{xxx}^3 u + mu\partial_x u = 3\partial_x u \partial_{xx}^2 u + u\partial_{xxx}^3 u, \\ (t, x) \in R_+ \times R, \end{aligned} \quad (3)$$

where $m > 0$ is a constant. Letting $y = u - \partial_{xx}^2 u$, $v = (m - \partial_{xx}^2)^{-1}u$ and using (3), we derive the conservation law

$$\begin{aligned} \int_R yv dx &= \int_R \frac{1 + \xi^2}{m + \xi^2} |\widehat{u}(\xi)|^2 d\xi \\ &= \int_R \frac{1 + \xi^2}{m + \xi^2} |\widehat{u}_0(\xi)|^2 d\xi \sim \|u_0\|_{L^2(R)}, \end{aligned} \quad (4)$$

where $u_0 = u(0, x)$. In fact, the conservation law (4) takes an important role in our further investigations of (3).

For $m = 4$, (3) reduces to the Degasperis-Procesi equation [2]

$$\begin{aligned} \partial_t u - \partial_{xxx}^3 u + 4u\partial_x u = 3\partial_x u \partial_{xx}^2 u + u\partial_{xxx}^3 u, \\ (t, x) \in R_+ \times R, \end{aligned} \quad (5)$$

which has been studied by many scholars (see [3–5]). Lundmark and Szmigielski [6] developed an inverse scattering approach for computing n -peakon solutions to (5). The traveling wave solutions of (5) were investigated in Vakhnenko and Parkes [5]. Holm and Staley [7] studied stability of solitons and peakons numerically. Lin and Liu [8] proved the stability of peakons for the Degasperis-Procesi equation (5) under certain assumptions. The precise blowup scenario result, a blowup result, and the global existence of strong solutions and global weak solutions to (5) can be found in [9]. Matsuno [10] studied multisoliton solutions and their peakon limits.

Analogous to the case of the Camassa-Holm equation, Henry [11] and Mustafa [12] showed that smooth solutions to (5) have infinite speed of propagation. For other methods to handle the problems relating to various dynamic properties of the Degasperis-Procesi equation and other shallow water equations, the reader is referred to [13–15] and the references therein.

Coclite and Karlsen [16] established the existence, uniqueness, and $L^1(R)$ stability of entropy weak solutions belonging to the class $L^1(R) \cap BV(R)$ for (5). They obtained existence of at least one weak solution satisfying a restricted set of entropy inequalities in the space $L^2(R) \cap L^4(R)$ and extended these results to a class of generalized Degasperis-Procesi equations in [16].

Motivated by the desire to extend parts of the results presented in Coclite and Karlsen [16], we consider (3) with its Cauchy problem in the form

$$\begin{aligned} \partial_t u - \partial_{txx}^3 u &= -\partial_x \left(\frac{m}{2} u^2 \right) + 3\partial_x u \partial_{xx}^2 u + u \partial_{xxx}^3 u \\ &= -\left(\frac{m}{2} u^2 \right)_x + \frac{1}{2} \partial_{xxx}^3 u^2, \\ u(0, x) &= u_0(x), \end{aligned} \quad (6)$$

which is equivalent to

$$\begin{aligned} u_t + uu_x &= -\frac{m-1}{2} \Lambda^{-2} (u^2)_x, \\ u(0, x) &= u_0(x), \end{aligned} \quad (7)$$

where $m > 0$ is a constant and $\Lambda = (1 - \partial_x^2)^{1/2}$.

The objective of this paper is to study (3). We investigate the existence of weak solutions in the space $L^1(R) \cap BV(R)$ under certain conditions. Several dynamical properties such as Oleinik type estimate and $L^{2N}(R)$ (N is a natural number) are obtained. As (3) includes the Degasperis-Procesi equation (5), parts of results presented in [16] are extended. Here we should mention that the generalized Degasperis-Procesi equation discussed in [16] does not include the model (3). We state that the ideas and approaches to prove our main results come from those in [16].

The rest of this paper is organized as follows. Section 2 establishes the L^2 , BV , and L^∞ estimates for the viscous approximations of problem (6). The main result is given in Section 3.

2. Viscous Approximations and Estimates

Firstly, we give some notations.

Set $R_+ = [0, +\infty)$. The space of all infinitely differentiable functions $\phi(t, x)$ with compact support in $R_+ \times R$ is denoted by C_0^∞ . We let $L^p = L^p(R)$ ($1 \leq p < +\infty$) be the space of all measurable functions h such that $\|h\|_{L^p}^p = \int_R |h(t, x)|^p dx < \infty$. We define $L^\infty = L^\infty(R)$ with the standard norm $\|h\|_{L^\infty} = \inf_{m(e)=0} \sup_{x \in R \setminus e} |h(t, x)|$. For any real number s ,

we let $H^s = H^s(R)$ denote the Sobolev space with the norm defined by

$$\|h\|_{H^s} = \left(\int_R (1 + |\xi|^2)^s |\hat{h}(t, \xi)|^2 d\xi \right)^{1/2} < \infty, \quad (8)$$

where $\hat{h}(t, \xi) = \int_R e^{-ix\xi} h(t, x) dx$.

For $T > 0$ and nonnegative number s , let $C([0, T]; H^s(R))$ denote the Frechet space of all continuous H^s -valued functions on $[0, T]$.

Defining

$$\phi(x) = \begin{cases} e^{1/(x^2-1)}, & |x| < 1, \\ 0, & |x| \geq 1 \end{cases} \quad (9)$$

and letting $\phi_\varepsilon(x) = \varepsilon^{-1/4} \phi(\varepsilon^{-1/4} x)$ with $0 < \varepsilon < 1/4$ and $u_{0,\varepsilon} = \phi_\varepsilon * u_0$, we know that $u_{0,\varepsilon} \in C^\infty$ for any $u_0 \in H^s$ with $s \geq 0$.

For simplicity, throughout this paper, we let c_0 denote any positive constant, which is independent of parameter ε and time t .

To establish the existence of solutions to the Cauchy problem (6), we will analyze the limiting behavior of a sequence of smooth functions $\{u_\varepsilon\}_{\varepsilon>0}$, where each function u_ε satisfies the viscous problem

$$\begin{aligned} \partial_t u_\varepsilon - \partial_{txx}^3 u_\varepsilon + m u_\varepsilon \partial_x u_\varepsilon &= 3\partial_x u_\varepsilon \partial_{xx}^2 u_\varepsilon + u_\varepsilon \partial_{xxx}^3 u_\varepsilon + \varepsilon \partial_{xx}^2 u_\varepsilon - \varepsilon \partial_{xxxx}^4 u_\varepsilon, \\ (t, x) &\in R_+ \times R, \\ u_\varepsilon(0, x) &= u_{0,\varepsilon}(x), \quad x \in R, \end{aligned} \quad (10)$$

which is equivalent to the parabolic-elliptic system

$$\begin{aligned} \partial_t u_\varepsilon + \partial_x \left(\frac{u_\varepsilon^2}{2} \right) + \partial_x P_\varepsilon &= \varepsilon \partial_{xx} u_\varepsilon, \\ P_\varepsilon - \partial_{xx}^2 P_\varepsilon &= \frac{m-1}{2} u_\varepsilon^2, \\ u_\varepsilon(0, x) &= u_{0,\varepsilon}(x). \end{aligned} \quad (11)$$

From the second identity of (11), we get

$$P_\varepsilon(t, x) = \frac{m-1}{4} \int_R e^{|x-y|} u_\varepsilon^2(t, y) dy. \quad (12)$$

2.1. L^2 Estimates and Several Consequences. Several properties for the smooth function $u_{0,\varepsilon}$ are given in the following Lemma.

Lemma 1. *The following estimates hold for any ε with $0 < \varepsilon < 1/4$ and $s \geq 0$:*

$$\begin{aligned}
\|u_{0,\varepsilon}\|_{L^2(R)} &\leq c_0 \|u_0\|_{L^2(R)}, \\
\|u_{0,\varepsilon}\|_{L^1(R)} &\leq c_0 \|u_0\|_{L^1(R)}, \\
\|u_{0,\varepsilon}\|_{BV(R)} &\leq c_0 \|u_0\|_{BV(R)}, \\
\|u_{0,\varepsilon}\|_{L^p(R)} &\leq c_0 \|u_0\|_{L^p(R)} \quad \text{for } 1 \leq p \leq \infty, \\
u_{0,\varepsilon} &\longrightarrow u_0 \quad (\varepsilon \longrightarrow 0) \text{ in } L^p(R) \text{ for } 1 \leq p \leq \infty, \\
\|u_{0,\varepsilon}\|_{H^q} &\leq c_0 \|u_0\|_{H^s}, \quad \text{if } q \leq s,
\end{aligned} \tag{13}$$

where c_0 is a constant independent of ε .

The proof of Lemma 1 is similar to that of Lemma 5 presented in [14]. Here we omit it.

Lemma 2. *Provided that $u_0 \in L^2(R)$, for any fixed $\varepsilon > 0$, there exists a unique global smooth solution $u_\varepsilon = u_\varepsilon(t, x)$ to the Cauchy problem (11) belonging to $C([0, \infty); H^s(R))$ with $s \geq 0$.*

Proof. We omit the proof since it is similar to the one found in [16] or [17] by using $u_{0,\varepsilon} \in C^\infty(R)$. $\square$

Lemma 3. *Assume that $u_0 \in L^2(R)$ holds and u_ε is a solution of problem (10). Then, the following bounds hold for any $t \geq 0$:*

$$\|u_\varepsilon\|_{L^2} \leq c_0 \|u_0\|_{L^2}, \tag{14}$$

$$\sqrt{\varepsilon} \|\partial_x u_\varepsilon\|_{L^2} \leq c_0 \|u_0\|_{L^2}, \tag{15}$$

where c_0 is a positive constant independent of ε and t .

Proof. Letting

$$mv_\varepsilon - \partial_{xx}^2 v_\varepsilon = u_\varepsilon \tag{16}$$

derives

$$v_\varepsilon = (m - \partial_{xx}^2)^{-1} u_\varepsilon. \tag{17}$$

Multiplying the first equation of problem (11) by $v_\varepsilon - \partial_{xx}^2 v_\varepsilon$ and integrating over R yield

$$\begin{aligned}
&\int_R \partial_t u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx - \varepsilon \int_R \partial_{xx}^2 u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&= - \int_R u_\varepsilon \partial_x u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&\quad - \int_R \partial_x P_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx.
\end{aligned} \tag{18}$$

For the left-hand side of this identity, using (16), we get

$$\begin{aligned}
&\int_R \partial_t u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx - \varepsilon \int_R \partial_{xx}^2 u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&= \int_R (m \partial_t v_\varepsilon - \partial_{txx}^3 v_\varepsilon) (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&\quad - \varepsilon \int_R (m \partial_{xx}^2 v_\varepsilon - \partial_{xxxx}^4 v_\varepsilon) (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&= \int_R (mv_\varepsilon \partial_t v_\varepsilon - v_\varepsilon \partial_{txx}^3 v_\varepsilon - m \partial_t v_\varepsilon \partial_{xx}^2 v_\varepsilon \\
&\quad + \partial_{xx}^2 v_\varepsilon \partial_{txx}^3 v_\varepsilon) dx \\
&\quad - \varepsilon \int_R (mv_\varepsilon \partial_{xx}^2 v_\varepsilon - m (\partial_{xx}^2 v_\varepsilon)^2 - v_\varepsilon \partial_{xxxx}^4 v_\varepsilon \\
&\quad + \partial_{xx}^2 v_\varepsilon \partial_{xxxx}^4 v_\varepsilon) dx \\
&= \int_R (mv_\varepsilon \partial_t v_\varepsilon - (m+1) v_\varepsilon \partial_{txx}^3 v_\varepsilon + \partial_{xx}^2 v_\varepsilon \partial_{txx}^3 v_\varepsilon) dx \\
&\quad - \varepsilon \int_R (mv_\varepsilon \partial_{xx}^2 v_\varepsilon - (m+1) v_\varepsilon \partial_{xxxx}^4 v_\varepsilon \\
&\quad + \partial_{xx}^2 v_\varepsilon \partial_{xxxx}^4 v_\varepsilon) dx \\
&= \int_R (mv_\varepsilon \partial_t v_\varepsilon + (m+1) \partial_x v_\varepsilon \partial_{tx}^2 v_\varepsilon + \partial_{xx}^2 v_\varepsilon \partial_{txx}^3 v_\varepsilon) dx \\
&\quad - \varepsilon \int_R (-m \partial_x v_\varepsilon \partial_x v_\varepsilon - (m+1) \partial_{xx}^2 v_\varepsilon \partial_{xx}^2 v_\varepsilon \\
&\quad - \partial_{xxx}^3 v_\varepsilon \partial_{xxx}^3 v_\varepsilon) dx \\
&= \frac{1}{2} \frac{d}{dt} \int_R (mv_\varepsilon^2 + (m+1) (\partial_x v_\varepsilon)^2 + (\partial_{xx}^2 v_\varepsilon)^2) dx \\
&\quad + \varepsilon \int_R (m (\partial_x v_\varepsilon)^2 + (m+1) (\partial_{xx}^2 v_\varepsilon)^2 + (\partial_{xxx}^3 v_\varepsilon)^2) dx.
\end{aligned} \tag{19}$$

For the right-hand side of (18), we conclude

$$\begin{aligned}
&- \int_R u_\varepsilon \partial_x u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx - \int_R \partial_x P_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&= - \int_R u_\varepsilon \partial_x u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
&\quad + \int_R (P_\varepsilon - \partial_{xx}^2 P_\varepsilon) \partial_x v_\varepsilon dx \\
&= - \int_R u_\varepsilon \partial_x u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx
\end{aligned}$$

$$\begin{aligned}
& + \frac{m-1}{2} \int_R u_\varepsilon^2 \partial_x v_\varepsilon dx \\
& = - \int_R u_\varepsilon \partial_x u_\varepsilon (v_\varepsilon - \partial_{xx}^2 v_\varepsilon) dx \\
& \quad - (m-1) \int_R u_\varepsilon \partial_x u_\varepsilon v_\varepsilon dx \\
& = \int_R (-m u_\varepsilon \partial_x u_\varepsilon v_\varepsilon + u_\varepsilon \partial_x u_\varepsilon \partial_{xx}^2 v_\varepsilon) dx \\
& = \int_R u_\varepsilon \partial_x u_\varepsilon (-m v_\varepsilon + \partial_{xx}^2 v_\varepsilon) dx \\
& = - \int_R u_\varepsilon^2 \partial_x u_\varepsilon dx \\
& = 0,
\end{aligned} \tag{20}$$

where we have used (16) and integration by parts.

From (18), (19), and (20), we have

$$\begin{aligned}
& m \|v_\varepsilon\|_{L^2}^2 + (m+1) \|\partial_x v_\varepsilon\|_{L^2}^2 + \|\partial_{xx}^2 v_\varepsilon\|_{L^2}^2 \\
& \quad + 2\varepsilon \int_R \left(m \|\partial_x v_\varepsilon\|_{L^2}^2 + (m+1) \|\partial_{xx}^2 v_\varepsilon\|_{L^2}^2 \right. \\
& \quad \quad \left. + \|\partial_{xxx}^3 v_\varepsilon\|_{L^2}^2 \right) dx \\
& = m \|v_\varepsilon(0, \cdot)\|_{L^2}^2 + (m+1) \|\partial_x v_\varepsilon(0, \cdot)\|_{L^2}^2 \\
& \quad + \|\partial_{xx}^2 v_\varepsilon(0, \cdot)\|_{L^2}^2.
\end{aligned} \tag{21}$$

From (17), we obtain

$$\begin{aligned}
& \|v_\varepsilon(0, \cdot)\|_{L^2}, \|\partial_x v_\varepsilon(0, \cdot)\|_{L^2}, \|\partial_{xx}^2 v_\varepsilon(0, \cdot)\|_{L^2} \\
& \leq c_0 \|u_{0,\varepsilon}\|_{L^2} \leq c_0 \|u_0\|_{L^2}.
\end{aligned} \tag{22}$$

It follows from (16) that

$$\begin{aligned}
& \|u_\varepsilon(t, \cdot)\|_{L^2(R)}^2 \\
& = \int_R (-\partial_{xx}^2 v_\varepsilon + m v_\varepsilon)^2 dx \\
& = \int_R (\partial_{xx}^2 v_\varepsilon)^2 dx - 2m \int_R v_\varepsilon \partial_{xx}^2 v_\varepsilon dx \\
& \quad + m^2 \int_R v_\varepsilon^2 dx \\
& = \int_R (\partial_{xx}^2 v_\varepsilon)^2 dx + 2m \int_R (\partial_x v_\varepsilon)^2 dx \\
& \quad + m^2 \int_R v_\varepsilon^2 dx.
\end{aligned} \tag{23}$$

Using (16), (21), and Lemma 1 derives that

$$\begin{aligned}
& \|u_\varepsilon\|_{L^2}^2 \\
& \leq 2 \|\partial_{xx}^2 v_\varepsilon\|_{L^2}^2 + 2m^2 \|v_\varepsilon\|_{L^2}^2 \\
& \leq \max(2, 2m) \left(m \|v_\varepsilon\|_{L^2}^2 + (m+1) \|\partial_x v_\varepsilon\|_{L^2}^2 \right. \\
& \quad \left. + \|\partial_{xx}^2 v_\varepsilon\|_{L^2}^2 \right) \\
& \leq \max(2, 2m) \left(m \|v_\varepsilon(0, \cdot)\|_{L^2}^2 + (m+1) \|\partial_x v_\varepsilon(0, \cdot)\|_{L^2}^2 \right. \\
& \quad \left. + \|\partial_{xx}^2 v_\varepsilon(0, \cdot)\|_{L^2}^2 \right) \\
& \leq c_0 \|u_{0,\varepsilon}\|_{L^2}^2 \\
& \leq c_0 \|u_0\|_{L^2}^2, \\
& \varepsilon \|\partial_x u_\varepsilon\|_{L^2}^2 \\
& \leq 2\varepsilon \|\partial_{xxx}^3 v_\varepsilon\|_{L^2}^2 + 2m^2 \varepsilon \|\partial_x v_\varepsilon\|_{L^2}^2 \\
& \leq \varepsilon \max(2, 2m) \left(m \|\partial_x v_\varepsilon\|_{L^2}^2 + (m+1) \|\partial_{xx}^2 v_\varepsilon\|_{L^2}^2 \right. \\
& \quad \left. + \|\partial_{xxx}^3 v_\varepsilon\|_{L^2}^2 \right) \\
& \leq \max(2, 2m) \left(m \|v_\varepsilon(0, \cdot)\|_{L^2}^2 + (m+1) \|\partial_x v_\varepsilon(0, \cdot)\|_{L^2}^2 \right. \\
& \quad \left. + \|\partial_{xx}^2 v_\varepsilon(0, \cdot)\|_{L^2}^2 \right) \\
& \leq c_0 \|u_{0,\varepsilon}\|_{L^2}^2 \\
& \leq c_0 \|u_0\|_{L^2}^2.
\end{aligned} \tag{24}$$

The proof of Lemma 3 follows from (24). $\square$

We give some bounds on the nonlocal term P_ε , in which all are consequences of the L^2 bound in Lemma 3.

Lemma 4. Assume that $u_0 \in L^2(R)$ holds. Then,

$$P_\varepsilon \geq 0 \quad \text{for } m \geq 1, \tag{25}$$

$$\|P_\varepsilon(t, \cdot)\|_{L^1(R)}, \quad \|\partial_x P_\varepsilon(t, \cdot)\|_{L^1(R)} \leq c_0 \frac{|m-1|}{2} \|u_0\|_{L^2}^2, \tag{26}$$

$$\|P_\varepsilon\|_{L^\infty(R_+ \times R)}, \quad \|\partial_x P_\varepsilon\|_{L^\infty(R_+ \times R)} \leq c_0 \frac{|m-1|}{2} \|u_0\|_{L^2}^2, \tag{27}$$

$$\|\partial_{xx}^2 P_\varepsilon(t, \cdot)\|_{L^1(R)} \leq c_0 \frac{|m-1|}{2} \|u_0\|_{L^2}^2, \tag{28}$$

where c_0 is a constant independent of ε and t .

Proof. Using (11), we get

$$P_\varepsilon(t, x) = \frac{m-1}{4} \int_R e^{-|x-y|} (u_\varepsilon(t, y))^2 dy, \quad (29)$$

$$\partial_x P_\varepsilon(t, x) = \frac{m-1}{4} \int_R e^{-|x-y|} \text{sign}(y-x) (u_\varepsilon(t, y))^2 dy. \quad (30)$$

From (29), we obtain (25). Using (14) and the Tonelli theorem, we have

$$\begin{aligned} & \int_R |P_\varepsilon(t, x)| dx \\ & \leq \frac{|m-1|}{4} \int_R \left(\int_R e^{-|x-y|} dx \right) (u_\varepsilon(t, y))^2 dy \\ & \leq c_0 \frac{|m-1|}{2} \int_R (u_\varepsilon(t, y))^2 dy \\ & \leq c_0 \frac{|m-1|}{2} \|u_0\|_{L^2}^2, \\ & \int_R |\partial_x P_\varepsilon(t, x)| dx \\ & \leq \frac{|m-1|}{4} \int_R \left(\int_R e^{-|x-y|} dx \right) (u_\varepsilon(t, y))^2 dy \\ & \leq c_0 \frac{|m-1|}{2} \int_R (u_\varepsilon(t, y))^2 dy \\ & \leq c_0 \frac{|m-1|}{2} \|u_0\|_{L^2}^2, \\ & |P_\varepsilon(t, x)| \leq \frac{|m-1|}{4} \int_R (u_\varepsilon(t, y))^2 dy \\ & \leq c_0 \frac{|m-1|}{4} \|u_0\|_{L^2}^2, \\ & |\partial_x P_\varepsilon(t, x)| \leq \frac{|m-1|}{4} \int_R (u_\varepsilon(t, y))^2 dy \\ & \leq c_0 \frac{|m-1|}{4} \|u_0\|_{L^2}^2. \end{aligned} \quad (31)$$

It follows from (31) that (26) and (27) hold. Using the second identity of problem (11), Lemma 3, and (26), we obtain (28). $\square$

Lemma 5. If $u_0 \in L^1(R) \cap L^2(R)$, it holds that

$$\|u_\varepsilon(t, \cdot)\|_{L^1} \leq \|u_0\|_{L^1} + c_0 t \|u_0\|_{L^2}^2. \quad (32)$$

Proof. Let functions η and $q : R \rightarrow R$ be such that $q'(u) = u\eta'(u)$. Multiplying the first equation in (11) with $\eta'(u_\varepsilon)$ gives rise to

$$\begin{aligned} & \partial_t \eta(u_\varepsilon) + \partial_x q(u_\varepsilon) + \eta'(u_\varepsilon) \partial_x P_\varepsilon \\ & = \varepsilon \eta'(u_\varepsilon) \partial_{xx}^2 u_\varepsilon \\ & = \varepsilon \partial_{xx}^2 \eta(u_\varepsilon) - \varepsilon \eta''(u_\varepsilon) (\partial_x u_\varepsilon)^2. \end{aligned} \quad (33)$$

Choosing $\eta(u) = |u|$ (modulo an approximation argument, see [16]) and then integrating the resulting equation over R yield

$$\frac{d}{dt} \int_R |u_\varepsilon| dx \leq \int_R |\text{sign}(u_\varepsilon) \partial_x P_\varepsilon| dx. \quad (34)$$

Using (26), we get

$$\frac{d}{dt} \|u_\varepsilon(t, \cdot)\|_{L^1} \leq c_0 \frac{|m-1|}{2} \|u_0\|_{L^2}^2, \quad (35)$$

from which we have (32). $\square$

2.2. BV and L^∞ Estimates. In this subsection we establish several supplementary estimates for the viscous approximations, which also are consequences of the L^2 bound in Lemma 3. In particular, we prove that the sequence $\{u_\varepsilon\}_{\varepsilon>0}$ is bounded in BV, which yields strong compactness of this sequence. To this end, we need to assume that $u_0 \in L^2(R)$ and $u_0 \in BV(R)$.

Lemma 6. Assume that $u_0 \in L^2(R)$ and $u_0 \in BV(R)$ hold. Then,

$$\|\partial_x u_\varepsilon(t, \cdot)\|_{L^1(R)} \leq c_0 (|u_0|_{BV} + t \|u_0\|_{L^2}^2), \quad t \geq 0. \quad (36)$$

Proof. Setting $q_\varepsilon := \partial_x u_\varepsilon$, we know that q_ε satisfies the equation

$$\partial_t q_\varepsilon + u_\varepsilon \partial_x q_\varepsilon + q_\varepsilon^2 + \partial_{xx}^2 P_\varepsilon = \varepsilon \partial_{xx}^2 q_\varepsilon. \quad (37)$$

If $\eta = \eta(u)$ and $q : R \rightarrow R$ satisfies $q'(u) = u\eta'(u)$, using the chain rule yields

$$\begin{aligned} & \partial_t \eta(q_\varepsilon) + \partial_x (u_\varepsilon q_\varepsilon) - q_\varepsilon \eta(q_\varepsilon) \\ & + \eta'(q_\varepsilon) q_\varepsilon^2 + \eta'(u_\varepsilon) \partial_{xx}^2 P_\varepsilon \\ & = \varepsilon \partial_{xx}^2 \eta(q_\varepsilon) - \varepsilon \eta''(q_\varepsilon) (\partial_x q_\varepsilon)^2. \end{aligned} \quad (38)$$

Choosing $\eta(u) = |u|$ (modulo an approximation argument) and then integrating the resulting equation over R give rise to

$$\frac{d}{dt} \int_R |\partial_x u_\varepsilon| dx \leq \int_R |\partial_{xx}^2 P_\varepsilon| dx. \quad (39)$$

Using (28), we have

$$\frac{d}{dt} \int_R |\partial_x u_\varepsilon| dx \leq c_0 \|u_0\|_{L^2(R)}^2, \quad (40)$$

from which we obtain (36). $\square$

Lemma 7. Assume that $u_0 \in BV(R) \cap L^2(R)$ holds. Then,

$$\begin{aligned} & \|u_\varepsilon(t, \cdot)\|_{L^\infty(R)} \leq c_0 (|u_0|_{BV(R)} + t \|u_0\|_{L^2(R)}^2), \\ & t \geq 0, \end{aligned} \quad (41)$$

$$\begin{aligned} & \|\partial_{xx}^2 P_\varepsilon(t, \cdot)\|_{L^\infty} \\ & \leq c_0 \left[\|u_0\|_{L^2(R)}^2 + (|u_0|_{BV(R)} + t \|u_0\|_{L^2(R)}^2)^2 \right], \end{aligned} \quad (42)$$

where c_0 is independent of ε and t .

Proof. Using

$$|u_\varepsilon(t, \cdot)| \leq \int_R |\partial_x u_\varepsilon(t, y)| dy = \|u_\varepsilon(t, \cdot)\|_{BV(R)} \quad (43)$$

and Lemma 6 derives (41). Using (27), (41), and the second equation of problem (11), we obtain inequality (42). $\square$

2.3. L^{2N} Estimate for Nature Number $N \geq 2$. Next we prove that the viscous approximations are bounded in $L^{2N}(R)$ for any nature number $N \geq 2$. From Lemmas 6 and 7, if $u_0 \in BV(R) \cap L^2(R)$, we have the inequality $u_{0,\varepsilon} \in L^\infty(R)$ from which we derive

$$u_{0,\varepsilon} \in L^2(R) \cap L^{2N}(R). \quad (44)$$

Lemma 8. Assume that $u_0 \in BV(R) \cap L^2(R)$ holds. For any $0 < \varepsilon < 1/4$, it has

$$\begin{aligned} \|u_\varepsilon(t, \cdot)\|_{L^{2N}(R)} &\leq c_0(1 + \|u_0\|_{L^2}^2)^2 t \\ &\quad + \|u_\varepsilon(0, \cdot)\|_{L^{2N}(R)}^{2N} e^{c_0(1 + \|u_0\|_{L^2}^2)^2 t}. \end{aligned} \quad (45)$$

Proof. Choosing $\eta(u) = (1/2N)u^{2N}$ in (33), writing

$$\begin{aligned} \partial_t \eta(u_\varepsilon) + \partial_x q(u_\varepsilon) + \eta'(u_\varepsilon) \partial_x P_\varepsilon \\ = \varepsilon \partial_{xx}^2 \eta(u_\varepsilon) - \varepsilon \eta''(u_\varepsilon) (\partial_x u_\varepsilon)^2 \\ = \varepsilon \eta'(u_\varepsilon) \partial_{xx}^2 u_\varepsilon \\ = \varepsilon u_\varepsilon^{2N-1} \partial_{xx}^2 u_\varepsilon, \end{aligned} \quad (46)$$

and integrating (46) over R gives rise to

$$\begin{aligned} \frac{1}{2N} \frac{d}{dt} \|u_\varepsilon(t, \cdot)\|_{L^{2N}}^{2N} &= - \int_R u_\varepsilon^{2N-1} \partial_x P_\varepsilon dx \\ &\quad + \varepsilon \int_R u_\varepsilon^{2N-1} \partial_{xx}^2 u_\varepsilon dx. \end{aligned} \quad (47)$$

Integration by parts shows

$$\begin{aligned} \varepsilon \int_R u_\varepsilon^{2N-1} \partial_{xx}^2 u_\varepsilon dx \\ = -(2N-1) \varepsilon \int_R u_\varepsilon^{(2N-2)} (\partial_x u_\varepsilon)^2 dx \leq 0. \end{aligned} \quad (48)$$

Letting $p_1 = 2N/(2N-1)$, $q_1 = 2N$, we have $1/p_1 + 1/q_1 = 1$. Using Hölder's inequality, (26), and (27), we obtain

$$\begin{aligned} \left| \int_R u_\varepsilon^{2N-1} \partial_x P_\varepsilon dx \right| \\ \leq \left(\int_R u_\varepsilon^{2N} dx \right)^{(2N-1)/2N} \left(\int_R (\partial_x P_\varepsilon)^{2N} dx \right)^{1/2N} \\ \leq \left(1 + \int_R u_\varepsilon^{2N} dx \right)^{(2N-1)/2N} \\ \times \left(\|\partial_x P_\varepsilon\|_{L^\infty(R_+ \times R)}^{2N-1} \|\partial_x P_\varepsilon\|_{L^1(R)}^{1/2N} \right) \\ \leq \left(1 + \int_R u_\varepsilon^{2N} dx \right) (1 + \|\partial_x P_\varepsilon\|_{L^\infty(R_+ \times R)}) \|\partial_x P_\varepsilon\|_{L^1(R)}^{1/2N} \\ \leq c_0 (1 + \|u_0\|_{L^2}^2) (1 + \|u_0\|_{L^2}^2)^{1/2N} \left(1 + \int_R u_\varepsilon^{2N} dx \right) \\ \leq c_0 (1 + \|u_0\|_{L^2}^2)^2 \left(1 + \int_R u_\varepsilon^{2N} dx \right). \end{aligned} \quad (49)$$

Using (47)–(49) gives rise to

$$\begin{aligned} \frac{1}{2N} \frac{d}{dt} \|u_\varepsilon(t, \cdot)\|_{L^{2N}}^{2N} \\ \leq c_0 (1 + \|u_0\|_{L^2}^2)^2 (1 + \|u_\varepsilon(t, \cdot)\|_{L^{2N}}^{2N}), \end{aligned} \quad (50)$$

from which we obtain (45) by Gronwall's inequality. $\square$

2.4. Oleinik Type Estimate

Lemma 9 (Oleinik type estimate). Assume that $u_0 \in L^2(R) \cap BV(R)$. Then, for each $t \in (0, T]$ with T being fixed,

$$\partial_x u_\varepsilon(t, x) \leq \frac{1}{t} + K_T, \quad x \in R, \quad (51)$$

where $K_T = c_0 \{ \|u_0\|_{L^2(R)}^2 + (\|u_0\|_{BV(R)} + T \|u_0\|_{L^2(R)}^2)^{1/2} \}$.

Proof. Setting $q_\varepsilon := \partial_x u_\varepsilon$, it follows from (11) and (42) that

$$\partial_t q_\varepsilon + u_\varepsilon \partial_x q_\varepsilon + q_\varepsilon^2 - \varepsilon \partial_{xx}^2 q_\varepsilon = -\partial_{xx}^2 P_\varepsilon \leq K_T^2. \quad (52)$$

Considering the ordinary differential equation

$$\frac{df}{dt} + f^2 = K_T^2 \quad (53)$$

and using comparing theorem, we have

$$\partial_x u_\varepsilon(t, x) \leq \frac{1}{t} + K_T, \quad x \in R, \quad (54)$$

which completes the proof. $\square$

3. Existence in $L^1(R) \cap BV(R)$

Using the estimates established in Section 2, we will show the existence of weak solutions to problem (6) under the assumption $u_0 \in L^1(R) \cap BV(R)$.

We state the concepts of weak solutions.

Definition 10 (weak solution). We call a function $u : R_+ \times R \rightarrow R$ a weak solution of the Cauchy problem (6) provided that

- (i) $u \in L^\infty(R_+; L^2(R))$ and
- (ii) $\partial_t u + \partial_x(u^2/2) + \partial_x P^u(t, x) = 0$ in $D'([0, \infty) \times R)$; that is, for all $\phi \in C_0^\infty([0, \infty) \times R)$, there holds the identity

$$\begin{aligned} \int_{R^+} \int_R \left(u \partial_t \phi + \frac{u^2}{2} \partial_x \phi - \partial_x P^u \phi \right) dx dt \\ + \int_R u_0(x) \phi(0, x) dx = 0, \end{aligned} \quad (55)$$

where

$$\begin{aligned} P^u(t, x) &= G_1 * \left(\frac{m-1}{2} u^2 \right)(t, x) \\ &= \frac{m-1}{4} \int_R e^{-|x-y|} (u(t, y))^2 dy. \end{aligned} \quad (56)$$

Remark 11. It follows from part (i) of Definition 10 that $u \in L^1((0, T) \times R)$ for any $T > 0$ and $\partial_x P^u \in L^\infty(R_+ \times R)$ (see Lemma 5). Therefore, (55) makes sense.

We assume that

$$\begin{aligned} \|u(t_2, \cdot) - u(t_1, \cdot)\|_{L^1(R)} &\leq C_T |t_2 - t_1| \\ \text{for } \forall t_1, t_2 &\in [0, T], \end{aligned} \quad (57)$$

where $C_T = c_0 [\|u_0\|_{L^2(R)}^2 + (\|u_0\|_{BV(R)} + t\|u_0\|_{L^2(R)})^2]$. Therefore, we have

$$\lim_{t \rightarrow 0^+} \|u(t, \cdot) - u_0\|_{L^1(R)} = 0. \quad (58)$$

Our main results are summarized in the following Theorem.

Theorem 12. *Provided that $u_0 \in L^1(R) \cap L^2(R) \cap BV(R)$ and the solution $u(t, x)$ satisfies (57), then there exists a weak solution to the Cauchy problem (6). The weak solution u satisfies the following estimates for any $t \in (0, T)$:*

$$\begin{aligned} \|u(t, \cdot)\|_{L^1(R)} &\leq \|u_0\|_{L^1} + c_0 t \|u_0\|_{L^2}^2, \\ \|u(t, \cdot)\|_{BV}, \quad \|u(t, x)\|_{L^\infty} &\leq c_0 (\|u_0\|_{BV(R)} + t\|u_0\|_{L^2(R)}^2), \\ \|u_\varepsilon(t, \cdot)\|_{L^{2N}(R)} &\leq c_0 (1 + \|u_0\|_{L^2}^2) t \\ &\quad + \|u(0, \cdot)\|_{L^{2N}(R)}^{2N} e^{c_0(1 + \|u_0\|_{L^2}^2)t}, \end{aligned} \quad (59)$$

where c_0 is a positive constant independent of t and ε .

The following Oleinik type estimate holds for a.e. $(t, x) \in (0, T] \times R$:

$$\partial_x u(t, x) \leq \frac{1}{t} + K_T, \quad (60)$$

where

$$\begin{aligned} K_T &= c_0 \left\{ \|u_0\|_{L^2(R)}^2 \right. \\ &\quad \left. + (\|u_0\|_{BV(R)} + T\|u_0\|_{L^2(R)}^2)^2 \right\}^{1/2}. \end{aligned} \quad (61)$$

This theorem is an immediate consequence of Theorem 13 and results are presented in Section 2.

Theorem 13 (existence). *Assume that $u_0 \in L^1(R) \cap L^2(R) \cap BV(R)$ and the solution $u(t, x)$ satisfies (57). Then, there exists at least one weak solution to problem (6).*

Proof. Using the estimates obtained in Section 2, we take a standard argument to see that there exists a sequence of strictly positive numbers $\{\varepsilon_k\}_{k=1}^\infty$ tending to zero such that as $k \rightarrow \infty$

$$\begin{aligned} u_{\varepsilon_k} &\longrightarrow u \quad \text{a.e. in } R_+ \times R, \\ u_{\varepsilon_k} &\longrightarrow u \quad \text{in } L^1_{\text{loc}}(R_+ \times R). \end{aligned} \quad (62)$$

The previous estimates in Section 2 imply immediately that the limit function u satisfies (59)-(60).

Let us now prove that as $k \rightarrow \infty$

$$P_{\varepsilon_k} \longrightarrow P^u, \quad \partial_x P_{\varepsilon_k} \longrightarrow \partial_x P^u \quad \text{in } L^1((0, T) \times R) \quad \forall T > 0, \quad (63)$$

which follows from the following calculations:

$$\begin{aligned} \|P_{\varepsilon_k} - P^u\|_{L^1((0, T) \times R)} &\leq \left(\frac{|m-1|}{4} \right) \\ &\quad \times \int_0^T \int_R \left(\int_R e^{-|x-y|} \right. \\ &\quad \left. \times \left| (u_{\varepsilon_k}(t, y))^2 - (u(t, y))^2 \right| dy \right) dx dt \\ &\leq \left(\frac{|m-1|}{2} \right) \int_0^T \left(\int_R |u_{\varepsilon_k}(t, y) - u(t, y)| \right. \\ &\quad \left. \times |u_{\varepsilon_k}(t, y) + u(t, y)| dy \right) dt \\ &\leq C_T \int_0^T \int_R |u_{\varepsilon_k}(t, y) - u(t, y)| dy dt \longrightarrow 0 \\ &\quad \text{as } k \longrightarrow \infty, \\ \|\partial_x P_{\varepsilon_k} - \partial_x P^u\|_{L^1((0, T) \times R)} &\leq \left(\frac{|m-1|}{4} \right) \\ &\quad \times \int_0^T \int_R \left(\int_R e^{-|x-y|} \right. \\ &\quad \left. \times \left| (u_{\varepsilon_k}(t, y))^2 - (u(t, y))^2 \right| dy \right) dx dt \end{aligned}$$

$$\begin{aligned}
&\leq \left(\frac{|m-1|}{2}\right) \int_0^T \left(\int_R |u_{\varepsilon_k}(t, y) - u(t, y)| \right. \\
&\quad \left. \times |u_{\varepsilon_k}(t, y) + u(t, y)| dy \right) dt \\
&\leq C_T \int_0^T \int_R |u_{\varepsilon_k}(t, y) - u(t, y)| dy dt \longrightarrow 0 \\
&\quad \text{as } k \longrightarrow \infty.
\end{aligned} \tag{64}$$

From (62) to (64), we complete the proof. $\square$

Conflict of Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

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References

- [1] A. Constantin and D. Lannes, "The hydrodynamical relevance of the Camassa-Holm and Degasperis-Procesi equations," *Archive for Rational Mechanics & Analysis*, vol. 192, no. 1, pp. 165–186, 2009.
- [2] A. Degasperis and M. Procesi, "Asymptotic integrability," in *Symmetry & Perturbation Theory*, vol. 1, no. 1, pp. 23–37, World Scientific, Singapore, 1999.
- [3] J. Escher, Y. Liu, and Z. Y. Yin, "Global weak solutions and blow-up structure for the Degasperis-Procesi equation," *Journal of Functional Analysis*, vol. 241, no. 2, pp. 457–485, 2006.
- [4] Y. Fu, Y. Liu, and C. Z. Qu, "On the blow-up structure for the generalized periodic Camassa-Holm and Degasperis-Procesi equations," *Journal of Functional Analysis*, vol. 262, no. 7, pp. 3125–3158, 2012.
- [5] V. O. Vakhnenko and E. J. Parkes, "Periodic and solitary-wave solutions of the Degasperis-Procesi equation," *Chaos, Solitons & Fractals*, vol. 20, no. 5, pp. 1059–1073, 2004.
- [6] H. Lundmark and J. Szmigielski, "Multi-peakon solutions of the Degasperis-Procesi equation," *Inverse Problems*, vol. 19, no. 6, pp. 1241–1245, 2003.
- [7] D. D. Holm and M. F. Staley, "Wave structure and nonlinear balances in a family of evolutionary PDEs," *SIAM Journal on Applied Dynamical Systems*, vol. 2, no. 3, pp. 323–380, 2003.
- [8] Z. W. Lin and Y. Liu, "Stability of peakons for the Degasperis-Procesi equation," *Communications on Pure & Applied Mathematics*, vol. 62, no. 1, pp. 125–146, 2009.
- [9] Z. Y. Yin, "Global weak solutions for a new periodic integrable equation with peakon solutions," *Journal of Functional Analysis*, vol. 212, no. 1, pp. 182–194, 2004.
- [10] Y. Matsuno, "Multisoliton solutions of the Degasperis-Procesi equation and their peakon limit," *Inverse Problems*, vol. 21, no. 5, pp. 1553–1570, 2005.
- [11] D. Henry, "Infinite propagation speed for the Degasperis-Procesi equation," *Journal of Mathematical Analysis & Applications*, vol. 311, no. 2, pp. 755–759, 2005.
- [12] O. G. Mustafa, "A note on the Degasperis-Procesi equation," *Journal of Nonlinear Mathematical Physics*, vol. 12, no. 1, pp. 10–14, 2005.
- [13] Y. Liu and Z. Y. Yin, "Global existence and blow-up phenomena for the Degasperis-Procesi equation," *Communications in Mathematical Physics*, vol. 267, no. 3, pp. 801–820, 2006.
- [14] S. Y. Lai and Y. H. Wu, "The local well-posedness and existence of weak solutions for a generalized Camassa-Holm equation," *Journal of Differential Equations*, vol. 248, no. 8, pp. 2038–2063, 2010.
- [15] Y. Zhou, "Blow-up of solutions to the DGH equation," *Journal of Functional Analysis*, vol. 250, no. 1, pp. 227–248, 2007.
- [16] G. M. Coclite and K. H. Karlsen, "On the well-posedness of the Degasperis-Procesi equation," *Journal of Functional Analysis*, vol. 233, no. 1, pp. 60–91, 2006.
- [17] G. M. Coclite, K. H. Karlsen, and H. Holden, "Wellposedness for a parabolic-elliptic system," *Discrete & Continuous Dynamical Systems*, vol. 13, no. 3, pp. 659–682, 2005.