Research Article

Manufacturing Data Uncertainties Propagation Method in Burn-Up Problems

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A nuclear data-based uncertainty propagation methodology is extended to enable propagation of manufacturing/technological data (TD) uncertainties in a burn-up calculation problem, taking into account correlation terms between Boltzmann and Bateman terms. The methodology is applied to reactivity and power distributions in a Material Testing Reactor benchmark. Due to the inherent statistical behavior of manufacturing tolerances, Monte Carlo sampling method is used for determining output perturbations on integral quantities. A global sensitivity analysis (GSA) is performed for each manufacturing parameter and allows identifying and ranking the influential parameters whose tolerances need to be better controlled. We show that the overall impact of some TD uncertainties, such as uranium enrichment, or fuel plate thickness, on the reactivity is negligible because the different core areas induce compensating effects on the global quantity. However, local quantities, such as power distributions, are strongly impacted by TD uncertainty propagations. For isotopic concentrations, no clear trends appear on the results.

1. Introduction

Sensitivity analysis (SA) methods are invaluable tools allowing the study of how the uncertainty in the model output relies on the different sources of uncertainty in the model inputs [1]. As a ranking method, it can be used to determine the most contributing input variables to an output behavior as the noninfluential inputs or clarify some correlated effects within the model. The objectives of SA are numerous; one can mention model verification and understanding, model simplification, and factor prioritization [2]. Finally, the SA helps in the validation of a computer code, guidance research efforts, or the justification in terms of system design safety. There is a high amount of literature on procedures and techniques for SA. The main outcomes can be found in [3, 4]. There are many possible uses of SA, described within the categories of decision support, communication, increased understanding, or system quantification and model development. Many different approaches to SA are described

elsewhere, varying in the experimental design used and in the way results are processed. An example of manufacturing uncertainties propagation is described in [5].

Tolerance analysis is also becoming an important tool for nuclear engineering design. This seemingly arbitrary task of assigning tolerances can have a large effect on the cost and performance of manufactured products, such as fuel design and fabrication. However, the fact of propagating tolerances instead of uncertainties does not lead to a representative approach of the errors because; in this case, only a bias is taken into account. It is then imperative to understand what kind of physical data creates and propagates uncertainties on the neutronics parameters for both safety and performance reasons. In Material Testing Reactors (MTR), the performance parameters can be core fuel cycle or isotope production. Therefore, it is necessary to calculate isotopic concentrations uncertainties in the reactor core.

We focus in this paper on technological data propagation, with a special attention to uranium enrichment and plate

thickness, as example of manufacturing uncertainties propagation. In a general *engineering* framework much broader than nuclear engineering only, as tolerances affect both cost and quality of a product, tolerancing is now considered as being a critical engineering design function. As such, tolerance allocation is a significant task that deserves considerable attention. The current situation is a compromise between designers, who usually specify tight tolerances to ensure high quality, and manufacturers, who prefer loose tolerances to reduce manufacturing cost [6]. So, adequate (i.e., reasonable or well-balanced) tolerances must be achieved in order to both ensure the desired performance and ease the fabrication process.

In general, there are no specific guidelines for allocating tolerances for any component, but [4] quotes the following paragraph [7]: “The most common practice is to allocate some tolerance that seems appropriate on the basis of experience or intuition, and then conduct an analysis to ensure that the allocated tolerance suits the desired design function. In order to do this, the designer must be able to realize all possible effects of the tolerances specified, especially if universal interchangeability is one of the design goals. The effects of specified tolerances are generally analyzed by creating an analytical model that can predict the accumulation of tolerances in an assembly. Prediction of tolerance accumulation is necessary because critical fits, clearances, etc. are usually controlled by the accumulation of several component tolerances”.

After a reminder of the theoretical approach and the implementation of tolerance analysis in the MC propagation methodology and UQ in coupled Boltzmann/Bateman problem, a practical example is given for complete depletion calculation, based on a Material Testing Reactor (MTR) benchmark. This latter is described, and the associated tolerance data, based on an actual series of manufacturing feedback, are detailed. One will focus on two main technological parameters: uranium enrichment of the plates as well as their thicknesses.

The uncertainty propagation will be performed for two different integral quantities: the reactivity, that is, a more global parameter, and the power factor (i.e., the plate fission rate distribution), more sensitive to local variations. A particular focus on the concentrations of some important isotopes will also be made.

2. Evaluation of the Technological Uncertainties

The method used to evaluate the uncertainties comes from complete work performed in [8, 9].

The complete evaluation of propagated uncertainties on neutronics parameters requires a precise knowledge of both nuclear data and manufacturing uncertainties. If the primers are relatively well known and characterized through consistent covariance matrices, such as the latest ENDF/B-VII.1 [10] or COMAC [11], manufacturing uncertainties are sometimes sparse and often not taken into account in the UQ process. However, those values can be built by considering (supposed known) tolerances.

The statistical nature of uncertainty analysis naturally relies on the use of Monte Carlo sampling methodology. Monte Carlo sampling methods can be used to perform uncertainty propagation throughout the whole core calculation process. The manufacture of a technological item is simulated, for example, by creating a set of component dimensions with small random changes to simulate natural process variations. In this case, a Gaussian model can be selected as a statistical distribution of uncertainties, and tolerances can be chosen as variances values at 3σ , built by expert elicitation.

Next, the resulting assembly dimensions are calculated from the simulated set of component dimensions. The numbers of outliers that fall outside the specification limits are then counted. Sample sizes generally range between 5,000 and 100,000, based on the required accuracy of the simulation. The accuracy of Monte Carlo sampling increases with larger sample sizes. Obviously, the computational effort of large sample sizes can be significant, but Monte Carlo sampling offers many advantages because of its flexibility. It also allows the generation of a sample of uncertain inputs. We then obtain a sample corresponding to the outputs of the calculation code.

Of course, the best and more rigorous way is to get actual measurement of each series of manufacturing parameters that would allow building the propagated bias on integral parameters between the theoretical core (i.e., without tolerance) and the actual (i.e., *as built*) core. The measurement of each sample enables postulating a statistical model of its manufacturing uncertainty. This is the methodology used in the present study.

2.1. Benchmark Description. The benchmark used is the present paper is a Material Testing Reactor based on 20% enriched $^{235}\text{U}\text{U}_3\text{Si}_2\text{Al}$ fuel plates [9]. A unique type of assembly has been modelled to build the whole core. The benchmark does not contain absorbing assembly in order to simplify the calculation: the goal being to give orders of magnitude of the propagated uncertainties.

A fuel assembly is made of 22 1.27 mm thick Zircaloy plates (in green). Each plate contains a 0.51 mm thick $\text{U}_3\text{Si}_2\text{Al}$ fuel blade, called “the meat.” The blue elements of Figure 1 represent the water. The assembly stiffeners are made of aluminium.

The benchmark study is performed in 15 energy groups using the APOLLO2.8.3 deterministic code [12] based on a MOC (Method of Characteristics) calculation scheme [13]. The calculation is performed on a 2D quarter of core with ad hoc symmetries. The full Boltzmann/Bateman calculation scheme is described in [8]. Each plate is discretized in 8 sectors, to get a more precise estimation of the local power factor and the concentrations.

For the sake of the present work, two technological parameters and their associated tolerances will be studied. In the following, they will be noted UO2MB for “UO₂ mass balance” and PTh for “plate thickness.”

Perturbation at the beginning of the calculation enables assessing global sensitivities.

FIGURE 1: Geometric representation of the benchmark.

FIGURE 2: Statistical model of manufacturing uncertainty.

2.2. Statistical Distributions of Manufacturing Parameters. The statistical distributions are built by calculating standard deviations and average values of manufacturing parameters. We consider, in the following, that measurements of technological parameters are available. An example of this kind of measurement is presented on Figure 2. The means mentioned under each plot correspond to the ratio between theoretical and measurement averaged values. In some cases, inconsistencies are observed. One way for circumvent this inconsistency is to perform a calculation of both average and theoretical cores and compare the results on relevant integral values (k_{eff} , power distributions, etc.). This comparison gives a manufacturing bias (not calculated in this work) which can

further be taken into account in the uncertainty tabulation of the core. The red and green curves represent, respectively, the probability density function and the cumulative distribution function, normalized. The histogram reproduces the frequency of the values observed by measurement. If the distribution function of a manufacturing parameter is a known law, the best is to make the propagation of this law in the propagation process. If it is not, we can perform the propagation, considering that the parameter follows a Gaussian law. As a Gaussian represents the maximum entropy law, the propagation is then conservative. In our example, the UO2MB parameter is not exactly a Gaussian but we will consider it as such. In these two parameters, the uncertainty

on PTh is more important (2.5% at 1σ) than uncertainty on UO2MB (0.2% at 1σ). However, as we will see in the following, the impact of UO2MB is higher on neutronics uncertainties.

As shown in [8, 9], the response function is modelled by

$$f: \begin{cases} \mathcal{R}^n \longrightarrow \mathcal{R}^p \\ \underline{\kappa} = (\kappa_1, \dots, \kappa_n) \longrightarrow \underline{Y} = (Y_1, \dots, Y_p) = f(\underline{\kappa}), \end{cases} \quad (1)$$

where

$$\delta f(t) = \delta h(t) + \sum_{i \in \mathcal{S}} [\delta(h \circ N_i)(t)] \quad (2)$$

and different methods can be used to calculate sensitivity indices and uncertainties. The method developed in the following does not allow getting the derivative from (7) as $\partial h / \partial \kappa_k$ which is the sensitivity of the interest response h to the input κ_k . In fact, this would be possible if large CPU time is available on HPC clusters. For example, to calculate the sensitivity of reactivity to PTh would require 1078 full calculations (i.e., coupled Boltzmann/Bateman from 0 to 100,000 MW-d/T). The 49 assemblies of the benchmark are composed each of 22 plates. For UO2MB only, 8624 full calculations should be performed. This kind of calculation is impossible today in reasonable time. Then, calculations are performed for each parameter by sampling their distribution law. Five hundred calculations are performed for each technological parameter, enabling a reconstruction of an estimator of the output distribution laws. Hence, parameters that impact the neutronics values can be identified and ranked. Another difficulty arises here through the coupled calculation and the associated uncertainty propagation coming from direct and transmutation terms. Transmutation and direct effects can be decorrelated by doubling the number of calculations, with the methods coming from [9]. In this paper, we will use the MC method described therein with the goal to get a global sensitivity analysis. However, its major drawback is that it does not take into account correlations between input parameters. These correlations could be calculated using Pearson coefficients if the measurements are correlated. For example, if the manufacturer simultaneously measures both fuel size and mass, correlations between these two parameters can be extracted. The geometrical perturbations are done for each assembly independently, using the Salome tools (<http://www.salome-platform.org/>).

The sampling can be performed using a multidimensional Gaussian law whose probability function is given by

$$g(\underline{\zeta}) = \frac{1}{\sqrt{2\pi}} \frac{1}{\det(\underline{M}(\underline{\kappa}))} \cdot \exp \left[-\frac{1}{2} (\underline{\kappa} - \underline{\mu})^T [\underline{M}(\underline{\kappa})]^{-1} (\underline{\kappa} - \underline{\mu}) \right], \quad (3)$$

where $\text{IE}(\underline{\kappa}) = \underline{\mu} = (\mu_1, \dots, \mu_n)$ is the manufacturing data mathematical expectancies and $\underline{M}(\underline{\kappa})$ is the covariance matrix of inputs $\underline{\kappa}$.

R (<https://www.r-project.org/>) has here been used as the tool to perform the sampling, made independently of each technological parameter.

The ℓ simulations performed by sampling enable building the uncertainties of different quantities Y of interest from

$$\begin{aligned} \hat{Y}(t) - Y_{\text{ref}}(t) &= \underbrace{\frac{1}{\ell} \sum_{\ell} Y_{\ell}(t) - Y_{\text{ref}}(t)}_{\text{Bias}} \\ &\pm \underbrace{\sqrt{\frac{1}{\ell-1} \sum_{\ell} \left(Y_{\ell}(t) - \frac{1}{\ell} \sum_{\ell} Y_{\ell}(t) \right)^2}}_{\text{Uncertainty}} \xrightarrow{\ell \mapsto +\infty} \\ &\pm \varepsilon_Y(t), \quad \forall t. \end{aligned} \quad (4)$$

The equations from [9] for local isotopic concentrations uncertainties are still valid then, and we resume here the theory for the sake of completeness.

During the calculation process, it is possible to extract concentrations values at each evolution step.

Using (4) and if Y represents an isotopic concentration, it is then possible to build the uncertainty, comparing to a reference calculation.

We write $\underline{M}_N(A, t)$, the isotopic covariance matrix in the mesh element A at step t (unknown), as follows:

$$\underline{M}_N(A, t) = \underline{Z}^T(A, t) \underline{\Omega} \underline{Z}(A, t), \quad (5)$$

where $\underline{Z}(A, t)$ is a matrix representing sensitivities of isotopic concentrations $((\partial N_i / \partial \kappa_k)(A, t))_{i,k}$ and $\underline{\Omega}$ is the correlation matrix $(r_{k1,k2})_{k1,k2}$ between input manufacturing uncertainties.

As $\underline{\Omega}$ is a correlation matrix, only ones are on its diagonal. Then we can get the uncertainty to each isotopic concentration i with the formulation:

$$Z_{i,i}(A, t) = \left(\sqrt{\sum_k \frac{\partial N_i}{\partial \kappa_k}(A, t) \frac{\partial N_i}{\partial \kappa_k}(A, t) \varepsilon_k \varepsilon_k} \right)_{i,i} \quad (6)$$

with $Z_{i,i}(A, t)$ the uncertainty of isotopic concentration in a core mesh. The MC estimator of $Z_{i,i}$ is then written as

$$\begin{aligned} \hat{N}_i(A, t) - N_{\text{ref}}(A, t) &= \underbrace{\frac{1}{\ell} \sum_{\ell} N_{i,\ell}(A, t) - N_{\text{ref}}(A, t)}_{\text{Bias}} \\ &\pm \underbrace{\sqrt{\frac{1}{\ell-1} \sum_{\ell} \left(N_{i,\ell}(A, t) - \frac{1}{\ell} \sum_{\ell} N_{i,\ell}(A, t) \right)^2}}_{\text{Uncertainty}} \xrightarrow{\ell \mapsto +\infty} \\ &\pm \varepsilon_{N_i}(t) = Z_{i,i}(A, t), \quad \forall t. \end{aligned} \quad (7)$$

FIGURE 3: Uncertainty on reactivity (pcm) induced by PTh.

3. Results on Propagation of Two Physical Parameters

In this paragraph, we discuss the results obtained on both manufacturing parameters propagated in the calculation code. In the following, the shapes of the outputs follow a normal law. Its variance is here called “uncertainty” and mean is here called “bias.”

3.1. Propagation of Plates Thickness. Calculations are performed with perturbations of the plate thickness. For each calculation, the thickness is sampled in its distribution law.

3.1.1. Propagation on Reactivity. The results related to reactivity uncertainty are shown in Figure 3. We observe a weak impact on the propagated uncertainty (blue curve). The impact is around 5 pcm at 1σ during all the irradiation. This low value is justified by the fact that the perturbations are almost compensated on the whole core geometry. For some plates, the perturbation increases the thickness, as, for others, it decreases. So, because of the important number of plates (more than a thousand), the average of all the perturbations cancels out. Moreover, the perturbations are made with a constant mass balance in the plate to get only the geometrical perturbation. If the uncertainty on plate mass balance is known, another calculation can be performed to get this particular impact.

However, we observe an important impact on the bias (compared to the standard deviation). This impact is a model bias, coming from the mesh perturbations and from the average of measured values compared to theoretical value (Figure 2). In fact, when the thickness of the plates is modified, the perturbation of the calculation mesh is

FIGURE 4: Uncertainty on power factors (%) induced by PTh.

automatically produced. Another part of this bias comes from the convergence of the MC estimator.

3.1.2. Propagation on Power Factors. As for the reactivity, the perturbation of plate thickness has a relatively weak impact on the power factors (Figure 4). The propagated uncertainty is of the order of 0.16% at 1σ for the maximal value. At the beginning of life, the uncertainties are concentrated principally in the centre of the assemblies. In fact, the thickness perturbation modifies the moderation ratio. This effect is most important in the centre of the assemblies, where the power factor is lower. During depletion, the uncertainties decrease with uranium consumption. So, after relatively large burn-up values, the total uncertainty on the perturbed core converges to the one of the reference core. The uncertainties induced by technological uncertainties on plate thicknesses are then compensated by the fuel consumption.

FIGURE 5: Uncertainty on reactivity (pcm) induced by UO2MB.

3.2. Propagation of UO2 Mass Balance. For the UO2MB parameter, the results are slightly different and are detailed in the next 2 paragraphs.

3.2.1. Propagation on Reactivity. The propagated uncertainty values for reactivity, presented on Figure 5, remain weak for the same reasons as previously. At the beginning of life, the uncertainty on reactivity is close to 3 pcm at 1σ . During irradiation, this value increases to reach 15 pcm at 1σ . The bias observed on the reactivity is only the result of the convergence of MC estimator. The bias is weak compared to the total uncertainty of the reactivity (including nuclear data [8, 9]). Increasing the number of samples is possible but not efficient because of the large CPU time needed for each calculation. The uncertainty increase comes from a flux displacement in the core during the depletion process. During irradiation, the areas with higher mass balance uncertainties are consumed at first, hence inducing a flux displacement to the outer regions. The elements in the outer regions become more important in the reactivity uncertainty, which is then increased consequently.

3.2.2. Propagation on Power Factors. For the power factors, we observe a random map (Figure 6) because of the absence of strong space dependency for mass balance uncertainty. Hence, no trend on power factors uncertainties can be observed or pointed out during depletion. We can only observe the results of a lot of sampling. The maximal value of uncertainty is 1% at 1σ . With this value being higher compared to uncertainty coming from PTh parameter and PTh uncertainty being ten times lower than uncertainty on UO2MB, this parameter has a higher impact on power distribution. During irradiation, the uranium consumption

FIGURE 6: Uncertainty on power factors (%) coming from UO2MB.

in the core centre is responsible for the decrease of the associated UO2MB propagated uncertainty in this area.

4. Computation of Simultaneous Nuclear Data and Manufacturing Data Uncertainty Propagation

The total propagated uncertainty coming from both nuclear data (fission yield and cross sections) and manufacturing data is obtained by applying the process previously described in [8, 9]. The covariance matrices used are the same as those used in [9] (read [14, 15] for fission yields and [11, 16–18] for cross sections). The results obtained by the propagation are described for the three parameters: reactivity, power factors, and isotopic concentrations.

FIGURE 7: Uncertainty on the reactivity during depletion (pcm at 1σ).

4.1. Reactivity. For the reactivity (cf. Figure 7), the total uncertainty is mainly the uncertainty propagated from nuclear data as the manufacturing uncertainties propagation is very low. So, the uncertainty is the result of fission yield and cross section uncertainty propagation. The total uncertainty coming from nuclear data and technological data is here starting around 400 pcm at the beginning of life of the reactor core. It increases during the life to reach 600 pcm at 1σ at 100,000 MW·d/t. As detailed in [9] the uncertainty on reactivity coming from fission yield is around 120–170 pcm at 1σ from some tenth of MW·d/t. This has the particularity to modify the bent shape of the curve at the beginning of life. The observed bias is negligible and can come from the mesh differences between the input data or the Monte Carlo estimator convergence. At the end of life, the cross section uncertainties are the major contributors to the total reactivity uncertainty, the fission yield uncertainty being responsible for about 20% of the total.

4.2. Power Factors. The power factors map uncertainties are principally caused by the UO2MB parameter except at the core-reflector interface. We observe in Figure 8 at the beginning of life an uncertainty of 1.8% at 1σ for the fuel at the interface, which corresponds to the value already found in [9], that is, without technological data. When one moves away from this interface, the uncertainty drops significantly for reaching 1.0% at 1σ , which is the value coming from UO2MB only, the propagated uncertainty from nuclear data being almost zero in this area. We showed in [9] that nuclear data uncertainty propagation gives 0.6% at 1σ in the core centre and that fission yields were negligible for power factors uncertainties. During the depletion process, the power factor uncertainty decreases everywhere and converges to the UO2MB value.

FIGURE 8: Uncertainty on power factors (% at 1σ).

4.3. Isotopic Concentrations. For the isotopic concentrations, no clear trend can be extracted from the results. The next figures reproduce cumulated propagated uncertainties for several isotopic concentrations. Some present an uncertainty coming principally from nuclear data while others come from manufacturing data. We showed in [9] that propagated nuclear data uncertainties on xenon concentration reached 8% at 1σ . In Figure 9, the total uncertainty (taking into account nuclear data and manufacturing data) is increased, demonstrating that manufacturing uncertainties induce an additional uncertainty on xenon isotopic concentration. However, for ^{239}Pu and ^{149}Sm build-ups (Figure 10), the uncertainties are slightly reduced, indicating that negative correlations are generated in the calculation by manufacturing data and nuclear data perturbation. For ^{235}U , Figure 11 shows an important heterogeneous uncertainty, essentially coming from UO2MB, hence from manufacturing

FIGURE 9: Uncertainty on local isotopic concentrations for some Pu241 and Xe135 at 48000 MW·d/t (% at 1σ).

uncertainty. The uncertainties on ^{155}Gd and ^{241}Pu are also presented as they are important.

Table 1 summarizes some major results obtained for isotopic concentration uncertainties as a function of the location.

5. Conclusions

A methodology for manufacturing uncertainty propagation based on Monte Carlo sampling has been presented and tested on a MTR benchmark. An adequate use of the manufacturer information enables simulating different realisations of the core. Those realisations enable getting estimators of standard deviation for different neutronics parameters such as reactivity, power factors, and local isotopic concentrations. Two manufacturing parameters are propagated in the study.

FIGURE 10: Uncertainty on local isotopic concentrations for some Sm149 and Pu239 at 48000 MW·d/t (% at 1σ).

TABLE 1: Summary of isotopic concentrations uncertainties.

Position isotope	Core centre	Core periphery
Pu241	2,0%	4,0–4,5%
Xe135	11,5%	10,0–11,5%
Sm149	3,5%	1,0–1,5%
Pu239	1,0–1,1%	1,0–1,1%
Gd155	12,0–15,0%	8,0–12,0%
U235	0,8–1,2%	0,8–1,2%

One is the UO_2 mass balance in the fuel and the other is a geometric characteristic, the fuel plate thickness.

It is shown that the propagation of these manufacturing parameters uncertainties on reactivity is negligible. This comes from the fact that the perturbed core elements are numerous. So, globally for the core, compensations arise and

FIGURE 11: Uncertainty on local isotopic concentrations for some Gd155 and U235 at 48000 MW-d/t (% at 1σ).

contribute to reducing the reactivity perturbation, which is a global parameter. As an example, the propagation of a 2,5% uncertainty on the fuel plate thickness, coupled with a 0.2% uncertainty on the UO_2 mass balance, induces a combined effect on the reactivity of less than 20 pcm ($10^{-5} \Delta k/k$) at 1σ .

For local parameters, the results are different. The mass balance strongly impacts the power factors map during all the depletion cycle. Although the uncertainties decrease during depletion, because of the fuel consumption, it remains important (approximately 1% at 1σ for local power factors). More surprisingly, the plate thickness has almost no impact on the power factor. It is principally the induced perturbation on the moderation ratio coming from the thickness perturbation, which impacts the power map. This explains the more significant uncertainty at the centre of the fuel assembly.

In a second step, a complete uncertainty propagation of all uncertain data (nuclear data such as cross sections and fission

yields and manufacturing data) through the depletion cycle has been performed to observe the combination of all these uncertainties. It is shown that even if these manufacturing data uncertainties can be neglected for reactivity, there is no clear trend for isotopic concentrations due to the coupling of different sources of uncertainties. For example, contributions to the ^{135}Xe uncertainty during depletion come from both nuclear data and manufacturing data. From [9], nuclear data propagated uncertainties were responsible for 8% on ^{135}Xe concentration, for a total propagated uncertainty from both nuclear data and technological data of 11% at 1σ . On the opposite, ^{235}U is mostly sensitive to manufacturing data uncertainties and less sensitive to nuclear data. We got (see [9]) less than 0.1% of propagated uncertainty on ^{235}U concentration from nuclear data and the present work exhibits a 1.0% at 1σ , due to both nuclear data and technological data uncertainties.

Competing Interests

The authors declare that they have no competing interests.

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