

The Investigation of Dissociation of Trifluoroacetyl Fluoride in the Field of Pulsed CO₂ Laser

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CF₃COF dissociation channels in the pressure region from 0.2 to 4 Torr under the pulsed CO₂ laser radiation are investigated. The dissociation products COF₂, C₂F₆, and CF₄ are identified according to the infrared spectra, and their relative output is measured over the whole range of the initial substance pressures. A scheme of laserochemical reactions describing the trifluoroacetyl fluoride dissociation channels is presented. The correctness of the chosen scheme is confirmed on the basis of the elementary theory of active collisions. The measurements of CF₃COF dissociation selectivity for ¹⁸O and ¹³C are given.

KEY WORDS: Multiphoton infrared dissociation; Trifluoroacetyl; isotope selectivity.

The interaction of complicated organic molecules with the infrared laser radiation has been investigated by numerous authors.¹ Initial interest was aroused in particular by the demonstration that the process was isotopically selective. However, in spite of a sufficiently large amount of experimental and theoretical papers, it is not always possible to describe unambiguously the dissociation channels. The investigation of the dissociation of new compounds makes a certain contribution into this problem.

In the present paper the channels of trifluoroacetyl fluoride dissociation are investigated at the excitation by the pulsed CO₂ laser. This substance can be considered as an initial material for carbon and oxygen isotopes separation. CF₃COF represents a gaseous substance

with the boiling temperature of 214 K. The molecule structure corresponds to the C_s symmetry group. The infrared absorption spectrum contains 15 principal vibrational bands:² $\nu_1 = 1890 \text{ cm}^{-1}$, $\nu_2 = 1340 \text{ cm}^{-1}$, $\nu_3 = 1254 \text{ cm}^{-1}$, $\nu_4 = 1099 \text{ cm}^{-1}$ (C—F), $\nu_5 = 806 \text{ cm}^{-1}$, $\nu_6 = 761 \text{ cm}^{-1}$, $\nu_7 = 692 \text{ cm}^{-1}$, $\nu_8 = 595 \text{ cm}^{-1}$, $\nu_9 = 390 \text{ cm}^{-1}$, $\nu_{10} = 228 \text{ cm}^{-1}$, $\nu_{11} = 1214 \text{ cm}^{-1}$, $\nu_{12} = 519 \text{ cm}^{-1}$, $\nu_{13} = 426 \text{ cm}^{-1}$, $\nu_{14} = 242 \text{ cm}^{-1}$, $\nu_{15} = 50 \text{ cm}^{-1}$. At room temperature, CF_3COF is a stable gaseous substance. The substance may stay in the reactor for several hours, which results in the high stability of the measurements.

CF_3COF irradiation at room temperature was carried out by a pulsed CO_2 laser with pre-ionization at the cathode. The duration of the laser pulse at its half-height was 120–150 ns, the pulse energy was 6 J at the wavelength R(16) 1075.99 cm^{-1} , and 3 J at the wavelength R(40) 1090 cm^{-1} . The area of the laser beam at the inlet and outlet of the experimental reactor was approximately the same, equal to 250 mm^2 . The experimental reactor represents a cylindrical cell 100 mm long and 30 mm in diameter made of stainless steel. The inlet and outlet windows are closed with plane-parallel NaCl plates, and the seal is made of fluoroplastic with an indium spacer. The reactor was prepassivated with fluorine. The absorption spectra have been recorded by the spectrophotometer IKS-29 over the range from 4000 to 400 cm^{-1} . To record the absorption spectra of the dissociation products at the pressure below 4 Torr, the number of experiments was increased with decreasing pressure of the working substance introduced per one experiment. The dissociation products were frozen out into the measuring reactor, and the recording of the IR absorption spectra was carried out at their constant pressure of 4 Torr (4 experiments for the pressure of 1 Torr in the reactor, 8 experiments for that of 0.5 Torr, etc.). The threshold CF_3COF dissociation energy was $\sim 0.8 \text{ J/cm}^2$ in the spectrum maximum, at R(20) 1078.6 cm^{-1} .

The analysis of the IR absorption spectra of CF_3COF dissociation products suggest the presence of the following substances: COF_2 (absorption bands at 1929 cm^{-1} , 1243.7 cm^{-1} , 965 cm^{-1}), C_2F_4 (absorption bands at 1340 cm^{-1} and 1190 cm^{-1}), C_2F_6 (absorption bands at 1250 cm^{-1} and 1116 cm^{-1}), CF_4 (absorption bands at 1280 cm^{-1} and 2174 cm^{-1}).

To carry out the quantitative analysis of the dissociation products, the intensities of pure COF_2 and C_2F_4 absorption as the functions of the pressure in the experimental reactor were measured. Figure 1

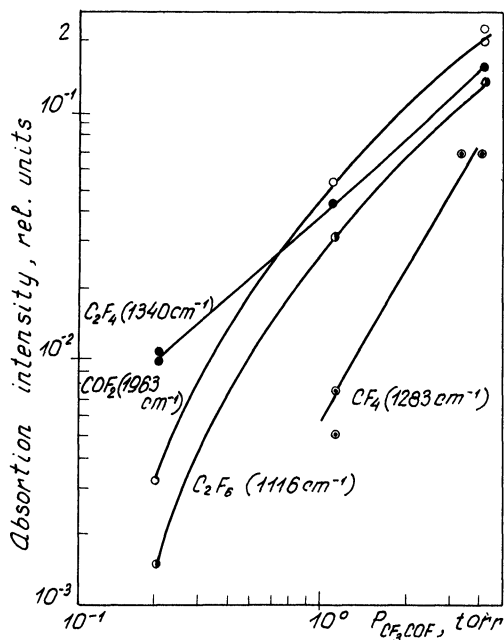
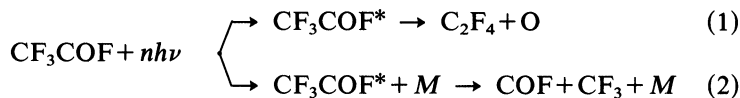


FIGURE 1 The dependence of the output of CF_3COF dissociation products on its pressure ($\nu = 1082.3 \text{ cm}^{-1}$; $E = 6 \text{ J}$).

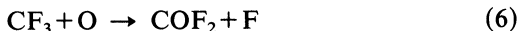
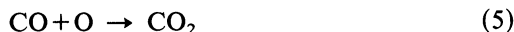
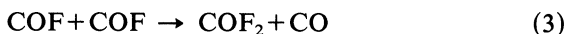
shows the experimentally observed dependences. The intensities of the spectral lines of all dissociation products are normalized to the same scale of relative units.

On the basis of the obtained data, we can come to the following conclusions: in the pressure region below 0.6 Torr, the reaction product is mainly C_2F_4 , while the amount of C_2F_6 and COF_2 is small. At the pressure above 1 Torr, the production rates of all the substances become equal and, in any case, remain constant up to 4 Torr (the pressures above 4 Torr were not investigated). Besides, at the pressure of 1 Torr, CF_4 appears among the products. The obtained data have allowed us to suggest that the formation of CF_3COF dissociation

products proceeds according to the following mechanism:



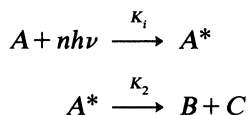
i.e., the initial molecule decomposition may proceed by two channels. Then the following reactions occur:

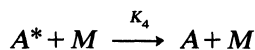
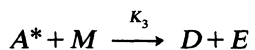


The interaction of COF and CF₃ radicals with the initial molecule may serve as a possible channel of C₂F₆ and COF₂ formation. The increase of CF₃COF pressure should lead to the dependence observed in the experiment (see Figure 1). To obtain the answer concerning the role of such processes, the buffer gas (He or Xe) was added to CF₃COF. The increase of the buffer pressure up to 40 Torr did not lead to any qualitative changes of the dependence obtained before. Therefore, we may conclude that the role of such secondary processes is slight.

To confirm the scheme presented above, we apply the elementary theory of active collisions. It is well-known³ that this theory is based on the supposition that bimolecular reactions are realized at the collision of two molecules. The reaction rate is proportional to the collision number in a unit volume during a unit of time. The molecule decomposition proceeds through an activated complex representing a certain instantaneous state of the active molecule, the excitation energy being concentrated at certain degrees of freedom and, finally, leading to the dissociation.

Let us consider the dissociation of the initial molecule *A* located in the thermal reservoir consisting of the molecules *M*, under the action of the laser radiation:





where B , C , D and E are the dissociation products.

We may apply the stationarity principle to the activated molecule A^* serving as an unstable intermediate complex:

$$\frac{d[A^*]}{dt} = k_1[A] - k_2[A^*] - k_3[A^*][M] - k_4[A^*][M] = 0$$

hence,

$$[A^*] = \frac{k_1[A]}{k_2 + (k_3 + k_4)[M]}$$

The observed decomposition rate $[A]$ is determined from the expression

$$-\frac{d[A]}{dt} = k_2[A^*] + k_3[A^*][M]$$

Substituting the concentration $[A^*]$ from the previous equation, we obtain

$$-\frac{d[A]}{dt} = k_2 \frac{k_1[A]}{k_2 + (k_3 + k_4)[M]} + k_3 \frac{k_1[A][M]}{k_2 + (k_3 + k_4)[M]}$$

Two limiting cases are possible depending on the initial substance pressure.

1. At low pressures

$$k_2 \gg (k_3 + k_4)[M]$$

$$-\frac{d[A]}{dt} = k_1[A] + k_1 \frac{k_3}{k_2} [A][M]$$

The analysis of the obtained expression shows that at small $[A]$ and $[M]$, the decomposition of $[A]$ occurs in the first order fashion. The increase of the initial substance pressure leads to the primary decomposition in the second order fashion (the second term makes the main contribution).

2. At high pressures

$$\begin{aligned} (k_3 + k_4)[M] &\gg k_2 \\ -\frac{d[A]}{dt} &= k_2 \frac{k_1[A]}{(k_3 + k_4)[M]} + \frac{k_1 k_3 [A][M]}{(k_3 + k_4)[M]} \\ &= \frac{k_1 k_2 [A]}{k_3 k_4 [M]} + \frac{k_1 k_3}{k_3 + k_4} [A] \end{aligned}$$

In case of the pure substance dissociation $[A] = [M]$, and $-d[A]/dt = k'[A] + b$ is a first-order reaction where

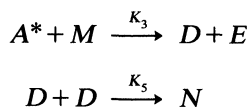
$$k' = \frac{k_1 k_2}{k_3 + k_4}, \quad b = \frac{K_1 K_2}{K_3 + K_4}$$

The above schematic representation of the reaction confirms the validity of the above scheme (1)-(4).

The production rates of all the products become equal at the pressure of 0.5 Torr. In the general case, the expression of the reaction rate is the following:

$$\begin{aligned} w_1 &= k'_1[A] && \text{in the first order} \\ w_2 &= k'_2[A][M] && \text{in the second order} \end{aligned}$$

CF_3COF decomposition is experimentally observed by the output of laserochemical reaction products C_2F_4 , COF_2 and C_2F_6 . Now we show that the final products can unambiguously define the initial compound decomposition in the case under consideration. We derive the production of C_2F_6 out of CF_3COF on the basis of the notion of two subsequent second order reactions:



The kinetics of the process under consideration will be described by the following equations:

$$\begin{aligned} \frac{d[N]}{dt} &= k_5[D]^2 \\ \frac{d[D]}{dt} &= k_3[A^*][M] - k_5[D]^2 \end{aligned}$$

According to the stationarity condition,

$$\frac{d[D]}{dt} = 0$$

then

$$[D]^2 = \frac{k_3}{k_5}[A^*][M]$$

$$\frac{d[N]}{dt} = k_3[A^*][M]$$

Hence, the production rate of C_2F_6 is determined by CF_3COF dissociation rate.

The parallel reactions (5) and (6) are competing with respect to the oxygen atoms. The reaction (6) is the least probable, since it should initiate the subsequent reaction



The absence of CF_4 among the products at low pressures of the initial substance confirms the above supposition.

CF_4 formation at CF_3COF pressure above 1 Torr may be, generally, described through several mechanisms. However, the experimental data do not allow us to describe the observed situation unambiguously. It should be noted that CF_4 formation is observed at the dissociation of numerous molecules (the pressure exceeding several Torr) containing CF_3 groups (CF_3COCF_3 , CF_3OF , etc.). Therefore, we may believe that the mechanism should be of a sufficiently general character, i.e., it should be described by an equation of the type of Eq. (7).

A pure CF_3COF of natural isotopic composition was used (at pressure 1 Torr) for preparing the dissociation products isotopically modified in ^{13}C and ^{18}O . Spectral dependence of the degree of COF_2 enrichment in the lasing frequency range from 980 to 1072 cm^{-1} was determined. The best results at the P(30) 1037.4 cm^{-1} line were as follows: 6% for ^{13}C and 0.53% for ^{18}O which corresponds to a selectivity 10 for both the isotopes. Isotopic analyses were done on the mass spectrometer MI-1201 on the mass lines 66, 67, 68 (COF_2^+) and 116, 117, 118 (CF_3COF^+). The C_2F_6 enrichment in ^{13}C , measured in terms of $C_2F_5^+$ ions at pressures ~ 1 Torr represents the enrichment for COF_2 . According to the dissociation mechanism proposed and

Eqs. (1) and (2) the highest enrichment in ^{13}C should be expected in tetrafluoroethylene resulting from the monomolecular decay of the initial material. The C_2F_4 produced directly in the mass spectrometer from the residual CF_3COF however prevents from isolating the mass lines of laser dissociation products.

CONCLUSIONS

It is shown that the CF_3COF dissociation depending on its pressure in the field of a pulsed CO_2 laser light can be described by a reaction as a result of active collisions of particles having some adequate energy. Low dissociation threshold and selectivity ~ 10 with the dissociation yield $\sim 8\%$ allow trifluoroacetyl fluoride to be considered as rather a promising material for oxygen isotopes separation.

References

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