## Research Article

# Ground State Solution for a Fourth Order Elliptic Equation of Kirchhoff Type with Critical Growth in $\mathbb{R}^{N}$ 

Li Zhou ${ }^{1 / 2}$ and Chuanxi Zhu $\mathbb{1 0}^{\mathbf{1 , 3}}$<br>${ }^{1}$ Zhejiang University of Science \& Technology, Hangzhou, Zhejiang 310023, China<br>${ }^{2}$ Department of Mathematics, Nanchang University, Nanchang, Jiangxi 330031, China<br>${ }^{3}$ School of Mathematics, Dalian University of Technology, Dalian, Liaoning 116024, China<br>Correspondence should be addressed to Chuanxi Zhu; chuanxizhu@126.com

Received 16 July 2022; Accepted 26 August 2022; Published 24 September 2022
Academic Editor: Sergey Shmarev
Copyright © 2022 Li Zhou and Chuanxi Zhu. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

In this paper, we consider the following fourth order elliptic Kirchhoff-type equation involving the critical growth of the form $\left\{\begin{array}{l}\Delta^{2} u-\left(a+b \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u+V(x) u=\left(I_{\alpha} * F(u)\right) f(u)+\lambda|u|^{2^{* *}-2} u, \text { in } \mathbb{R}^{N}, \\ u \in H^{2}\left(\mathbb{R}^{N}\right),\end{array}\right.$ where $a>0, b \geq 0, \lambda$ is a positive parameter, $\alpha \in(N$ $-2, N), 5 \leq N \leq 8, V: \mathbb{R}^{N} \longrightarrow \mathbb{R}$ is a potential function, and $I_{\alpha}$ is a Riesz potential of order $\alpha$. Here, $2^{* *}=2 N /(n-4)$ with $N$ $\geq 5$ is the Sobolev critical exponent, and $\Delta^{2} u=\Delta(\Delta u)$ is the biharmonic operator. Under certain assumptions on $V(x)$ and $f(u)$, we prove that the equation has ground state solutions by variational methods.

## 1. Introduction

In this article, we study the following fourth-order elliptic Kirchhoff-type equation involving the critical growth of the form:
$\left\{\begin{array}{l}\Delta^{2} u-\left(a+b \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u+V(x) u=\left(I_{\alpha} * F(u)\right) f(u)+\lambda|u|^{2^{* *-2}} u, \text { in } \mathbb{R}^{N}, \\ u \in H^{2}\left(\mathbb{R}^{N}\right),\end{array}\right.$
where $a>0, b \geq 0, \lambda$ is a positive parameter, $5 \leq N \leq 8, V$ $: \mathbb{R}^{N} \longrightarrow \mathbb{R}$ is a potential function, and $I_{\alpha}$ is a Riesz potential of order $\alpha \in(N-2, N)$ defined by $I_{\alpha}=(\Gamma((N-\alpha) / 2)) /(\Gamma(\alpha /$ 2) $\left.\pi^{N / 2} 2^{\alpha}|x|^{N-\alpha}\right)$. Here, $2^{* *}=2 N /(N-4)$ with $N \geq 5$ is the Sobolev critical exponent, and $\Delta^{2} u=\Delta(\Delta u)$ is the biharmonic operator, that is, $\Delta^{2} u=\sum_{i=1}^{N}\left(\partial^{4} / \partial x_{i}^{4}\right) u+\sum_{i \neq j}^{N}\left(\partial^{4} / \partial x_{i}^{2} \partial x_{j}^{2}\right) u$. Besides, $V(x): \mathbb{R}^{N} \longrightarrow \mathbb{R}$ is a potential function satisfying
(V1) $V \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$ and $\inf _{x \in \mathbb{R}^{N}} V(x):=V_{0}>0$
(V2) meas $\left\{x \in \mathbb{R}^{N}: V(x) \leq M\right\}<\infty$, where meas denotes the Lebesegue measure in $\mathbb{R}^{N}$ and $V_{0}$ and $M$ are positive constants.

Furthermore, we suppose that the function $f \in C^{1}(\mathbb{R}, \mathbb{R})$ satisfies
(f1) $f(t)=o\left(t^{\alpha / N}\right)$ as $t \longrightarrow 0$
(f2) $\lim _{|t| \longrightarrow+\infty}\left(f(t) / t^{(\alpha+2) /(N-2)}\right)=0$
(f3) $f(t) / t$ is increasing on $(0,+\infty)$ and decreasing on $(-\infty, 0)$
(f4) $f(t)$ is increasing on $\mathbb{R}$
On the one hand, in 2012, by the variational methods, Wang and An [1] studied the following fourth-order equation of Kirchhoff type:

$$
\begin{cases}\Delta^{2} u-M\left(\int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=f(x, u), & x \in \Omega  \tag{2}\\ u=\nabla u=0, & x \in \partial \Omega\end{cases}
$$

and obtained the existence and multiplicity of solutions. Later, Cabada and Figueiredo [2] considered a class of generalized extensible beam equations with critical growth in $\mathbb{R}^{N}$ as follows:

$$
\left\{\begin{array}{l}
\Delta^{2} u-M\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u+u=\lambda f(u)+|u|^{2^{* *}-2} u, x \text { in } \mathbb{R}^{N}  \tag{3}\\
u \in H^{2}\left(\mathbb{R}^{N}\right)
\end{array}\right.
$$

where $M: \mathbb{R}^{+} \longrightarrow \mathbb{R}^{+}$is continuous, $f \in C(\mathbb{R}, \mathbb{R}), N \geq 5$, and $\lambda>0$ is a parameter. With the help of the minimax theorem and the truncation technique, the existence of nontrivial solutions of equation (3) is proved for $\lambda$ sufficiently large. Recently, Song and Shi [3] proved the multiplicity of solutions for the following fourth-order elliptic equation with critical exponent

$$
\begin{cases}\Delta^{2} u-g\left(\int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=\lambda h(x, u)+|u|^{2^{* *}-2} u, & x \in \Omega  \tag{4}\\ u=\nabla u=0, & x \in \partial \Omega\end{cases}
$$

where $\Omega \subset \mathbb{R}^{N}$ is an open bounded domain with smooth boundary, $g: \mathbb{R}^{+} \longrightarrow \mathbb{R}^{+}$is continuous, $h \in C(\bar{\Omega} \times \mathbb{R}, \mathbb{R})$, $N \geq 5$, and $\lambda$ is a positive parameter. Soon after that, Liang et al. [4] obtained multiplicity of solutions to the following generalized extensible beam equation with critical growth:

$$
\begin{align*}
& \Delta^{2} u-M\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u+V(x) u  \tag{5}\\
& \quad=k(x)|u|^{q-2} u+\lambda|u|^{2^{* *}-2} u, x \text { in } \mathbb{R}^{N}
\end{align*}
$$

In fact, in the earlier time, Ma has already applied the variational methods to study the existence and multiplicity of solutions for the following fourth-order boundary value problem of Kirchhoff type:

$$
\begin{gather*}
u^{\prime \prime \prime \prime}-M\left(\int_{0}^{1}\left|u^{\prime}\right|^{2} d x\right) u^{\prime}=f(x, u), \quad x \in(0, L), \\
\left\{\begin{array}{l}
u^{\prime \prime \prime \prime}-M\left(\int_{0}^{1}\left|u^{\prime}\right|^{2} d x\right) u^{\prime \prime}=q(x) f\left(x, u, u^{\prime}\right) x \in(0,1), \\
u(0)=u(1)=u^{\prime \prime}(0)=u^{\prime \prime}(1)=0 .
\end{array}\right. \tag{6}
\end{gather*}
$$

For more details, readers can refer to $[5,6]$ and the references therein.

Actually, for the special case of problem (2) with $M(t)$ $=a+b t$, then problem (2) is reduced to the following
fourth-order elliptic equations of Kirchhoff type:

$$
\begin{cases}\Delta^{2} u-\left(a+b \int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=f(x, u), & x \in \Omega  \tag{7}\\ u=\nabla u=0, & x \in \partial \Omega\end{cases}
$$

This problem is related to the stationary analog of the evolution equation of Kirchhoff type:

$$
\begin{equation*}
u_{t t}+\Delta^{2} u-\left(a+b \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u=f(x, t) \tag{8}
\end{equation*}
$$

Dimensions one and two are relevant from the point of view of physics, engineering, and other sciences, because in those situations model (8) is considered a good approximation for describing nonlinear vibrations of beams or plates (see [7, 8]). Different approaches have been taken to attack this problem under various hypotheses on the nonlinearity. For example, very recently, Wang et al. [9] concentrated on the following Navier BVPs:

$$
\begin{cases}\Delta^{2} u-\lambda\left(a+b \int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=f(x, u), & x \in \Omega  \tag{9}\\ u=\nabla u=0, & x \in \partial \Omega\end{cases}
$$

where $\Omega \subset \mathbb{R}^{N}$ is a smooth bounded domain and $\lambda, a, b>0$. Applying mountain pass techniques and the truncation method, they obtained the existence of nontrivial solution to equation (9) for $\lambda$ small enough when $f(x, u)$ satisfies some superlinear assumptions. For whole space $\mathbb{R}^{N}$, Song and Chen [10] studied the class of Schrödinger-Kirchhofftype biharmonic problems:

$$
\left\{\begin{array}{l}
\Delta^{2} u-\left(a+b \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u+V(x) u=f(x, u), x \in \mathbb{R}^{N}  \tag{10}\\
u \in H^{2}\left(\mathbb{R}^{N}\right)
\end{array}\right.
$$

where $f$ satisfies the Ambrosetti-Rabinowitz type conditions. Under appropriate assumptions on $V$ and $f$, the existence of infinitely many solutions is proved by using the symmetric mountain pass theorem.

On the other hand, in the past decades, many scholars have studied the following problem:

$$
\begin{equation*}
-\Delta u+V(x) u=\left(I_{\alpha} *|u|^{p}\right)|u|^{p-2} u \tag{11}
\end{equation*}
$$

which is called nonlinear Choquard type equation. For the physical background, we refer to [11-13] and the references therein. Mathematically, the existence and qualitative properties of solutions to equation (5) have been studied for decades by variational methods. See [11, 14, 15] for earlier results and [16-26] for a recent work.

Motivated by the works we mentioned above, especially by $[4,10]$, we consider the combination of equations (7) and (11) and extend to the general convolution case in $\mathbb{R}^{N}$. In our paper, we get the ground state solution of problem (1).

Our main results are as follows:

Theorem 1. If $V$ satisfies (V1)-(V2) and $f \in C^{1}(\mathbb{R}, \mathbb{R})$ verifies (f1)-(f4), then problem (1) has a ground state solution.

For the convenience of expression, hereafter, we use the following notations:
$H:=H^{2}\left(\mathbb{R}^{N}\right)=\left\{u \in L^{2}\left(\mathbb{R}^{N}\right):|\nabla u|, \Delta u \in L^{2}\left(\mathbb{R}^{N}\right)\right\}$ equipped with scalar product $\langle u, v\rangle_{H}=\int_{\mathbb{R}^{N}}(\Delta u \Delta v+\nabla u \nabla v+$ $u v) d x$
$E:=\left\{u \in H: \int_{\mathbb{R}^{N}} V(x) u^{2} d x<\infty\right\}$ equipped with scalar product $\langle u, v\rangle=\int_{\mathbb{R}^{N}}(\Delta u \Delta v+\nabla u \nabla v+V(x) u v) d x$, therefore, $\|u\|=\left[\int_{\mathbb{R}^{N}}\left(|\Delta u|^{2}+|\nabla u|^{2}+V(x) u^{2}\right) d x\right]^{1 / 2}$
(i) $L^{s}\left(\mathbb{R}^{N}\right)(1 \leq s \leq \infty)$ denotes the Lebesgue space with the norm $|u|_{s}=\left(\int_{\mathbb{R}^{N}}|u|^{s} d x\right)^{1 / s}$
(ii) For any $u \in H^{2}\left(\mathbb{R}^{N}\right) \backslash\{0\}$, $u_{t}$ is denoted as

$$
u_{t}= \begin{cases}0, & t=0  \tag{12}\\ \sqrt{t} u\left(\frac{x}{t}\right), & t>0\end{cases}
$$

(iii) $C, C_{1}, C_{2}, \cdots$ denote positive constants possibly different in different lines

Remark 2. By the assumptions of $V$, it is obvious that the embedding $E^{\circ} H^{2}\left(\mathbb{R}^{N}\right)$ is continuous. Furthermore, the embedding $E^{\circ} L^{r}\left(\mathbb{R}^{N}\right)$ is continuous for $r \in\left[1,2^{* *}\right)$, and compact for $r \in\left[1,2^{*}\right)$, where $2^{*}=2 N /(N-2)$ if $N \geq 3$ and $2^{*}$ $=\infty$ if $N=1$ or 2 . Thus, for each $s \in\left[2,2^{* *}\right]$, there is a $C_{\varepsilon}$ $>0$ such that $|u|_{s} \leq C_{\varepsilon}\|u\|$ for all $u \in E$.

## 2. Preliminaries

In this section, we will give some very important lemmas.

Lemma 3. Assume (f1)-(f4) hold, then we have the following:
(1) For all $\varepsilon>0$, there is a $C_{\varepsilon}>0$ such that $|f(t)| \leq \varepsilon$ $|t|^{\alpha / N}+C_{\varepsilon}|t|^{(\alpha+2) /(N-2)}$ and $|F(t)| \leq \varepsilon|t|^{(N+\alpha) / N}+C_{\varepsilon}$ $|t|^{(N+\alpha) /(N-2)}$
(2) For all $\varepsilon>0$, there is a $C_{\varepsilon}>0$ such that for every $p$ $\in\left(2,2^{*}\right),|F(t)| \leq \varepsilon\left(|t|^{(N+\alpha) / N}+|t|^{(N+\alpha) /(N-2)}\right)+C_{\varepsilon}$ $|t|^{p(N+\alpha) / 2 N}$, and $|F(t)|^{2 N /(N+\alpha)} \leq \varepsilon\left(|t|^{2}+|t|^{2 N /(N-\alpha)}\right)$ $+C_{\varepsilon}|t|^{p}$
(3) For any $s \neq 0, s f(s)>2 F(s)$ and $F(s)>0$

Proof. One can easily obtain the results by elementary calculation.

Lemma 4 (Hardy-Littlewood-Sobolev inequality [27]). Let $0<\alpha<N, p, q>1$ and $1 \leq r<s<\infty$ be such that

$$
\begin{equation*}
\frac{1}{p}+\frac{1}{q}=1+\frac{\alpha}{N}, \frac{1}{r}-\frac{1}{s}=\frac{\alpha}{N} . \tag{13}
\end{equation*}
$$

(1) For any $f \in L^{p}\left(\mathbb{R}^{N}\right)$ and $g \in L^{q}\left(\mathbb{R}^{N}\right)$, one has

$$
\begin{equation*}
\left|\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{f(x) g(y)}{|x-y|^{N-\alpha}} d x d y\right| \leq C(N, \alpha, p)\|f\|_{L^{p}\left(\mathbb{R}^{N}\right)}\|g\|_{L^{q}\left(\mathbb{R}^{N}\right)} \tag{14}
\end{equation*}
$$

(2) For any $f \in L^{r}\left(\mathbb{R}^{N}\right)$, one has

$$
\|l / \mid \cdot\|^{N-\alpha} * f\left\|_{L^{s}\left(\mathbb{R}^{N}\right)} \leq C(N, \alpha, r)\right\| f \|_{L^{r}\left(\mathbb{R}^{N}\right)}
$$

Remark 5. By Lemma 3(1), Lemma 4(1) and Sobolev imbedding theorem, we can get

$$
\begin{align*}
& \left|\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) F(u) d x\right| \\
& \quad \leq C|F(u)|_{2 N /(N+\alpha)}^{2} \\
& \quad \leq C\left[\int_{\mathbb{R}^{N}}\left(|u|^{(N+\alpha) / N}+|u|^{(N+\alpha) /(N-2)}\right)^{2 N /(N+\alpha)} d x\right]^{(N+\alpha) / N} \\
& \quad \leq C\left[\int_{\mathbb{R}^{N}}\left(|u|^{2}+|u|^{2 N /(N-2)}\right) d x\right]^{(N+\alpha) / N} \\
& \quad \leq C\left(\|u\|^{(2 N+2 \alpha) / N}+\|u\|^{(2 N+2 \alpha) /(N-2)}\right) \tag{15}
\end{align*}
$$

## 3. Variational Formulation

The associated energy function of problem (1) is given by

$$
\begin{align*}
I(u)= & \frac{1}{2} \int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+\frac{1}{2} \int_{\mathbb{R}^{N}}\left[a|\nabla u|^{2}+V(x) u^{2}\right] d x \\
& +\frac{b}{4}\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right)^{2}-\frac{1}{2} \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) F(u) d x  \tag{16}\\
& -\frac{\lambda}{2^{* *}} \int_{\mathbb{R}^{N}}|u|^{2^{* *}} d x
\end{align*}
$$

i.e. the critical points of the functional $I(u)$ are weak solutions of problem (1). Under the assumptions, $I \in C^{1}(E$, $\mathbb{R}$ ), and for all $u, v \in E$, it holds that

$$
\begin{align*}
\left\langle I^{\prime}(u), v\right\rangle= & \int_{\mathbb{R}^{N}} \Delta u \Delta v d x+a \int_{\mathbb{R}^{N}} \nabla u \nabla v d x+\int_{\mathbb{R}^{N}} V(x) u v d x \\
& +b \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x \int_{\mathbb{R}^{N}} \nabla u \nabla v d x \\
& -\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) f(u) v d x-\lambda \int_{\mathbb{R}^{N}}|u|^{2^{* *}-2} u v d x . \tag{17}
\end{align*}
$$

Thus,

$$
\begin{align*}
\left\langle I^{\prime}(u), u\right\rangle= & \int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+\int_{\mathbb{R}^{N}}\left[a|\nabla u|^{2}+V(x) u^{2}\right] d x \\
& +b\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right)^{2}-\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) f(u) u d x \\
& -\lambda \int_{\mathbb{R}^{N}}|u|^{2^{* *}} d x . \tag{18}
\end{align*}
$$

In this section, we prove the following results.
Lemma 6. The functional $I(u)$ possesses the mountain-pass geometry, i.e.,
(1) There exists $\rho, \delta>0$ such that $I(u) \geq \delta$ for all $\|u\|=\rho$
(2) There exists $e \in H^{2}\left(\mathbb{R}^{N}\right)$ such that $\|e\|>\rho$ and $I(e)$ $<0$

Proof. (1) By Lemma 3(1) and Lemma 4, we have

$$
\begin{align*}
I(u)= & \frac{1}{2} \int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+\frac{1}{2} \int_{\mathbb{R}^{N}}\left[a|\nabla u|^{2}\right. \\
& \left.+V(x) u^{2}\right] d x+\frac{b}{4}\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right)^{2} \\
& -\frac{1}{2} \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) F(u) d x-\frac{\lambda}{2^{* *}} \int_{\mathbb{R}^{N}}|u|^{2^{* *}} d x  \tag{19}\\
\geq & C_{1}\|u\|^{2}-C_{2}\left(\|u\|^{(2 N+2 \alpha) / N}+\|u\|^{(2 N+2 \alpha) /(N-2)}\right) \\
& -C_{3}\|u\|^{2^{* *}} .
\end{align*}
$$

Thus, there exists $\rho, \delta>0$ such that $I_{\infty} \geq \delta$ for all $\|u\|$ $=\rho>0$ small enough.
(2) For any $u \in E \backslash\{0\}$,

$$
\begin{aligned}
I\left(u_{t}\right)= & \frac{t^{N-3}}{2} \int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+\frac{a t^{N-1}}{2} \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x \\
& +\frac{t^{N+1}}{2} \int_{\mathbb{R}^{N}} V(x) u^{2} d x+\frac{b t^{2 N-2}}{4}\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right)^{2} \\
& -\frac{t^{N+\alpha}}{2} \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) F(u) d x \\
& -\frac{\lambda}{2^{* *}} t^{\left(N^{2}+3 N\right) /(N-4)} \int_{\mathbb{R}^{N}}|u|^{2^{* *}} d x \longrightarrow-\infty
\end{aligned}
$$

as $t \longrightarrow \infty$, since $\alpha>N-2$; thus, we see $I_{\infty}\left(u_{t}\right)<0$ for $t>0$ large. Note that

$$
\begin{align*}
\left\|u_{t}\right\|^{2}= & t^{N-3} \int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+t^{N-1} \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x  \tag{21}\\
& +t^{N+1} \int_{\mathbb{R}^{N}} V(x) u^{2} d x
\end{align*}
$$

Taking $e=u_{t_{0}}$, with $t_{0}>0$ large, we have $\|e\|>\rho$ and $I_{\infty}(e)<0$.

Hence, we define the mountain-pass level of $I$ :

$$
\begin{equation*}
c=\inf _{\gamma \in \Gamma} \max _{t \in[0,1]} I(\gamma(t))>0, \tag{22}
\end{equation*}
$$

where $\Gamma=\{\gamma \in C([0,1], E): \gamma(0)=0, I(\gamma(1))<0\}$.
Lemma 7. If $\left\{u_{n}\right\} \subset E$ is a $(P S)_{c}$ sequence of $I$, then $\left\{u_{n}\right\}$ is bounded in $E$.

Proof. Let $\left\{u_{n}\right\} \subset E$ be $(P S)_{c}$ sequence, i.e., $I\left(u_{n}\right) \longrightarrow c$ and $I^{\prime}\left(u_{n}\right) \longrightarrow 0$; then, we have

$$
\begin{align*}
c= & I\left(u_{n}\right)-\frac{1}{4}\left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle+o(1) \\
= & \frac{1}{4} \int_{\mathbb{R}^{N}}\left|\Delta u_{n}\right|^{2} d x+\frac{1}{4} \int_{\mathbb{R}^{N}}\left[a\left|\nabla u_{n}\right|^{2}+V(x) u_{n}^{2}\right] d x \\
& +\frac{1}{4} \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right)\left[f\left(u_{n}\right) u_{n}-2 F\left(u_{n}\right)\right] d x  \tag{23}\\
& +\left(\frac{\lambda}{4}-\frac{\lambda}{2^{* *}}\right) \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x+o(1) \\
& \geq C_{4}\left\|u_{n}\right\|^{2} .
\end{align*}
$$

Consequently, $\left\{u_{n}\right\}$ is bounded in $E$.
Remark 8. By Lemma 7, we can assume that there exists a $u$ such that

$$
\left\{\begin{array}{l}
u_{n} u \text { in } E,  \tag{24}\\
u_{n} \longrightarrow u \text { in } L_{l o c}^{s}\left(\mathbb{R}^{N}\right), \quad \forall s \in\left[2,2^{* *}\right), \\
u_{n} \longrightarrow \text { ua.e.on } \mathbb{R}^{N} .
\end{array}\right.
$$

Then, by the similar method as Lemma 3.3 in [4], we can obtain $\lim _{n \longrightarrow \infty} \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x=\int_{\mathbb{R}^{N}}|u|^{2^{* *}} d x$.

Lemma 9. $I(u)$ satisfies $(P S)$ condition.
Proof. Let $\left\{u_{n}\right\} \subset E$ be $(P S)_{c}$ sequence, i.e., $I\left(u_{n}\right) \longrightarrow c$ and $I^{\prime}\left(u_{n}\right) \longrightarrow 0$; by Lemma $7,\left\{u_{n}\right\}$ is bounded in E. Hence,
up to a subsequence, we may assume that there exists a $u$ such that

$$
\left\{\begin{array}{l}
u_{n} u \text { in } E,  \tag{25}\\
u_{n} \longrightarrow u \text { in } L_{l o c}^{s}\left(\mathbb{R}^{N}\right), \quad \forall s \in\left[2,2^{* *}\right), \\
u_{n} \longrightarrow \text { ua.e.on } \mathbb{R}^{N}
\end{array}\right.
$$

Then, using the lower semicontinuity of the norm, Brezis-Lieb Lemma [28] and Remark 8, we have

$$
\begin{align*}
o(1)\left\|u_{n}\right\|= & \left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle=\int_{\mathbb{R}^{N}}\left|\Delta u_{n}\right|^{2} d x \\
& +\int_{\mathbb{R}^{N}}\left[a\left|\nabla u_{n}\right|^{2}+V(x) u_{n}^{2}\right] d x+b\left(\int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2} d x\right)^{2} \\
& -\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x-\lambda \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x \\
= & \int_{\mathbb{R}^{N}}\left|\Delta u_{n}-\Delta u\right|^{2} d x+\int_{\mathbb{R}^{N}}\left[a\left|\nabla\left(u_{n}-u\right)\right|^{2}\right. \\
& \left.+V(x)\left|u_{n}-u\right|^{2}\right] d x+\int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+\int_{\mathbb{R}^{N}}\left[a|\nabla u|^{2}\right. \\
& \left.+V(x) u^{2}\right] d x+b\left(\int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2} d x\right)^{2} \\
& -\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x \\
& -\lambda \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x+o(1) \\
\geq & C_{5}\left\|u_{n}-u\right\|^{2}+\int_{\mathbb{R}^{N}}|\Delta u|^{2} d x+\int_{\mathbb{R}^{N}}\left[a|\nabla u|^{2}\right. \\
& \left.+V(x) u^{2}\right] d x+b\left(\int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right)^{2} \\
& -\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F(u)\right) f(u) u d x-\lambda \int_{\mathbb{R}^{N}}|u|^{2^{* *}} d x+o(1) . \tag{26}
\end{align*}
$$

Thus, we can have $\left\|u_{n}-u\right\|^{2} \longrightarrow 0$, which implies that $\left\{u_{n}\right\}$ strongly converges to $u$ in $E$. This completes the proof of Lemma 9.

Remark 10. Now, we recall the Nehari manifold

$$
\begin{equation*}
\mathscr{M}:=\left\{u \in E \backslash\{0\}:\left\langle I^{\prime}(u), u\right\rangle=0\right\} . \tag{27}
\end{equation*}
$$

Let $m=\inf _{u \in \mathscr{M}} I(u)$; then for any $u \in \mathscr{M}$, we have

$$
\begin{equation*}
I(u)=I(u)-\frac{1}{4}\left\langle I^{\prime}(u), u\right\rangle \geq C_{6}\left\|u_{n}\right\|^{2} \geq 0 \tag{28}
\end{equation*}
$$

Hence, $m$ is well defined. Moreover, by the similar argument as Chapter 4 [29], we have the following characterization:

$$
\begin{equation*}
c=\inf _{\gamma \in \Gamma} \max _{t \in[0,1]} I(\gamma(t))=m=\inf _{u \in \mathscr{M}} I(u) . \tag{29}
\end{equation*}
$$

## 4. Ground State Solution for Problem (1)

In this section, we prove the main theorem.
Proof of Theorem 11. From Lemmas 6 and 7, we know that there exists a bounded $(P S)_{c}$ sequence $\left\{u_{n}\right\}$, that is, $I\left(u_{n}\right)$ $\longrightarrow c=m, I^{\prime}\left(u_{n}\right) \longrightarrow 0$. Next, let $\delta:=\limsup _{n \longrightarrow \infty} \sup _{y \in \mathbb{R}^{N}} \int_{B_{1}(y)}\left|u_{n}\right|^{2}$ $\mathrm{d} x$. We claim $\delta>0$. On the contrary, by similar argument as Lions' concentration compactness principle, we can proof $u_{n} \longrightarrow 0$ in $L^{p}\left(\mathbb{R}^{N}\right)$ for $2<p<2^{* *}$. By Lemma 3(2), for any $\varepsilon>0$, there exists a constant $C_{\varepsilon}>0$ such that

$$
\begin{align*}
& \underset{n \longrightarrow \infty}{\limsup } \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x \\
& \quad \leq \operatorname{Climsup} \\
& n \longrightarrow \infty  \tag{30}\\
& \left.\quad+\int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2 N /(N-2)} d x\right)+\int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2} d x \\
& \left.\quad \leq C\left[\varepsilon C_{\mathbb{R}^{N}}+\left.C_{\varepsilon} \limsup _{n \longrightarrow \infty}\right|^{p} d x\right]_{\mathbb{R}^{N}}\left|u_{n}\right|^{p} d x\right]^{(N+\alpha) / N} \\
& \quad=C\left(\varepsilon C_{8}\right)^{(N+\alpha) / N} .
\end{align*}
$$

Note that $\varepsilon$ is arbitrary; we get

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x=o(1) \tag{31}
\end{equation*}
$$

Combining with $I^{\prime}\left(u_{n}\right) \longrightarrow 0$ and Remark 8, we can get

$$
\begin{align*}
o(1)= & \left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle=\int_{\mathbb{R}^{N}}\left|\Delta u_{n}\right|^{2} d x+a \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2} d x \\
& +b\left(\int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2} d x\right)^{2}+\int_{\mathbb{R}^{N}} V(x)\left|u_{n}\right|^{2} d x \\
& -\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x-\lambda \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x  \tag{32}\\
\geq & C_{9}\left\|u_{n}\right\|^{2}-\int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x \\
& -\lambda \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x,
\end{align*}
$$

which implies that

$$
\begin{equation*}
C_{9}\left\|u_{n}\right\|^{2} \leq \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(u_{n}\right)\right) f\left(u_{n}\right) u_{n} d x+\lambda \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{2^{* *}} d x+o(1) . \tag{33}
\end{equation*}
$$

Then, we have $\left\|u_{n}\right\| \longrightarrow 0$, which implies $u_{n} \longrightarrow 0$ in $E$. We deduce that $c=0$, which contradicts to the fact that $c>0$. Hence, $\delta>0$, and there exists $\left\{y_{n}\right\} \subset \mathbb{R}^{N}$ such that $\int_{B_{1}\left(y_{n}\right)}$ $\left|u_{n}\right|^{p} d x \geq(\delta / 2)>0$. We set $v_{n}(x)=u_{n}\left(x+y_{n}\right)$; then, $\left\|u_{n}\right\|=$
$\left\|v_{n}\right\|, \int_{B_{1}(0)}\left|v_{n}\right|^{p} d x>\delta / 2$ and $I\left(v_{n}\right) \longrightarrow c=m, I^{\prime}\left(v_{n}\right) \longrightarrow 0$. Thus, there exists a $v_{0} \neq 0$ such that

$$
\left\{\begin{array}{l}
v_{n} v_{0} \text { in } E,  \tag{34}\\
v_{n} \longrightarrow v_{0} \text { in } L_{l o c}^{s}\left(\mathbb{R}^{N}\right), \quad \forall s \in\left[2,2^{* *}\right) \\
v_{n} \longrightarrow v_{0} \text { a.e.on } \mathbb{R}^{N}
\end{array}\right.
$$

Then, for any $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$, we have $0=\left\langle I^{\prime}\left(v_{n}\right), \varphi\right\rangle+$ $o(1)=\left\langle I^{\prime}\left(v_{0}\right), \varphi\right\rangle$, which means $v_{0}$ is a solution of equation (1).

On the other hand, combining with the Fatou Lemma, we can obtain

$$
\begin{align*}
m= & I\left(v_{n}\right)-\frac{1}{4}\left\langle I^{\prime}\left(v_{n}\right), v_{n}\right\rangle+o(1) \\
= & \frac{1}{4} \int_{\mathbb{R}^{N}}\left|\Delta v_{n}\right|^{2} d x+\frac{1}{4} \int_{\mathbb{R}^{N}}\left[a\left|\nabla v_{n}\right|^{2}+V(x) v_{n}^{2}\right] d x \\
& +\frac{1}{4} \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(v_{n}\right)\right)\left[f\left(v_{n}\right) v_{n}-2 F\left(v_{n}\right)\right] d x \\
& +\left(\frac{\lambda}{4}-\frac{\lambda}{2^{* *}}\right) \int_{\mathbb{R}^{N}}\left|v_{n}\right|^{2^{* *}} d x+o(1)  \tag{35}\\
\geq & \frac{1}{4} \int_{\mathbb{R}^{N}}\left|\Delta v_{0}\right|^{2} d x+\frac{1}{4} \int_{\mathbb{R}^{N}}\left[a\left|\nabla v_{0}\right|^{2}+V(x) v_{0}^{2}\right] d x \\
& +\frac{1}{4} \int_{\mathbb{R}^{N}}\left(I_{\alpha} * F\left(v_{0}\right)\right)\left[f\left(v_{0}\right) v_{0}-2 F\left(v_{0}\right)\right] d x \\
& +\left(\frac{\lambda}{4}-\frac{\lambda}{2^{* *}}\right) \int_{\mathbb{R}^{N}}\left|v_{0}\right|^{2^{* *}} d x+o(1) \\
= & I\left(v_{0}\right)-\frac{1}{4}\left\langle I^{\prime}\left(v_{0}\right), v_{0}\right\rangle+o(1)=I\left(v_{0}\right)+o(1)
\end{align*}
$$

At the same time, we know $m=c \leq I\left(v_{0}\right)$ by the definition of $m$. Then, we can deduce that $v_{0}$ is a ground state solution of equation (1). Thus, we complete the proof of Theorem 11.

## Data Availability

Data is not available.

## Conflicts of Interest

The authors declare that they have no conflicts of interest.

## Authors' Contributions

LZ curated the data and is responsible for the writing of the original draft. CZ acquired funding and is assigned to the validation. LZ and CZ did the writing, review, and editing.

## Acknowledgments

This work was supported by the National Natural Science Foundation of China (Grant Nos. 11771198 and 11901276) and Science and Technology Project of Education Department of Jiangxi Province (Grant No. GJJ218406).

## References

[1] F. Wang and Y. An, "Existence and multiplicity of solutions for a fourth-order elliptic equation," Boundary Value Problems, vol. 2012, no. 1, 2012.
[2] A. Cabada and G. M. Figueiredo, "A generalization of an extensible beam equation with critical growth in $\mathbb{R}^{N}$," Nonlinear Analysis: Real World Applications, vol. 20, pp. 134-142, 2014.
[3] Y. Song and S. Shi, "Multiplicity of solutions for fourth-order elliptic equations of Kirchhoff type with critical exponent," Journal of Dynamical and Control Systems, vol. 23, no. 2, pp. 375-386, 2017.
[4] S. H. Liang, Z. Y. Liu, and H. L. Pu, "Multiplicity of solutions to the generalized extensible beam equations with critical growth," Nonlinear Analysis, vol. 197, article 111835, 2020.
[5] T. F. Ma, "Existence results and numerical solutions for a beam equation with nonlinear boundary conditions," Applied Numerical Mathematics, vol. 47, pp. 189-196, 2003.
[6] T. F. Ma, "Positive solutions for a nonlocal fourth order equation of Kirchhoff type," Conference Publications, vol. 2007, pp. 694-703, 2007.
[7] J. M. Ball, "Initial-boundary value problems for an extensible beam," Journal of Mathematical Analysis and Applications, vol. 42, no. 1, pp. 61-90, 1973.
[8] H. M. Berger, "A new approach to the analysis of large deflections of plates," Journal of Applied Mechanics, vol. 22, no. 4, pp. 465-472, 1955.
[9] F. Wang, M. Avci, and Y. An, "Existence of solutions for fourth order elliptic equations of Kirchhoff type,"Journal of Mathematical Analysis and Applications, vol. 409, no. 1, pp. 140146, 2014.
[10] H. Song and C. Chen, "Infinitely many solutions for Schrödinger-Kirchhoff-type fourth order elliptic equations," Proceedings of the Edinburgh Mathematical Society, vol. 60, no. 4, pp. 1003-1020, 2017.
[11] E. H. Lieb, "Existence and uniqueness of minimizing solution of Choquard's nonlinear equation," Studies in Applied Mathematics, vol. 57, pp. 93-105, 1977.
[12] I. M. Moroz, R. Penrose, and P. Tod, "Spherically-symmetric solutions of the Schrödinger-Newton equations," Classical and Quantum Gravity, vol. 15, no. 9, pp. 2733-2742, 1998.
[13] S. Pekar, Untersuchung über die Elektronentheorie der Kristalle, Akademie Verlag, Berlin, 1954.
[14] P. L. Lions, "The Choquard equation and related questions," Nonlinear Analysis, vol. 4, no. 6, pp. 1063-1072, 1980.
[15] P. L. Lions, "The concentration-compactness principle in the Calculus of Variations. The locally compact case, part 1," Annales de l'Institut Henri Poincaré C, Analyse non linéaire, vol. 1, no. 2, pp. 109-145, 1984.
[16] V. Moroz and J. Van Schaftingen, "Ground states of nonlinear Choquard equations: existence, qualitative properties and decay asymptotics," Journal of Functional Analysis, vol. 265, pp. 153-184, 2013.
[17] V. Moroz and J. Van Schaftingen, "Existence of ground states for a class of nonlinear Choquard equations," Transactions of the American Mathematical Society, vol. 367, pp. 6557-6579, 2015.
[18] V. Moroz and J. Van Schaftingen, "Nonexistence and optimal decay of supersolutions to Choquard equations in exterior
domains," Journal of Differential Equations, vol. 254, no. 8, pp. 3089-3145, 2013.
[19] V. Moroz and J. Van Schaftingen, "Groundstates of nonlinear Choquard equations: hardy-Littlewood-Sobolev critical exponent," Communications in Contemporary Mathematics, vol. 17, no. 5, article 1550005, 2015.
[20] V. Moroz and J. Van Schaftingen, "A guide to the Choquard equation," Journal of Fixed Point Theory and Applications, vol. 19, no. 1, pp. 773-813, 2017.
[21] M. Chimenti and J. Van Schaftingen, "Nodal solutions for the Choquard equation," Journal of Functional Analysis, vol. 271, no. 1, pp. 107-135, 2016.
[22] L. Ma and Z. Lin, "Classification of positive solitary solutions of the nonlinear Choquard equation," Archive for Rational Mechanics and Analysis, vol. 195, no. 2, pp. 455-467, 2010.
[23] P. Chen and X. C. Liu, "Ground states for Kirchhoff equation with Hartree-type nonlinearities," Journal of Mathematical Analysis and Applications, vol. 473, no. 1, pp. 587-608, 2019.
[24] W. Zhang, S. Yuan, and L. Wen, "Existence and concentration of ground-states for fractional Choquard equation with indefinite potential," Advances in Nonlinear Analysis, vol. 11, no. 1, pp. 1552-1578, 2022.
[25] W. Zhang and J. Zhang, "Multiplicity and concentration of positive solutions for fractional unbalanced double phase problems," Journal of Geometric Analysis, vol. 32, no. 9, p. 235, 2022.
[26] J. Zhang, W. Zhang, and V. D. Rădulescu, "Double phase problems with competing potentials: concentration and multiplication of ground states," Mathematische Zeitschrift, vol. 301, pp. 4037-4078, 2022.
[27] E. H. Lieb and M. Loss, Analysis, American Mathematical Scoiety. Province, RL, 2nd edition, 2001.
[28] H. Brézis and E. H. Lieb, "A relation between pointwise convergence of functions and convergence of functionals," Proceedings of the American Mathematical Society, vol. 88, no. 3, pp. 486-490, 1983.
[29] M. Willem, Minimax Theorems, Proress in Nonlinear Differential Equations and Their Applications 24, Birkhäuser, Boston, MA, 1996.

