

Research Article

Oscillatory Behaviour of the Nonlinear Damped Fractional Partial Dynamic Equation

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In this paper, a partial dynamic equation of fractional order is considered with Neumann and Dirichlet boundary conditions, and we studied the oscillation properties of the fractional partial dynamic equation on time scales. Riccati transformation technique is used to establish oscillation criteria for the fractional partial dynamic equation. The obtained results are verified with examples.

1. Introduction

In recent years, the importance of fractional-order calculus has been much motivated by researchers rather than the integer order because of increasing applications in signal processing, neural networks, and electrical and mechanical engineering. Definitions for fractional derivatives in continuous and discrete cases using various operators and functions have been given by many such as Hadamard, Euler, Riemann, and Grunwald, but fundamental properties of derivatives were not satisfied by these derivatives, and with these derivatives, solving the differential equations of fractional order is not easy. In [1], the definition of the conformable fractional derivative has been given, and it has satisfied some fundamental properties of derivatives. The properties and application of the conformable derivative have been presented in [2, 3] and the references cited therein. Meanwhile, the oscillation properties of solutions are an important qualitative tool to study the solutions of dynamical systems. The oscillation of different types of both integer-order and fractional-order differential, difference, and dynamic equations has been investigated by many authors [4–20]. The time-scale calculus was introduced by Stefan Hilger to unify the theory of continuous and discrete cases. The time-scale analysis for multivariable cases is discussed in [21–23], and the fractional time-scale calculus is studied in [24]. To the best of our knowledge, we observe that the oscillation of solutions of the fractional partial dynamic equation on the time scale was not considered so far.

This motivates the authors to establish oscillation results of the following fractional partial dynamic equation with the damped term:

$$\left(\ell(\eta)\chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta)\right)^{(\Delta^{\alpha})_{\eta}} + \hbar(\eta)\chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) + \nu(\xi,\eta)f(\chi(\xi,\eta)) = \tau(\eta)\nabla^{2}_{\xi}\chi(\xi,\eta),$$
(1)

where $(\xi, \eta) \in \Omega \times [\eta_0, \infty)_{\Theta}$ with the Neumann boundary condition

$$\frac{\partial \chi(\xi,\eta)}{\partial \nu} = 0, \quad (\xi,\eta) \in \partial \Omega \times [\eta_0,\infty)_{\Theta}, \tag{2}$$

and the Dirichlet boundary condition

$$\chi(\xi,\eta) = 0, \quad (\xi,\eta) \in \partial\Omega \times [\eta_0,\infty)_{\Theta}, \tag{3}$$

where $(\Delta^{\alpha})_{\eta}$ is the conformable fractional partial dynamic operator with respect to the time variable η , Ω is a bounded domain in \mathbb{R}^n with piecewise smooth boundary $\partial\Omega$, ν is the unit exterior normal vector to $\partial\Omega$ and $\alpha \in (0, 1]$, and $\nabla_{\xi}^2 =$ $\sum_{i=1}^n \partial_{\xi_i \xi_i}^2$ is the Laplacian operator. $\ell(\eta), \hbar(\eta)$, and $\tau(\eta)$ are real-valued and rd-continuous functions on $[0, \infty)_{\Theta}, \ell(\eta) > 0$ is conformable fractional differentiable with $\ell(\eta)^{(\Delta^{\alpha})_{\eta}}$ continuous, and $v(\xi, \eta) \in C(\overline{\Omega} \times [\eta_0, \infty)_{\Theta}, \mathbb{R})$ with $v(\eta) = \min_{\xi \in \Omega}$ $v(\xi, \eta)$. The continuous function f satisfies $\chi f(\chi) > 0$ for $\chi \neq 0$.

A time scale Θ is a nonempty closed subset of the real numbers \mathbb{R} which is unbounded above. The time-scale interval of the form $[\eta_0, \infty)_{\Theta}$ is denoted by $[\eta_0, \infty)_{\Theta} = [\eta_0, \infty) \cap \Theta$ for $\eta_0 \in \Theta$, $\eta_0 > 0$.

A nontrivial solution $\chi(\xi, \eta)$ which satisfies (1) on $\Omega \times [\eta_0, \infty)_{\Theta}$ along with boundary condition (2) (or (3)) is called oscillatory if it has neither eventually positive solution (EPS) nor eventually negative solution (ENS). Equation (1) is oscillatory if all the solutions of it are oscillatory.

2. Preliminaries

The following definitions can be found in [25], where there is a detailed introduction to time-scale calculus. A time scale Θ is a nonempty closed subset of the real numbers \mathbb{R} . We will use intervals of the form $[t_0, \infty)_{\Theta} = [t_0, \infty) \cap \Theta$ for $t_0 \in \Theta$. For a point $t \in \Theta$, we have the following definitions: the forward jump operator is defined as $\sigma(t) = \inf\{s \in \Theta, s > t\}$. The backward jump operator is defined as $\rho(t) = \sup\{s \in \Theta, s < t\}$. The graininess is defined as $\mu(t) = \sigma(t) - t$. A point $t \in \Theta$ is said to be right dense if $\sigma(t) = t$. A function $g: \Theta \longrightarrow \mathbb{R}$ is said to be rd-continuous if it is continuous at each right-dense point and there exists a finite left limit of g at all left-dense points. The set of rd-continuous functions is denoted by $C_{rd}(\Theta, \mathbb{R})$.

To define derivatives, we introduce

$$\Theta^{k} = \begin{cases} \Theta \setminus (\rho (\sup \Theta), \sup \Theta) & \text{if } \sup \Theta < \infty, \\ \Theta & \text{if } \sup \Theta < \infty. \end{cases}$$
(4)

Definition 1 (see [24]). Let $f: \Theta \longrightarrow \mathbb{R}$; $\eta \in \Theta^k$, and $\alpha \in (0, 1]$. For $\eta > 0$, we define $f^{(\Delta^\alpha)}(\eta)$ to be the number (provided it exists) with the property that, given any $\epsilon > 0$, there is a δ-neighborhood $\mathcal{N}_{\eta} \in \Theta$ of $\eta, \delta > 0$ such that

$$\left| \left[f^{\sigma}(\eta) - f(\kappa) \right] \eta^{1-\alpha} - f^{(\Delta^{\alpha})}(\eta) \left[\sigma(\eta) - \kappa \right] \right| \le \varepsilon |\sigma(\eta) - \kappa| \text{ for all } \kappa \in \mathcal{N}_{\eta},$$
(5)

 $f^{(\Delta^{\alpha})}(\eta)$ is called the conformable fractional derivative of f of order α at η , and the conformable fractional derivative at 0 is $f^{(\Delta^{\alpha})}(0) = \lim_{\eta \longrightarrow 0^+} f^{(\Delta^{\alpha})}(\eta)$.

Lemma 1. Let $a \in \Theta^k$ and $b \in \Theta$, and assume that $f: \Theta \times \Theta^k \longrightarrow \mathbb{R}$ is continuous at (η, η) where $\eta \in \Theta^k$ with $a < \eta$. Suppose that, for all $\epsilon > 0$, there is a δ -neighborhood U of η_i such that

$$\left| \left[f^{\sigma_i}(\eta) - f^{\kappa}_i(\eta) \right] \eta_i^{1-\alpha} - f^{(\Delta^{\alpha})_{\eta_i}}(\eta) \left[\sigma_i(\eta) - \kappa \right] \right| \le \varepsilon \left| \sigma_i(\eta) - \kappa \right| \text{ for all } \kappa \in U,$$
(6)

where $f^{(\Delta^{\alpha})\eta_i}$ denotes the partial derivative of f with respect to η_i . Then,

$$g(\eta) = \int_{a}^{\eta} f(\eta, \kappa) \Delta \kappa \Rightarrow g^{(\Delta^{\alpha})_{\eta_{i}}}(\eta) = \int_{a}^{\eta} f^{(\Delta^{\alpha})_{\eta_{i}}}(\eta, \kappa) \Delta \kappa + f(\sigma_{i}(\eta), \eta),$$

$$h(\eta) = \int_{\eta}^{b} f(\eta, \kappa) \Delta \kappa \Rightarrow h^{(\Delta^{\alpha})_{\eta_{i}}}(\eta) = \int_{\eta}^{b} f^{(\Delta^{\alpha})_{\eta_{i}}}(\eta, \kappa) \Delta \kappa - f(\sigma_{i}(\eta), \eta).$$
(7)

Proof. The proof is analogous to the proof of Theorem 1.117 in [25]. \Box

for $\eta \ge \eta_0$, $\Lambda(\eta, \kappa) > 0$ for $\eta > \kappa \ge \eta_0$, and Λ has a nonpositive continuous α -partial fractional derivative $\Lambda^{(\Delta^{\alpha})_{\kappa}}(\eta, \kappa)$.

Definition 2. Let $\mathbb{B} = \{(\eta, \kappa) | \eta \ge \kappa \ge \eta_0\}$. Then, the class Λ is a collection of functions $\Lambda \in C_{rd}(\mathbb{B}, \mathbb{R})$ such that $\Lambda(\eta, \eta) = 0$

A function $g: \Theta \longrightarrow \mathbb{R}$ is said to be regressive provided $1 + \mu(t)g(t) \neq 0$ for each $t \in \Theta^k$. Let \mathscr{R} be the set of functions that are rd-continuous and regressive. Also, we define

 $\mathscr{R}^+ = \{g \in \mathscr{R}: 1 + \mu(t)g(t) > 0, t \in \Theta\}$. For $\gamma \in \mathscr{R}$ and $s, t \in \Theta$, the generalized exponential function is defined by

$$e_{\gamma}(t,s) = \exp\left(\int_{s}^{t} \zeta_{\mu(\tau)}(\gamma(\tau))\Delta\tau\right), \quad \zeta_{\mu}(z) = \begin{cases} \frac{1}{\mu}\log(1+\mu z), & \mu \neq 0, \\ \\ z, & \mu = 0. \end{cases}$$
(8)

For convenience, we use the following:

$$Y(\eta) = \int_{\Omega} \chi(\xi, \eta) d\xi,$$

$$R(\eta) = -\frac{\eta^{\alpha - 1} \hbar(\eta)}{\ell(\eta)}.$$
(9)

Lemma 2. If the fractional dynamic inequalities

$$\left[\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)\right]^{(\Delta^{\alpha})} + \hbar(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta) + \upsilon(\eta)f(\Upsilon(\eta)) \le 0, \eta \ge \eta_1,$$
(10)

$$\left[\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)\right]^{(\Delta^{\alpha})} + \hbar(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta) + \upsilon(\eta)f(\Upsilon(\eta)) \ge 0, \eta \ge \eta_1,$$
(11)

have no EPS and ENS, respectively, then every solution $\chi(\xi, \eta)$ of (1) and (2) is oscillatory.

Proof. Suppose that $\chi(\xi, \eta)$ is an EPS of (1) and (2). Then, by the definition of EPS, there exists $\eta_1 \ge \eta_0$ such that $\chi(\xi, \eta) > 0$ for $\eta \ge \eta_1$, which gives $\Upsilon(\eta) > 0$ for $\eta \ge \eta_1$. Now, taking integration to (1) in connection with ξ over Ω , we get

$$\int_{\Omega} \left[\ell(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) \right]^{(\Delta^{\alpha})_{\eta}} d\xi + \int_{\Omega} \hbar(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) d\xi + \int_{\Omega} v(\xi,\eta) f(\chi(\xi,\eta)) d\xi$$

$$= \int_{\Omega} \tau(\eta) \nabla_{\xi}^{2} \chi(\xi,\eta) d\xi.$$
(12)

Applying Green's formula and (2), we get

$$\int_{\Omega} \nabla_{\xi}^{2} \chi(\xi, \eta) d\xi = \int_{\partial \Omega} \frac{\partial \chi(\xi, \eta)}{\partial \nu} dS = 0.$$
(13)

By using the condition on v and applying Jensen's inequality, we obtain

$$\int_{\Omega} v(\xi,\eta) f(\chi(\xi,\eta)) d\xi \ge v(\eta) \int_{\Omega} f(\chi(\xi,\eta)) d\xi \ge v(\eta) f(\Upsilon(\eta)).$$
(14)

From (12)–(14), we have

$$\int_{\Omega} \left[\ell(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) \right]^{(\Delta^{\alpha})_{\eta}} d\xi$$

$$+ \hbar(\eta) \int_{\Omega} \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) d\xi + v(\eta) f(\Upsilon(\eta)) \leq 0.$$
(15)

Using that $[\ell(\eta)\chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta)]^{(\Delta^{\alpha})_{\eta}}$ is continuous in connection with ξ on the closed and bounded set $\overline{G} \subseteq \mathscr{R}^{n}$, we have δ independent of η in Lemma 1. Therefore, the conditions of Lemma 1 are satisfied. Thus, $\int_{G} [\chi(\xi,\eta)d\xi]^{(\Delta^{\alpha})_{\eta}} = \int_{G} \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta)d\xi$. Hence, from (15), we obtain (10). A similar argument is used for the ENS $\chi(\xi,\eta)$ of (1) and (2) to obtain (11). \Box

Lemma 3. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$ and $R(\eta) \in \mathcal{R}_+$. If (10) has an EPS $\Upsilon(\eta)$, then there exists sufficiently large $\eta_1 \in [\eta_0, \infty)_{\Theta}$ such that $[\ell(\eta) \Upsilon^{(\Delta^{\alpha})}(\eta) / e_{R(\eta)}(\eta, \eta_0)]^{(\Delta^{\alpha})} < 0$ and $\Upsilon^{(\Delta^{\alpha})}(\eta) > 0$ on $[\eta, \infty)_{\Theta}$.

Proof. Given $\Upsilon(\eta)$ is an EPS of (10), there is sufficiently large η_2 so that $\Upsilon(\eta) > 0$ on $[\eta, \infty)_{\Theta}$. Now, we have

$$\left[\frac{\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)}{e_{R(\eta)}(\eta,\eta_{0})}\right]^{(\Delta^{\alpha})} = \frac{e_{R(\eta)}(\eta,\eta_{0})\left[\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)\right]^{(\Delta^{\alpha})} - \ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)\left[e_{R(\eta)}(\eta,\eta_{0})\right]^{(\Delta^{\alpha})}}{e_{R(\eta)}(\eta,\eta_{0})e_{R(\eta)}(\sigma(\eta),\eta_{0})}
= \frac{e_{R(\eta)}(\eta,\eta_{0})\left[\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)\right]^{(\Delta^{\alpha})} + \hbar(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)e_{R(\eta)}(\eta,\eta_{0})}{e_{R(\eta)}(\eta,\eta_{0})e_{R(\eta)}(\sigma(\eta),\eta_{0})}$$
(16)
$$= \frac{-v(\eta)f(\Upsilon(\eta))}{e_{R(\eta)}(\eta,\eta_{0})e_{R(\eta)}(\sigma(\eta),\eta_{0})}$$

$$= \frac{v(\eta) f(1(\eta))}{e_{R(\eta)} \left(\sigma(\eta), \eta_0\right)} < 0$$

Since $R(\eta) \in \mathbb{R}_+$, we get $e_{R(\eta)}(\sigma(\eta), \eta_0) > 0$, and $\ell(\eta)$ $\Upsilon^{(\Delta^{\alpha})}(\eta)/e_{R(\eta)}(\eta, \eta_0)$ is decreasing on $[\eta_2, \infty)_{\Theta}$. Therefore, we get that $\Upsilon^{(\Delta^{\alpha})}(\eta)$ is the EPS or ENS. Suppose that $\Upsilon^{(\Delta^{\alpha})}(\eta) < 0$ on $[\eta_2, \infty)_{\Theta}$ for sufficiently large $\eta_2 > \eta_1$. Then,

$$\Upsilon(\eta) - \Upsilon(\eta_3) = \int_{\eta_2}^{\eta} \frac{\ell(\kappa) \Upsilon^{(\Delta^{\alpha})}(\kappa)}{\ell(\kappa)} \Delta^{\alpha} \kappa \le \ell(\eta_2) \Upsilon^{(\Delta^{\alpha})}(\eta_2) \int_{\eta_2}^{\eta} \frac{\Delta^{\alpha} \kappa}{\ell(\kappa)}.$$
(17)

Letting $\eta \longrightarrow \infty$ and using the condition that $\int_{\eta}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$, we get $\lim_{\eta \longrightarrow \infty} \Upsilon(\eta) = -\infty$ which is a contradiction to $\Upsilon(\eta) > 0$, and hence, $\Upsilon^{(\Delta^{\alpha})}(\eta) > 0$. \Box

The following lemma holds if we proceed as in the above lemma with an ENS $\Upsilon(\eta)$.

Lemma 4. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$ and $R(\eta) \in \mathcal{R}_+$. If (11) has an ENS $\Upsilon(\eta)$, then there exists sufficiently large $\eta_1 \in [\eta_0, \infty)_{\Theta}$ such that $[\ell(\eta) \Upsilon^{(\Delta^{\alpha})}(\eta) / e_{R(\eta)}(\eta, \eta_0)]^{(\Delta^{\alpha})} > 0$ and $\Upsilon^{(\Delta^{\alpha})}(\eta) < 0$ on $[\eta, \infty)_{\Theta}$. **Theorem 1.** Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$, $R(\eta) \in \mathcal{R}_+$, and f' exists with $f'(\xi) \ge N$ for some N > 0, $\forall \xi \ne 0$. If there exists $z(\eta) > 0$ on $[H, \infty)_{\Theta}$ such that

$$\limsup_{\eta \to \infty} \int_{H}^{\eta} \left[\frac{v(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_{0})} - \frac{\ell(\kappa)(z^{(\Delta^{\alpha})}(\kappa))^{2}}{4Nz(\kappa)e_{R(\eta)}(\kappa,\eta_{0})} \right] \Delta^{\alpha} \kappa = \infty,$$
(18)

then (1) and (2) are oscillatory.

Proof. Suppose that the solution $\chi(\xi, \eta)$ of (1) and (2) is the EPS. Define a generalized Riccati function as

$$w(\eta) = \left[\frac{z(\eta)\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)}{f[\Upsilon(\eta)]e_{R(\eta)}(\eta,\eta_0)}\right].$$
(19)

By Lemma 3, it is lucid that $\Upsilon > 0$ and $\Upsilon^{(\Delta^{\alpha})} \ge 0$, respectively. Hence, $w(\eta) \ge 0$ for $\eta \ge \eta_1$. Then,

$$w^{(\Delta^{\alpha})}(\eta) = \frac{z(\eta)}{f[\Upsilon(\eta)]} \left[\frac{e_{R(\eta)}(\eta, \eta_{0}) \left[\ell(\eta) \Upsilon^{(\Delta^{\alpha})}(\eta) \right]^{(\Delta^{\alpha})} - \left(e_{R(\eta)}(\eta, \eta_{0}) \right)^{(\Delta^{\alpha})} \ell(\eta) \Upsilon^{(\Delta^{\alpha})}(\eta)}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} \right] \\ + \frac{\ell(\sigma(\eta)) \Upsilon^{(\Delta^{\alpha})}(\sigma(\eta))}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} \left[\frac{f(\Upsilon(\eta) z^{(\Delta^{\alpha})}(\eta) - z(\eta) [f(\Upsilon(\eta))]^{(\Delta^{\alpha})}}{f[\Upsilon(\eta)] f[\Upsilon(\sigma(\eta)]]} \right] \\ \leq \frac{-v(\eta) z(\eta)}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} + \frac{z^{(\Delta^{\alpha})}(\eta) w(\sigma(\eta))}{z(\sigma(\eta))} - \frac{N z(\eta) e_{R(\eta)}(\eta, \eta_{0})}{\ell(\eta)} \left[\frac{w(\sigma(\eta))}{z(\sigma(\eta))} \right]^{2} \\ \leq \frac{-v(\eta) z(\eta)}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} + \frac{\ell(\eta) \left(z^{(\Delta^{\alpha})}(\eta) \right)^{2}}{4N z(\eta) e_{R(\eta)}(\eta, \eta_{0})}.$$
(20)

By taking α -integration from η_1 to η , we obtain

$$\limsup_{\eta \to \infty} \int_{\eta_1}^{\eta} \left[\frac{v(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_0)} - \frac{\ell(\kappa)(z^{(\Delta^{\alpha})}(\kappa))^2}{4Nz(\kappa)e_{R(\eta)}(\kappa,\eta_0)} \right] \Delta^{\alpha} \kappa$$

$$\leq w(\eta_1) - w(\eta) \leq w(\eta_1) \leq \infty,$$
(21)

which is a contradiction to (18). Similarly, we get a contradiction while assuming that $\chi(\xi, \eta)$ is the ENS of (1) and (2). \Box

Theorem 2. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$, $R(\eta) \in \mathcal{R}_+$, and $f(\xi) \ge Ex$ for some E > 0, $\forall \xi \ne 0$. If there exists $z(\eta) > 0$ on $[H, \infty)_{\Theta}$ such that

$$\limsup_{\eta \to \infty} \int_{H}^{\eta} \left[\frac{Ev(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_{0})} - \frac{\ell(\kappa)(z^{(\Delta^{\alpha})}(\kappa))^{2}}{4z(\kappa)e_{R(\eta)}(\kappa,\eta_{0})} \right] \Delta^{\alpha} \kappa = \infty,$$
(22)

then (1) and (2) are oscillatory.

Proof. Suppose that the solution $\chi(\xi, \eta)$ of (1) and (2) is the EPS. Define a generalized Riccati function as

$$w(\eta) = \left[\frac{z(\eta)\ell(\eta)\Upsilon^{(\Delta^{\alpha})}(\eta)}{\Upsilon(\eta)e_{R(\eta)}(\eta,\eta_0)}\right].$$
 (23)

By Lemma 3, it is lucid that $\Upsilon > 0$ and $\Upsilon^{(\Delta^{\alpha})} \ge 0$, respectively. Hence, $w(\eta) \ge 0$ for $\eta \ge \eta_1$. Then,

$$w^{(\Delta^{\alpha})}(\eta) = \frac{z(\eta)}{Y(\eta)} \left[\frac{e_{R(\eta)}(\eta, \eta_{0}) \left[\ell(\eta) Y^{(\Delta^{\alpha})}(\eta) \right]^{(\Delta^{\alpha})} - \left(e_{R(\eta)}(\eta, \eta_{0}) \right)^{(\Delta^{\alpha})} \ell(\eta) Y^{(\Delta^{\alpha})}(\eta)}{e_{R(\eta)}(\sigma(\eta), \eta_{0}) e_{R(\eta)}(\sigma(\eta), \eta_{0})} \right] + \frac{\ell(\sigma(\eta)) Y^{(\Delta^{\alpha})}(\sigma(\eta))}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} \left[\frac{Y(\eta) z^{(\Delta^{\alpha})}(\eta) - z(\eta) [Y(\eta)]^{(\Delta^{\alpha})}}{Y(\eta) Y(\sigma(\eta)} \right] \right]$$

$$\leq \frac{-Ev(\eta) z(\eta)}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} + \frac{z^{(\Delta^{\alpha})}(\eta) w(\sigma(\eta))}{z(\sigma(\eta))} - \frac{z(\eta) e_{R(\eta)}(\eta, \eta_{0})}{\ell(\eta)} \left[\frac{w(\sigma(\eta))}{z(\sigma(\eta))} \right]^{2}$$

$$\leq \frac{-Ev(\eta) z(\eta)}{e_{R(\eta)}(\sigma(\eta), \eta_{0})} + \frac{\ell(\eta) (z^{(\Delta^{\alpha})}(\eta))^{2}}{4z(\eta) e_{R(\eta)}(\eta, \eta_{0})}.$$
(24)

By taking α -integration from η_1 to η , we obtain

$$\limsup_{\eta \to \infty} \int_{\eta_{1}}^{\eta} \left[\frac{Ev(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_{0})} - \frac{\ell(\kappa)(z^{(\Delta^{\alpha})}(\kappa))^{2}}{4z(\kappa)e_{R(\eta)}(\kappa,\eta_{0})} \right] \Delta^{\alpha} \kappa$$

$$\leq w(\eta_{1}) - w(\eta) \leq w(\eta_{1}) \leq \infty,$$
(25)

which is a contradiction to (22). Similarly, we get a contradiction while assuming that $\chi(\xi, \eta)$ is the ENS of (1) and (2). \Box

Remark 1. In Theorem 1, (18) can be replaced by

$$\limsup_{\eta \to \infty} \frac{1}{\Lambda(\eta, H)} \int_{H}^{\eta} \Lambda(\eta, \kappa) \\ \cdot \left[\frac{v(\kappa) z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa), \eta_{0})} - \frac{\ell(\kappa) (z^{(\Delta^{\alpha})}(\kappa))^{2}}{4N z(\kappa) e_{R(\eta)}(\kappa, \eta_{0})} \right] \Delta^{\alpha} \kappa = \infty,$$
(26)

for $\Lambda \in \Lambda$.

By Theorem 1, we have

$$w^{(\Delta^{\alpha})}(\eta) \leq \frac{-v(\eta)z(\eta)}{e_{R(\eta)}(\sigma(\eta),\eta_{0})} + \frac{\ell(\eta)(z^{(\Delta^{\alpha})}(\eta))^{2}}{4Nz(\eta)e_{R(\eta)}(\eta,\eta_{0})}.$$
 (27)

Multiplying the above inequality by $\Lambda(\eta, \kappa)$ and taking α -integration from *H* to η , we get

$$\int_{-H}^{\eta} \Lambda(\eta,\kappa) w^{(\Delta^{\alpha})}(\kappa) \Delta^{\alpha} \kappa \leq \int_{-H}^{\eta} \Lambda(\eta,\kappa) \\ \cdot \left[\frac{-v(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_0)} + \frac{\ell(\kappa)(z^{(\Delta^{\alpha})}(\kappa))^2}{4Nz(\kappa)e_{R(\eta)}(\kappa,\eta_0)} \right] \Delta^{\alpha} \kappa.$$
(28)

Using integration by parts, we obtain

$$\begin{split} \left[\Lambda(\eta,\kappa)w(\eta)\right]_{H}^{\eta} &- \int_{H}^{\eta} \Lambda^{(\Delta^{\alpha})}(\eta,\kappa)w(\sigma(\kappa))\Delta^{\alpha}\kappa \\ &\leq \int_{H}^{\eta} \Lambda(\eta,\kappa) \left[\frac{-v(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_{0})} + \frac{\ell(\kappa)\left(z^{(\Delta^{\alpha})}(\kappa)\right)^{2}}{4Nz(\kappa)e_{R(\eta)}(\kappa,\eta_{0})} \right] \Delta^{\alpha}\kappa. \end{split}$$

$$(29)$$

By using the definition of the function Λ and rearranging the terms in the above inequality, we have

$$\frac{1}{\Lambda(\eta,H)} \int_{H}^{\eta} \Lambda(\eta,\kappa) \left[\frac{v(\kappa)z(\kappa)}{e_{R(\kappa)} \left(\sigma(\kappa),\eta_{0}\right)} - \frac{\ell(\kappa) \left(z^{(\Delta^{\alpha})}(\kappa)\right)^{2}}{4Nz(\kappa)e_{R(\eta)} \left(\kappa,\eta_{0}\right)} \right] \Delta^{\alpha} \kappa \leq w(H) \leq \infty, \tag{30}$$

which contradicts (26). Likewise, we obtain a contradiction when we take $\chi(\xi, \eta)$ is the *ENS* of (1) and (2).

Similarly, (22) can be replaced by the following in Theorem 2:

$$\limsup_{\eta \to \infty} \frac{1}{\Lambda(\eta, H)} \int_{-H}^{\eta} \Lambda(\eta, \kappa) \left[\frac{Ev(\kappa)z(\kappa)}{e_{R(\kappa)}(\sigma(\kappa), \eta_0)} - \frac{\ell(\kappa) (z^{(\Delta^{\alpha})}(\kappa))^2}{4z(\kappa) e_{R(\eta)}(\kappa, \eta_0)} \right] \Delta^{\alpha} \kappa = \infty,$$
(31)

for $\Lambda \in \Lambda$.

4. Oscillation of (1) with Boundary Condition (3)

In this section, we use the eigenvalues of the Laplacian, $-\nabla^2 \phi(\xi) = \lambda \phi(\xi)$ in *G*, $\phi(\xi) = 0$ on ∂G . The principal eigenvalue λ_1 is positive, and the corresponding eigenfunction ϕ_1 is also positive in the interior of *G*; see [26], Theorem 2, page 356. Furthermore, we normalize this eigenvector so that $\int_G \phi_1 = 1$.

Lemma 5. If $\chi(\xi, \eta)$ is an EPS of (1) and (3), then there exists $\eta_1 \ge \eta_0$ such that $\Upsilon_1(\eta) = \int_G \chi(\xi, \eta) \phi_1(\xi) d\xi > 0$ and

$$\begin{bmatrix} \ell(\eta) \Upsilon_1^{(\Delta^{\alpha})}(\eta) \end{bmatrix}^{(\Delta^{\alpha})} + \hbar(\eta) \Upsilon_1^{(\Delta^{\alpha})}(\eta) + v(\eta) f(\Upsilon_1(\eta)) \le 0, \quad \forall \eta \ge \eta_1.$$
 (32)

Also, if $\chi(\xi, \eta)$ is an ENS, then there exists $\eta_1 \ge \eta_0$ such that $\Upsilon_1(\eta) < 0$ and

$$\begin{bmatrix} \ell(\eta) \Upsilon_1^{(\Delta^{\alpha})}(\eta) \end{bmatrix}^{(\Delta^{\alpha})} + \hbar(\eta) \Upsilon_1^{(\Delta^{\alpha})}(\eta) + v(\eta) f(\Upsilon_1(\eta)) \ge 0, \quad \forall \eta \ge \eta_1.$$

$$(33)$$

Proof. Since $\chi(\xi, \eta)$ is an EPS of (1) satisfying (3), there exists $\eta_1 \ge \eta_0$ such that $\chi(\xi, \eta) \ge 0$. $\phi_1(\xi) > 0$ in the interior of Ω , which implies that $\Upsilon_1(\eta) = \int_{\Omega} \chi(\xi, \eta) \phi_1(\xi) d\xi > 0$. Multiply (1) by $\phi_1(\xi)$, and integrate in connection with ξ over Ω ; then,

$$\int_{\Omega} \left(\ell(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) \right)^{(\Delta^{\alpha})_{\eta}} \phi_{1}(\xi) d\xi + \int_{\Omega} \hbar(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) \phi_{1}(\xi) d\xi + \int_{\Omega} \upsilon(\xi,\eta) f(\chi(\xi,\eta)) \phi_{1}(\xi) d\xi = \int_{\Omega} \tau(\eta) \nabla_{\xi}^{2} \chi(\xi,\eta) \phi_{1}(\xi) d\xi .$$
(34)

On the right-hand side of (34), we have

$$\int_{\Omega} \tau(\eta) \nabla_{\xi}^{2} \chi(\xi, \eta) \phi_{1}(\xi) d\xi$$

$$= -\lambda \tau(\eta) \int_{\Omega} \chi(\xi, \eta) \phi_{1}(\xi) d\xi = -\lambda \tau(\eta) \Upsilon_{1}(\eta) \leq 0,$$
(35)

which implies that

$$\int_{\Omega} \left(\ell(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) \right)^{(\Delta^{\alpha})_{\eta}} \phi_{1}(\xi) d\xi$$

+
$$\int_{\Omega} \hbar(\eta) \chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta) \phi_{1}(\xi) d\xi$$
(36)
+
$$\int_{\Omega} v(\xi,\eta) f(\chi(\xi,\eta)) \phi_{1}(\xi) d\xi \leq 0.$$

Applying Jensen's inequality and condition on v to the third term of the above inequality and using Lemma 1 and continuity of $(\ell(\eta)\chi^{(\Delta^{\alpha})_{\eta}}(\xi,\eta))^{(\Delta^{\alpha})_{\eta}}$ in the first term of the above inequality, we obtain

$$\left[\ell(\eta)\Upsilon_{1}^{(\Delta^{\alpha})}(\eta)\right]^{(\Delta^{\alpha})} + \hbar(\eta)\Upsilon_{1}^{(\Delta^{\alpha})}(\eta) + \upsilon(\eta)f(\Upsilon_{1}(\eta)) \leq 0.$$
(37)

A similar proof will be used for the ENS to obtain (33). \Box

The following lemmas and theorems consider the solution $\Upsilon_1(\eta)$ of (1) and (3), and the proofs of the following are similar to those of lemmas and theorems proved for the solution $\Upsilon(\eta)$ of (1) and (2). Thus, we omit the proofs.

Lemma 6. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$ and $R(\eta) \in \mathcal{R}_+$. If (32) has an EPS $\Upsilon(\eta)$, then there is sufficiently large $\eta_1 \in [\eta_0, \infty)_{\Theta}$ such that $[\ell(\eta) \Upsilon_1^{(\Delta^{\alpha})}(\eta) / e_{R(\eta)}(\eta, \eta_0)]^{(\Delta^{\alpha})} < 0$ and $\Upsilon_1^{(\Delta^{\alpha})}(\eta) > 0$ on $[\eta_1, \infty)_{\Theta}$.

Lemma 7. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$ and $R(\eta) \in \mathcal{R}_+$. If (33) has an ENS $Y_1(\eta)$, then there is sufficiently large $\eta_1 \in [\eta_0, \infty)_{\Theta}$ such that $[\ell(\eta) Y_1^{(\Delta^{\alpha})}(\eta) / e_{R(\eta)}(\eta, \eta_0)]^{(\Delta^{\alpha})} > 0$ and $Y_1^{(\Delta^{\alpha})}(\eta) < 0$ on $[\eta_1, \infty)_{\Theta}$.

Theorem 3. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$, $R(\eta) \in \mathcal{R}_+$, and f' exists with $f'(\eta) \ge N$ for some $N \ge 0$, $\forall \xi \ne 0$. If there exists $\psi(\eta) > 0$ on $[H, \infty)_{\Theta}$ such that

$$\limsup_{\eta \to \infty} \int_{H}^{\eta} \left[\frac{v(\kappa)\psi(\kappa)}{e_{R(\kappa)}(\sigma(\kappa),\eta_0)} - \frac{\ell(\kappa)(\psi^{(\Delta^{\alpha})}(\kappa))^2}{4N\psi(\kappa)e_{R(\eta)}(\kappa,\eta_0)} \right] \Delta^{\alpha}\kappa = \infty,$$
(38)

then (1) and (3) are oscillatory.

Theorem 4. Assume that $\int_{\eta_0}^{\infty} \Delta^{\alpha} \kappa / \ell(\kappa) = \infty$, $R(\eta) \in \mathcal{R}_+$, and $f(\xi) \ge Ex$ for some E > 0, $\forall \xi \ne 0$. If there exists $\psi(\eta) > 0$ on $[H, \infty)_{\Theta}$ such that

$$\limsup_{\eta \longrightarrow \infty} \int_{H}^{\eta} \left[\frac{E v(\kappa) \psi(\kappa)}{e_{R(\kappa)} (\sigma(\kappa), \eta_0)} - \frac{\ell(\kappa) (\psi^{(\Delta^{\alpha})}(\kappa))^2}{4 \psi(\kappa) e_{R(\eta)} (\kappa, \eta_0)} \right] \Delta^{\alpha} \kappa = \infty,$$
(39)

then (1) and (3) are oscillatory.

Remark 2. In Theorem 3, (38) can be replaced by

$$\limsup_{\eta \to \infty} \frac{1}{\Lambda(\eta, H)} \int_{H}^{\eta} \Lambda(\eta, \kappa) \\ \cdot \left[\frac{v(\kappa)\psi(\kappa)}{e_{R(\kappa)}(\sigma(\kappa), \eta_0)} - \frac{\ell(\kappa)(\psi^{(\Delta^{\alpha})}(\kappa))^2}{4N\psi(\kappa)e_{R(\eta)}(\kappa, \eta_0)} \right] \Delta^{\alpha} \kappa$$
(40)

 $=\infty$, for $\Lambda \in \Lambda$.

Similarly, (39) can be replaced by the following in Theorem 4:

$$\limsup_{\eta \to \infty} \frac{1}{\Lambda(\eta, H)} \int_{H}^{\eta} \Lambda(\eta, \kappa) \\ \cdot \left[\frac{Ev(\kappa)\psi(\kappa)}{e_{R(\kappa)}(\sigma(\kappa), \eta_0)} - \frac{\ell(\kappa)(\psi^{(\Delta^{\alpha})}(\kappa))^2}{4N\psi(\kappa)e_{R(\eta)}(\kappa, \eta_0)} \right] \Delta^{\alpha} \kappa$$
(41)
= ∞ , for $\Lambda \in \Lambda$.

5. Examples

The equations considered below are an illustration to validate our established results.

Example 1. Consider

$$\left(\eta^{-1} \chi^{\left(\Delta^{0.5}\right)_{\eta}}(\xi,\eta) \right)^{\left(\Delta^{0.5}\right)_{\eta}} + \eta^{-3} \chi^{\left(\Delta^{0.5}\right)_{\eta}}(\xi,\eta)$$

$$+ \eta^{-1} \chi(\xi,\eta) \left(1 + |\chi(\xi,\eta)| \right) = \eta \nabla_{\xi}^{2} \chi(\xi,\eta),$$

$$(42)$$

with $\chi_x(0,\eta) = \chi_x(\pi,\eta) = 0$. Here, $\Theta = \mathbb{R}$; $\Omega = (0,\pi)$; $\alpha = 0.5$; $\ell(\eta) = \eta^{-1}$; $\hbar(\eta) = \eta^{-3}$; $f(\chi(\xi,\eta)) = \chi(\xi,\eta)(1 + |\chi(\xi,\eta)|)$; $v(\eta) = \eta^{-1}$; and $\eta_0 = 1$. It is lucid that $f(\chi) \ge \chi \Longrightarrow E = 1$, $\sigma(\eta) = \eta, \mu(\eta) = 0, R(\eta) = -\eta^{-(5/2)}$,

$$\int_{\eta_0}^{\infty} \frac{\Delta^{\alpha} \kappa}{\ell(\kappa)} = \int_{\eta_0}^{\infty} \frac{\kappa^{-(1/2)} d\kappa}{\kappa^{-1}} = \int_{\eta_0}^{\infty} \kappa^{1/2} d\kappa = \infty,$$

$$1 \ge e_{R(\eta)}(\eta, \eta_0) = \exp\left(\int_{\eta_0}^{\eta} R(\kappa) d\kappa\right) \ge 1 + \int_{1}^{\eta} R(\kappa) d\kappa \quad (43)$$

$$= 1 - \int_{1}^{\eta} \kappa^{-(5/2)} d\kappa = 1 + \frac{2}{3} \left(\eta^{-(3/2)} - 1\right) \ge \frac{1}{3}.$$

Let $z(\eta) = \eta^{1/2}$. Then, every condition of Theorem 2 is fulfilled. Hence, (42) with the given boundary condition is oscillatory.

Example 2. Consider

$$\left(\eta^{-1/2} \chi^{(\Delta^{0.5})_{\eta}}(\xi,\eta)\right)^{(\Delta^{0.5})_{\eta}} + \eta^{-5/2} \chi^{(\Delta^{0.5})_{\eta}}(\xi,\eta)$$

$$+ \eta^{-1} \chi(\xi,\eta) \left(E + e^{\chi(\xi,\eta)}\right) = \eta \nabla_{\xi}^{2} \chi(\xi,\eta),$$
(44)

with $\chi(0,\eta) = \chi(\pi,\eta) = 0$. Here, $\eta \in [2,\infty)_{\mathbb{Z}}$; $\Theta = \mathbb{Z}$; $\Omega = (0,\pi)$; $\alpha = 0.5$; $\ell(\eta) = \eta^{-(1/2)}$; $\hbar(\eta) = \eta^{-(5/2)}$; $\nu(\eta) = \eta^{-1}$; $\eta_0 = 2$; and $f(\chi(\xi,\eta)) = (\chi(\xi,\eta))^{\gamma} (E + e^{\chi(\xi,\eta)})$.

Therefore, we get $f(\chi) \ge E\chi$, $\sigma(\eta) = \eta + 1$, $\mu(\eta) = 1$, $R(\eta) = -\eta^{-(5/2)}$, and $1 + \mu(\eta)R(\eta) = 1 - \eta^{-(5/2)} \ge 1 - (1/2) > 0$ which implies $R(\eta) \in \mathcal{R}^+$. Also,

$$\int_{\eta_0}^{\infty} \frac{\Delta^{\alpha} \kappa}{\ell(\kappa)} = \int_{2}^{\infty} \frac{\kappa^{\alpha-1} \Delta \kappa}{\ell(\kappa)} = \sum_{\kappa=2}^{\infty} \frac{\kappa^{\alpha-1}}{\kappa^{-(1/2)}} = \sum_{\kappa=2}^{\infty} 1 = \infty,$$

$$1 > e_{R(\eta)}(\eta, \eta_0) = \exp\left(\int_{\eta_0}^{\eta} R(\kappa) \Delta \kappa\right) \ge 1 + \int_{2}^{\eta} R(\kappa) \Delta \kappa = 1 - \sum_{\kappa=2}^{\eta-1} \kappa^{-(5/2)}$$

$$\ge 1 - \int_{1}^{\eta-1} \kappa^{-(5/2)} d\kappa = 1 + \frac{2}{3} \left((\eta - 1)^{-(3/2)} - 1\right) > \frac{1}{3}.$$
(45)

Let $\psi(\eta) = \eta$. Then, all the conditions of Theorem 4 are fulfilled. Hence, (44) with the given boundary condition is oscillatory.

6. Conclusion

In this article, we established the oscillation criteria for the fractional partial dynamic equation on time scales. The obtained results are improved in the sense that which provide sufficient criteria for the oscillation of the considered equation with two boundary conditions. Finally, numerical examples are also presented to validate the theoretical results of this study. Additionally, more complex systems, including the fractional dynamic equation with time delay, are interesting research topics and will be examined in our future works.

Data Availability

No data were used to support this study.

Conflicts of Interest

The authors declare that there are no conflicts of interest regarding the publication of this paper.

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