

## Research Article

# Phase Portraits and Traveling Wave Solutions of Fokas System in Monomode Optical Fibers

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The main purpose of this paper is to investigate the bifurcation and traveling wave solution of the Fokas system in monomode optical fibers by using the method of planar dynamical system. Firstly, the Fokas systems are reduced to two-dimensional planar dynamic system by using the traveling wave transformation. Secondly, by selecting fixed parameters, the phase portraits are drawn by using the Maple software. Finally, the Jacobi elliptic function solutions, the hyperbolic function solutions, and trigonometric function solutions of the Fokas system are obtained.

## 1. Introduction

In recent years, the complex nonlinear partial differential equations [1–4] have been widely used in nonlinear optics, fluid mechanics, quantum mechanics, biology, communication, control, and other fields [5–9], which mainly include the well-known Schrödinger equation [10], the Radhakrishnan-Kundu-Lakshmanan equation [11], the Kundu-Mukherjee-Naskar equation [12], the Lakshmanan-Porsezian-Daniel equation [13], the Triki-Biswas equation [14], the Fokas-Lenells equation [15], the Gerdjikov-Ivanov equation [16], the Ginzburg-Landau equation [17], and the Biswas-Arshed equation [18]. The study of dynamic behavior and exact traveling wave solutions of complex nonlinear partial differential equations has always been a very important hot topic.

The Fokas system first proposed by Fokas [19] is a very important class of complex nonlinear partial differential equations, which is a generalization of nonlinear Schrödinger equation. In this paper, we consider the following Fokas system [20, 21]:

$$\begin{cases} ip_t + a_1 p_{xx} + a_2 pq = 0, \\ a_3 q_y - a_4 (|p|^2)_x = 0, \end{cases} \quad (1)$$

where  $p = p(t, x, y)$  and  $q = q(t, x, y)$  are the complex valued functions, which represent the propagation of nonlinear pulse in monomode optical fibers. The parameters  $a_1$ ,  $a_2$ ,  $a_3$ , and  $a_4$  are the nonzero constants. In recent years, many experts and scholars have studied the exact solution of Fokas system, many important methods have been proposed, such as the Jacobi elliptic function expansion method [22], the Hirota's bilinear method [23–26], the Painlevé analysis method [27], the bilinear Bäcklund transformation method [28], the  $\exp(-\psi(k))$ -expansion method [29], the improved F-expansion method [30], the extended rational sine-cosine method [31], and the Exp-function method [32]. Although some important methods of constructing Fokas system have been established, the bifurcation of the dynamic system of Fokas system and the more general traveling wave solution have not been studied. In the paper, our main purpose is to study the bifurcation and traveling wave solutions of the Fokas system by using the theory of the planar dynamic system [33, 34].

The article is organized as follows: in Section 2, the phase portraits are drawn. Moreover, the traveling wave solution of (1) is constructed by using the method of planar dynamic system. In Section 3, a conclusion is presented.

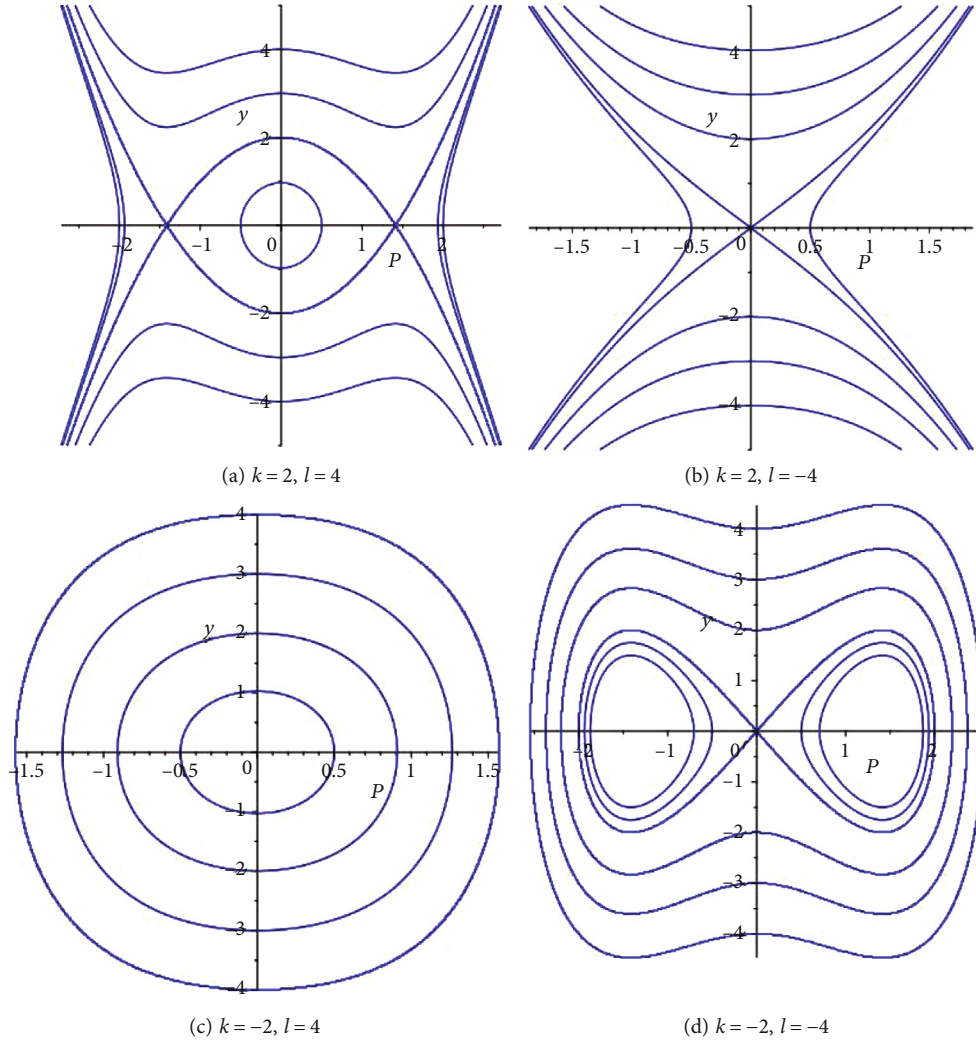


FIGURE 1: The phase portraits of system (8).

## 2. Phase Portraits and Traveling Wave Solution of (1)

First, let us make a traveling wave transformation as follows:

$$p(t, x, y) = P(\xi)e^{i(k_1x+k_2y+k_3t+k_4)}, q(t, x, y) = Q(\xi), \xi = x + y - vt, \quad (2)$$

where  $k_1, k_2,$  and  $k_3$  are real parameters and  $v$  stands for the wave transformation.

Applying the traveling wave transformation, we obtain

$$(-v + 2a_1k_1) iP' - k_3P + a_1P'' - a_1k_1^2P + a_2PQ = 0, \quad (3)$$

$$a_3Q' - 2a_4PP' = 0. \quad (4)$$

Integrating both sides of Equation (4) and making the integral constant be zero, we obtain

$$Q(\xi) = \frac{a_4}{a_3} P^2(\xi). \quad (5)$$

Then, substituting Equation (5) into Equation (3) and setting  $v = 2a_1k_1$ , we have

$$P'' + lP - kP^3 = 0, \quad (6)$$

where  $l = -(k_3 + a_1k_1^2/a_1)$  and  $k = -(a_2a_4/a_1a_3)$ .

Let  $dP(\xi)/d\xi = y$ . System (6) can be transformed into a two-dimensional planar dynamic system:

$$\begin{cases} \frac{dP(\xi)}{d\xi} = y, \\ \frac{dy}{d\xi} = kP^3 - lP, \end{cases} \quad (7)$$

and the Hamiltonian system as follows:

$$H(P, y) = \frac{1}{2}y^2 - \frac{k}{4}P^4 + \frac{l}{2}P^2 = h, \quad h \in \mathbb{R}. \quad (8)$$

Assume that  $M_i(P_i, 0)$  be the equilibrium points of system (7). Suppose further that  $F(P) = kP^3 - lP$ ,  $\lambda_{1,2} = \pm$

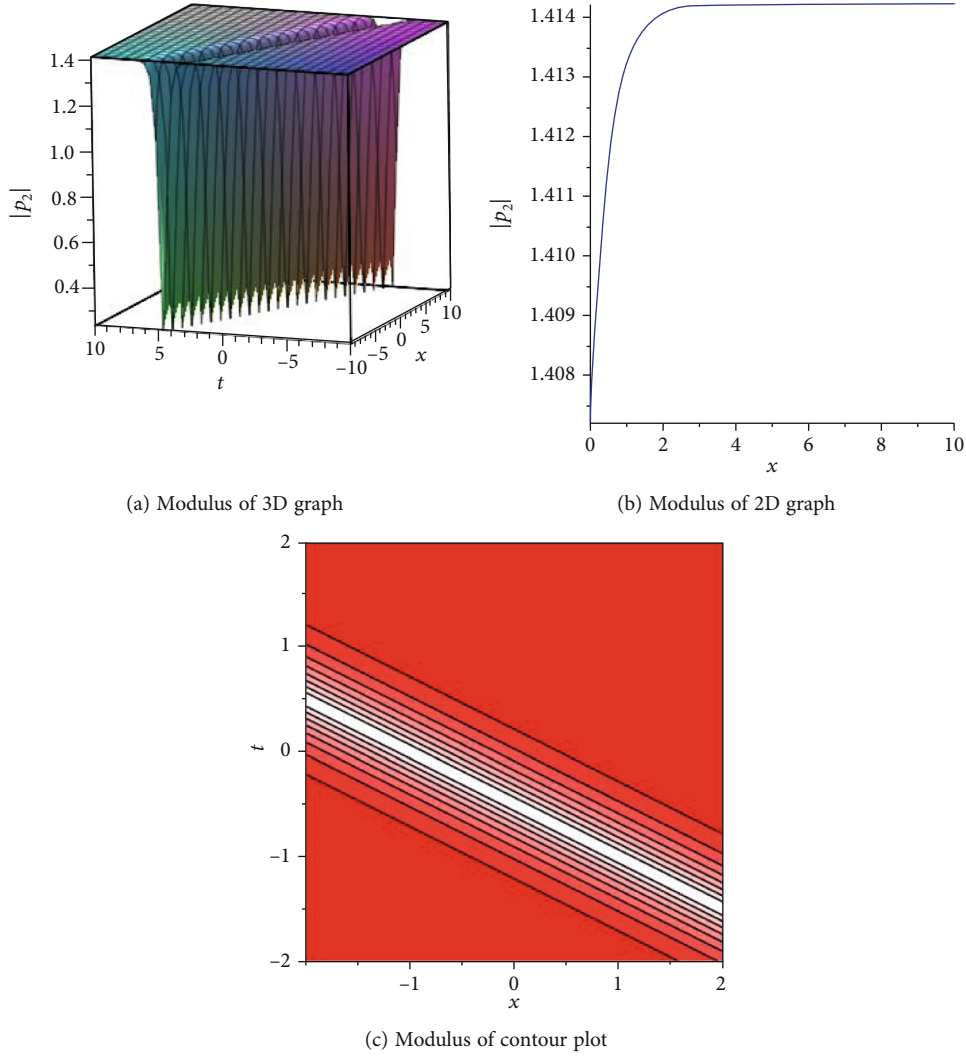


FIGURE 2: Traveling wave solution of  $p_2(t, x, y)$  with  $a_1 = -1, a_2 = 2, a_3 = 1, a_4 = 1/2, k_1 = 1, k_3 = 3, \xi_0 = 0, y = 1$ .

$\sqrt{F'(P)}$ ,  $h_0 = H(0, 0) = 0$ , and  $h_1 = H(\pm\sqrt{l/k}, 0) = l^2/4k$ . If  $k > 0$ , system (8) has three equilibrium points  $M_1(0, 0)$ ,  $M_2(\sqrt{l/k}, 0)$ , and  $M_3(-\sqrt{l/k}, 0)$ . Similarly, if  $kl < 0$ , system (8) has one equilibrium point  $M_4(0, 0)$ . If  $F'(P_i) > 0$ ,  $F'(P_i) = 0$ , and  $F'(P_i) < 0$ , we can obtain that the equilibrium point  $M_i(P_i, 0)$  is saddle point, degraded saddle point, and center point, respectively. The phase portraits of system (8) are drawn as shown in Figure 1.

Case 1  $k > 0, l > 0$ .

(i) If  $h \in (0, l^2/4k)$ , Equation (8) can be rewritten as

$$y^2 = \frac{k}{2}P^4 - lP^2 + 2h = \frac{k}{2}(P^2 - \mu_1^2)(P^2 - \mu_2^2), \quad (9)$$

where  $\mu_1 = \sqrt{(l + \sqrt{l^2 - 4kh})/k}$  and  $\mu_2 = \sqrt{(l - \sqrt{l^2 - 4kh})/k}$ .

Substituting (9) into  $dP(\xi)/d\xi = y$ , we can obtain

$$\int_P^{\mu_2} \frac{dP}{\sqrt{(P^2 - \mu_1^2)(P^2 - \mu_2^2)}} = \pm \sqrt{\frac{k}{2}}(\xi - \xi_0). \quad (10)$$

Integrating Equation (10) and combining Equation (2), we can obtain the Jacobi elliptic function solutions of Equation (1).

$$p_1(t, x, y) = \pm \mu_1 \operatorname{sn} \left( \mu_2 \sqrt{\frac{k}{2}}(x + y - 2a_1 k_1 t - \xi_0), \frac{\mu_1}{\mu_2} \right) e^{i(k_1 x + k_2 y + k_3 t + k_4)},$$

$$q_1(t, x, y) = \frac{a_4 \mu_1^2}{a_3} \left[ \operatorname{sn} \left( \mu_2 \sqrt{\frac{k}{2}}(x + y - 2a_1 k_1 t - \xi_0), \frac{\mu_1}{\mu_2} \right) \right]^2 e^{i(k_1 x + k_2 y + k_3 t + k_4)}. \quad (11)$$

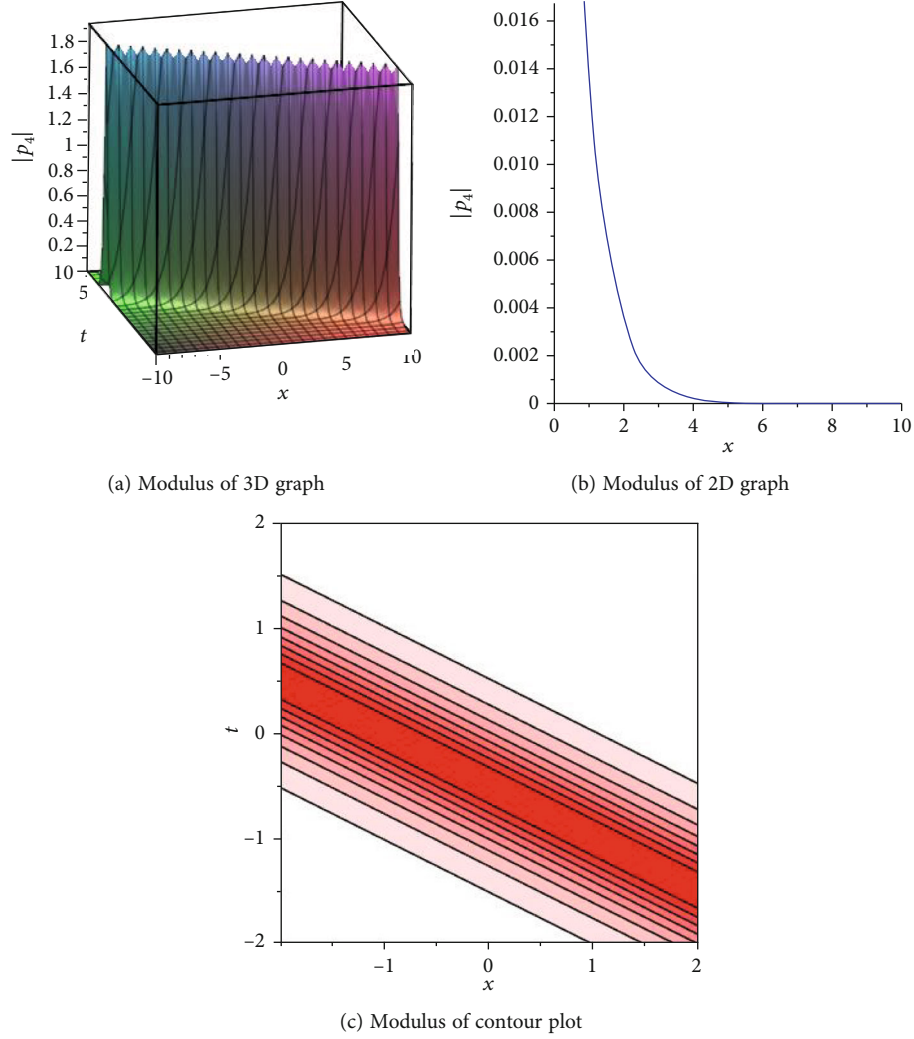


FIGURE 3: Traveling wave solution of  $p_4(t, x, y)$  with  $a_1 = 1, a_2 = 2, a_3 = 1, a_4 = 1/2, k_1 = -1, k_3 = 1, \xi_0 = 0, y = 1$ .

(ii) If  $h = l^2/4k$ , we obtain  $\mu_1 = \mu_2 = \pm\sqrt{lk}$ . Then, we obtain the hyperbolic function solution of Equation (1).

$$p_2(t, x, y) = \pm\sqrt{\frac{l}{k}} \tanh\left(\sqrt{\frac{l}{2}}(x+y-2a_1k_1t-\xi_0)\right) e^{i(k_1x+k_2y+k_3t+k_4)},$$

$$q_2(t, x, y) = \frac{a_4l}{a_3k} \left[ \tanh\left(\sqrt{\frac{l}{2}}(x+y-2a_1k_1t-\xi_0)\right) \right]^2 e^{i(k_1x+k_2y+k_3t+k_4)}. \quad (12)$$

When  $a_1 = -1, a_2 = 2, a_3 = 1, a_4 = 1/2, k_1 = 1, k_3 = 3, \xi_0 = 0, y = 1$ , the modulus of solution  $p_2(t, x, y)$  is plotted as shown in Figure 2.

Case 2  $k < 0, l < 0$ .

(i) If  $h \in (-l^2/4k, 0)$ , Equation (8) can be rewritten as

$$y^2 = \frac{k}{2}P^4 - lP^2 + 2h = -\frac{k}{2}(P^2 - Q_1^2)(Q_2^2 - P^2), \quad (13)$$

where  $Q_1 = \sqrt{(l + \sqrt{l^2 - 4kh})/k}$  and  $Q_2 = \sqrt{(l - \sqrt{l^2 - 4kh})/k}$ . Substituting (9) into  $dP(\xi)/d\xi = y$ , we can obtain

$$\int_P^{Q_2} \frac{dP}{\sqrt{(P^2 - Q_1^2)(Q_2^2 - P^2)}} = \pm\sqrt{-\frac{k}{2}}(\xi - \xi_0), \quad (14)$$

$$\int_{-Q_2}^P \frac{dP}{\sqrt{(P^2 - Q_1^2)(Q_2^2 - P^2)}} = \pm\sqrt{-\frac{k}{2}}(\xi - \xi_0). \quad (15)$$

Integrating Equation (14) and Equation (15), we can

obtain the Jacobi elliptic function solutions of Equation (1).

$$\begin{aligned} p_3(t, x, y) &= \pm Q_2 \operatorname{dn} \left( Q_2 \sqrt{-\frac{k}{2}} (x + y - 2a_1 k_1 t - \xi_0), \frac{\sqrt{Q_2^2 - Q_1^2}}{Q_2} \right) e^{i(k_1 x + k_2 y + k_3 t + k_4)}, \\ q_3(t, x, y) &= \frac{a_4 Q_2^2}{a_3} \left[ \operatorname{dn} \left( Q_2 \sqrt{-\frac{k}{2}} (x + y - 2a_1 k_1 t - \xi_0), \frac{\sqrt{Q_2^2 - Q_1^2}}{Q_2} \right) \right]^2 e^{i(k_1 x + k_2 y + k_3 t + k_4)}. \end{aligned} \quad (16)$$

(ii) If  $h = 0$ , we obtain  $Q_1 = 0$  and  $Q_2 = \sqrt{2l/k}$ . Then, we obtain the solution of Equation (1).

$$\begin{aligned} p_4(t, x, y) &= \pm \sqrt{\frac{2l}{k}} \operatorname{sech} \left( \sqrt{-l} (x + y - 2a_1 k_1 t - \xi_0) \right) e^{i(k_1 x + k_2 y + k_3 t + k_4)}, \\ q_4(t, x, y) &= \frac{2la_4}{ka_3} \left[ \operatorname{sech} \left( \sqrt{-l} (x + y - 2a_1 k_1 t - \xi_0) \right) \right]^2 e^{i(k_1 x + k_2 y + k_3 t + k_4)}. \end{aligned} \quad (17)$$

When  $a_1 = 1, a_2 = 2, a_3 = 1, a_4 = 1/2, k_1 = -1, k_3 = 1, \xi_0 = 0, y = 1$ , the modulus of solution  $p_4(t, x, y)$  is plotted as shown in Figure 3.

(iii) If  $0 < h < +\infty$ , Equation (8) can be rewritten as

$$y^2 = \frac{k}{2} P^4 - l P^2 + 2h = -\frac{k}{2} (Q_3^2 + P^2) (Q_4^2 - P^2), \quad (18)$$

$$\text{where } Q_3 = \sqrt{-(l + \sqrt{l^2 - 4kh})/k} \quad \text{and} \quad Q_4 = \sqrt{(l - \sqrt{l^2 - 4kh})/k}.$$

Substituting (18) into  $dP(\xi)/d\xi = y$ , we can obtain

$$\int_0^P \frac{dP}{\sqrt{(Q_3^2 + P^2)(Q_4^2 - P^2)}} = \pm \sqrt{-\frac{k}{2}} (\xi - \xi_0). \quad (19)$$

Integrating Equation (19), we can obtain the Jacobi elliptic function solutions of Equation (1).

$$\begin{aligned} p_5(t, x, y) &= \pm Q_3 \operatorname{cn} \left( Q_2 \sqrt{-\frac{k(Q_3^2 + Q_4^2)}{2}} (x + y - 2a_1 k_1 t - \xi_0), \frac{Q_4}{\sqrt{Q_3^2 + Q_4^2}} \right) e^{i(k_1 x + k_2 y + k_3 t + k_4)}, \\ q_5(t, x, y) &= \frac{a_4 Q_3^2}{a_3} \left[ \operatorname{cn} \left( Q_2 \sqrt{-\frac{k(Q_3^2 + Q_4^2)}{2}} (x + y - 2a_1 k_1 t - \xi_0), \frac{Q_4}{\sqrt{Q_3^2 + Q_4^2}} \right) \right]^2 e^{i(k_1 x + k_2 y + k_3 t + k_4)}. \end{aligned} \quad (20)$$

### 3. Conclusion

In this paper, we study the bifurcation and traveling wave solution of the Fokas system in monomode optical fibers by using the method of planar dynamical system. By select-

ing fixed parameters, the phase portraits, three-dimensional graphs, two-dimensional graphs, and contour plot are drawn by using the Maple software, which explains the propagation of optical solitons in nonlinear optical fibers from different angles. Compared with the published literature [20, 21], the solutions obtained in this paper are more abundant; we also get the Jacobi elliptic function solutions, the hyperbolic function solutions, and trigonometric function solutions of the Fokas system.

### Data Availability

No underlying data was collected or produced in this study.

### Conflicts of Interest

The author declares no conflict of interest.

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### References

- [1] L. Tang, "Bifurcations and dispersive optical solitons for the cubic-quartic nonlinear Lakshmanan-Porsezian-Daniel equation in polarization-preserving fibers," *Optik*, vol. 270, p. 170000, 2022.
- [2] Z. Li, P. Li, and T. Y. Han, "Bifurcation, traveling wave solutions, and stability analysis of the fractional generalized Hirota-Satsuma coupled KdV equations," *Discrete Dynamics in Nature and Society*, vol. 2021, Article ID 5303295, 6 pages, 2021.
- [3] L. Tang, "Bifurcation analysis and multiple solitons in birefringent fibers with coupled Schrodinger-Hirota equation," *Chaos, Solitons & Fractals*, vol. 161, p. 112383, 2022.
- [4] L. Tang, "Bifurcations and multiple optical solitons for the dual-mode nonlinear Schrodinger equation with Kerr law nonlinearity," *Optik*, vol. 265, p. 169555, 2022.
- [5] J. X. Zuo and X. C. Lin, "High-power laser systems," *Laser & Photonics Reviews*, vol. 16, no. 5, p. 2100741, 2022.
- [6] X. D. Cai, R. Tang, H. Y. Zhou et al., "Dynamically controlling terahertz wavefronts with cascaded metasurfaces," *Advanced Photonics*, vol. 3, p. 036003, 2021.
- [7] T. Zhong, M. Cheng, S. P. Lu, X. T. Dong, and Y. Li, "RCEN: a deep-learning-based background noise suppression method for DAS-VSP records," *IEEE Geoscience and Remote Sensing Letters*, vol. 19, p. 342, 2022.
- [8] D. Yang, T. Y. Zhu, S. Wang, S. Z. Wang, and Z. Xiong, "LFRSNet: a robust light field semantic segmentation network combining contextual and geometric features," *Frontiers in environmental Science*, vol. 10, p. 996513, 2022.
- [9] B. Zhu, Q. S. Zhong, Y. S. Chen et al., "A novel reconstruction method for temperature distribution measurement based on ultrasonic tomography," *IEEE Transactions on Ultrasonics, Ferroelectrics, and Frequency Control*, vol. 69, no. 7, pp. 2352–2370, 2022.
- [10] T. Y. Han, Z. Li, and X. Zhang, "Bifurcation and new exact traveling wave solutions to time-space coupled fractional

- nonlinear Schrodinger equation," *Physics Letters A*, vol. 395, p. 127217, 2021.
- [11] K. Zhang, X. Y. He, and Z. Li, "Bifurcation analysis and classification of all single traveling wave solution in fiber Bragg gratings with Radhakrishnan-Kundu-Lakshmanan equation," *AIMS Mathematics*, vol. 7, no. 9, pp. 16733–16740, 2022.
- [12] C. Peng, F. R. Zhang, H. W. Zhao, and Z. Li, "New optical solitons in Bragg grating fibers for the nonlinear coupled (-)-dimensional Kundu-Mukherjee-Naskar system via complete discrimination system method," *Adv. Math. Phys.*, vol. 2022, article 8184270, pp. 1–13, 2022.
- [13] C. Peng, Z. Li, and H. W. Zhao, "New exact solutions to the Lakshmanan-Porsezian-Daniel equation with Kerr law of nonlinearity," *Mathematical Problems in Engineering*, vol. 2022, Article ID 7340373, 10 pages, 2022.
- [14] Z. Li and Z. G. Lian, "Optical solitons and single traveling wave solutions for the Triki-Biswas equation describing monomode optical fibers," *Optik*, vol. 258, p. 168835, 2022.
- [15] M. M. A. Khater, A. E. S. Ahmed, S. H. Alfalqi, J. F. Alzaidi, S. Elbendary, and A. M. Alabdali, "Computational and approximate solutions of complex nonlinear Fokas-Lenells equation arising in optical fiber," *Results in Physics*, vol. 25, p. 104322, 2021.
- [16] Z. Li and T. Y. Han, "Classification of all single traveling wave solutions of fractional perturbed Gerdjikov-Ivanov equation," *Mathematical Problems in Engineering*, vol. 2021, Article ID 1283083, 7 pages, 2021.
- [17] Z. Li and T. Y. Han, "New exact traveling wave solutions of the time fractional complex Ginzburg-Landau equation via the conformable fractional derivative," *Advances in Mathematical Physics*, vol. 2021, Article ID 8887512, 12 pages, 2021.
- [18] Z. Li, "Bifurcation and traveling wave solution to fractional Biswas-Arshed equation with the beta time derivative," *Chaos, Solitons & Fractals*, vol. 160, p. 112249, 2022.
- [19] A. S. Fokas, "Integrable nonlinear evolution partial differential equations in 4+2 and 3+1 dimensions," *Physical Review Letters*, vol. 96, p. 190201, 2016.
- [20] S. Tarla, K. K. Ali, T. C. Sun, R. Yilmazer, and M. S. Osman, "Nonlinear pulse propagation for novel optical solitons modeled by Fokas system in monomode optical fibers," *Results in Physics*, vol. 36, p. 105381, 2022.
- [21] K. Zhang, T. Y. Han, and Z. Li, "New single traveling wave solution of the Fokas system via complete discrimination system for polynomial method," *AIMS Mathematics*, vol. 8, no. 1, pp. 1925–1936, 2023.
- [22] H. Khatri, M. S. Gautam, and A. Maik, "Localized and complex soliton solutions to the integrable (4+1)-dimensional Fokas equation," *SN Applied Sciences*, vol. 1, pp. 1–9, 2019.
- [23] S. Zhang, C. Tian, and W. Y. Qian, "Bilinearization and new multisoliton solutions for the (4+1)-dimensional Fokas equation," *Pramana-Journal of Physics*, vol. 86, no. 6, pp. 1259–1267, 2016.
- [24] Y. L. Cao, J. G. Rao, D. Mihalache, and J. S. He, "Semi-rational solutions for the (2+1)-dimensional nonlocal Fokas system," *Applied Mathematics Letters*, vol. 80, pp. 27–34, 2018.
- [25] Y. L. Cao, J. S. He, Y. Cheng, and D. Mihalache, "Reduction in the (4+1)-dimensional Fokas equation and their solutions," *Nonlinear Dynamics*, vol. 99, no. 4, pp. 3013–3028, 2020.
- [26] J. G. Rao, D. Mihalache, Y. Cheng, and J. S. He, "Lump-soliton solutions to the Fokas system," *Physics Letters A*, vol. 383, no. 11, pp. 1138–1142, 2019.
- [27] W. Tan, Z. D. Dai, and D. Q. Qiu, "Parameter limit method and its application in the (4+1)-dimensional Fokas equation," *Computers and Mathematics with Applications*, vol. 75, no. 12, pp. 4214–4220, 2018.
- [28] R. X. Yao, Y. L. Shen, and Z. B. Li, "Lump solutions and bilinear Bäcklund transformation for the (4+1)-dimensional Fokas equation," *Mathematical Sciences*, vol. 14, no. 4, pp. 301–308, 2020.
- [29] P. Verma and L. Kaur, "New exact solutions of the (4+1)-dimensional Fokas equation via extended version of  $(\exp(-\psi(k)))$ -expansion method," *International Journal of Computers and Applications*, vol. 7, no. 3, p. 104, 2021.
- [30] S. Sarwar, "New soliton wave structures of nonlinear (4 + 1)-dimensional Fokas dynamical model by using different methods," *Alexandria Engineering Journal*, vol. 60, no. 1, pp. 795–803, 2021.
- [31] K. J. Wang, "Abundant exact soliton solutions to the Fokas system," *Optik*, vol. 249, p. 168265, 2022.
- [32] K. J. Wang, J. Liu, and W. Jun, "Soliton solutions to the Fokas system arising in monomode optical fibers," *Optik*, vol. 251, p. 168319, 2022.
- [33] A. Saha and S. Banerjee, *Nonlinear Dynamics and Applications*, Springer Cham, Switzerland, 2022.
- [34] A. Saha and S. Banerjee, *Dynamical Systems and Nonlinear Waves in Plasmas*, Springer, Boca Raton, 2021.