Research Article

Electromagnetic Interference Shielding and Characterization of Ni\(^{2+}\) Substituted Cobalt Nanoferrites Prepared by Sol-Gel Auto Combustion Method

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Received 24 September 2022; Revised 25 October 2022; Accepted 14 November 2022; Published 28 November 2022

Academic Editor: Hom Kandel

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The structural, magnetic, and dielectric properties of a series of Ni\(^{2+}\) substituted cobalt nanoferrite particle samples with the composition \(\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4\) (where \(0.0 \leq x \leq 1.0\)) synthesized by using the sol-gel auto combustion route are presented in this report. The electromagnetic interference shielding of \(\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/\text{PVA}\) nanocomposite films has been determined in the microwave X-band (8.2–12.4 GHz) frequencies. X-ray analysis revealed the single-phase formation of nickel-substituted cobalt nanoferrite samples. The decreasing trend of lattice parameters with Ni\(^{2+}\) substitution indicates the incorporation of Ni\(^{2+}\) into the crystal structure, obeying Vegard’s law. FTIR showed the absorption bands at 560–590 cm\(^{-1}\) (\(\nu_1\)) and 390–400 cm\(^{-1}\) (\(\nu_2\)) were attributed to (A-site) tetrahedral and (B-site) octahedral groups complex, respectively which confirm the spinel structure of the samples. Field emission scanning electron microscopy showed agglomerated grains of different sizes and shapes in the morphological observation. EDS reveals the chemical composition of the prepared samples. TEM analysis revealed that the synthesized particles were nearly monodisperse, show to be roughly spherical in shape, and have a polycrystalline nature. The dielectric constant and loss tangent (\(\tan\delta\)) is found to decrease with increasing frequency which shows normal behavior for ferrimagnetic materials. The magnetic properties determined using VSM have substantially changed with the substitution of Ni\(^{2+}\) ions. The saturation magnetization and the experimentally magnetic moment are observed to decrease with an increase in Ni\(^{2+}\) content \(x\). A series of \(\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/\text{PVA}\) nanocomposite films are prepared by applying simple, rapid, and inexpensive methods for EMI shielding materials. The vector network analyzer data were used to evaluate the electromagnetic interference (EMI) shielding properties of the \(\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/\text{PVA}\) samples. At 9.2 GHz, a study of reflection loss showed a minimum reflection loss (RL) of ~32.08 dB. Also, the synthesized \(\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/\text{PVA}\) nanocomposite samples show improved performance for EMI efficiency which proves the utility of this doping. With this low RL value, the results and techniques also promise a simple, effective approach to achieve light-weight \(\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/\text{PVA}\) nanocomposite films and make it excellent microwave absorbers, capable of working at gigahertz frequencies for application potentials in EMI shielding material, communication, radar stealth technology, and electronic warfare.

1. Introduction

Ferrite is produced by mixing metallic elements including cobalt (Co), barium (Ba), magnesium (Mg), nickel (Ni), manganese (Mn), and zinc (Zn), in small amounts with iron oxide (\(\text{Fe}_2\text{O}_4\)) in large amounts [1]. The general chemical formula for spinel ferrites is \(M\text{Fe}_2\text{O}_4\) where \(M\) is the divalent metallic ion. The characteristic of spinel ferrites can be changed by substituting a different metal ion.

The synthesis and characterization of nanoferrite particles are taking great interest due to their wide range of applications in many areas such as electrical [2], electronic...
Advances in Materials Science and Engineering

Based on polyvinyl alcohol and Ni²⁺ substituted cobalt nanoferrite particles is reported in detail. The nanocomposite detail.

frequency range of X-band (8.2–12.4GHz) was discussed in shielding against the interference caused by EMI signal in the ration magnetization [9].

volume ratio, high magnetic permeability, and high satu-

frequency magnetic felds are good nanomaterials [22, 23].

resistance to decrease. Ferromagnetic materials for low-fre-

loop toward the magnetic resistance causing the magnetic

nanomaterials permeability value results in a magnetic fux

EMI shielding due to their electrical properties. Magnetic field shielding works by suppressing radiation from the external magnetic field, which is a low-frequency field. Increasing the nanomaterials permeability value results in a magnetic flux loop toward the magnetic resistance causing the magnetic resistance to decrease. Ferromagnetic materials for low-frequency magnetic fields are good nanomaterials [22, 23].

In this study, the effect of Ni⁷⁺ substitution on the structural, magnetic, and dielectric characteristics of cobalt nano ferrite nanoparticles is reported in detail. The nanocomposite based on polyvinyl alcohol and Ni²⁺ substituted cobalt nanoferrite particles is synthesized and tested to study shielding against the interference caused by EMI signal in the frequency range of X-band (8.2–12.4GHz) was discussed in detail.

2. Chemicals and Reagents

Analytical grade chemicals and reagents from Honeywell/Fluka and Biochem Chemopharma were obtained. These materials were used for the synthesis of Ni²⁺ substituted cobalt spinel nanoferrite particles with the composition of Co₁₋ₓNiₓFe₂O₄ (where x = 0.0 ≤ x ≤ 1.0) without further purification. Cobalt (II) nitrate hexahydrate (Co(NO₃)₂·6H₂O), nickel (II) nitrate hexahydrate (Ni(NO₃)₂·6H₂O), and iron (III) nitrate nonahydrate (Fe(NO₃)₃·9H₂O) were used as a metal ion source. Citric acid monohydrate (C₆H₈O₇·H₂O) acts as a chelating agent (fuel) and helps in the homogeneous distribution of metal ions. Dilute ammonia (NH₃) is used to maintain the solution neutral pH ~ 7. Ultrapure water (Milli-Q grade, 18 MΩ·cm⁻¹) is used during the experimentation for the synthesis of the materials. Details of chemicals and reagents are provided in Table 1.

3. Experimental Procedures

The experimental procedures are described in the following sections.

3.1. Synthesis of Co₁₋ₓNiₓFe₂O₄ Nanoparticles. Using the sol-gel auto combustion approach [24], nanocrystalline Ni²⁺ substituted cobalt nano ferrite of the chemical composition Co₁₋ₓNiₓFe₂O₄ (where x = 0.0 ≤ x ≤ 1.0) is synthesized as shown in Figure 1. By dissolving the desired proportions of cobalt (II) nitrate, nickel (II) nitrate, and iron (III) nitrate in ultrapure water, citric acid acts as an organic fuel during the calcination process and chelating agent for metal ions in sol-gel auto combustion in preparation of many oxides. Citric acid is separately dissolved in ultrapure water in a separate beaker. The molar ratio of metal nitrate ions to citric acid was 3:2:2. After completely mixing the two solutions by using ultrasonic for 10 minutes, continual stirring was carried out until the temperature reached 50°C and remained at this temperature for 30 minutes. After that, dilute ammonia was added dropwise to the solution to attain a pH of 7. The temperature was then raised to 90°C. The obtained solution was transformed into a viscous brown gel phase. The viscous gel was placed in an oven at 200°C to initiate an auto combustion reaction and produce a fluffy powder and dried in an oven for 30 minutes. The prepared fluffy powder was ground and the obtained as-burnt pow-

3.2. Pelletizing and Nanocomposite Preparation. A small amount of (2 wt%) PVA was added as a binder to the samples that were calcinated at 800°C before being pressed into circular pellets with a diameter of 12 mm and a thickness of around 2 mm using a hydraulic press at a
Contineous heating at 90 °C for 2hrs and heating by Mortar and Pestle Manual grinding under continuous stirring at 80°C. Ten, 10% weight of the sample was added to 40 mlof water then sintered at 900 °C for 3 hours to densify the samples and remove the excess binder. Slowly it is allowed to cool naturally to evaluate the dielectric characteristics.

4. Results and Discussion

Results and discussion are discussed in the following sections.

4.1. Crystallographic Analysis. High-intensity X-ray diffraction (XRD, model Panalytical (X’pert Pro, Netherlands)) equipped with a copper-κα radiation source (with an incident X-ray wavelength of λ = 0.154 nm, 40 mA, 40 kV) confirmed the successful fabrication, crystal structure, and phase identification of the synthesized as-burnt and calcined samples with the composition Co_{1−x}Ni_{x}Fe_{2}O_{4} (where x = 0.0 ≤ x ≤ 1.0), all measurements were taken in increments of 0.02° in the 2θ range (20°–70°). Figure 3 shows the X-ray diffraction pattern of the Co_{1−x}Ni_{x}Fe_{2}O_{4} (where x = 0.0 ≤ x ≤ 1.0) nanoparticles for as-burnt and calcined at 600°C, 700°C, and 800°C.

Calcination was carried out to remove any external phases and reduce internal stress, the size of the Co_{1−x}Ni_{x}Fe_{2}O_{4} (where x = 0.0 ≤ x ≤ 1.0) nanoparticles also increase which led to crystal growth prepared nanoparticles. All the samples were found to be face-centered cubic (FCC) with an Fd-3m space group. According to the standard ICSD cards [00-001-1121] and [00-010-0325], the product can be mainly indexed with the miller indices of the reflection planes of (111), (220), (311), (222), (400), (422), (511), and (440), which can be indexed to CoFe_{2}O_{4} and NiFe_{2}O_{4} respectively, and this pattern was also confirmed with previously published data [25–27]. These diffraction peaks confirmed the formation of nanometer-sized particles, and the substitution of the spinel crystal structure. The XRD result indicates both high crystallinity and purity of the Co_{1−x}Ni_{x}Fe_{2}O_{4} (where x = 0.0 ≤ x ≤ 1.0) nanoparticles samples.

The peak attributed to the (311) plane has the highest peak intensity, and it changes to a higher diffraction angle when nickel substitution in cobalt ferrites increases. The shifting of the (311) peak towards higher 2theta angles, which indicates a drop in lattice constant with Ni^{2+} substitution, is clearly shown in Figure 4. The result might be due to the difference in ionic radii of Ni^{2+} (0.69 Å) ion and Co^{2+} (0.745 Å) ion.

<table>
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<th>Molar mass</th>
<th>Role</th>
<th>Supplier</th>
</tr>
</thead>
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<td>Co(NO3)2.6H2O</td>
<td>291.04</td>
<td>Metal ion source</td>
<td>Honeywell/fluka</td>
</tr>
<tr>
<td>Nickel (II) nitrate</td>
<td>Ni(NO3)2.6H2O</td>
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<td>Metal ion source</td>
<td>Biochem chemopharma</td>
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<tr>
<td>Iron (III) nitrate</td>
<td>Fe(NO3)3.9H2O</td>
<td>404.06</td>
<td>Metal ion source</td>
<td>Honeywell/fluka</td>
</tr>
<tr>
<td>Ammonia</td>
<td>NH3</td>
<td>18.02</td>
<td>pH controller</td>
<td>Biochem chemopharma</td>
</tr>
<tr>
<td>Ultrapure water (milli-Q grade, 18 MΩ·cm⁻¹)</td>
<td>H2O</td>
<td>18.02</td>
<td>Reaction medium</td>
<td>Biochem chemopharma</td>
</tr>
</tbody>
</table>

| Table 1: Chemicals and reagents. |

![Figure 1: Schematic representation of the sol-gel autocombustion method for the Co_{1−x}Ni_{x}Fe_{2}O_{4} (where x = 0.0 ≤ x ≤ 1.0).](image-url)
The lattice parameter “a” can be determined using the following relation in the case of cubic crystal structure [28]:

$$\frac{1}{d_{hkl}^2} = \frac{h^2 + k^2 + l^2}{a^2},$$  \hspace{1cm} (1)

$$V = a^3,$$

where $d$ is the interplanar spacing and $h$, $k$, and $l$ are the Miller indices of the crystal planes. Figure 5 shows that the lattice parameter value decreases linearly with nickel concentration, which can be described by the decreased ionic radius of Ni$^{2+}$ compared to Co$^{2+}$ this variation can be explained by Vegard’s law [29]. The unit cell dimensions of pure nickel ferrites and pure cobalt ferrites are quite similar to the values obtained in this report.

Because of the smaller crystallite size, the peak intensity decreases as nickel substitution increases from $x=0.0$ to $x=1.0$. From the most intense peak (311), the average crystallite size of the prepared Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x=0.0 \leq x \leq 1.0$) nano ferrite particles were calculated using the Scherrer formula [30].

$$D = \frac{kl}{\beta \cos \theta},$$  \hspace{1cm} (2)

where $D$ represents the average crystallite size, $k$ corresponds to the Scherrer constant (0.90), $\beta$ is the full width at half-maximum, $\lambda$ is the radiation wavelength of the X-ray, and $\theta$ is the diffraction angle corresponding to (311) plane, respectively. This method produces nano-crystals with a size of 49–65 nm, which is suitable for practical applications such as wireless communication, electronic warfare, radiation medical exposure, sensing, and aerospace that need small-volume, light, and very efficient screening systems to protect a specific component and its surroundings [31]. The peaks in the XRD pattern for all the samples get sharper and narrower when the calcined temperature increases to 800°C, with a decrease in their FWHM (full width at half maxima). This shows that crystallinity and particle size have increased. The observed behavior indicates that the present samples have nanocrystalline nature. It is detected that with the calcined temperature the particle size also increases as shown in Figure 6.

Calcining the samples will increase the crystallinity. By increasing the temperature, the estimated average crystallite size increases so, particle size and grain size will increase and the dominant peak shape becomes sharper by calcining because the same wavelength is hitting one particle instead of hitting small particles of the same orientation. Additionally, we observed a decrease in the FWHM which is a narrowing of the peaks corresponding to an increase in crystallization of the nano-materials with calcination temperature, if full width at half maximum decreases, which means that the crystallinity also increases [32, 33].

Since each primitive unit cell of the spinel structure contains 8 molecules, the value of the X-ray density $\rho_x$ was calculated according to [34]

$$\rho_x = \frac{8M}{Na^3},$$  \hspace{1cm} (3)

where $M$ represents the molecular weight of the sample, $N$ is Avogadro’s number ($= 6.0225 \times 10^{23}$ atom/mole), and $a$ is the experimental lattice parameter. The measured X-ray density is tabulated in Table 2. With increasing the nickel concentration, X-ray density also increases as in Figure 5, subsequently, the cobalt atom is lighter than the nickel atom. The bulk density $\rho_b$ was calculated from the following formula [35]:

$$\rho_b = \frac{m}{\pi r^2 d},$$  \hspace{1cm} (4)

where $r$ is the radius, $m$ is the mass, and $d$ represents the thickness of the sample. The obtained results are tabulated in Table 2.

Because of the existence of pores in the sample formed during the synthesis process, which is further characterized by the term porosity, bulk density is often lower than X-ray density, and the percentage bulk porosity of the sample was estimated using the following relation [36]:

$$P = \left(1 - \frac{\rho_b}{\rho_x}\right) \times 100,$$  \hspace{1cm} (5)

There is an inverse relationship between bulk porosity and bulk density.
The ferrite systems' physical properties are influenced by hopping length $L$, the distance between the magnetic ions, the hopping length in $A$-site ($L_A$-tetrahedral) and $B$-site ($L_B$-octahedral) were estimated using the following equations [37]:

\[
L_A = \frac{a \sqrt{3}}{4},
\]

\[
L_B = \frac{a \sqrt{2}}{4}.
\]
Table 2 shows the computed values of the hopping length $L_A$ and $L_B$ of the various compositions. As the Ni$^{2+}$-substitution varies, the hopping length changes as well. The amount in the tetrahedral and octahedral sites at which Ni$^{2+}$ substitution changes with increased Ni$^{2+}$ concentration, has a great effect in changing the value of hopping length, which is connected to the fact that Ni$^{2+}$ ion has lower ionic radii than Co$^{2+}$ ion.

Specific surface area ($S$) was estimated by using the following relation:

$$S = \frac{6000}{D \cdot \rho_x}.$$  \hfill (7)

Figure 4: XRD pattern of all the compositions Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) calcinated at 800°C along with the magnified view of (311) peaks.

Figure 5: Variation of lattice parameter ($a$) and X-ray density ($\rho_x$) as functions of Ni$^{2+}$ content.

The calculated values of the hopping factor for the Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) nanoferrite particle samples were presented in Table 2. It was observed that with increasing Ni$^{2+}$ concentration, the crystallite size decreases and a similar trend is observed for the packing factor values.

The structural refinement was carried out using the Rietveld refinement approach [40–42] with the fullprof program after the computation of structural parameters was carried out. Wyckoff positions for cobalt, nickel, and iron,

Figure 6: Variation of particle size ($D$) with calcination temperature.

$$P = \frac{D}{d_{(440)}},$$  \hfill (8)

where $D$ is the crystallite size and $\rho_x$ is the X-ray density [38].
4.2. Vibrational Spectral Analysis. Typical room-temperature Fourier transform infrared (FTIR) spectra, for (Co₁₋ₓNiₓFe₂O₄ (where x = 0.0 ≤ x ≤ 1.0)) nanoferrite particle samples calcinated at 800°C recorded by using (Shimadzu, IRAffinity-1, Japan) by means of KBr pellets is depicted in Figure 9. The spectra obtained for the samples in the frequency range of 300–4000 cm⁻¹. All of the samples’ functional groups have two separate absorption bands below 700 cm⁻¹, which support ferrite crystallization and also confirm the spinel structure of the samples. All the samples show an absorption band around 560–590 cm⁻¹ and 390–400 cm⁻¹ which are found to agree with the previously reported values [25, 27].

Wadron and Hafner have attributed that the higher absorption band (ν₁) around 560–590 cm⁻¹ is appointed as the tetrahedral metal complex groups stretching vibration (Fe³⁺-O²⁻), which consists of bonding between A-site metal cation and oxygen anion. The lower absorption band (ν₂) around 390–400 cm⁻¹ is attributed to the octahedral metal complex groups stretching vibration (Fe³⁺-O²⁻), which is regarded as bonding between B-site metal cation and oxygen anion [45]. Changes in metal-oxygen bond length (Fe³⁺-O²⁻) at both tetrahedral and octahedral coordination are attributed to the formation of two main absorption bands below 700 cm⁻¹ [45]. The Debye temperature is the temperature at which the lattice exhibits the maximum vibration. For all the samples, the following equation is used to calculate the Debye temperature [45]:

\[ \theta_D = \frac{hc\nu_{12}}{k_B}, \]  

where \( c \) stands for the velocity of light (3 × 10⁸ m/s), \( k_B \) denotes Boltzmann’s constant (1.38 × 10⁻²³ J/K), \( h \) represents Plank’s constant (6.624 × 10⁻³⁴ J s), and \( \nu_{12} \) corresponds to the average wavenumber of absorption bands expressed as \( \nu_{12} = (\nu_1 + \nu_2)/2 \), \( \nu_1 \) and \( \nu_2 \) are the frequency of absorption bands related to A-site and B-site. With increasing calcination temperature, the Debye temperature increases, which corresponds to an increase in the normal vibration mode of the crystal. The calculated values of the absorption bands, average absorption bands, and Debye temperature are given in Table 4.

4.3. Field-Emission Scanning Electron Microscopy (FESEM). The cross-sectional morphology of the Co₁₋ₓNiₓFe₂O₄ (where x = 0.0 ≤ x ≤ 1.0) nanoparticles was analyzed by the field-emission scanning electron microscopy (FESEM) (Mira3-XMU, TESCAN, Japan). FESEM micrographs of the CoFe₂O₄, Co₀.₆Nio₀.₄Fe₂O₄, and NiFe₂O₄ nanoferrite particle samples calcinated at 800°C are shown in Figure 10. The ferrites were investigated using FESEM to better understand their morphological microstructures. It illustrates that most of the nanoparticles have a nearly spherical shape. The nanoparticles have a nonuniform and inhomogeneous size distribution. The estimated average grain size diameter of spherical particles for CoFe₂O₄, Co₀.₆Nio₀.₄Fe₂O₄, and NiFe₂O₄ calcinated at 800°C is found to be 139.4, 107.4, and 113.5 nm along with a standard deviation of 67.57, 56.28, and 61.97 nm, this decrease might be due to the lower ionic radii of nickel ions than cobalt ions.

cations were taken at 8a (1/8, 1/8, 1/8), 16d (1/2, 1/2, 1/2), respectively, and 32e (x, y, z) for oxygen for the Fd-3m space group in cubic spinel ferrites [43]. Results of refinement for typical samples (CoFe₂O₄, Co₀.₆Nio₀.₄Fe₂O₄, and NiFe₂O₄) are shown in Figure 7. Based on the visual difference between observed-calculated intensities and the χ² parameter the goodness of fit was decided.

The refined values of the reliability parameters along with the goodness of fit factor (χ²) index, lattice parameter, and volume of unit cells are listed in Table 3. Due to the difference in the ionic radii of Ni²⁺ ions and Co²⁺ ions lattice constant of all the samples decreases with increasing Ni²⁺ content. Figure 8 illustrates the computational modeling structure of the typical sample Co₁₋₉NiₓFe₂O₄ (x = 0.4) created with VESTA software, a program in which 2D and 3D structures of atoms and molecules of the crystal structure can be visualized and modeled. The crystal structure is represented by the ball and stick model, where oxygen is represented by small spheres and either cobalt and nickel or iron is represented by large spheres, and the polyhedral model, where octahedral sites and tetrahedral sites are presented within the crystal structure [44].

<table>
<thead>
<tr>
<th>Parameters</th>
<th>0.0</th>
<th>0.2</th>
<th>0.4</th>
<th>0.6</th>
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<td>Lattice parameter a (Å)</td>
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<td>8.361</td>
<td>8.357</td>
<td>8.327</td>
<td>8.332</td>
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<td>583.669</td>
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<td>578.56</td>
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<td>2.956</td>
<td>2.954</td>
<td>2.944</td>
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<tr>
<td>Surface area (m²/gm)</td>
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<td>18.45</td>
<td>16.917</td>
<td>20.159</td>
<td>22.396</td>
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Table 2: Calculated values of various structural parameters for Co₁₋₉NiₓFe₂O₄ (where x = 0.0 ≤ x ≤ 1.0) nanoferrites, for dominant peak (311) at T = 800°C.
Figure 7: Rietveld refinement XRD plot for the typical (a) CoFe$_2$O$_4$, (b) Co$_{0.6}$Ni$_{0.4}$Fe$_2$O$_4$, and (c) NiFe$_2$O$_4$ samples. The measured data points ($Y_{obs}$) are represented by red solid circles. The black solid lines correspond to calculated patterns ($Y_{calc}$). The green vertical lines are the Bragg positions. The difference between observed and calculated ($Y_{obs}-Y_{calc}$) plots is also shown as lines in blue color at the bottom of each pattern.

Table 3: Parameters obtained from Rietveld refinement for Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) samples.

<table>
<thead>
<tr>
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<td>Profile factor $R_p$ (%)</td>
<td>47.4</td>
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<td>49</td>
<td>96.8</td>
<td>29.6</td>
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<td>Weighted residual factor $R_{wp}$ (%)</td>
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<td>Expected residual factor $R_{exp}$ (%)</td>
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<td>13.17</td>
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<td>Bragg $R$ factor $R_B$ (%)</td>
<td>14.2</td>
<td>21.7</td>
<td>15.7</td>
<td>39</td>
<td>3.9</td>
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<tr>
<td>$R_F$ (%)</td>
<td>15</td>
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<td>12.6</td>
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Teagglomeration is an indication of the prepared sample’s high reactivity during heat treatment, and it could possibly be due to magnetostatics interaction between the nanoparticles as the magnetic nanoparticles attract each other by Van der Waals forces and magnetic dipolar interactions [46]. The formation of agglomerated grain structure is a characteristic feature of the sol-gel auto-combustion process.

Figure 8: Typical polyhedral model crystal structure of $\text{Co}_{0.6}\text{Ni}_{0.4}\text{Fe}_2\text{O}_4$: (a) ball and stick model and (b) polyhedral model.

Figure 9: FTIR spectra of the prepared $\text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4$ (where $x = 0.0 \leq x \leq 1.0$) nanoferrite particle samples calcinated at $800^\circ$C.
4.4. Energy-Dispersive Spectrum (EDS). The attached energy-dispersive spectrum (EDS) model Aztec (Oxford instrument PLC, UK) was used to investigate the qualitative chemical composition of the samples which was carried out to accurately indicate and confirm the existence of Ni, Co, Fe, and O present in the compositions. The energy dispersive spectrum shows that nearly the same stoichiometric amount of the elements is present in the chemical composition without any major impurity. Typical EDS patterns of samples calcined at 800°C are depicted in Figure 11. Table 5 summarizes the weight and atomic percentages of individual Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) nanoferrite particle samples calcined at 800°C. The theoretical weight percentages from the stoichiometric formula compared to the experimentally established percentages are listed in Table 5.

4.5. Transmission Electron Microscopy (TEM). Transmission electron microscopy (TEM) measurements were recorded at operating voltages of 200 kV. Typical high-resolution TEM micrographs for the CoFe$_2$O$_4$ and Co$_{0.8}$Ni$_{0.2}$Fe$_2$O$_4$ nanoparticle samples calcined at 800°C are depicted in Figure 12. The majority of particles in the TEM images show to be spherical in shape and have a polycrystalline nature. Because of the sample’s strong magnetic nature, by interfacial forces, a large number of small particles are held together and compose these bigger spherical (agglomerate) shape particles. The estimated average particle size is 82.51 and 52.31 nm along with a standard deviation of 41.81 nm. An overall increase is noticed in the value of the magnetic hysteresis loop obtained at room temperature for nanocrystalline Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) since a high proportion of Ni$^{2+}$ ions choose A-sites and contribute to hopping transport. Because the amount of charge carriers, i.e., electrons, has increased, the conductivity has increased as well.

4.6. Dielectric Properties. The LCR meter (Keysight E4980A/AL Precision LCR meter) device was used to study the dielectric behavior at different frequencies (20 Hz–2 MHz) regions. The complex permittivity of all the Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) nanoferrite particle samples calcined at 800°C was calculated. For that purpose, the dielectric constants’ ($\varepsilon'$) values (also known as the real permittivity) of all samples are measured using the following formula by using their known capacitance at room temperature (Figure 13(a)): $\varepsilon' = \frac{Cd}{A\varepsilon_0}$,

where $\varepsilon_0$ is the permittivity of the vacuum, $A$ represents the area of the sample, and $d$ is the thickness of the sample.

It exhibits dielectric dispersion at lower frequencies and achieves a stabilized relaxing value at higher frequencies, based on the Maxwell–Wagner model of interfacial polarization and exactly corresponding to Koop’s phenomenological theory of dispersion [47]. Conducting grains and insulating grain boundaries are the two layers of the dielectric structure according to this theory.

The average grain size of specimens is related to the dielectric constant of the same composition. As a result, at lower frequencies, the decrease is rapid. Because the electron exchange between ferric and ions ferrous is not followed by the alternating field at higher frequencies, the rate of decrease in dielectric constant with respect to frequency slows down [48]. An overall increase is noticed in the value of the dielectric constant with respect to the substitution of Ni$^{2+}$ in place of cobalt except for a slight decrease for compositions $(x = 0.2, 0.4)$. The maximum values are possessed by $(x = 1.0)$ compositions.

Also, for the calculation of dielectric loss (\(\varepsilon''\)) along with AC conductivity ($\sigma_{AC}$) of the samples, the dielectric constant (\(\varepsilon'\)) and the dissipation factor (\(\tan\delta\)) were used. The relations used for these measurements are as follows [49]:

\[
\varepsilon'' = \tan\delta \times \varepsilon',
\]

\[
\sigma_{AC} = \varepsilon_0 \varepsilon' \tan\delta,
\]

where $\omega$ represents the angular frequency. The variation of the dielectric constant with respect to frequency is similar to that of the dielectric loss as shown in Figure 13(b). At higher frequencies, the decrease in space charge polarization causes a decrease in dielectric loss [50].

The value of AC conductivity increases linearly with frequency for all the samples, which is the expected characteristic of ferrites Figure 13(c). It can be explained by Verwey’s hopping mechanism [51], which states that electrical conductivity in ferrites is mostly caused by electron hopping between ions of the same element in more than one valence state [52]. Conductive grains become increasingly active at higher frequencies of the applied field, also the conductivity increases due to an increase in the hopping frequency [53]. The AC conductivity increases when doping is increased from $(x = 0.0$ to $x = 1.0$) since a high proportion of Ni$^{2+}$ ions choose A-sites and contribute to hopping transport. Because the amount of charge carriers, i.e., electrons, has increased, the conductivity has increased as well.

4.7. Magnetic Properties Using Vibrating Sample Magnetometer (VSM). The magnetic properties were investigated at room temperature using a vibrating sample magnetometer (LBKFB model Meghnatis Daghig Havig Kavir Company) in the applied field range of $-15$ to $+15$ kOe. Figure 14 shows a typical magnetic hysteresis loop obtained at room temperature for nanocrystalline Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) nanoferrite particle samples calcined at

<table>
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</tr>
</thead>
<tbody>
<tr>
<td>$\gamma_1$ (cm$^{-1}$)</td>
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</tr>
<tr>
<td>$\gamma_2$ (cm$^{-1}$)</td>
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</tr>
<tr>
<td>$\gamma_{12}$ (cm$^{-1}$)</td>
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</tr>
<tr>
<td>$\theta_D$ (K)</td>
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</tr>
</tbody>
</table>

Table 4: FTIR absorption data for Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) calcined at 800°C.
Figure 10: FESEM micrograph of the (a) CoFe$_2$O$_4$, (b) Co$_{0.6}$Ni$_{0.4}$Fe$_2$O$_4$, and (c) NiFe$_2$O$_4$ nanoferrite particle samples calcinated at 800°C, the (Inset) corresponding histogram.
$800^\circ$C, and magnified magnetization curves are shown as an inset in the figure under low applied magnetic field. All of the samples had ferrimagnetic properties with S-shaped hysteresis loops.

The decrease in saturation magnetization with decreasing cobalt concentration can be explained by the fact that due to their relatively strong orbital contribution to magnetic moment Co$^{2+}$ ions are known to have considerable
induced anisotropy [54]. Increasing Ni\(^{2+}\) concentration causes a decrease in coercivity, and this may be due to the decrease in the anisotropy field, which in turn decreases the domain wall energy [55, 56].

In the present study, the saturation magnetization values are very high, and close to that of the bulk cobalt ferrite indicating that the sol-gel auto-combustion technique is a good synthesis method. An increase in particle size as the calcined temperature increases may also cause the saturation magnetization to increase. However, CoFe\(_2\)O\(_4\) has higher saturation magnetization than NiFe\(_2\)O\(_4\) due to the high ionic magnetic moment of Co than Ni. Similar high saturation magnetization values have been reported in the literature by several authors [57–60]. The magnetization increases with an increase in grain size as the surface-to-volume ratio decreases [61, 62]. The magnetic studies revealed that another reason for a high value of saturation magnetization is due to canting of spins occurrence as explained by the Yafet–Kittel model. The smaller canting angle of magnetic ions shows the increase in the overlap of the wave functions between the two nearest neighboring magnetic ions and also superexchange interactions between the magnetic ions and oxygen anions lead to a higher saturation magnetization [63].

For the spinel ferrites, the cations in tetrahedral (A-site) and octahedral (B-site) have opposite aligned magnetic

<table>
<thead>
<tr>
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<th>Theoretical</th>
<th>% atomic</th>
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<td>O</td>
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<td>27.28</td>
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<td>14.3</td>
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<td>Ni</td>
<td></td>
<td></td>
<td></td>
</tr>
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</tr>
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<td>15.02</td>
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<tr>
<td></td>
<td>Ni</td>
<td>28.5</td>
<td>25.05</td>
<td>14.3</td>
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</tbody>
</table>

Figure 12: TEM micrographs and (inset) particle distribution histogram of the (a) CoFe\(_2\)O\(_4\) and (b) Co\(_{0.6}\)Ni\(_{0.4}\)Fe\(_2\)O\(_4\) ferrite nanoparticle samples calcinated at 800°C.
moments, where the theoretical magnetic moment values \( n^\text{th}_B \) of \( \text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4 \) ferrite nanoparticles were calculated using the proposed cation distribution and the ionic magnetic moment [64].

\[
   n^\text{th}_B = M_B - M_A, \tag{12}
\]

where \( M_B \) and \( M_A \) are the Bohr magneton on the A-site and B-site, respectively.

The experimental value of magnetic moment (\( n^\epsilon_B \), in units of Bohr magneton) was computed using the following relation [65]:

\[
   n^\epsilon_B = \frac{M_{\text{spinel ferrite}} \times M_s}{5585}, \tag{13}
\]

where \( M_{\text{spinel ferrite}} \) is the molecular weight of the synthesized ferrite in g/mol, 5585 is the magnetic factor, and \( M_s \) is saturation magnetization in emu/g.

Because the experimental values of the magnetic moment are smaller than the theoretical magnetic moment which means that Neel’s two-sub-lattice collinear model is not suitable for the obtained samples and the magnetic order is not governed by the Neel-type magnetic order. This difference between the theoretical and experimental magnetic moment showed that we need to invoke Yafet–Kittel three sublattice model. This suggests that the magnetic order in all the nickel-substituted cobalt ferrite samples shows a Y-K type of magnetic ordering.

According to the Neel model when the canting angle is zero means that the sample shows a Neel-type of magnetic ordering indicating that magnetization can be explained on the basis of the Neels two sublattice theory. While, according to the Y-K model, the B lattice can be divided into two sublattices, \( B_1 \) and \( B_2 \), each having magnetic moments equal in magnitude and each oppositely canted at the same angle.

Figure 13: Dielectric parameters of all the prepared \( \text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4 \) (where \( x = 0.0 \leq x \leq 1.0 \)) nanoparticle samples calcinated at 800°C.
\[ a_{Y-K} = \cos^{-1}\left(\frac{n_B^r + M_A}{M_B}\right). \tag{14} \]

Moreover, there is a significant triangular (or canted) noncollinear type spin arrangement in \( B \)-site, since for all the samples the calculated values of \( Y-K \) angles are nonzero, which strengthens the \( B-B \) interaction and in turn, reduces the \( A-B \) interaction. The system can be explained according to the Yafet and Kittel three sublattice model. The increase in spin canting angles for the samples with an increase in \( \text{Ni}^{2+} \) content suggests the increased favor for triangular spin arrangements on \( B \) sites resulting in the decrease in the \( A-B \) exchange interaction and thus enhancing the \( B-B \) interaction.

From experimental results, it is observed that the values of \( M_r \) and hence \( n_B \) goes on decreasing as the concentration of \( \text{Ni}^{2+} \) is increased. The decrease in \( M_r \) and hence \( n_B \) for all the samples is caused by nonzero \( Y-K \) angles.

In the present study, the calculated squareness ratio for compositions \( x = 0.0, 0.2, 0.4, 0.6 \) (high \( H_J \)) has been found to be nearly 0.5, indicating that these samples have a virtually single domain structure and correspond to uniaxial anisotropy, while the other compositions \( x = 0.8, \) and 1.0 (lower \( H_J \)) possess a ratio of 0.43 and 0.32 respectively, attributed to that these ferrites have a multidomain structure, also according to the Stoner–Wohlfarth model squareness ratio of an assembly of noninteracting 3D random particles is 0.5, it can be concluded that all the samples in this study have uniaxial anisotropy because the squareness values are nearly equal or lower than 0.5 [68].

The calculated values obtained from Figure 14 of saturation magnetization \( M_s \), coercivity \( H_c \), remanent magnetization \( M_r \), squareness ratio \( M_s/M_r \), experimental \( n_B^r \) and theoretical \( n_B^{th} \) magnetic moment, and \( Y-K \) angles \( (a_{Y-K}) \) with respect of \( \text{Ni}^{2+} \) concentration for all the samples calcined at 800°C are listed in Table 6. For \( \text{(NiFe}_2\text{O}_4) \) composition the values of all the magnetic parameters are much lower than other compositions. A noticeable decrease in \( M_r \) from 42.89 emu/g to 15.49 emu/g is found with an increase in \( x \) from (0.0 to 1.0).

4.8. Vector Network Analyzer (VNA). The blocking of incident electromagnetic radiation is known as electromagnetic interference (EMI). Possible EMI shielding based on \( \text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/PVA \) nanocomposites has been proposed and illustrated in Figure 15. When EM radiation incident on the shielding surface because of multiple/internal reflections through the \( \text{Ni}^{2+} \) substituted cobalt ferrite/PVA interfaces that exist in the shielding nanocomposite film, from the outer side of the nanocomposite film some EM radiation is reflected, a small portion is transmitted from the film, and within the nanocomposite, the remaining EM radiation is absorbed. Many different ways can be used to measure EMI shielding such as impedance parameters [69], electrical and magnetic parameters [70], and VNA [71, 72].

A novel and promising application of the electromagnetic \( \text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/PVA \) nanocomposites is their ability to shield electromagnetic radiation. The mechanisms of energy loss in magnetic materials are due to dielectric and magnetic properties, which depend on the imaginary part of the complex permittivity and complex permeability. The EM shielding performances of the \( \text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/PVA \) nanocomposites films were analyzed by vector network analyzer an Agilent N5230 A (Keysight Technologies, Inc. USA) in the X-band (8.2–12.4 GHz), by calculating the scattering parameters (\( S \)-parameters) for a rectangular strip with dimensions (10.16 * 22.86 mm) of nanocomposite film samples.

EMI shielding was calculated from scattering \( S \)-parameters and then complex values of permittivity and permeability were calculated using the Nicholson–Ross method to investigate the underlying shielding mechanisms of \( \text{Co}_{1-x}\text{Ni}_x\text{Fe}_2\text{O}_4/PVA \) nanocomposites films and understand the shielding mechanism. The method used to measure reflection loss (RL) using complex electromagnetic parameters obtained experimentally is a short-circuit approach, in which RL values are simulated [73]. The electromagnetic parameters, that is, complex dielectric permittivity \( (\varepsilon_r = \varepsilon_r' - j\varepsilon_r'') \) and magnetic permeability

![Figure 14: Magnetic hysteresis loops for Co$_{1-x}$Ni$_x$Fe$_2$O$_4$ (where $x = 0.0 \leq x \leq 1.0$) nanoferite particle samples calcinated at 800°C; inset showing magnified hysteresis loops under low magnetic field.](image-url)
When EM waves pass from a material, SE is expressed as the logarithmic ratio of incident $P_i$ and transmitted $P_r$ power of EM wave, mathematically expressed as in the following [77]:

$$SE_{\text{total}}(\text{dB}) = 10 \log_{10} \left( \frac{P_i}{P_r} \right).$$  \hspace{1cm} (15)

The performance of shielding material to attenuate EMI can be expressed by the EMI shielding effectiveness and calculated in decibels (dB) units.

As the incident EM wave falls on the material, three basic interaction phenomena occur; transmittance ($T$), reflectance ($R$), and absorbance ($A$) [78]. The transmission coefficient and reflection coefficient are related to $S$-parameters as follows:

$$T = |S_{21}|^2 = |S_{11}|^2,$$

$$R = |S_{22}|^2 = |S_{12}|^2.$$  \hspace{1cm} (16)

The scattering parameters represent $S_{11}$ as the coefficient of forwarding reflection and $S_{12}$ as the coefficient of reverse transmittance. The power of transmission and reflection characterizes the scattering parameter. In turn, from the above equations, the absorption coefficient is measured which is mathematically expressed as follows:

$$A = 1 - R - T.$$  \hspace{1cm} (17)

The samples were bombarded by electromagnetic waves from both sides and scattering parameters were obtained from VNA and shielding effectiveness through absorbance ($SE_A$), shielding effectiveness through reflection ($SE_R$), and total shielding effectiveness ($SE_T$) were calculated using the following equations:

$$SE_A = -10 \times \log_{10} \left( \frac{T}{1 - R} \right),$$

$$SE_R (\text{dB}) = -10 \times \log_{10} \left( \frac{R}{1 - R} \right),$$

$$SE_T = SE_A + SE_R.$$  \hspace{1cm} (18)

According to Schelkunoff’s hypothesis, a material’s total shielding effectiveness ($SE_T$) is determined by its electrical conductivity and is proportional to its reflection, absorption, and multiple reflections.

When the absorption loss is $\geq 10 \text{dB}$, multiple/internal reflection is neglected for total SE. Because of absorption, while moving from one boundary to another the magnitude of EM waves is neglected at high frequencies. Materials with a high ability of absorption and high thickness can neglect safely multiple/internal reflections. As a result, only absorption loss ($SE_A$) and reflection loss ($SE_R$) contribute to $SE_T$ [79]. The total EMI SE (measured directly from S-parameters) can be obtained as follows:

$$\text{EMI SE} = 10 \times \log_{10} \left( \frac{1}{|S_{12}|^2} \right) = 10 \times \log_{10} \left( \frac{1}{|S_{21}|^2} \right).$$  \hspace{1cm} (19)

Figure 17 shows the EMI shielding effectiveness values of prepared nanocomposite PF1, PF2, and PF3 samples in the frequency range X-band (8.2–12.4) GHz. The obtained result of EMI SE of PF2 is 27 dB. This result can be enhanced for the composites with higher magnetic permeability, electrical conductivity, and larger thicknesses [56]. Pubby et al. [28]
also reported a significant enhancement and increase in absorption shielding effectiveness parameters with the addition of cobalt in nickel ferrites, so they concluded that the mixed nickel-cobalt ferrite has potential for usage in microwave shielding applications.

The oscillatory behavior of absorption in the Ni$^{2+}$ substituted cobalt nanoferrites samples are due to the hopping of electrons, diverse relaxation frequencies of various dipoles formed in the ferrite structure, and the relaxation due to interfacial polarization.

The Ni$^{2+}$ substituted cobalt nanoferrites act as an absorbing material and show improved EMI shielding values due to their dielectric and magnetic losses in the microwave frequency band. The dielectric properties perfectly matched the magnetic properties. The magnetic loss of these magnetic materials results from their spin relaxation in the high-frequency alternating electromagnetic fields, ferrimagnetism’s, and the resonance absorption of moving magnetic domains wall.

The advantage of using polymer nanocomposite is that a large number of nanoferrite particles can accommodate in small film thickness because of the small filler particle size, resulting in more EMR attenuation due to a large number of interfaces. All communication devices operate in the microwave range and emit electromagnetic radiation into the environment which can be shielded using the EMI shielding
The theory of EMI is based on essential mechanisms such as reflection loss and absorption loss. The transmission line theory method is used to calculate the frequency dependence of reflection loss (RL) at a thickness \((d)\) based on data of complex dielectric permittivity and complex magnetic permeability, which characterizes the electromagnetic wave absorption properties [80].

\[
Z_{\text{in}} = \sqrt{\frac{\mu_r}{\varepsilon_r}} \tan \left( \frac{i2\pi f d \sqrt{\mu_r \varepsilon_r}}{c} \right),
\]

\[
R_L = 20 \log \left( \frac{Z_{\text{in}} - 1}{Z_{\text{in}} + 1} \right) \text{ (dB)},
\]
where \( Z_{in} \) is the absorber’s input impedance, \( f \) represents the frequency, \( d \) is the absorber thickness, and \( c \) is the velocity of light. Figure 18 characterizes the measured absorption spectra of PF1, PF2, and PF3. The maximum RL for the prepared nanocomposites in the X-band frequency range is \((-32.08\,\text{dB})\) for the PF2 sample.

To rationalize the discussion, we compared our results with those reported in the literature, and the data are comprehensively presented in Table 7.

To rationalize the discussion, we compared our results with those reported in the literature, and the data are comprehensively presented in Table 7.

### 4. Conclusions

Nanostructured Ni\(^{2+}\) substituted cobalt spinel ferrite with a composition of Co\(_{1-x}\)Ni\(_x\)Fe\(_2\)O\(_4\) (where \( x = 0.0 \leq x \leq 1.0 \)) and the corresponding composite film was synthesized successfully. The obtained result demonstrates that variation in calcination temperature could turn effectively the structural, dielectric, and magnetic properties of Co\(_{1-x}\)Ni\(_x\)Fe\(_2\)O\(_4\) (where \( x = 0.0 \leq x \leq 1.0 \)) nanoferrite particles. From XRD, the crystallite size of the synthesized samples was calculated and is found to be in the nano-range between 49 and 65 nm. Due to the substitution of Ni\(^{2+}\) ions, a decrease in the lattice parameters values and a shift of the main peak (311) towards a higher angle occurs. With increasing Ni\(^{2+}\) content a decrease in the grain size is observed. The nanocomposite films were successfully prepared through the solution casting technique. Shielding effectiveness of \(-32.08\,\text{dB}\) was observed in a broad X-band of frequency regions (8.2–12.4) GHz.

### Data Availability

All data used to support the findings of this study are included within the article.

### Conflicts of Interest

The authors declare that they have no conflicts of interest.

### References


### Table 7: Comparison of EMI shielding properties obtained in the present study with that of the various ferrites reported in the literature.

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<th>Year</th>
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<td>Sol-gel citrate</td>
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<td>8.2–12.4</td>
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<td>[83]</td>
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<tr>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.2})Fe(</em>{2-2x})O(_4)</td>
<td>Sol-gel citrate</td>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.2})Fe(</em>{2-2x})O(_4)/PVA</td>
<td>−44.92</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.25})Fe(</em>{2-2x})O(_4)</td>
<td>Sol-gel citrate</td>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.25})Fe(</em>{2-2x})O(_4)/PVA</td>
<td>−40.43</td>
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<tr>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.3})Fe(</em>{2-2x})O(_4)</td>
<td>Sol-gel citrate</td>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.3})Fe(</em>{2-2x})O(_4)/PVA</td>
<td>−32.08</td>
<td>8.2–12.4</td>
<td>2020</td>
<td>[84]</td>
</tr>
<tr>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.35})Fe(</em>{2-2x})O(_4)</td>
<td>Sol-gel citrate</td>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.35})Fe(</em>{2-2x})O(_4)/PVA</td>
<td>−36.1</td>
<td>0.1–20</td>
<td>2020</td>
<td>[85]</td>
</tr>
<tr>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.4})Fe(</em>{2-2x})O(_4)</td>
<td>Sol-gel citrate</td>
<td>Ni(<em>{0.6})Zn(</em>{0.4})(CoZr)(<em>{0.4})Fe(</em>{2-2x})O(_4)/PVA</td>
<td>−39.2</td>
<td>0.1–20</td>
<td>2020</td>
<td>[86]</td>
</tr>
</tbody>
</table>

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[76] F. Movassagh-Alanagh, A. Bordbar-Khiabani, and A. Ahangari-Asl, "Three-phase PANI@ nano-Fe3O4@ CFs heterostructure: fabrication, characterization and investigation of microwave absorption and EMI shielding of PANI@ nano-Fe3O4@ CFs/epoxy hybrid composite," *Composites Science and Technology*, vol. 150, pp. 65–78, 2017.


