

## SHORT COMMUNICATION

### Variations of Thin Metallic Zinc Film Resistances with Sputtering Rate

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Variations of resistance of evaporated<sup>1</sup> or sputtered<sup>2</sup> films during deposition have been studied for deposition times lower than five minutes and for resistances higher than  $100 \Omega \text{ sq}^{-1}$ . As our main objective is to study the conduction mechanisms<sup>3</sup> in sputtered films over the thickness range from 200 to 1500 Å (i.e. in the  $10$  to  $100 \Omega \text{ sq}^{-1}$  sheet resistance range), we report in this note our investigations about the electrical resistance  $R(T)$  of zinc films for deposition time,  $T$  greater than two minutes and for four average deposition rates.

Preparation of films has been described in a previous paper;<sup>4</sup> they are deposited by d.c. diode sputtering of a zinc target (99.9% purity) in an atmosphere of  $U$  grade argon. As broken sputtering is equivalent to continuous deposition<sup>4</sup>, sputtering was stopped every minute to measure the resistance  $R$  with a multimeter. The average sputtering rate was determined from the deposition time and the film thickness measured by an optical method;<sup>5</sup> this method is adequate for we have observed very slight anisotropic effects,<sup>6</sup> which seem more important for higher sputtering rates;<sup>7</sup> it varied with the intensity  $I_e$  of the glow discharge current, the voltage  $U_e$  remaining constant (curves 1, 2, 3 on Figure 2). Variations of the deposition rate as a function of the intensity  $I_e$  of glow discharge current for a voltage equal to 1500 V are shown in Figure 1. This curve is in good agreement with the results of Laville Saint-Martin<sup>8</sup> who established that sputtering rate  $\nu$  is given by

$$\nu = I_e U_e \exp \{-AU_e^{-1} - BU_e\} \quad (1)$$

where  $A, B$  are constants.

For high voltage values in the range 1000 to 1750 V, eq. (1) may be expressed as

$$\nu \propto I_e \quad (2)$$

Variations of zinc film resistance  $R(T)$  versus

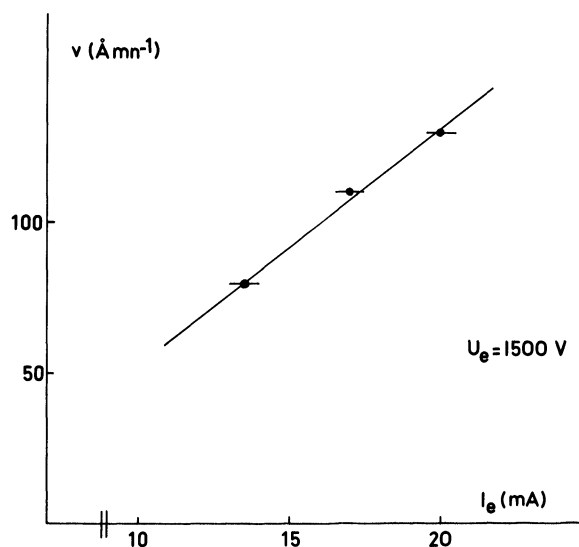


FIGURE 1 Variations of the average sputtering rate  $\nu$  with the intensity  $I_e$  of the glow discharge.

deposition time  $T$  are plotted in Figure 2. Attempts have been made<sup>1,2</sup> to fit these experimental variations to an empirical equation in the form

$$R(T) = R_{\infty} \exp \frac{1}{K_1 + K_2 T} \quad (3)$$

where  $R_{\infty}$  is the limiting value of  $R(T)$  when  $T$  is large and  $K_1$  and  $K_2$  are constants.

As we have established that films thicker than 5000 Å exhibit bulk properties,<sup>4</sup> their resistance  $R_b \approx 2.5 \Omega \text{ sq}^{-1}$  is assumed equal to  $R_{\infty}$ . Substituting for  $R_{\infty}$  in Eq. 3, this yields

$$R(T) = R_b \exp \frac{1}{K_1 + K_2 T} \quad (4)$$

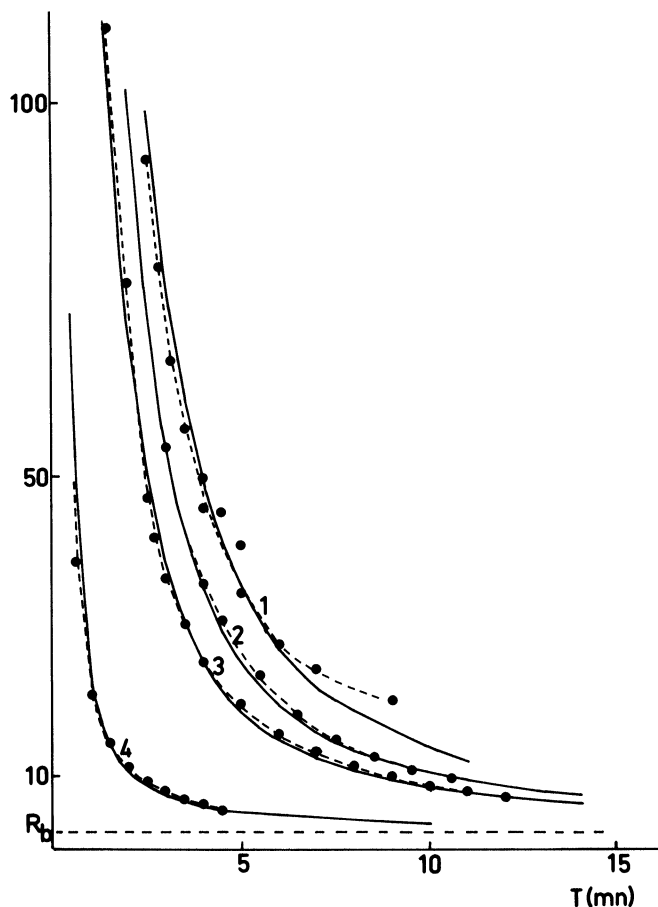


FIGURE 2 Experimental (dotted lines) and theoretical (full lines) variations of zinc film resistance  $R$ , with sputtering rate  $\nu$ , equal to: (1)  $80 \text{ \AA mn}^{-1}$ , (2)  $110 \text{ \AA mn}^{-1}$ , (3)  $130 \text{ \AA mn}^{-1}$ , (4)  $550 \text{ \AA mn}^{-1}$ .

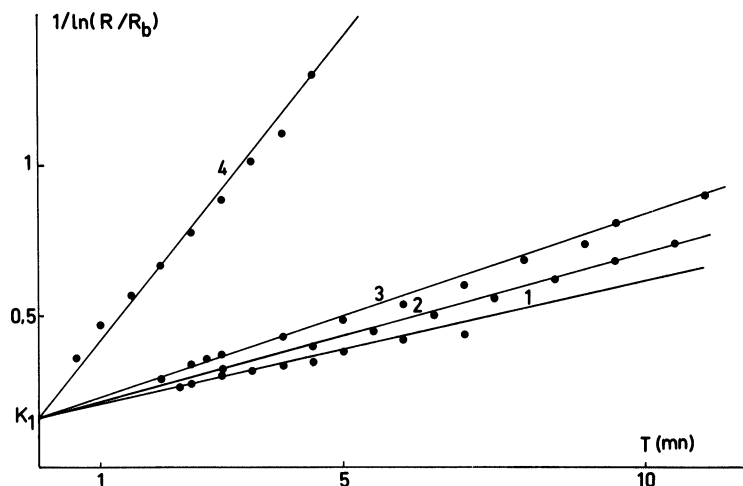


FIGURE 3  $1/\text{Ln}[R(T)/R_b]$  versus  $T$  (deposition time) with sputtering rate,  $\nu$ , equal to (1)  $80 \text{ \AA mn}^{-1}$ , (2)  $110 \text{ \AA mn}^{-1}$ , (3)  $130 \text{ \AA mn}^{-1}$ , (4)  $550 \text{ \AA mn}^{-1}$ .

$K_1$  and  $K_2$  are determined by plotting  $1/\text{Ln}[R(T)/R_b]$  versus deposition time  $T$  (Figure 3). The slopes of the best fit straight lines determine the values of  $K_2$ , while  $K_1$  is calculated from the intercept with the vertical axis.

From Figure 3 it can be seen that  $K_1$  is independent of the deposition rate as indicated by Eq. 4. Thus,

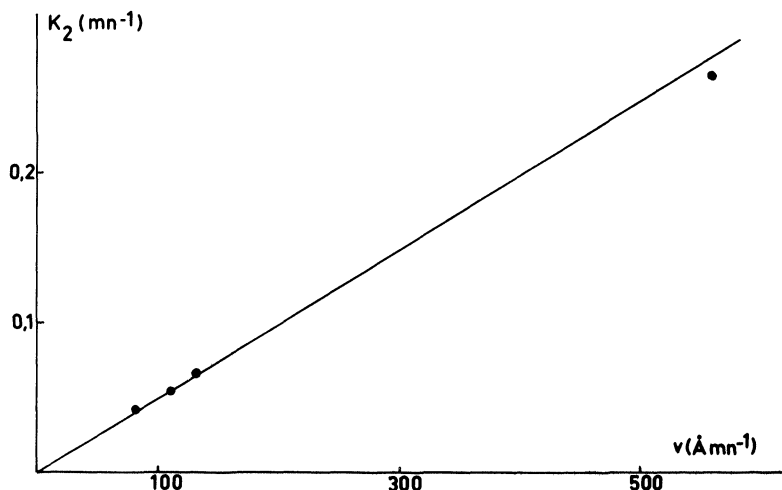
$$K_1 \approx 1/\text{Ln}[R_o/R] \quad (5)$$

where  $R_o$  is the substrate resistance.

However the value of  $R_o$  calculated from the experimental value of  $K_1$  is low (about  $6 \text{ k}\Omega$ ). As quasi-linear growth occurs only above the first critical thickness<sup>9</sup> we assume that this value corresponds to a smaller thickness for which the granular structure consists of a large number of empty channels distributed throughout the film. It has been shown that the value of the resistance of such a film depends essentially on the geometrical arrangement of the sputtering chamber which remained unchanged in our experiments.

Calculations allow one to determine suitable values of  $K_2$ , leading to a good agreement between the experimental and theoretical curves (Figure 2) in the resistance range  $10$  to  $100 \text{ }\Omega \text{ sq}^{-1}$ . A discrepancy of less than 10% is observed except for the lower sputtering rate but experimental accuracy is low in this case (15%).

The observed slight departures from the theoretical resistance at low  $R$  values occur at low sputtering rates,<sup>3</sup> whereas considerable departures have previously been observed by Laville Saint-Martin<sup>7</sup> for higher sputtering rate.

FIGURE 4  $K_2$  versus deposition rate,  $\nu$ .

We observe (Figure 4) that the magnitude of constant  $K_2$  differs markedly for different deposition rates. Constant  $K_2$  has been defined as a velocity constant related to the sputtering rate  $\nu$  and may be given approximately by

$$K_2 \approx \alpha \nu \quad \text{with} \quad \alpha = 5.10^{-4} \text{ Å}^{-1}$$

It may be concluded that in these experimental conditions the measured resistance fits the theoretical formulae obtained with the deposition rate as a parameter in the range  $80 \text{ Å mn}^{-1}$  to  $600 \text{ Å mn}^{-1}$ ; a simple way to predetermine thin film resistances is thus available.

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