

Research Article

Higher-Order Uniformly Convergent Numerical Scheme for Singularly Perturbed Differential Difference Equations with Mixed Small Shifts

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This paper deals with numerical treatment of singularly perturbed differential difference equations involving mixed small shifts on the reaction terms. The highest-order derivative term in the equation is multiplied by a small perturbation parameter ε taking arbitrary values in the interval $(0, 1]$. For small values of ε , the solution of the problem exhibits exponential boundary layer on the left or right side of the domain and the derivatives of the solution behave boundlessly large. The terms having the shifts are treated using Taylor's series approximation. The resulting singularly perturbed boundary value problem is solved using exponentially fitted operator FDM. Uniform stability of the scheme is investigated and analysed using comparison principle and solution bound. The formulated scheme converges uniformly with linear order before Richardson extrapolation and quadratic order after Richardson extrapolation. The theoretical analysis of the scheme is validated using numerical test examples for different values of ε and mesh number N .

1. Introduction

Differential equations play a prominent role in many disciplines including engineering, physics, economics, and biology. Currently different authors are working on analytical and numerical solutions of differential equations using different techniques [1, 2]. Differential difference equations (DDEs) are differential equations where the evolution of the system not only depends on the present state of the system but also depends on the past history. Singularly perturbed differential difference equations are differential equations in which the highest-order derivative term is multiplied by a small perturbation parameter ε and involves at least one term with delay. In general, when the perturbation parameter tends to zero, the smoothness of the solution of the singularly perturbed differential difference equations (SPDDEs) deteriorates and it forms boundary layer [3]. Such type of equations has applications in the study

of variational problems in control theory [4] and in modelling of neuronal variability [5].

The presence of singular perturbation parameter ε in the equation leads to oscillation in the computed solution, while using standard numerical methods like FDM, FEM, and spline method [6]. To avoid this oscillation, an unacceptably large number of mesh points are required when ε is very small. This is not practical and leads to round-off error. So, to overcome the drawbacks associated with standard numerical methods, different authors have developed schemes that converge uniformly.

Numerical treatments of a class of SPDDEs have received a great deal of attention recently because of their wide applications. It is of theoretical and practical interest to consider numerical methods for such problems. Owing to this, here we present some prior studies on numerical solution of the considered problem. Lange and Miura in [7–10] studied a class of second-order DDEs in which the second

derivative term is multiplied by a small parameter. The authors extend the method of matched asymptotic expansions initially developed for solving boundary value problems to obtain approximate solution for SPDDEs. In a series of papers [11–14], Kadalbajoo and Sharma developed uniformly convergent numerical methods using fitted mesh FDMs techniques. Swamy et al. [15–17] considered the problem and developed a numerical scheme using fitted operator finite difference techniques. Melesse et al. [18] applied initial value technique for treating the considered SPDDEs. Ranjan and Prasad [19] used modified fitted operator FDM for solving the problem. Sirisha et al. [20] developed fitted operator finite difference scheme using the procedure of domain decomposition. A number of authors have developed numerical scheme using exponentially fitted method for solving SPDDEs. To the authors' knowledge, none of them show the uniform convergence of their

schemes. This motivates to treat the considered SPDDEs and formulate the uniform convergence analysis of the scheme. Our contribution in this paper is to develop higher-order uniformly convergent numerical scheme using exponentially fitted FDM and to analyse the uniform convergence of the proposed scheme.

Notation 1. The symbol C is used to denote positive constant independent of ε and N . The norm $\|\cdot\|$ denotes the maximum norm.

2. Statement of the Problem

A class of second-order singularly perturbed differential difference equations having mixed shift on reaction terms is given by

$$-\varepsilon u''(x) + a(x)u'(x) + \alpha(x)u(x - \delta) + \omega(x)u(x) + \beta(x)u(x + \eta) = f(x), \quad x \in \Omega = (0, 1), \quad (1)$$

with the interval conditions

$$\begin{aligned} u(x) &= \phi(x), & x \in [-\delta, 0], \\ u(x) &= \psi(x), & x \in [1, 1 + \eta], \end{aligned} \quad (2)$$

where $\varepsilon \in (0, 1]$ is singular perturbation parameter and δ and η are delay and advance parameters satisfying $\delta, \eta < \varepsilon$. The functions $a(x), \alpha(x), \omega(x), \beta(x), f(x), \phi(x)$, and $\psi(x)$ are assumed to be sufficiently smooth, bounded, and independent of ε for guaranteeing the existence of unique solution. The coefficient functions $\alpha(x), \omega(x)$, and $\beta(x)$ are assumed to satisfy

$$\alpha(x) + \omega(x) + \beta(x) \geq q^* > 0, \quad \forall x \in \overline{\Omega}, \quad (3)$$

for some constant q^* .

For the case where $\delta, \eta < \varepsilon$, using Taylor's series approximation for the terms with shift is appropriate [21]. So, we approximate $u(x - \delta)$ and $u(x + \eta)$ as

$$u(x - \delta) \approx u(x) - \delta u'(x) + \frac{\delta^2}{2} u''(x) + O(\delta^3), \quad (4)$$

$$u(x + \eta) \approx u(x) + \eta u'(x) + \frac{\eta^2}{2} u''(x) + O(\eta^3).$$

Using the approximation in (4) into (1)-(2) gives

$$-c_\varepsilon u''(x) + p(x)u'(x) + q(x)u(x) = f(x), \quad x \in (0, 1), \quad (5)$$

with the boundary conditions

$$\begin{aligned} u(0) &= \phi(0), \\ u(1) &= \psi(1), \end{aligned} \quad (6)$$

where $c_\varepsilon = \varepsilon^2 - (\delta^2/2)\alpha - (\eta^2/2)\beta$, $p(x) = a(x) - \delta\alpha(x) + \eta\beta(x)$, and $q(x) = \alpha(x) + \beta(x) + \omega(x)$ for α and β are lower

bounds of $\alpha(x)$ and $\beta(x)$, respectively. For small values of δ, η , equations (1)-(2) and (5)-(6) are asymptotically equivalent, since the truncated term in (5)-(6) is order of $O(\delta^3, \eta^3)$.

Let us denote differential operator L for the differential equation in (5) as

$$Lu(x) = -c_\varepsilon u''(x) + p(x)u'(x) + q(x)u(x). \quad (7)$$

The solution of the problem in (5)-(6) exhibits regular boundary layer of thickness $O(c_\varepsilon)$ and the position of the boundary layer depends on the sign of $p(x)$. If $p(x) < 0$, left boundary layer exists and, for $p(x) > 0$, right boundary layer exists. If $p(x), x \in [0, 1]$ changes sign interior layer will occur [22].

Setting $c_\varepsilon = 0$ in equations (5)-(6) gives the reduced problem. For the case where $p(x) > 0$, it is given by

$$\begin{cases} p(x)u_0'(x) + q(x)u_0(x) = f(x), & \forall x \in \Omega, \\ u_0(0) = \phi(0), \end{cases} \quad (8)$$

and, for the case where $p(x) < 0$, it is given by

$$\begin{cases} p(x)u_0'(x) + q(x)u_0(x) = f(x), & \forall x \in \Omega, \\ u_0(1) = \psi(1). \end{cases} \quad (9)$$

It is a first-order initial value problem; for small values of c_ε , the solution of (5)-(6) is very close to the solution of (8) or (9).

2.1. Properties of the Analytical Solution

Lemma 1 (The maximum principle). *Let z be a sufficiently smooth function defined on Ω , which satisfies $z(0) \geq 0$ and $z(1) \geq 0$. Then $Lz(x) > 0, \forall x \in \Omega$, implies that $z(x) \geq 0, \forall x \in \overline{\Omega}$.*

Proof. Let $x^* \in \overline{\Omega}$ be such that $z(x^*) = \min_{(x) \in \overline{\Omega}} z(x)$ and suppose that $z(x^*) < 0$. It is clear that $x^* \notin \{0, 1\}$. Since $z(x^*) = \min_{(x) \in \overline{\Omega}} z(x)$ form extrema values in calculus we have $z'(x^*) = 0$ and $z''(x^*) \geq 0$ giving that $Lz(x^*) < 0$, which is contradiction to the assumption made above: $Lz(x^*) > 0, \forall x \in \Omega$. Therefore, $z(x) \geq 0, \forall x \in \Omega$. \square

Lemma 2 (Stability estimate). *Let u be the solution of (5)-(6); then, it satisfies the bound*

$$|u(x)| \leq \frac{\|f\|}{q_*} + \max\{|\phi(0)|, |\psi(1)|\}, \quad (10)$$

where q^* is lower bound of $q(x)$.

$$\begin{aligned} L\vartheta_{\pm}(x) &= -c_{\varepsilon}\vartheta_{\pm}''(x) + p(x)\vartheta_{\pm}'(x) + q(x)\vartheta_{\pm}(x) \\ &= -c_{\varepsilon}(0 \pm u''(x)) + p(x)(0 \pm u'(x)) + q(x)\left(\frac{\|f\|}{q_*} + \max\{\phi(0), \psi(1)\} \pm u(x)\right) \\ &= q(x)\left(\frac{\|f\|}{q_*} + \max\{\phi(0), \psi(1)\}\right) \pm f(x) \\ &\geq 0, \quad \text{since } q(x) \geq q^* > 0, \end{aligned} \quad (12)$$

which implies that $L\vartheta_{\pm}(x) \geq 0$. Hence, using the maximum principle in Lemma 1, we obtain $\vartheta_{\pm}(x) \geq 0, \forall x \in \overline{\Omega}$. \square

Lemma 3. *The derivatives of the solutions of (5)-(6) satisfy the bound*

$$|u^{(k)}(x)| \leq C \left(1 + c_{\varepsilon}^{-k} \exp\left(-\frac{p^* x}{c_{\varepsilon}}\right)\right), \quad x \in \overline{\Omega}, 0 \leq k \leq 4, \quad (13)$$

for left boundary layer problems and

$$|u^{(k)}(x)| \leq C \left(1 + c_{\varepsilon}^{-k} \exp\left(-\frac{p^*(1-x)}{c_{\varepsilon}}\right)\right), \quad x \in \overline{\Omega}, 0 \leq k \leq 4, \quad (14)$$

Proof. Let us define barrier functions $\vartheta_{\pm}(x)$ as $\vartheta_{\pm}(x) = (\|f\|/q^*) + \max\{\phi(0), \psi(1)\} \pm u(x)$ and apply the maximum principle to obtain the required bound.

On the boundary points, we have

$$\vartheta_{\pm}(0) = \frac{\|f\|}{q_*} + \max\{\phi(0), \psi(1)\} \pm u(0) \geq 0, \quad (11)$$

$$\vartheta_{\pm}(1) = \frac{\|f\|}{q_*} + \max\{\phi(0), \psi(1)\} \pm u(1) \geq 0.$$

On the differential operator,

for right boundary layer problems.

Proof. See [23]. \square

3. Numerical Scheme

First, let us discretize the domain $\overline{\Omega} = [0, 1]$ into N equal number of subintervals with mesh length $h = (1/N)$ as $\overline{\Omega}^N = \{x_i = ih\}_{i=0}^N$. Let $u(x)$ be smooth function on the domain $\overline{\Omega} = [0, 1]$; then, using Taylor series approximation, we have

$$\begin{aligned} u(x_{i+1}) &= u_{i+1} = u_i + hu'_i + \frac{h^2}{2!}u''_i + \frac{h^3}{3!}u^{(3)}_i + \frac{h^4}{4!}u^{(4)}_i + \frac{h^5}{5!}u^{(5)}_i + \frac{h^6}{6!}u^{(6)}_i + O(h^7), \\ u(x_{i-1}) &= u_{i-1} = u_i - hu'_i + \frac{h^2}{2!}u''_i - \frac{h^3}{3!}u^{(3)}_i + \frac{h^4}{4!}u^{(4)}_i - \frac{h^5}{5!}u^{(5)}_i + \frac{h^6}{6!}u^{(6)}_i + O(h^7). \end{aligned} \quad (15)$$

Taking the difference in (15), we obtain

$$u_{i-1} - 2u_i + u_{i+1} = \frac{2h^2}{2!}u''_i + \frac{2h^4}{4!}u^{(4)}_i + \frac{2h^6}{6!}u^{(6)}_i + O(h^8). \quad (16)$$

Differentiating (16) two times gives

$$u''_{i-1} - 2u''_i + u''_{i+1} = \frac{2h^2}{2!}u^{(4)}_i + \frac{2h^4}{4!}u^{(6)}_i + \frac{2h^6}{6!}u^{(8)}_i + O(h^8). \quad (17)$$

Now, multiplying (17) by $-(h^2/12)$ and adding with (16) to eliminate the term with $u^{(4)}_i$ gives

$$u_{i-1} - 2u_i + u_{i+1} = \frac{h^2}{12} [u_{i-1}'' + 10u_i'' + u_{i+1}''] + \tau, \quad (18)$$

where $\tau = O(h^6)$. Evaluating (5) at x_{i-1} , x_i , and x_{i+1} , respectively, we obtain

$$\begin{aligned} -c_\epsilon u_{i-1}'' &= -p(x_{i-1})u_{i-1}' - q(x_{i-1})u_{i-1} + f(x_{i-1}), \\ -c_\epsilon u_i'' &= -p(x_i)u_i' - q(x_i)u_i + f(x_i), \\ -c_\epsilon u_{i+1}'' &= -p(x_{i+1})u_{i+1}' - q(x_{i+1})u_{i+1} + f(x_{i+1}). \end{aligned} \quad (19)$$

Next, approximate the first derivative terms u_{i-1}' , u_i' , and u_{i+1}' in (19), using the right shifted, central, and left shifted finite difference approximations as

$$\begin{aligned} u_{i-1}' &= \frac{u_{i-1} - 4u_i + 3u_{i+1}}{2h} + hu_i'' + O(h^2), \\ u_i' &= \frac{u_{i+1} - u_{i-1}}{2h} + O(h^2), \\ u_{i+1}' &= \frac{-3u_{i-1} + 4u_i - u_{i+1}}{2h} - hu_i'' + O(h^2). \end{aligned} \quad (20)$$

Substituting (20) into (19) and then (19) into (18) gives

$$\begin{aligned} & -\left(c_\epsilon - \frac{hp_{i-1}}{12} + \frac{hp_{i+1}}{12}\right) \left(\frac{U_{i-1} - 2U_i + U_{i+1}}{h^2}\right) + \frac{p_{i-1}}{24h} (-3U_{i-1} + 4U_i - U_{i+1}) \\ & + \frac{10p_i}{24h} (U_{i+1} - U_{i-1}) + \frac{p_{i+1}}{24h} (U_{i-1} - 4U_i + 3U_{i+1}) + \frac{q_{i-1}}{12} U_{i-1} + \frac{10q_i}{12} U_i + \frac{q_{i+1}}{12} U_{i+1} \\ & = \frac{1}{12} [f_{i-1} + 10f_i + f_{i+1}] + O(h^2), \quad i = 1, 2, \dots, N-1, \end{aligned} \quad (21)$$

where p_{i-1} , p_i , and p_{i+1} are denoted for $p(x_{i-1})$, $p(x_i)$, and $p(x_{i+1})$, respectively. We denote U_i for the approximate solution of $u(x_i)$ in the above discretization.

To get small truncation error in boundary layer region, we apply exponentially fitted operator finite difference method (FOFDM). For developing the FOFDM, we use the theory developed in asymptotic method for treating

singularly perturbed BVPs. Let us consider and treat the left and the right boundary layer cases separately.

3.1. Left Boundary Layer Problems. In this case, the boundary layer occurs near $x = 0$. From the theory of singular perturbation given in [24], the zeroth-order asymptotic solution of (5)-(6) is given by

$$u(x) = u_0(x) + \frac{p(0)}{p(x)} (\phi(0) - u_0(0)) \exp\left(-\int_0^x \left(\frac{p(x)}{c_\epsilon} - \frac{q(x)}{p(x)}\right) dx\right) + O(c_\epsilon), \quad (22)$$

where u_0 is the solution of the reduced problem. Using Taylor's series approximation for $u_0(x)$, $p(x)$, and $q(x)$ centred at $x_i = ih$ up to first order and considering $c_\epsilon \rightarrow 0$, the discretized form of (22) becomes

$$u(ih) = u_0(ih) + (\phi(0) - u_0(0)) \exp(-\rho p_i), \quad (23)$$

where $\rho = h/c_\epsilon$, $h = 1/N$. Similarly, we write

$$\begin{aligned} u((i+1)h) &= u_0(ih) + (\phi(0) - u_0(0)) \exp(-\rho p_i(i+1)), \\ u((i-1)h) &= u_0(ih) + (\phi(0) - u_0(0)) \exp(-\rho p_i(i-1)). \end{aligned} \quad (24)$$

To handle the effect of the perturbation parameter, exponential fitting factor $\sigma_1(\rho)$ is multiplied on the term containing the perturbation parameter as

$$\begin{aligned} & -\left[\sigma_1(\rho)c_\epsilon - \frac{hp_{i-1}}{12} + \frac{hp_{i+1}}{12}\right] \left(\frac{U_{i-1} - 2U_i + U_{i+1}}{h^2}\right) + \frac{p_{i-1}}{24h} (-3U_{i-1} + 4U_i - U_{i+1}) \\ & + \frac{10p_i}{24h} (U_{i+1} - U_{i-1}) + \frac{p_{i+1}}{24h} (U_{i-1} - 4U_i + 3U_{i+1}) + \frac{q_{i-1}}{12} U_{i-1} + \frac{10q_i}{12} U_i + \frac{q_{i+1}}{12} U_{i+1} \\ & = \frac{1}{12} [f_{i-1} + 10f_i + f_{i+1}], \quad i = 1, 2, \dots, N-1. \end{aligned} \quad (25)$$

Multiplying both sides of (25) by h and substituting ρ for h/c_ε and taking the limit as $h \rightarrow 0$ give

$$\begin{aligned} & -\lim_{h \rightarrow 0} \frac{\sigma_1(\rho)}{\rho} (U_{i-1} - 2U_i + U_{i+1}) + \frac{p_0}{24} \lim_{h \rightarrow 0} (-3U_{i-1} + 4U_i - U_{i+1}) + \frac{10p_0}{24} \lim_{h \rightarrow 0} (U_{i+1} - U_{i-1}) \\ & + \frac{p_0}{24} \lim_{h \rightarrow 0} (U_{i-1} - 4U_i + 3U_{i+1}) = 0. \end{aligned} \quad (26)$$

From (23) and (24), we obtain

$$\begin{cases} \lim_{h \rightarrow 0} (U_{i-1} - 2U_i + U_{i+1}) = (\phi(0) - u_0(0))e^{(-p_0 i \rho)} [e^{(p_0 \rho)} - 2 + e^{-(p_0 \rho)}], \\ \lim_{h \rightarrow 0} (-3U_{i-1} + 4U_i - U_{i+1}) = (\phi(0) - u_0(0))e^{(-p_0 i \rho)} [-3e^{(p_0 \rho)} + 4 - e^{-(p_0 \rho)}], \\ \lim_{h \rightarrow 0} (U_{i-1} - 4U_i + 3U_{i+1}) = (\phi(0) - u_0(0))e^{(-p_0 i \rho)} [e^{(p_0 \rho)} - 4 + 3e^{-(p_0 \rho)}], \\ \lim_{h \rightarrow 0} (U_{i+1} - U_{i-1}) = (\phi(0) - u_0(0))e^{(-p_0 i \rho)} [e^{(p_0 \rho)} - e^{-(p_0 \rho)}]. \end{cases} \quad (27)$$

Substituting (27) into (26) and simplifying give

$$\frac{\sigma_1(\rho)}{\rho} [e^{(p_0 \rho)} - 2 + e^{-(p_0 \rho)}] = \frac{p_0}{24} \left([-3e^{(p_0 \rho)} - e^{-(p_0 \rho)}] + 10[e^{(-p_0 \rho)} - e^{(p_0 \rho)}] + [e^{(p_0 \rho)} + 3e^{-(p_0 \rho)}] \right). \quad (28)$$

The exponential fitting factor is obtained as

$$\sigma_1(\rho) = \frac{\rho p(0)}{2} \coth\left(\frac{\rho p(0)}{2}\right). \quad (29)$$

$$L_L^h U_i = \frac{1}{12} [f_{i-1} + 10f_i + f_{i+1}], \quad i = 1, 2, \dots, N-1, \quad (30)$$

where

Hence, the required finite difference scheme becomes

$$\begin{aligned} L_L^h U_i = & -\left[\sigma_1(\rho) c_\varepsilon - \frac{hp_{i-1}}{12} + \frac{hp_{i+1}}{12} \right] \left(\frac{U_{i-1} - 2U_i + U_{i+1}}{h^2} \right) + \frac{p_{i-1}}{24h} (-3U_{i-1} + 4U_i - U_{i+1}) \\ & + \frac{10p_i}{24h} (U_{i+1} - U_{i-1}) + \frac{p_{i+1}}{24h} (U_{i-1} - 4U_i + 3U_{i+1}) + \frac{q_{i-1}}{12} U_{i-1} + \frac{10q_i}{12} U_i + \frac{q_{i+1}}{12} U_{i+1}, \end{aligned} \quad (31)$$

with the boundary values $U_0 = \phi(0)$ and $U_N = \psi(1)$.

3.2. Right Boundary Layer Problems. In this case, the boundary layer occurs near $x = 1$. From [24], the zeroth-order asymptotic solution of (5)-(6) is given by

$$u(x) = u_0(x) + \frac{p(1)}{p(x)} (\psi(1) - u_0(1)) \exp\left(-\int_x^1 \left(\frac{p(x)}{c_\varepsilon} - \frac{q(x)}{p(x)}\right) dx\right) + O(c_\varepsilon), \quad (32)$$

where u_0 is the solution of the reduced problem.

Using Taylor's series approximation for $u_0(x)$, $p(x)$, and $q(x)$ centred at $x_i = ih$ up to first order and considering $c_\varepsilon \rightarrow 0$, the discretized form of (24) becomes

$$u(ih) = u_0(ih) + (\phi(0) - u_0(0)) \exp\left(-p_i \left(\frac{1}{c_\varepsilon - \rho}\right)\right). \quad (33)$$

Using similar procedures as the left boundary layer case, the exponential fitting factor is obtained as

$$\sigma_2(\rho) = \frac{\rho p(1)}{2} \coth\left(\frac{\rho p(1)}{2}\right). \quad (34)$$

Hence, the required finite difference scheme becomes

$$L_R^h U_i = \frac{1}{12} [f_{i-1} + 10f_i + f_{i+1}], \quad i = 1, 2, \dots, N-1, \quad (35)$$

where

$$\begin{aligned} L_R^h U_i = & -\left[\sigma_2(\rho)c_\varepsilon - \frac{hp_{i-1}}{12} + \frac{hp_{i+1}}{12}\right] \left(\frac{U_{i-1} - 2U_i + U_{i+1}}{h^2}\right) + \frac{p_{i-1}}{24h} (-3U_{i-1} + 4U_i - U_{i+1}) \\ & + \frac{10p_i}{24h} (U_{i+1} - U_{i-1}) + \frac{p_{i+1}}{24h} (U_{i-1} - 4U_i + 3U_{i+1}) + \frac{q_{i-1}}{12} U_{i-1} + \frac{10q_i}{12} U_i + \frac{q_{i+1}}{12} U_{i+1}. \end{aligned} \quad (36)$$

3.3. Convergence Analysis. In this section, we show the stability and convergence analysis for the right boundary layer problems. In similar manner, it is proved for the left boundary layer case. First, we need to prove the discrete comparison principle for the scheme in (35) for guaranteeing existence of unique discrete solution.

Lemma 4 (Discrete comparison principle). *Assume that, for mesh function U_i there exists a comparison function V_i such that $L_R^h U_i \leq L_R^h V_i$, $i = 1, 2, \dots, N-1$ and if $U_0 \leq V_0$ and $U_N \leq V_N$, then $U_i \leq V_i$, $\forall i, i = 0, 1, 2, \dots, N$.*

Proof. The matrix associated with operator L_R^h is of size $(N+1) \times (N+1)$ and satisfies the property of M -matrix. See the detailed proof in [23].

This lemma gives guarantee for the existence of unique discrete solution. In the next lemma, we discuss the uniform stability of the discrete solution. \square

Lemma 5 (Discrete stability estimate). *The solution U_i of the discrete scheme in (35) satisfies the following bound:*

$$|U_i| \leq \frac{\|L_R^h U_i\|}{q^*} + \max\{|U_0|, |U_N|\}. \quad (37)$$

Proof. Let $Q = (\|L_R^h U_i\|/q^*) + \max\{U_0, U_N\}$ and define barrier functions \mathfrak{g}_i^\pm by $\mathfrak{g}_i^\pm = Q \pm U_i$.

On the boundary points, we obtain

$$\mathfrak{g}_0^\pm = Q \pm U_0 = \frac{\|L_R^h U_i\|}{q^*} + \max\{U_0, U_N\} \pm \phi(0) \geq 0,$$

$$\mathfrak{g}_N^\pm = Q \pm U_N = \frac{\|L_R^h U_i\|}{q^*} + \max\{U_0, U_N\} \pm \psi(1) \geq 0. \quad (38)$$

On the discretized spatial domain x_i , $1 < i < N - 1$, we obtain

$$\begin{aligned}
 L_R^h \vartheta_i^\pm &= - \left(c_\varepsilon \sigma(\rho) - \frac{hp_{i-1}}{12} + \frac{hp_{i+1}}{12} \right) \left(\frac{Q \pm U_{i+1} - 2(Q \pm U_i) + Q \pm U_{i-1}}{h^2} \right) \\
 &\quad + \frac{p_{i-1}}{24h} (-3(Q + U_{i-1}) + 4(Q + U_i) - (Q + U_{i+1})) + \frac{10p_i}{24h} ((Q + U_{i+1}) - (Q + U_{i-1})) \\
 &\quad + \frac{p_{i+1}}{24h} ((Q + U_{i-1}) - 4(Q + U_i) + 3(Q + U_{i+1})) \\
 &\quad + \frac{q_{i-1}}{12} (Q + U_{i-1}) + \frac{10q_i}{12} (Q + U_i) + \frac{q_{i+1}}{12} (Q + U_{i+1}) \\
 &= \mp \left(c_\varepsilon \sigma(\rho) - \frac{hp_{i-1}}{12} + \frac{hp_{i+1}}{12} \right) \left(\frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} \right) \pm \frac{p_{i-1}}{24h} (-3U_{i-1} + 4U_i - U_{i+1}) \\
 &\quad \pm \frac{10p_i}{24h} (U_{i+1} - U_{i-1}) \pm \frac{p_{i+1}}{24h} (U_{i-1} - 4U_i + 3U_{i+1}) \pm \frac{q_{i-1}}{12} U_{i-1} \pm \frac{10q_i}{12} U_i \pm \frac{q_{i+1}}{12} U_{i+1} \\
 &\quad + \left(\frac{q_{i-1}}{12} + \frac{10q_i}{12} + \frac{q_{i+1}}{12} \right) Q \\
 &= \left(\frac{q_{i-1}}{12} + \frac{10q_i}{12} + \frac{q_{i+1}}{12} \right) \left(\frac{\|L_R^h U_i\|}{q^*} + \max\{U_0, U_N\} \right) \pm \frac{1}{12} [f_{i-1} + 10f_i + f_{i+1}] \\
 &\geq 0, \quad \text{since } q(x_i) \geq q^* > 0.
 \end{aligned} \tag{39}$$

By the discrete comparison principle in Lemma 4, we obtain $\vartheta_i^\pm \geq 0$, $\forall x_i \in \overline{\Omega}^N$. Hence, the required bound is satisfied.

Now, let us denote the right shifted, centred, and left shifted finite differences, respectively, as

$$\begin{aligned}
 \Delta_R u(x_i) &= \frac{u_{i-1} - 4u_i + 3u_{i+1}}{2h} + hu_i'', \\
 \Delta_C u(x_i) &= \frac{u_{i+1} - u_{i-1}}{2h}, \\
 \Delta_L u(x_i) &= \frac{-3u_{i-1} + 4u_i - u_{i+1}}{2h} - hu_i''.
 \end{aligned} \tag{40}$$

Using Taylor's series approximation, we obtain the bound

$$\begin{aligned}
 \|u_{i-1}' - \Delta_L u(x_i)\| &\leq CN^{-2} \|u^{(4)}(x_i)\|, \\
 \|u_i' - \Delta_C u(x_i)\| &\leq CN^{-2} \|u''(x_i)\|, \\
 \|u_{i+1}' - \Delta_R u(x_i)\| &\leq CN^{-2} \|u^{(4)}(x_i)\|,
 \end{aligned} \tag{41}$$

where $\|u^{(k)}(x_i)\| = \max_{x_0 \leq x_i \leq x_N} |u^{(k)}(x_i)|$, $k = 2, 4$. Similarly, we have

$$\begin{aligned}
 \|u''(x_i) - D^+ D^- u(x_i)\| &\leq CN^{-2} \|u^{(4)}(x_i)\|, \\
 \|D^+ D^- u(x_i)\| &\leq C \|u''(x_i)\|.
 \end{aligned} \tag{42}$$

Now, for $\rho > 0$, C_1 and C_2 are constants, and we have $|\rho \coth(\rho) - 1| \leq C_1 \rho^2$, for $\rho \leq 1$. For $\rho \rightarrow \infty$, since $\lim_{\rho \rightarrow \infty} \coth(\rho) = 1$, $|\rho \coth(\rho) - 1| \leq C_1 \rho$ is given. In general, for all $\rho > 0$, we write

$$C_1 \frac{\rho^2}{\rho + 1} \leq \rho \cot h(\rho) - 1 \leq C_2 \frac{\rho^2}{\rho + 1}, \tag{43}$$

implying that

$$c_\varepsilon \left[p(x_i) \frac{\rho}{2} \cot h\left(p(1) \frac{\rho}{2}\right) - 1 \right] \leq c_\varepsilon \frac{(N^{-1}/c_\varepsilon)^2}{(N^{-1}/c_\varepsilon) + 1} = \frac{N^{-2}}{N^{-1} + c_\varepsilon}. \tag{44}$$

The following theorem gives truncation error bound of the proposed scheme. \square

Theorem 1. Let $u(x_i)$ and U_i be solutions of (5)-(6) and (35), respectively. Then the following error estimate holds:

$$\|L_R^h(u(x_i) - U_i)\| \leq \frac{CN^{-2}}{N^{-1} + c_\varepsilon} \left(1 + c_\varepsilon^{-4} \max_{x_1 \leq x_i \leq x_{N-1}} \exp\left(-\frac{p^*(1-x_i)}{c_\varepsilon}\right) \right). \quad (45)$$

Proof. Consider the truncation error that is given by

$$\begin{aligned} \|L^h u(x_i) - L_R^h U_i\| &\leq \|-c_\varepsilon(u''(x_i) - \sigma(\rho)D^+ D^- u(x_i))\| + \left\|\frac{P_{i-1}}{12}(u'_{i-1} - \Delta_L u(x_i))\right\| \\ &\quad + \left\|\frac{10P_i}{12}(u'_i - \Delta_C u(x_i))\right\| + \left\|\frac{P_{i+1}}{12}(u'_{i+1} - \Delta_R u(x_i))\right\| \\ &\leq \|-c_\varepsilon\left[p(1)\frac{\rho}{2}\cot h\left(p(1)\frac{\rho}{2}\right) - 1\right]D^+ D^- u(x_i)\| + \|c_\varepsilon(u''(x_i) - D^+ D^- u(x_i))\| \\ &\quad + \left\|\frac{P_{i-1}}{12}(u'_{i-1} - \Delta_L u(x_i))\right\| + \left\|\frac{10P_i}{12}(u'_i - \Delta_C u(x_i))\right\| + \left\|\frac{P_{i+1}}{12}(u'_{i+1} - \Delta_R u(x_i))\right\|. \end{aligned} \quad (46)$$

Using the bounds in (44), (41), and (42) gives

$$\begin{aligned} \|L^h(u(x_i) - U_i)\| &\leq \frac{CN^{-2}}{N^{-1} + c_\varepsilon} \|u''(x_i)\| + c_\varepsilon CN^{-2} \|u^{(4)}(x_i)\| + CN^{-2} \|u^{(4)}(x_i)\| \\ &\leq \frac{CN^{-2}}{N^{-1} + c_\varepsilon} \|u''(x_i)\| + CN^{-2} \|u^{(4)}(x_i)\|. \end{aligned} \quad (47)$$

Substituting the bounds for the derivatives of the solution in Lemma 3, we obtain

$$\begin{aligned} \|L^h(u(x_i) - U_i)\| &\leq \frac{CN^{-2}}{N^{-1} + c_\varepsilon} \left(1 + c_\varepsilon^{-2} \exp\left(-\frac{p^*(1-x_i)}{c_\varepsilon}\right) \right) \\ &\quad + CN^{-2} \left(1 + c_\varepsilon^{-4} \exp\left(-\frac{p^*(1-x_i)}{c_\varepsilon}\right) \right). \end{aligned} \quad (48)$$

Since $c_\varepsilon^{-2} \leq c_\varepsilon^{-4}$, we obtain

$$\begin{aligned} \|L^h(u(x_i) - U_i)\| &\leq \frac{CN^{-2}}{N^{-1} + c_\varepsilon} \\ &\quad \cdot \left(1 + c_\varepsilon^{-4} \max_{x_1 \leq x_i \leq x_{N-1}} \exp\left(-\frac{p^*(1-x_i)}{c_\varepsilon}\right) \right). \end{aligned} \quad (49)$$

□

Lemma 6. For $c_\varepsilon \rightarrow 0$, and for given fixed N , we obtain

$$\begin{aligned} \lim_{c_\varepsilon \rightarrow 0} \max_{x_j} \frac{\exp(-p^* x_j / c_\varepsilon)}{c_\varepsilon^m} &= 0, \\ \lim_{c_\varepsilon \rightarrow 0} \max_{x_j} \frac{\exp(-p^* x_j / c_\varepsilon)}{c_\varepsilon^m} &= 0, \quad m = 1, 2, 3, \dots, \end{aligned} \quad (50)$$

where $x_j = jh$, $h = 1/N$, $\forall j = 1, 2, \dots, N-1$.

Proof. See [25–27].

□

Theorem 2. Under the hypothesis of boundedness of discrete solution, the solution of the discrete schemes in (30) satisfies the following uniform error bound:

$$\sup_{0 < c_\varepsilon \ll 1} \|u(x_i) - U_i\| \leq CN^{-1}. \quad (51)$$

Proof. Substituting the results in Lemma 6 into Theorem 1, applying Lemma 5 gives the required bound. □

3.4. Richardson Extrapolation. We apply the Richardson extrapolation technique to accelerate the rate of convergence of the proposed scheme. Richardson extrapolation is a convergence acceleration technique that involves combination of two computed approximations of solution. Interested reader can see the details of Richardson extrapolation in [28]. From (34) and Lemma 6, we obtain

$$\|L_R^h(u(x_i) - U_i)\| \leq C(N^{-1} + N^{-2}), \quad (52)$$

where $u(x_i)$ and U_i are the exact and approximate solutions of (5)-(6), respectively. Applying Lemma 5 in (52) gives

$$|u(x_i) - U_i| \leq CN^{-1} + O(N^{-2}). \quad (53)$$

Let us denote U_i^{2N} for an approximate solution on $2N$ number of mesh points by including the mid-points. Applying the same procedure as in (52) and (53) on $2N$ number of mesh points, we obtain

$$\left| u(x_i) - U_i^{2N} \right| \leq C(2N)^{-1} + O(N^{-2}). \quad (54)$$

Combining (53) and (54) for removing the term CN^{-1} results in

$$\left| u(x_i) - (2U_i^{2N} - U_i) \right| \leq CN^{-2}, \quad (55)$$

giving

$$U_i^{\text{ext}} = 2U_i^{2N} - U_i, \quad (56)$$

as the extrapolated solution. The error bound for the extrapolated solution in (56) becomes

$$\sup_{0 < c_\varepsilon \ll 1} \|u(x_i) - U_i^{\text{ext}}\| \leq CN^{-2}. \quad (57)$$

4. Examples and Numerical Results

In this section, we consider numerical examples to illustrate the theoretical findings of the developed schemes.

Example 1. Consider the problem

$$-\varepsilon u''(x) + u'(x) + 2u(x - \delta) + 5u(x) - u(x + \eta) = 0, \quad (58)$$

with interval conditions $u(x) = 1$, $-\delta \leq x \leq 0$, $u(x) = -1$, $1 \leq x \leq 1 + \eta$.

Example 2. Consider the problem

$$\begin{aligned} -\varepsilon u''(x) + (1 + \exp(x^2))u'(x) + xu(x - \delta) \\ - x^2u(x) + (1 - \exp(-x))u(x + \eta) = -1, \end{aligned} \quad (59)$$

with interval conditions $u(x) = 1$, $-\delta \leq x \leq 0$, $u(x) = -1$, $1 \leq x \leq 1 + \eta$.

Example 3. Consider the problem

$$-\varepsilon u''(x) - 0.5u'(x) + 3u(x - \delta) + 2u(x) - 2u(x + \eta) = -1, \quad (60)$$

with interval conditions $u(x) = 1$, $-\delta \leq x \leq 0$, $u(x) = 0$, $1 \leq x \leq 1 + \eta$.

Example 4. Consider the problem

$$\varepsilon u''(x) + 2.5u'(x) - 2\exp(x)u(x - \delta) - u(x) - xu(x + \eta) = 1, \quad (61)$$

with interval conditions $u(x) = 1$, $-\delta \leq x \leq 0$, $u(x) = 1$, $1 \leq x \leq 1 + \eta$.

The exact solution of the constant coefficient interval value problem

$$\varepsilon u''(x) + au'(x) + \alpha u(x - \delta) + \omega u(x) + \beta u(x + \eta) = f(x), \quad (62)$$

on $\Omega = (0, 1)$, with the interval conditions $u(x) = \phi(x)$, $x \in [-\delta, 0]$, $u(x) = \psi(x)$, $x \in [1, 1 + \eta]$, is given by

$$u(x) = c_1 e^{m_1 x} + c_2 e^{m_2 x} + \frac{f}{c}, \quad (63)$$

where $c = \alpha + \omega + \beta$, $c_1 = (-f + \psi(1)c + e^{m_2}(f - \phi(0)c))/c(e^{m_1} - e^{m_2})$, $c_2 = (-f + \psi(1)c + e^{m_1}(f - \phi(0)c))/c(e^{m_1} - e^{m_2})$, $m_1 = (-(a - \delta\alpha + \eta\beta) + \sqrt{(a - \delta\alpha + \eta\beta)^2 - 4c\varepsilon})/2\varepsilon$, and $m_2 = (-(a - \delta\alpha + \eta\beta) - \sqrt{(a - \delta\alpha + \eta\beta)^2 - 4c\varepsilon})/2\varepsilon$.

The exact solutions of the variable coefficient problems are not known. So, we use the procedure of the double mesh technique to calculate maximum absolute error. The maximum absolute error is defined as

$$E_\varepsilon^N = \max_{0 \leq i \leq N} |U_i^N - U_i^{2N}|, \quad (64)$$

where U_i^N denotes the solution of the problem on N number of mesh points and U_i^{2N} denotes the numerical solution on $2N$ number of mesh points by including the mid-points $x_{(i+1)/2}$ into the mesh numbers. The uniform error estimate is defined as

$$E_\varepsilon^N = \max_\varepsilon |E_\varepsilon^N|. \quad (65)$$

The rate of convergence of the scheme is given by

$$R_\varepsilon^N = \log_2 \left(\frac{E_\varepsilon^N}{E_\varepsilon^{2N}} \right), \quad (66)$$

and the uniform rate of convergence is given as

$$R^N = \log_2 \left(\frac{E^N}{E^{2N}} \right). \quad (67)$$

In Tables 1–8, the maximum absolute error of Examples 1–4 using the proposed scheme is given. In Tables 1, 3, 5, and 7, the maximum absolute error before the Richardson extrapolation is given, and in Tables 2, 4, 6, and 8, the maximum absolute error after the Richardson extrapolation is given. As one observes in the tables, for each number of mesh interval N as $\varepsilon \rightarrow 0$, the maximum absolute error becomes stable and uniform. This indicates that the proposed scheme convergence is independent of the perturbation parameter ε . In the last two rows of each table, we observe the ε -uniform error and the ε -uniform rate of convergence of the scheme. The scheme before the extrapolation gives first-order uniform convergence and the extrapolated scheme gives second-order uniform convergence. In Tables 9 and 10, we compare the maximum absolute error of the proposed scheme with recently published papers in [14, 15, 19]. As one observes, the proposed scheme gives more accurate result.

TABLE 1: Maximum absolute error of Example 1 before Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$1.3808e-03$	$3.4045e-04$	$8.4818e-05$	$2.1193e-05$	$5.2970e-06$	$1.3242e-06$
2^{-4}	$1.7200e-02$	$5.0151e-03$	$1.4739e-03$	$3.5993e-04$	$9.0018e-05$	$2.2464e-05$
2^{-8}	$4.1828e-02$	$2.6301e-02$	$1.4357e-02$	$6.0620e-03$	$1.8873e-03$	$5.0407e-04$
2^{-12}	$4.1711e-02$	$2.6194e-02$	$1.4838e-02$	$7.9779e-03$	$4.1434e-03$	$2.1118e-03$
2^{-16}	$4.1704e-02$	$2.6186e-02$	$1.4834e-02$	$7.9750e-03$	$4.1418e-03$	$2.1124e-03$
2^{-20}	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
2^{-24}	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
2^{-28}	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
2^{-32}	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
2^{-36}	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
2^{-40}	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
E^N	$4.1703e-02$	$2.6186e-02$	$1.4833e-02$	$7.9748e-03$	$4.1417e-03$	$2.1123e-03$
R^N	0.6714	0.8200	0.8953	0.9452	0.9714	–

TABLE 2: Maximum absolute error of Example 1 after Richardson extrapolation for $\delta = 0.6\epsilon$ $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$3.4045e-04$	$2.1193e-05$	$1.3242e-06$	$8.2744e-08$	$5.3296e-09$	$1.3016e-09$
2^{-4}	$5.0151e-03$	$3.5993e-04$	$2.2464e-05$	$1.4036e-06$	$8.7720e-08$	$2.2310e-08$
2^{-8}	$2.6301e-02$	$6.0620e-03$	$5.0407e-04$	$3.3479e-05$	$2.0797e-06$	$4.2147e-07$
2^{-12}	$2.6194e-02$	$7.9779e-03$	$2.1118e-03$	$4.0826e-04$	$3.2897e-05$	$5.2346e-06$
2^{-16}	$2.6186e-02$	$7.9750e-03$	$2.1124e-03$	$5.3615e-04$	$1.3446e-04$	$3.2541e-05$
2^{-20}	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
2^{-24}	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
2^{-28}	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
2^{-32}	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
2^{-36}	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
2^{-40}	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
E^N	$2.6186e-02$	$7.9748e-03$	$2.1123e-03$	$5.3614e-04$	$1.3446e-04$	$3.2541e-05$
R^N	1.7153	1.9166	1.9781	1.9954	2.0468	–

TABLE 3: Maximum absolute error of Example 2 before Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$3.4492e-03$	$1.6925e-03$	$8.3747e-04$	$4.1601e-04$	$2.0724e-04$	$1.0343e-04$
2^{-4}	$2.2763e-02$	$7.6376e-03$	$1.4739e-03$	$1.5941e-03$	$7.8163e-04$	$3.8784e-04$
2^{-8}	$2.3670e-02$	$1.1054e-02$	$5.2737e-03$	$2.5716e-03$	$1.2314e-03$	$4.9403e-04$
2^{-12}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-16}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-20}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-24}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-28}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-32}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-36}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
2^{-40}	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
E^N	$2.3670e-02$	$1.1054e-02$	$5.2735e-03$	$2.5654e-03$	$1.2638e-03$	$6.2704e-04$
R^N	1.0985	1.0677	1.0396	1.0214	1.0111	–

TABLE 4: Maximum absolute error of Example 2 after Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$1.6925e-03$	$4.1601e-04$	$1.0343e-04$	$2.5819e-05$	$6.4524e-06$	$1.1321e-06$
2^{-4}	$7.6376e-03$	$1.5941e-03$	$3.8784e-04$	$9.6460e-05$	$2.4085e-05$	$7.6001e-06$
2^{-8}	$1.1054e-02$	$2.5716e-03$	$4.9403e-04$	$1.0586e-04$	$2.6498e-05$	$8.0231e-06$
2^{-12}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5625e-04$	$3.0982e-05$	$8.2543e-06$
2^{-16}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
2^{-20}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
2^{-24}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
2^{-28}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
2^{-32}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
2^{-36}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
2^{-40}	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
E^N	$1.1054e-02$	$2.5654e-03$	$6.2704e-04$	$1.5584e-04$	$3.8901e-05$	$9.2541e-06$
R^N	2.1073	2.0326	2.0085	2.0022	2.0716	—

TABLE 5: Maximum absolute error of Example 3 before Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$1.1318e-03$	$2.7968e-04$	$6.9551e-05$	$1.7366e-05$	$4.3405e-06$	$1.0850e-06$
2^{-4}	$1.4786e-02$	$3.6872e-03$	$8.5968e-04$	$2.1118e-04$	$5.2578e-05$	$1.3152e-05$
2^{-8}	$4.7613e-02$	$2.9531e-02$	$1.4879e-02$	$5.5516e-03$	$1.4704e-03$	$3.3684e-04$
2^{-12}	$4.8375e-02$	$3.0553e-02$	$1.7528e-02$	$9.4597e-03$	$4.9243e-03$	$2.4698e-03$
2^{-16}	$4.8423e-02$	$3.0586e-02$	$1.7549e-02$	$9.4718e-03$	$4.9323e-03$	$2.5185e-03$
2^{-20}	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9327e-03$	$2.5187e-03$
2^{-24}	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9328e-03$	$2.5187e-03$
2^{-28}	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9328e-03$	$2.5187e-03$
2^{-32}	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9328e-03$	$2.5187e-03$
2^{-36}	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9328e-03$	$2.5187e-03$
2^{-40}	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9328e-03$	$2.5187e-03$
E^N	$4.8426e-02$	$3.0589e-02$	$1.7550e-02$	$9.4725e-03$	$4.9328e-03$	$2.5187e-03$
R^N	0.6628	0.8015	0.8897	0.9413	0.9697	—

TABLE 6: Maximum absolute error of Example 3 after Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$2.7968e-04$	$1.7366e-05$	$1.0850e-06$	$6.7812e-08$	$4.2386e-09$	$1.0101e-09$
2^{-4}	$3.6872e-03$	$2.1118e-04$	$1.3152e-05$	$8.2173e-07$	$5.0203e-08$	$1.2341e-08$
2^{-8}	$2.9531e-02$	$5.5516e-03$	$3.3684e-04$	$2.0489e-05$	$1.2809e-06$	$3.7125e-07$
2^{-12}	$3.0553e-02$	$9.4597e-03$	$2.4698e-03$	$3.8494e-04$	$2.2756e-05$	$4.8133e-06$
2^{-16}	$3.0586e-02$	$9.4718e-03$	$2.5185e-03$	$6.3984e-04$	$1.5770e-04$	$3.8042e-05$
2^{-20}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
2^{-24}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
2^{-28}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
2^{-32}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
2^{-36}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
2^{-40}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
E^N	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
R^N	1.6912	1.9111	1.9768	1.9942	2.0421	—

TABLE 7: Maximum absolute error of Example 4 before Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$9.0320e-03$	$4.5788e-03$	$2.2918e-03$	$1.1469e-03$	$5.7338e-04$	$2.8666e-04$
2^{-4}	$3.7979e-02$	$1.6150e-02$	$6.1639e-03$	$2.4525e-03$	$1.0562e-03$	$4.8456e-04$
2^{-8}	$4.3631e-02$	$2.4943e-02$	$1.3421e-02$	$6.9160e-03$	$3.1649e-03$	$1.2055e-03$
2^{-12}	$4.3723e-02$	$2.4996e-02$	$1.3451e-02$	$7.0011e-03$	$3.5749e-03$	$1.8068e-03$
2^{-16}	$4.3728e-02$	$2.4999e-02$	$1.3452e-02$	$7.0020e-03$	$3.5755e-03$	$1.8071e-03$
2^{-20}	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
2^{-24}	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
2^{-28}	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
2^{-32}	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
2^{-36}	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
2^{-40}	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
E^N	$4.3729e-02$	$2.4999e-02$	$1.3452e-02$	$7.0021e-03$	$3.5755e-03$	$1.8071e-03$
R^N	0.8067	0.8940	0.9420	0.9696	0.9845	–

TABLE 8: Maximum absolute error of Example 4 after Richardson extrapolation for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
2^0	$4.5788e-03$	$1.1469e-03$	$2.8666e-04$	$7.1658e-05$	$1.7914e-05$	$4.4784e-06$
2^{-4}	$1.6150e-02$	$2.4525e-03$	$4.8456e-04$	$1.1290e-04$	$2.7710e-05$	$6.5434e-06$
2^{-8}	$2.4943e-02$	$6.9160e-03$	$1.2055e-03$	$1.6886e-04$	$3.2527e-05$	$8.9632e-06$
2^{-12}	$2.4996e-02$	$7.0011e-03$	$1.8068e-03$	$4.5088e-04$	$7.6356e-05$	$1.2456e-05$
2^{-16}	$2.4999e-02$	$7.0020e-03$	$1.8071e-03$	$4.5547e-04$	$1.1410e-04$	$2.8541e-05$
2^{-20}	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
2^{-24}	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
2^{-28}	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
2^{-32}	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
2^{-36}	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
2^{-40}	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
E^N	$2.4999e-02$	$7.0021e-03$	$1.8071e-03$	$4.5548e-04$	$1.1410e-04$	$2.8541e-05$
R^N	1.8360	1.9541	1.9882	1.9971	1.9992	–

TABLE 9: Example 3: comparison of maximum absolute error of proposed scheme and result in [15] for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 8$	16	32	64	128	256
Proposed scheme						
10^{-4}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
10^{-5}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
10^{-6}	$3.0589e-02$	$9.4725e-03$	$2.5187e-03$	$6.3989e-04$	$1.6062e-04$	$3.9001e-05$
Result in [15]						
10^{-4}	0.10207612	0.06281638	0.03556178	0.01906740	0.00989895	0.00504914
10^{-5}	0.10210236	0.06283382	0.03557070	0.01907016	0.00989794	0.00504604
10^{-6}	0.10210499	0.06283557	0.03557159	0.01907044	0.00989798	0.00504673

TABLE 10: Example 4: comparison of maximum absolute error of proposed scheme and results in [14, 19] for $\delta = 0.6\epsilon$ and $\eta = 0.5\epsilon$.

$\epsilon \downarrow$	$N \longrightarrow 10$	100	10	100	10	100
	Result in [19]		Result in [14]		Proposed scheme	
10^{-1}	$1.5339e-02$	$1.9170e-04$	0.01596700	0.00017254	$9.2052e-03$	$5.8912e-04$
10^{-2}	$2.8175e-02$	$1.8656e-03$	0.08016582	0.00199505	$2.0378e-02$	$1.0991e-03$
10^{-3}	$2.8534e-02$	$3.3897e-03$	0.10168369	0.00999987	$2.0543e-02$	$2.2978e-03$
10^{-4}	$2.8570e-02$	$3.3954e-03$	0.10406804	0.01230494	$2.0554e-02$	$2.3059e-03$

In Figure 1, the influence of the delay parameter on the behaviour of the solution of Examples 3 and 4 is shown for $\epsilon = 2^{-3}$ and $\delta = 0.1\epsilon, 0.5\epsilon$, and 0.9ϵ . From Figure 2, we observe

the numerical solution of Examples 1–4 for different values of the perturbation parameter $\epsilon = 2^{-5}, 2^{-6}$, and 2^{-7} . As observed in the figures, for ϵ going small, strong boundary layer is created.

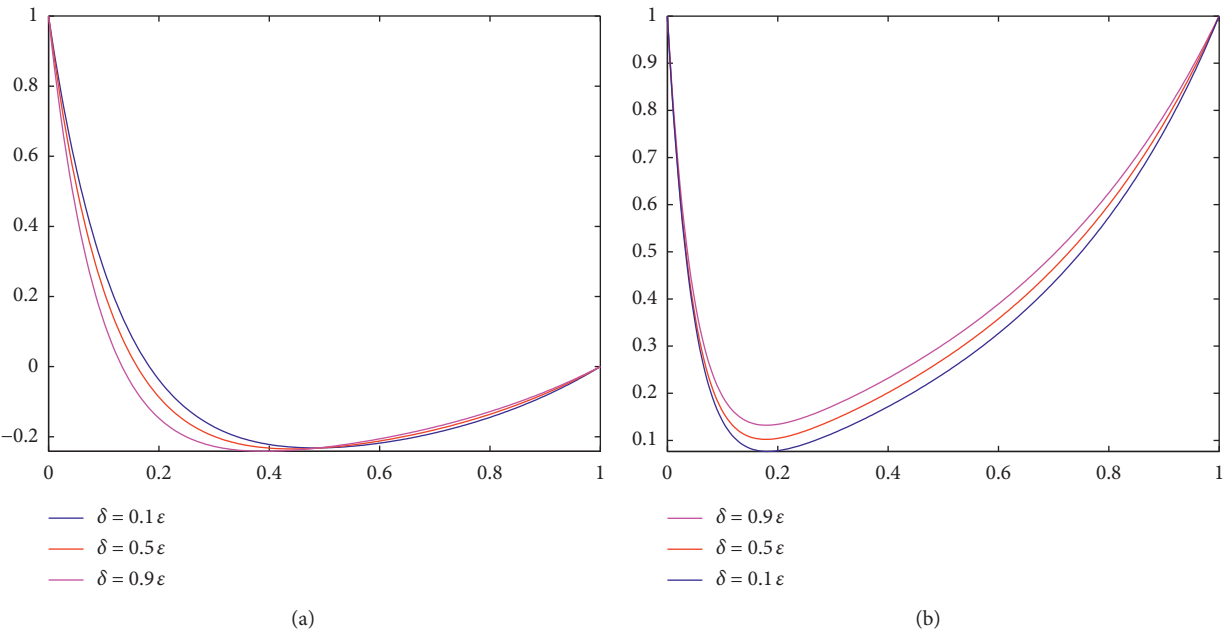


FIGURE 1: Effect of delay on the solution behaviour for $\varepsilon = 2^{-3}$. (a) Example 3. (b) Example 4.

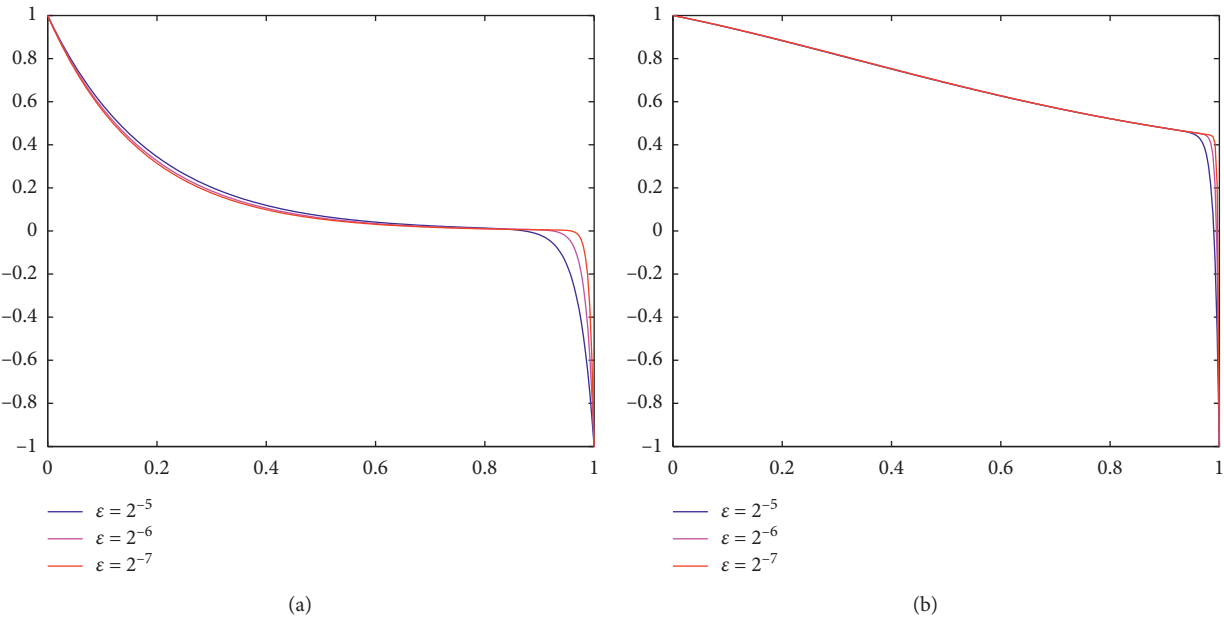


FIGURE 2: Continued.

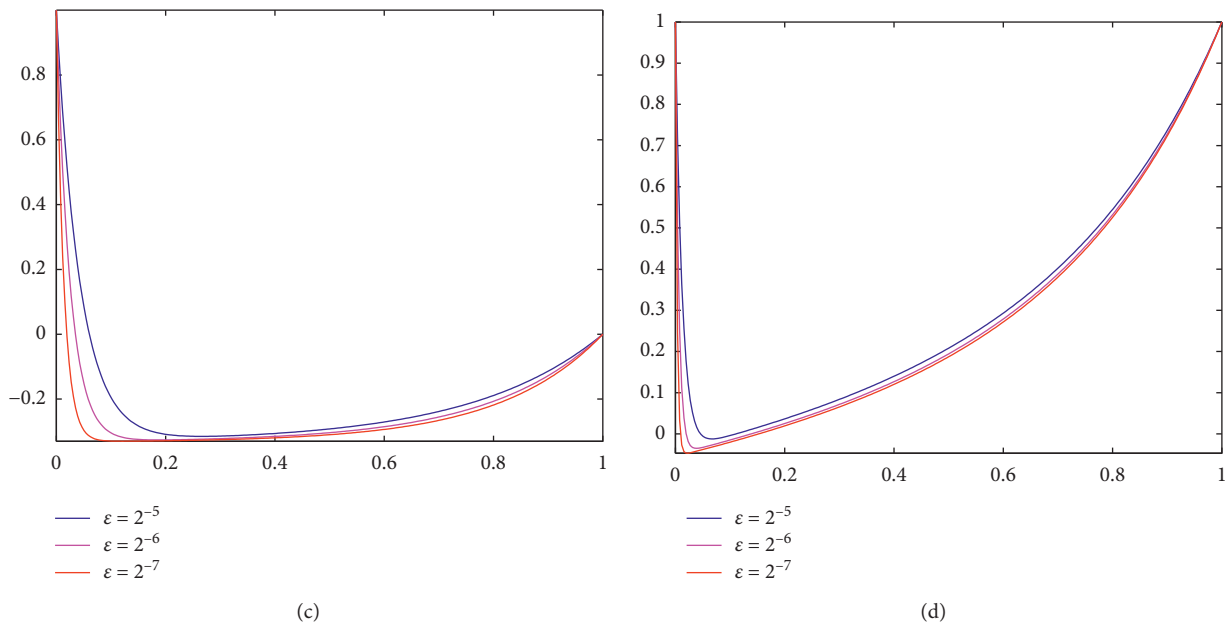


FIGURE 2: Boundary layer formation as ε goes small. (a) Example 1. (b) Example 2. (c) Example 3. (d) Example 4.

5. Conclusion

In this paper, singularly perturbed differential difference equations having mixed small shifts on reaction terms of the equation are considered. The considered problem exhibits boundary layer for small values of the perturbation parameter. The bounds and the behaviour of the continuous solution are discussed. Numerical scheme is developed using the technique of exponentially fitted finite difference method. Stability of the scheme is investigated using comparison principle and solution bound. The proposed scheme converges uniformly with rate of convergence of one before Richardson extrapolation and of two after Richardson extrapolation is applied. Test examples exhibiting boundary layers are considered to validate the theoretical finding. The finding in the computation agrees well with the theoretical findings. The proposed scheme gives more accurate results than existing research findings in the literature. In future works, we extend the proposed scheme for singularly perturbed parabolic differential difference equations and singularly perturbed problems with degenerate coefficients.

Data Availability

No data were used to support the study.

Conflicts of Interest

The authors declare no conflicts of interest.

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