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A NON-LINEAR HYPERBOLIC EQUATION

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<u>ABSTRACT</u>. In this paper the following Cauchy problem, in a Hilbert space H, is considered:

$$(I + \lambda A)u'' + A^{2}u + [\alpha + M(|A^{\frac{1}{2}}u|^{2})]Au = f$$
$$u(0) = u_{0}$$
$$u'(0) = u_{1}$$

M and f are given functions, A an operator in H, satisfying convenient hypothesis, $\lambda \ge 0$ and α is a real number.

For u_0 in the domain of A and u_1 in the domain of $A^{\frac{1}{2}}$, if $\lambda > 0$, and u_1 in H, when $\lambda = 0$, a theorem of existence and uniqueness of weak solution is proved. <u>KEY WORDS AND PHRASES</u>. Nonlinear Wave Equation, Cauchy Problem, Existence and Uniqueness.

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1. INTRODUCTION.

The physical origin of the problem here considered lies in the theory of vibrations of an extensible beam of length ℓ , whose ends are held a fixed distance apart, hinged or clamped, and is either stretched or compressed by an axial force, taking into account the fact that, during vibration, the elements of a beam perform not only a translatory motion, but also rotate; see Timoshenko [9].

A mathematical model for this problem is an initial-boundary value problem for the non-linear hyperbolic equation

$$\frac{\partial^2 u}{\partial t^2} - \lambda \frac{\partial^4 u}{\partial t^2 \partial t^2} + \frac{\partial^4 u}{\partial t^4} - [\alpha + \int_0^\ell [\frac{\partial u}{\partial s} (s,t)]^2 ds] \frac{\partial^2 u}{\partial t^2} = 0, \quad (1.1)$$

where u(l,t) is the deflection of pointlat time t, α is a real constant, proportional to the axial force acting on the beam when it is constrained to lie along the l axis, and λ is a nonnegative constant ($y^{\lambda} = 0$ means neglecting the rotatory inertia, while $\lambda > 0$ means considering it). The non-linearity of the equation is due to considering the extensibility of the beam.

This model, when $\lambda = 0$, was treated by Dickey [2], Ball [1] and, in a Hilbert space formulation, by Medeiros [5]. For related problems, see Pohozaev [7], Lions [4], Menzala [6] and Rivera [8].

In this paper, a theorem of existence and uniqueness of weak solution for a Cauchy problem in a Hilbert space H, is proved for the equation

$$(I + \lambda A)u'' + A^{2}u + [\alpha + M(|A^{2}u|^{2})]Au = f,$$
 (1.2)

with suitable conditions on the operator A and the given functions M and f.

This paper is divided in three parts. In Part 1, the theorem is stated and existence of a weak solution is proved. In Part 2, its uniqueness is established. Finally, an application is given, in Part 3, when H is $L^2(\Omega)$, Ω a bounded open set with regular boundary in Rⁿ, and A is the Laplace operator - Δ .

2. EXISTENCE OF WEAK SOLUTION.

Let H be a real Hilbert space, with inner product (,) and norm ||.

Let A be a linear operator in H, with domain D(A) = V dense in H. With the graph norm of A, denoted || ||, i.e.

$$||v||^2 = |v|^2 + |Av|^2$$
, for v in V,

V is a real Hilbert space and its injection in H is continuous. We assume this injection compact.

Suppose A self-adjoint and positive, i.e., there is a constant
$$k > 0$$
 such that
 $(Av,v) \ge k |v|^2$, for v in V. (2.1)

Let V' be the dual of V and <,> denote the pairing between V' and V. Identifying H and H', it follows that $V \subset H \subset V'$. Injections being continuous and dense, it is known that, for h in H and v in V, <h,v> = (h,v).

Define
$$A^2$$
: $V \rightarrow V'$ by
< A^2u , $v > = (Au, Av)$, for u, v in V. (2.2)

It follows that A^2 is a bounded linear operator from V into V'. Let a(u,v) denote the bilinear form in $D(A^{\frac{1}{2}})$ associated to A, i.e.,

$$a(u,v) = (A^{\frac{1}{2}}u, A^{\frac{1}{2}}v), \text{ for } u, v \text{ in } D(A^{\frac{1}{2}})$$

a(u) means a(u,u).

Given $\lambda \ge 0$, consider in W = D((λA)²) the graph norm of (λA)², denoted $| |_{\lambda}$, i.e.,

$$|w|_{\lambda}^{2} = |w|^{2} + \lambda |A^{2}w|^{2}$$
, for w in W

Note that W = H, if λ = 0, and W = D($A^{\frac{1}{2}}$), if λ > 0; hence V is dense in W. Let α be a real number, M a real C¹ function, with M'(s) \geq 0, for s \geq 0. Assume the existence of positive constants m₀ and m₁ such that M(s) \geq m₀ + m₁s, for $s \ge 0$. Notice that, should M be the identity function, replacement of $\alpha + s$ by $(\alpha - m_0) + (m_0 + s)$, with arbitrary $m_0 > 0$, ensures the fulfilment of the above condition on M.

The theorem can now be stated.

THEOREM. Given f in $L^2(0,T;H)$, u_0 in V, u_1 in W, there is a unique function u = u(t), $0 \le t < T$, such that:

- a) u $\in L^{\infty}(0,T;V)$
- b) $u' \in L^{\infty}(0,T;W)$
- c) u is a weak solution of

$$(I + \lambda A)u'' + A^2u + [\alpha + M(|A^2u|^2)] Au = f,$$
 (2.3a)

i.e., for every v in V, u satisfies in \mathcal{D}' (0,T):

$$\frac{d}{dt} [(u'(t),v) + \lambda a(u'(t),v)] + (Au(t),Av) + [\alpha + M(a(u(t)))] a(u(t),v) = (f(t),v) , \qquad (2.3b)$$

d) The following initial conditions hold:

$$u(0) = u_0, u'(0) = u_1$$
 (2.4ab)

Before proving the theorem, some remarks are pertinent.

Equation (2.3a) makes sense, because (a) and (b) above imply that u, $A^{\frac{1}{2}}u$, Au, u', $(\lambda A)^{\frac{1}{2}}u'$ belong to $L^{\infty}(0,T;H)$.

Initial condition (2.4a) makes sense, because it is known, (see Lions [3]) that if u and u' are in $L^{\infty}(0,T;H)$, then

Now, initial condition (2.4b) must be understood.

Remember $u' \in L^{\infty}(0,T;W)$ implies that $(I+\lambda A)u' \in L^{\infty}(0,T;V')$, because

$$\langle (I + \lambda A)u', v \rangle = (u', v) + \lambda a(u', v), \text{ for } v \text{ in } V$$

From (2.3a), it follows that $(I + \lambda A)u'' \in L^2(0,T;V')$. The fact that both $(I + \lambda A)u'$ and $(I + \lambda A)u''$ belong to $L^2(0,T;V')$ ensures that

$$(I + \lambda A)u' \in C^{0}(0,T;V')$$
 (2.6)

Therefore $(I + \lambda A)u'(0)$ is defined. Given u_1 in W, set $(I + \lambda A)u'(0) =$

 $(I + \lambda A)u_1$, in V'. It follows that u'(0) = u₁, because, it will be proved below,

$$(I + \lambda A)w = 0$$
, for w in W, implies $w = 0$. (2.7)

Indeed, V being dense in W, there is a sequence $(v_j)_{j \in \mathbb{N}}$ in V that converges to w in W, i.e., as $j \rightarrow \infty$,

$$|\mathbf{w} - \mathbf{v}_{j}|_{\lambda}^{2} = |\mathbf{w} - \mathbf{v}_{j}|^{2} + \lambda \mathbf{a}(\mathbf{w} - \mathbf{v}_{j}) \longrightarrow 0$$

and

$$0 = \langle (\mathbf{I} + \lambda \mathbf{A}) \mathbf{w}, \mathbf{v}_{j} \rangle = (\mathbf{w}, \mathbf{v}_{j}) + \lambda \mathbf{a}(\mathbf{w}, \mathbf{v}_{j})$$

tends to

$$(\mathbf{w},\mathbf{w}) + \lambda \mathbf{a}(\mathbf{w}) = |\mathbf{w}|_{\lambda}^{2}$$

Hence w = 0.

Proof of Existence:

It will follow Galerkin method. Suppose, for simplicity, v separable. Let, then, (w_j) be a sequence in V such that, for each m, the set $j \in \mathbb{N}$ w_1, \ldots, w_m is linearly independent and the finite linear combinations of w_1, w_2, \ldots are dense in V. Let V_m denote the finite subspace of V, spanned by w_1, \ldots, w_m .

<u>Approximate Solutions</u> Search for $u_m(t) = \sum_{j=1}^m g_{jm}(t)w_j$ in V_m , such that, for all v in V_m ,

$$((I + \lambda A)u_{m}^{"}(t), v) + (Au_{m}(t), Av) + [\alpha + M(a(u_{m}(t))](Au_{m}(t), v) = (f(t), v)$$
(2.8)

$$u_{m}(0) = u_{om}, u'_{m}(0) = u_{1m},$$
 (2.9)

where u_{om} converges to u_{o} in V and u_{1m} to u_{1} in W.

This system of ordinary differential equations with initial conditions has a solution $u_m(t)$, defined for $0 \le t < t_m \le T$. It is convenient to emphasize that the matrix $((I+\lambda A)w_j, w_i)$, i,j=1,...,m, is invertible, for otherwise the homogeneous system of linear equations

$$\sum_{j=1}^{m} ((I + \lambda A)w_{j}, w_{i})x_{j} = 0 , i = 1, ..., m_{i}$$

would have a non-trivial solution $\alpha_1, \ldots, \alpha_m$, hence

$$\Big|\sum_{j=1}^{m} \alpha_{j} w_{j}\Big|_{\lambda}^{2} = ((I + \lambda A) \sum_{j=1}^{m} \alpha_{j} w_{j}, \sum_{i=1}^{m} \alpha_{i} w_{i}) = 0,$$

a contradiction to the linear independence of w_1, \ldots, w_m .

(ii) <u>A Priori Estimates</u>

Set

For $v = 2u'_{m}(t)$, (2.8) becomes:

$$\frac{d}{dt} |u_{m}^{\prime}(t)|^{2} + \lambda \frac{d}{dt} a(u_{m}^{\prime}(t)) + \frac{d}{dt} |Au_{m}(t)|^{2} + \alpha \frac{d}{dt} a(u_{m}(t)) + M(a(u_{m}(t))) \frac{d}{dt} a(u_{m}(t)) = 2(f(t), u_{m}^{\prime}(t)) \qquad (2.10)$$

$$\overline{M}(\sigma) = \int_{0}^{\sigma} M(s) ds.$$

We integrate (2.10) from 0 to t < t_m and obtain:

$$\begin{aligned} |u_{m}^{\prime}(t)|^{2} + \lambda a(u_{m}^{\prime}(t)) + |Au_{m}(t)|^{2} + \overline{M}(a(u_{m}(t))) \\ &\leq K_{m} + |\alpha|a(u_{m}(t)) + \int_{0}^{t} |u_{m}^{\prime}(s)|^{2} ds, \end{aligned}$$
(2.11)
where $K_{m} = |u_{1m}|^{2} + \lambda a(u_{1m}) + |Au_{0m}|^{2} + M(a(u_{0m})) + \int_{0}^{T} |f(s)|^{2} ds.$

By choice, u_{om} and u_{lm} converge respectively to u_o in V and to u_l in W (remember that $|u_{lm}|_{\lambda}^2 = |u_{lm}|^2 + \lambda a(u_{lm})$).

Therefore, there is a constant $C_0 > 0$, independent of m and greater than K_m such that (2.11) still holds, with K_m replaced by C_0 .

Now,
$$M(s) \ge m_0 + m_1 s$$
 implies $\overline{M}(\sigma) \ge m_0 \sigma + \frac{m_1}{2} \sigma^2$ (2.12)

For $\sigma = a(u_m(t))$, from (2.11), (2.12) and $|\alpha|\sigma \leq \frac{|\alpha|}{2m_1} + \frac{m_1}{2}\sigma^2$

one obtains:

$$|u_{m}'(t)|^{2} + \lambda a(u_{m}'(t)) + |Au_{m}(t)|^{2} + m_{o}a(u_{m}(t)) \leq C + \int_{0}^{t} |u_{m}'(s)|^{2} ds,$$
 (2.13)

where $C = C_0 + \frac{|\alpha|}{2m_1}$, a constant independent of m.

In particular,

$$|u_{\underline{m}}'(t)|^{2} \leq C + \int_{0}^{t} |u_{\underline{m}}'(s)|^{2} ds.$$

Hence, applying Gronwall inequality

$$|u_{m}'(t)|^{2} \leq C e^{T}$$
, (2.14)

It follows from (2.13) and (2.14) that

$$|u'_{m}(t)|^{2} + \lambda a(u'_{m}(t)) + |Au_{m}(t)|^{2} + m_{o}a(u_{m}(t)) \leq K,$$
 (2.15)

where $K = C(1 + Te^{T})$, for all t in $[0, t_m]$ and all m.

In particular, as $k|u_m(t)|^2 \le a(u_m(t))$, it follows that $u_m(t)$ remains bounded; hence it can be extended to [0,T]. Therefore, (2.15) holds, in fact, for all m and t in [0,T].

(iii) Passage to the Limit

It follows that there is a sub-sequence of (u_m) , still dentoed (u_m) , for which, as $m \rightarrow \infty$, the following is true, in the weak star convergence of $L^{\infty}(0,T;H)$:

$$u_m \rightarrow u,$$
 (2.16)

$$a(u_{n}) \rightarrow a(u),$$
 (2.17)

$$Au_m \rightarrow Au,$$
 (2.18)

$$u'_{m} \rightarrow u',$$
 (2.19)

$$\lambda a(u'') \rightarrow \lambda a(u'),$$
 (2.20)

$$M(a(u_{m}))Au_{m} \rightarrow \psi \qquad (2.21)$$

It must still be proved that, in fact

$$\psi = M(a(u))Au$$
 (2.22)

(2.22) will be shown to follow from the Lemma below, whose proof, here reproduced, was given by J.L. Lions [3] and [4].

LEMMA. The mapping $v \rightarrow M(a(v))Av$ from V into H is monotonic. PROOF. The function $\overline{M}(\sigma) = \int_{0}^{\sigma} M(s)ds$ is non-decreasing (because $M'(\sigma) = M(\sigma) \ge 0$) and convex (because $\overline{M}''(\sigma) = M'(\sigma) \ge 0$). Take

$$\phi(\mathbf{v}) = \overline{\mathbf{M}}(\mathbf{a}(\mathbf{v})), \text{ for } \mathbf{v} \text{ in } \mathbf{V}$$

It is easy to see that ϕ has a Gateau derivative,

$$\phi'(v) = 2M(a(v))Av$$
, for v in V,

and that ϕ is convex, i.e., for $0 \le \rho \le 1$,

$$\phi(\rho v + (1-\rho)w) \leq \rho\phi(v) + (1-\rho)\phi(w)$$
, for v, w in V.

This inequality can be written in the form.

$$\frac{\phi(\mathbf{w} + \rho(\mathbf{v} - \mathbf{w})) - \phi(\mathbf{w})}{\rho} \leq \phi(\mathbf{v}) - \phi(\mathbf{w})$$

Passing to the limit, as $\rho \rightarrow 0$ it follows that

$$(\phi'(w), v-w) \leq \phi(v) - \phi(w)$$

and, interchanging the roles of v and w,

$$(\phi'(v), w-v) \leq \phi(w) - \phi(v)$$

Adding the two inequalities above, one obtains:

$$(\phi'(w) - \phi'(v), w-v) \ge 0,$$

This proves the Lemma.

It can now be shown that (2.22) holds.

Indeed, because of the Lemma, for all v in $L^{2}(0,T;V)$, it is true that

$$\int_{0}^{T} (M(a(u_{m}))Au_{m} - M(a(v))Av), u_{m} - v)dt \geq 0$$

Because (u_m) is bounded in $L^{\infty}(0,T;V)$ and (u'_m) in $L^{\infty}(0,T;H)$ and the injection of V in H is compact, (u_m) can, further, be supposed to converge to u strongly in $L^2(0,T;H)$. Hence, as $m \to \infty$:

$$\int_{0}^{T} (\psi - M(a(v))Av, u - v)dt \ge 0$$

Set $u - v = \rho w$, $\rho \ge 0$, divide the inequality by ρ and let $\rho \rightarrow 0$, to obtain:

$$\int_{0}^{T} (\psi - M(a(u))Au, w)dt \ge 0$$

This holds for all w in $L^{\infty}(0,T;V)$, hence $\psi = M(a(u))Au$.

In the following, let k be fixed, k < m; take v in V_k and let $m \rightarrow \infty$. (2.19) and (2.20) imply that, in $\mathcal{D}'(0,T)$,

$$\frac{d}{dt} (u''_m(t), v) \longrightarrow \frac{d}{dt} (u'(t), v) , \qquad (2.23)$$

$$\lambda \frac{d}{dt} a(u''_{m}(t), v) \rightarrow \lambda \frac{d}{dt} a(u'(t), v) \qquad (2.24)$$

Passing to the limit in (2.8), then (2.23) and (2.24), with (2.17), (2.18), (2.21) and (2.22) ensure that

$$\frac{d}{dt} [(u'(t),v) + \lambda a(u'(t),v)] + (Au(t),Av) + + [\alpha + M(a(u(t)))] a(u(t),v) = (f(t),v), \qquad (2.25)$$

holds in $\mathcal{D}'(0,T)$, for all v in V_k. By density, (2.25) holds in $\mathcal{D}'(0,T)$, for all v in V.

Therefore, u is, indeed, a weak solution of (2.3a).

It must still be shown, in order to complete the proof of existence, that u satisfies (2.4ab).

(iv) Initial Conditions

(2.19) means that, for v in V and θ in C'(0,T) such that $\theta(0) = 1$ and $\theta(T) = 0$, as $m \to \infty$

$$\int_{0}^{T} (u'_{m}(t), v)\theta(t)dt \rightarrow \int_{0}^{T} (u'(t), v)\theta(t)dt \qquad (2.26)$$

Because of (2.5) and (2.16), integrating (2.26) by parts, it follows that

$$(u_{om}, v) \rightarrow (u(0), v), \text{ for } v \text{ in } V.$$
 (2.27)

But $u_{om} \rightarrow u_{o}$ in V; hence (2.27) yields

$$(u_0, v) = (u(0), v)$$
, for v in V, i.e., u satisfies (2.4a).

To show that u satisfies (2.4b), consider equations (2.3b) and (2.8) for $v = w_j$, j = 1, 2.... It follows, using (2.17)-(2.18), (2.21) and (2.22), that, as $m \rightarrow \infty$

$$\frac{d}{dt} [(u'_{m}(t), w_{j}) + \lambda a(u'_{m}(t), w_{j})]$$
(2.28)

converges to $\frac{d}{dt} [(u'(t),w_j) + \lambda a(u'(t),w_j)]$ weak star in $L^{\infty}(0,T)$.

(2.28) means that for v in V, θ in C¹(0,T) such that $\theta(0)=1$, $\theta(T)=0$,

$$\int_{0}^{T} \frac{d}{dt} \left[\left(u'_{m}(t), w_{j} \right) + \lambda a(u'_{m}(t), w_{j}) \right] \theta(t) dt$$

$$\longrightarrow \int_{0}^{T} \frac{d}{dt} \left[\left(u'(t), w_{j} \right) + \lambda a(u'(t), w_{j}) \right] \theta(t) dt. \qquad (2.29)$$

Because of (2.6), (2.19) and (2.20), integrating (2.29) by parts, it follows that

$$(u_{lm}, w_j) + \lambda a(u_{lm}, w_j) \rightarrow (u'(0), w_j) + \lambda a(u'(0), w_j), \qquad (2.30)$$

But $u_{lm} \rightarrow u_{l}$ in W, hence the left-hand side of (2.30) converges also to $(u_{1}, w_{j}) + \lambda a(u_{1}, w_{j})$. Therefore

$$(u'(0), w_{j}) + \lambda a(u'(0), w_{j}) = (u_{1}, w_{j}) + \lambda a(u_{1}, w_{j})$$
(2.31)

As (2.31) holds for j = 1, 2, ..., it follows that, in fact, for all v in V:

$$(u'(0),v) + \lambda a(u'(0),v) = (u_1,v) + \lambda a(u_1,v).$$

In other words

$$(I + \lambda A)u'(0) = (I + \lambda A)u_1$$
 in V'

But this implies, (see [6]) $u'(0) = u_1$; i.e. u satisfies (2.4b).

3. UNIQUENESS

2

Let u and \overline{u} be two solutions of (2.3a) with the same initial conditions (2.4ab). Thus w = u - \overline{u} satisfies

$$(I + \lambda A)w'' + A^2w + \alpha Aw + M(a(u))w + [M(a(u)) - M(a(\bar{u}))] A\bar{u} = 0,$$
 (3.1)

$$w(0) = 0, w'(0) = 0,$$
 (3.2ab)

The standard energy method cannot be used to prove uniqueness, because, while the left-hand side of (3.1) belongs to $L^{2}(0,T;V')$, u' belongs to $L^{\infty}(0,T;W)$ and not to $L^{\infty}(0,T;V)$. A modification has to be made; this procedure can be found in Lions [3].

Consider:

$$z(t) = \begin{bmatrix} -\int_{t}^{s} w(\xi) d\xi & \text{for } t \leq s \\ 0 & \text{for } t > s \end{bmatrix}$$
(3.3)

$$w_{1}(t) = \int_{0}^{t} w(\xi) d\xi,$$
 (3.4)

and

Then

$$z(t) = w_1(t) - w_1(s),$$
 (3.5)

$$z(0) = -w_1(s)$$
, (3.6)

$$z(s) = 0$$
, (3.7)

and

$$z'(t) = w(t)$$
 . (3.8)

As $w \in L^{\infty}(0,T;V)$, it is clear [see (3.3) and (3.8)] that z and z' are in $L^{1}(0,T;V)$. Hence, it follows from (3.1) that

$$\int_{0}^{S} < (I + \lambda A)w''(t), z(t) > dt + \int_{0}^{S} (Aw(t), Az(t))dt +$$

$$+ \alpha \int_{0}^{S} (Aw(t), z(t))dt + \int_{0}^{S} M(a(u(t)))(Aw(t), z(t))dt +$$

$$+ \int_{0}^{S} [M(a(u(t))) - M(a(\bar{u}(t)))](A\bar{u}(t), z(t))dt = 0.$$
(3.9)

But, [see (3.8)]

<
$$(I+\lambda A)w''(t), z(t) > = \frac{d}{dt}((I+\lambda A)w'(t), z(t)) - ((I+\lambda A)w'(t), z'(t))$$

= $\frac{d}{dt}((I+\lambda A)w'(t), z(t)) - \frac{1}{2}\frac{d}{dt}((I+\lambda A)w(t), w(t))$

Therefore, using (3.2ab) and (3.7), it follows that (remember $|w|_{\lambda}^2 = |w|_{\lambda}^2 + \lambda a(w)$)

$$\int_{0}^{s} < (I + \lambda A) w''(t), z(t) > dt = -\frac{1}{2} |w(s)|^{2} . \qquad (3.10)$$

Now,[see (3.8)]

$$(Aw(t), Az(t)) = (Az'(t), Az(t)) = \frac{1}{2} \frac{d}{dt} |Az(t)|^2$$

Thus, [see (3.6) and (3.7)]

$$\int_{0}^{8} (Aw(t), Az(t))dt = -\frac{1}{2} |Aw_{1}(s)|^{2}$$
(3.11)

As
$$|w|_{\lambda} \ge |w|$$
, (3.9), (3.10) and (3.11) yield
 $|w(s)|^{2} + |Aw_{1}(s)|^{2} \le 2|\alpha| \int_{0}^{s} |(w(t), Az(t))|dt$
 $+ 2 \int_{0}^{s} M(a(u(t)))|(w(t), Az(t))|dt$
 $+ 2 \int_{0}^{s} |M(a(u(t))) - M(a(\bar{u}(t)))| |(\bar{u}(t), Az(t))|dt$ (3.12)

As u, $\overline{u} \in L^{\infty}(0,T;V)$ and, for $s \ge 0$, $M \ge 0$ is a C^1 function, with $M^1 \ge 0$, there is a constant C > 0 such that

$$2 \int_{0}^{s} M(a(u(t))) | (w(t), Az(t)) | dt \leq 2C_{0} \int_{0}^{s} |w(t)| | Az(t) | dt \qquad (3.13)$$

And

$$2 \int_{0}^{S} |M(a(u(t)) - M(a(\overline{u}(t)))| | (\overline{u}(t), Az(t))| dt$$

$$\leq 2C_{0} \int_{0}^{S} |a(u(t)) - a(\overline{u}(t))| | \overline{u}(t)| | Az(t)| dt$$

$$\leq 2C_{0}^{2} \int_{0}^{S} |(A(u(t) + \overline{u}(t)), w(t))| | Az(t)| dt$$

$$\leq 2C_{0}^{3} \int_{0}^{S} |w(t)| | Az(t)| dt \qquad (3.14)$$

Notice that, [see (3.5)]

$$2|w(t)| |Az(t)| \le 2[|w(t)|^2 + |Aw_1(t)|^2] + |Aw_1(s)|^2.$$
 (3.15)

Hence, it follows from (3.15) that

$$2|\alpha| \int_{0}^{s} |(w(t),Az(t))| dt \leq 2|\alpha| \int_{0}^{t} [|w(t)|^{2} + |Aw_{1}(t)|^{2}] dt + |\alpha|s|Aw_{1}(s)|_{2}^{2}(3.16)$$

Hence (3.13) and (3.15) give

$$2 \int_{0}^{s} M(a(u(t))) | (w(t), Az(t)) | dt$$

$$\leq 2C_{0} \int_{0}^{s} [|w(t)|^{2} + |Aw_{1}(t)|^{2}] dt + C_{0} s |Aw_{1}(s)|^{2}$$
(3.17)

Now (3.14) and (3.15) give

$$2 \int_{0}^{s} |M(a(u(t))) - M(a(\bar{u}(t)))| |(\bar{u}(t), Az(t))|dt$$

$$\leq 2c_{0}^{3} \int_{0}^{s} [|w(t)|^{2} + |Aw_{1}(t)|^{2}]dt + c_{0}^{3} s|Aw_{1}(s)|^{2}$$
(3.18)

It now follows from (3.12), with (3.16) - (3.17) that

$$|w(s)|^{2} + (1 - Cs)|Aw_{1}(s)|^{2} \le 2C \int_{0}^{s} [|w(t)|^{2} + |Aw_{1}(t)|^{2}]dt,$$
 (3.19)

where $C = |\alpha| + C_0 + C_0^3$.

Take s_o such that, for $0 \le s \le s_o$, $\frac{1}{2} \le 1 - Cs \le 1$. Hence (3.19) yields for $0 \le s \le s_o$:

$$|w(s)|^{2} + \frac{1}{2} |Aw_{1}(s)|^{2} \leq 2C \int_{0}^{s} [|w(t)|^{2} + |Aw_{1}(t)|^{2}]dt.$$

A fortiori, for $0 \le s \le s_0$,

$$|w(s)|^{2} + |Aw_{1}(s)|^{2} \leq 4C \int_{0}^{s} [|w(t)|^{2} + |Aw_{1}(t)|^{2}]dt.$$

Applying Gronwall inequality, it then follows that

$$w(s) = 0$$
, for $0 \le s \le s_0$.

Similarly, it is proved that $\hat{w}(s) = 0$, for $s \leq s \leq s + \tau$, with $\tau > 0$. It then follows that, in fact, w(s) = 0, for $0 \leq s < T$.

The proof of uniqueness is complete.

4. APPLICATION

For Ω a bounded open set in Rⁿ, with regular boundary, consider

$$H = L^{2}(\Omega), \quad V = H_{O}^{1}(\Omega) \bigcap H^{2}(\Omega)$$

Let Δ be the Laplace and ∇ the gradient operators in Rⁿ respectively. Take A = - Δ , hence $A^{\frac{1}{2}} = \nabla$. Hypothesis on A are satisfied. Notice that, in this case, the condition (Av,v) $\geq k |v|^2$, for v in V, is the Friedrichs - Poincaré inequality; the compactness of the injection of V in H is the Rellich theorem.

It is clear that

W =
$$L^{2}(\Omega)$$
, if $\lambda = 0$
W = $H^{1}(\Omega)$, if $\lambda > 0$

Now (,) and | | are respectively the inner product and the norm in $L^2(\Omega)$. Given

$$\begin{split} & u_{o} \in H_{o}^{1}(\Omega) \quad H^{2}(\Omega) \\ & u_{1} \in L^{2}(\Omega), \text{ if } \lambda = 0; \ u_{1} \in H^{1}(\Omega), \text{ if } \lambda > 0, \\ & f \in L^{2}(0,T; \ L^{2}(\Omega)), \end{split}$$

the theorem proved above ensures existence and uniqueness of weak solution for the non-linear hyperbolic equation

$$(\mathbf{I} - \lambda \Delta)\mathbf{u} + \Delta^2 \mathbf{u} - \bar{\iota} \alpha M(|\nabla \mathbf{u}|^2)] \Delta \mathbf{u} = \mathbf{f},$$

satisfying $u(0) = u_0, u'(0) = u_1$.

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