# STRONG OSCILLATIONS FOR SECOND ORDER NONLINEAR FUNCTIONAL DIFFERENTIAL EQUATIONS <br> JURANG YAN 

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#### Abstract

In this paper, we establish some strongly oscillation theorems for nonlinear second order functional differential equation


$$
x^{\prime \prime}(t)+p(t) f(x(t), x(g(t)))=0
$$

without assuming that $g(t)$ is retarded or advanced.

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1. INTRODUCTION. We consider the second order nonlinear functional differential equation

$$
\begin{equation*}
x^{\prime \prime}(t)+p(t) f(x(t), x(g(t)))=0 \tag{1.1}
\end{equation*}
$$

where $p(t), g(t) \varepsilon C\left(\left[t_{0}, \infty\right), R\right), g(t) \rightarrow \infty$ as $t \rightarrow \infty$ and $f(u, v) \varepsilon C(R, R)$ and has the sign of $u$ and $v$ when they have the same sign. We shall restrict our attention to solutions of (1.1) which exist on some positive half-1ine. A nontrivial solution $x(t)$ is called oscillatory if $x(t)$ has an unbounded set of zeros, and otherwise it is called nonoscillatory. Equation (1.1) is said to be oscillatory if every solution of (1.1) is oscillatory.

Oscillation theory for (1.1) has been developed by many authors. Bradley [1], Chiou [2], Erbe [3] Gol1witzer [4], Ladas [5], Travis [6], Waltman [7], Wong [8] and references therein. It is wellknown theorem of Wintner [9] and Leighton [10] that the linear equation

$$
x^{\prime \prime}(t)+p(t) \quad x(t)=0
$$

and Waltman [7] demonstrated that the equation

$$
\begin{equation*}
x^{\prime \prime}(t)+p(t) x(g(t))=0 \tag{1.2}
\end{equation*}
$$

is oscillatory if $p(t) \geqslant 0$ and $\int^{\infty} p(t) d t=\infty$. Travis [6] constructed a counterexample showing that the Leighton-Wintner oscillation theorem can not be extended to Equation (1.2) unless $p(t) \geqslant 0$. The author [11] extended Bradley-Waltman oscillation theorem to (1.1) i.e. if $p(t) \geqslant 0$ and $\int^{\infty} p(t) d t=\infty$, then (1.1) is ocillatory.

The purpose of this paper is to establish some strongly oscillation criteria for (1.1). We are primarily interested in the case when $p(t) \geqslant 0, \int^{\infty} p(t) d t<\infty$, are satisfied.

Considering the equation

$$
\begin{equation*}
x^{\prime \prime}(t)+\lambda p(t) x(t)=0 \tag{1.3}
\end{equation*}
$$

We shall call $p(t)$ a strongly oscillatory coefficient if (1.3) is oscillatory for all positive $\lambda$. If $p(t) \geqslant 0$, Nehari [12] shows that
$\lim \sup t \int_{t}^{\infty} p(s) d s=\infty$ $t \rightarrow \infty$
is a necessary and sufficient condition for $p(t)$ to be a strongly oscillatory coefficient. In general, motivated by Nehari, we define as follows: Fquation (1.1) is said to be strongly oscillatory if the related equation of (1.1)

$$
\begin{equation*}
x^{\prime \prime}(t)+\lambda p(t) f(x(t), x(g(t)))=0 \tag{1.4}
\end{equation*}
$$

is osillatory for all positive $\lambda$.

## 2. MAIN RESULTS.

For Equation (1.1) the following conditions are assumed to hold throughout the paper:
(i) $p(t) \geqslant 0$ and there exists $h(t) \leqslant m i n(g(t), t)$ such that $0<k \leqslant h^{\prime}(t)$ where $k$ is a constant.
(ii) there exists $m>0$ such that $|u| \geqslant m$ implies

$$
\left\lvert\, \lim _{\inf } \frac{|f(u, v)|}{|\phi(v)|} \geqslant \varepsilon_{1}>0\right.
$$

where $\phi(v) \varepsilon C^{\prime}(R), v \phi(v)>0$ and $\phi^{\prime}(v) \neq 0$ for $v \neq 0$, and $\quad \begin{aligned} & \text { lim } \\ & \text { where } \varepsilon_{1} \text { and } \delta \text { are constants. }\end{aligned} \quad \phi^{\prime}(v) \geqslant \delta>0$
We begin with a Lemma which needed in establishing our results.

LEMMA 2.1. Suppose that for $\lambda=\lambda_{0}>0$ Equation (1.1) has a nonoscillatory solution $x(t)$. Then the following inequality holds for all large $t$,

$$
\begin{equation*}
w(t) \geqslant \sigma \int_{t}^{\infty} w^{2}(s) d s+\lambda_{0} \varepsilon \int_{t}^{\infty} p(s) d s \tag{2.1}
\end{equation*}
$$

where $w(t)=x^{\prime}(t) / \phi(x(h(t))), \sigma$ and $\varepsilon$ are positive constants.
PROOF. Assume that Equation (1.4) at $\lambda=\lambda_{0}$, has a nonoscillatory solution $x(t)>0$ for $t \geqslant t_{0}>0$. A similar proof will hold if $x(t)<0$ for $t \geqslant t_{0}$. It is easy to verify that $x^{\prime \prime}(t)<0$ and $x^{\prime}(t)>0$ for all large $t$. Let $w(t)=x^{\prime}(t) / \phi(x(h(t)))$, then

$$
w^{\prime}(t)=-\lambda_{0} p(t) \frac{f(x(t), x(g(t)))}{\phi(x(h(t)))}-\frac{\phi^{\prime}(x(h(t))) x^{\prime}(h(t)) h^{\prime}(t)}{\phi(x(h(f)))} w(t) .
$$

 If $\underset{t \rightarrow \infty}{\lim } x(t)=\alpha$ is finite, then

$$
\lim _{t \rightarrow \infty} \frac{f(x(t), x(g(t)))}{p(x(g(t)))}=\frac{f(\alpha, \alpha)}{\phi(\alpha)}=\varepsilon_{2}>0 .
$$

If $\lim _{t \rightarrow \infty} x(t)=\infty$, then by (ii) we have that

$$
\frac{f(x(t), x(g(t)))}{\phi(x(g(t)))} \geqslant \varepsilon_{1}
$$

for large $t$. In either case, for sufficiently large $t$, we have that

$$
\begin{equation*}
\frac{f(x(t), x(g(t)))}{\phi(x(g(t)))} \geqslant \varepsilon, \text { where } \varepsilon=\min \left(\varepsilon_{1}, \varepsilon_{2}\right) . \tag{2.2}
\end{equation*}
$$

Since $x(t)$ is increasing, for large $t$ we have that

$$
p(t) \frac{f(x(t), x(g(t)))}{\phi(x(h(t)))} \geqslant p(t) \frac{f(x(t), x(g(t)))}{\phi(x(g(t)))} \geqslant \lambda_{0} \varepsilon_{p}(t)
$$

and in view of $x^{\prime \prime}(t)<0$ for large $t$ and (ii) we see that

$$
\frac{\phi^{\prime}(x(h(t))) x^{\prime}(h(t)) h^{\prime}(t)}{\dot{\varphi}(x(h(t)))} w(t) \geqslant \frac{\phi^{\prime}(x(h(t))) x^{\prime}(t) h^{\prime}(t)}{\phi(x(h(t)))} \geqslant k \delta w^{2}(t)=\sigma w^{2}(t)
$$

Thus for $t \geqslant t_{1} \geqslant t_{0}$,

$$
\begin{equation*}
w^{\prime}(t)+\sigma w^{2}(t)+\lambda_{0} \varepsilon p(t) \leqslant 0 \tag{2.3}
\end{equation*}
$$

and

$$
\begin{equation*}
w^{\prime}(t)+\sigma w^{2}(t) \leqslant 0 \tag{2.4}
\end{equation*}
$$

From (2.4), we have that

$$
\frac{d}{d t}\left(-\frac{1}{w(t)}+\sigma t\right) \leqslant 0 .
$$

Integrating the above inequality from $t_{1}$ to $t$ we obtain that

$$
\begin{equation*}
0<w(t) \leqslant \frac{1}{\sigma t+b} . \tag{2.5}
\end{equation*}
$$

where $b=\frac{1}{w\left(t_{1}\right)}-k t_{1}$. From (2.5) we have that

$$
\lim _{t \rightarrow \infty} w(t)=0
$$

Integrating (2.3) from $t_{1}$ to $t$, then letting $t \rightarrow \infty$ we obtain that (2.l) holds for all large.

We introduce the function sequence $\left\{p_{n}(t)\right\}, n=0,1,2, \ldots$,
defined as follows:

$$
\begin{align*}
& p_{0}(t)=p(t), \quad p_{1}(t)=\int_{t}^{\infty} p_{0}(s) d s \\
& \quad p_{n+1}(t)=\int_{t}^{\infty} p_{n}^{2}(s) d s, \quad n=1,2,3, \ldots \tag{2.7}
\end{align*}
$$

THEOREM 2.2. Assume that one of the following conditions holds, $\left(I_{1}\right)$ there is an integer $m \geqslant 1$ such that $p_{n}(t)$ is defined for $n=1,2, \ldots, m$, and

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup t p_{m}(t)=\infty ; \tag{2.8}
\end{equation*}
$$

$\left(I_{2}\right)$ there is an integer $m \geqslant 2$ such that $p_{n}(t)$ is defined for $n=1,2, \ldots, m$, but $p_{m}(t)$ does not exist, i.e.

$$
\int^{\infty} p_{m-1}^{2}(t) d t=\infty
$$

Then equation (1.1) is strongly oscillatory.
PROOF. Assume to the contrary that Equation (1.4) at $\lambda=\lambda_{0}>0$, has a nonoscillatory solution $x(t)>0$ for $t \geqslant t_{0} \geqslant 0$. A similar argument holds when $x(t)<0$ for $t \geqslant t_{0} \geqslant 0$. Let $w(t)=x^{\prime}(t) / \phi(x(h(t)))$. As in the proof of Lemma 2.1 , we can obtain

$$
\begin{equation*}
w^{\prime}(t)+\sigma w^{2}(t)+\lambda_{0} \varepsilon_{p_{0}}(t) \leqslant 0, \quad t \geqslant t_{1} \geqslant t_{0} \tag{2.9}
\end{equation*}
$$

Suppose $m=1$. Define $u_{o}(t)=\sigma w(t)$; then

$$
\begin{equation*}
u_{0}^{\prime}(t)+u_{0}^{2}(t)+\xi_{0} p_{0}(t) \leqslant 0, \quad t \geqslant t_{1} \geqslant t_{0} \tag{2.10}
\end{equation*}
$$

where $\xi_{0}=\lambda_{0} \varepsilon \sigma$. However, by a well-known theorem of Wintner [13] this implies the equation

$$
\begin{equation*}
y^{\prime \prime}(t)+\xi_{0} p_{0}(t) y(t)=0 \tag{2.11}
\end{equation*}
$$

is nonoscillatory. This contradicts the fact that the condition ( $I_{1}$ ) implies $p_{0}(t)$ is a strongly oscillatory coefficient.

If $m>1$, integrating (2.10) from $t_{1}$ to $t$ we obtain

$$
\begin{equation*}
u_{0}(t)-u\left(t_{1}\right)+\int_{t_{1}}^{t} u_{0}^{2}(s) d s+\xi_{0} \int_{t_{1}}^{t} p_{0}(s) d s<0 . \tag{2.12}
\end{equation*}
$$

Since $\lim _{t \rightarrow \infty} u_{0}(t)=1 \lim _{t} \sigma(t)=0$, from (2.12) we get

$$
\begin{equation*}
u_{0}(t) \geqslant u_{1}(t)+\xi_{0} p_{1}(t), \quad t \geqslant t_{1}, \tag{2.13}
\end{equation*}
$$

where $u_{1}(t)=\int_{t}^{\infty} u_{0}^{2}(s) d s$. Then (2.13) implies that

$$
u_{1}^{\prime}(t)=-u_{0}^{2}(t) \leqslant-u_{1}^{2}(t)-2 \xi_{0} u_{1}(t) p_{1}(t)-\xi_{0}^{2} p_{1}^{2}(t) .
$$

Hence

$$
\begin{equation*}
u_{1}^{\prime}(t)+u_{1}^{2}(t)+\xi_{1} p_{1}^{2}(t) \leqslant 0, \tag{2.14}
\end{equation*}
$$

where $\xi_{1}=\xi_{0}^{2}$.
If $m=2$, then by the Nehari theorem, the condition ( $I_{1}$ ) implies equation

$$
\begin{equation*}
y^{\prime \prime}(t)+\xi_{1} p_{1}^{2}(t) y(t)=0 \tag{2.15}
\end{equation*}
$$

is oscillatory, contradicting (2.14).
If the condition $\left(I_{2}\right)$ is satisfied, then by the Leighton-Wintner theorem we have that equation (2.15) is oscillatory, again a contradiction.

When $m>2$, we obtain inductively that

$$
\begin{equation*}
u_{m-1}^{\prime}(t)+u_{m-1}^{2}(t)+\xi_{m-1} p_{m-1}^{2}(t) \leqslant 0, \quad t \geqslant t^{*} \geqslant t_{1} . \tag{2.16}
\end{equation*}
$$

where

$$
u_{m-1}(t)=\int_{t}^{\infty} u_{m-2}^{2}(s) d s, \xi_{m-1} \text { is a constant and } p_{m-1}(t) \text { is }
$$

defined by (2.7). Applying the Wintner theorem again, it follows that equation

$$
\begin{equation*}
y^{\prime \prime}(t)+\xi_{m-1} p_{m-1}^{2}(t) y(t)=0 \tag{2.17}
\end{equation*}
$$

is nonoscillatory. But this contradicts again the fact that the condition $\left(I_{1}\right)$ or ( $I_{2}$ ) implies that equation (2.17) is oscillatory. The proof is thus complete.

REMARK. Theorem 2.2 includes Theorem 2.? in [6] as a special case, i.e. $\phi(u)=u$ and $m=1$.

Consider the function sequence $\left\{q_{n}(t, \xi, n)\right\}, n=1,2, \ldots$, which is defined as follows:

$$
\begin{align*}
& q_{0}(t, \eta)=n \int_{t}^{\infty} p(s) d s, q_{1}(t, \xi, \eta)=\xi \int_{t}^{\infty} q_{0}^{2}(s, \eta) d s+q_{0}(t, \eta), \\
& q_{n}(t, \xi, \eta)=\xi \int_{t}^{\infty} q_{n-1}^{2}(s, \xi, \eta) d s+q_{0}(t, \eta), \quad t \geqslant t_{0}, n=2,3, \ldots . \tag{2.18}
\end{align*}
$$

where $\zeta$ and $l$ are positive constants.

THEOREM 2.3. For all positive constants $\xi$ and $n$ assume that any one of following conditions is satisfied:
(II ${ }_{1}$ ) there is a positive integer $m$ such that $q_{n}(t, \xi, n)$ is defined for $\mathrm{n}=1,2, \ldots, \mathrm{~m}-1$, but $\mathrm{q}_{\mathrm{m}}(\mathrm{t}, \xi, n)$ does not exist;
( $I_{2}$ ) $q_{n}(t, \xi, n$ ) is defined for $n=1,2, \ldots$, but the function sequence (2.18) is not convergent for all large $t$.
( $\mathrm{II}_{3}$ ) the sequence (2.18) is convergent and $\lim q_{n}(t, \xi, \eta)=q(t, \zeta, \eta)$, but $q(t, \xi, n) \notin L^{2}\left(t_{0}, \infty\right)$. $n \rightarrow \infty$
Then equation (1.4) is strongly oscillatory.
PROOF. Assume that Equation(1.1) at $\lambda=\lambda_{0}>0$, has a solution $x(t)>0$ for $t \geqslant t_{0} \geqslant 0$. A similar proof will hold if $x(t)<0$ for $t \geqslant t_{0}$. Let
$w(t)=x(t) / \phi(x(h(t)))$. From Lemma 2.1 we can obtain (2.1). It follows that $w(t) \geqslant q_{0}\left(t, \eta_{0}\right)$, where $\eta_{0}=\lambda_{0} \varepsilon_{0}$. Hence

$$
\begin{equation*}
w^{2}(t) \geqslant q_{0}^{2}\left(t, \eta_{0}\right), \quad t \geqslant t_{1} \geqslant t_{0} \tag{2.19}
\end{equation*}
$$

Suppose that ( $\mathrm{II}_{1}$ ) holds. If $m=1$ from (2.1), (2.19) implies that
$\int_{t}^{\infty} q_{0}^{2}\left(s, n_{0}\right) d s \leqslant \int_{t}^{\infty} w^{2}(s) d s<\infty$. Thus

$$
q_{1}\left(t, \xi_{0}, \eta_{0}\right)=\xi_{0} \int_{t}^{\infty} q_{0}^{2}\left(s, \eta_{0}\right) d s+q_{0}\left(t, \eta_{0}\right) \leqslant w(t), \quad t \geqslant t_{1}
$$

where $\xi_{0}=\sigma$. This is in contradiciton to the nonexistence of $q_{1}$.
If $m>1$, from (2.1) and (2.19) we get that $q_{m-1} \leqslant w(t)$. Hence
$\int_{t}^{\infty} q_{m-1}\left(s, \xi_{0} n_{0}\right) d s \leqslant \int_{t}^{\infty} w^{2}(s) d s<\infty$. Applying (2.1) we have that

$$
q_{m}\left(t, \xi_{0}, \eta_{0}\right)=\xi \int_{t}^{\infty} q_{m-1}^{2}\left(s, \xi_{0}, \eta_{0}\right) d s+q_{0}\left(t, \eta_{0}\right) \leqslant w(t)
$$

and we arrive at a contradiction of ( $1 I_{1}$ ).
Suppose that ( $\mathrm{II}_{2}$ ) holds. From (2.18) and (2.19), we conclude that for all large $t$

$$
\begin{equation*}
q_{n-1}\left(t, \xi_{0}, \eta_{0}\right) \leqslant q_{n}\left(t, \xi_{0}, \eta_{0}\right) \leqslant w(t), n=1,2, \ldots \ldots \tag{2.20}
\end{equation*}
$$

Therefore $\lim _{n \rightarrow \infty} q_{n}\left(t, \xi_{0}, \eta_{0}\right)$ exists and has a finte limit. But this constradicts the fact that $q_{n}\left(t, \xi_{j} \eta_{0}\right)$ is not convergent.

Suppose that ( $\mathrm{II}_{3}$ ) holds. By (2.20),

$$
\begin{equation*}
\lim q_{n}\left(t, \xi_{0}, \eta_{0}\right)=q\left(t, \xi_{0}, \eta_{0}\right) \leqslant w(t) \tag{2.21}
\end{equation*}
$$

Using (2.21), we have that $\int_{t}^{\infty} q^{2}\left(s, \xi_{0}, \eta_{o}\right) d s \leqslant \int_{t}^{\infty} w^{2}(s) d s<\infty$ which contradicts the condition ( $\mathrm{II}_{3}$ ). This completes the proof.

THEOREM 2.4. Assume that

$$
\begin{equation*}
\int_{\alpha}^{\infty} \frac{d u}{\phi(u)}<\infty \quad \text { and } \int_{-\alpha}^{-\infty} \frac{d u}{\phi(u)}<\infty, \quad \alpha>0 \tag{2.22}
\end{equation*}
$$

Further assume that sequence (2.18) for all positive constants $\xi$ and $\eta$ satisfies any one of the following conditions:
( III $_{1}$ ) there is a positive integer $m$ such that $q_{n}(t, \xi, n)$ is defined for $n=0,1,2 \ldots m$, and

$$
\int_{t_{0}}^{\infty} q_{m}(s, \xi, \eta) d s=\infty
$$

( III $_{2}$ ) $q_{n}(t, \xi, n)$ is defined for $n=0,1,2, \ldots$, and $\lim _{n \rightarrow \infty} q_{n}(t, \xi, \eta)=q(t, \xi, \eta)$
exists and satisfies

$$
\int_{t_{0}}^{\infty} q(s, \xi, \eta) d s=\infty
$$

Then Equation (1.1) is strongly osillatory.
PROOF. Assume that Equation (1.4) at $\lambda=\lambda_{0}>0$, has a nonoscillatory solution
$x(t)>0$ for $t \geqslant t_{0}$. The case $x(t)<0$ is handled similarly. Let $w(t)=\frac{x^{\prime}(t)}{\phi(x(h(t))}$. By Lemmma 2.1, we find that (2.1) holds.

Suppose that ( $\mathrm{III}_{1}$ ) holds, then, as proof of Theorem 2.3,

$$
q_{m}\left(t, \xi_{0}, \eta_{0}\right)=\xi_{0} \int_{t}^{\infty} q_{m}^{2}\left(s, \xi_{0}, \eta_{0}\right) d s+q_{o}\left(t, \xi_{0}, \eta_{0}\right) \leqslant w(t)
$$

or

$$
\begin{equation*}
q_{m}\left(t, \xi_{0}, \eta_{0}\right) \leqslant \frac{x^{\prime}(t)}{\phi(x(h(t)))} \leqslant \frac{x^{\prime}\left(h(t) h^{\prime}(t)\right.}{k \phi(x(h(t)))} . \tag{2.23}
\end{equation*}
$$

From (2.22) and (2.23), we have that

$$
\lim _{t \rightarrow \infty} \int_{t_{0}}^{t} q_{m}\left(t, \xi_{o}, \eta_{o}\right) d s \leqslant \quad \lim _{t \rightarrow \infty} \int_{x\left(h\left(t_{0}\right)\right)}^{x(h(t))} \quad \frac{d u}{k \phi(u)}<\infty .
$$

This contradicts condition (III $)_{1}$.
Suppose that ( $\mathrm{III}_{2}$ ) holds, then it follows from (2.23) that
$\lim _{m \rightarrow \infty} q\left(t, \xi_{o}, \eta_{o}\right) \leqslant \frac{x^{\prime}\left(h(t) h^{\prime}(t)\right.}{k \phi(x(h(t)))}$, namely $q\left(t_{0} \xi_{0}, \eta_{o}\right) \leqslant \frac{x^{\prime}\left(h(t) h^{\prime}(t)\right.}{k \phi(x(h(t)))}$.
Hence
$\lim _{t \rightarrow \infty} \int_{t_{0}}^{t} q\left(s, \xi_{0}, \eta_{0}\right) d s \leqslant \lim _{t \rightarrow \infty} \int_{x\left(h\left(t_{0}\right)\right)}^{x(h(t))} \frac{d u}{\phi(u)}<\infty$.
which is again a contradiction, and the proof of the theorem is complete.
Equation (1.1) is said to be strongly bounded oscillatory if all bounded solutions of Equation (1.4) for any $\lambda \varepsilon(0, \infty)$ are oscillatory.

From the proof of Theorem (2.4), we see that the following result holds. COROLLARY 2.5. As sune that the condition (III ) or (III ${ }_{2}$ ) holds, then Equation (1.1) is strongly bounded oscillatory.

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