

Improvement of Stability of Rotor System by Introducing a Hydraulic Damper into an Active Journal Bearing

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Modelling and analysis of a rotor—bearing system with a new type of active oil bearing is presented. The active bearing is supplied with a flexible sleeve whose deformation can be changed during operation of the rotor. The flexible sleeve is also a part of a hydraulic damper whose parameters can be controlled during operation as well.

Finite Element Method (FEM) and the Guyan condensation technique was utilised to create mathematical model of both, the rotor and the flexible sleeve. The hydrodynamic pressure distribution in the oil film, for the instantaneous position of the flexible sleeve and rotor, was approximated by Reynolds equation.

The mathematical model of motion of a rotor system with the described active bearing developed in this paper allowed the influence of the introduced hydraulic damper on stability of the rotor-bearing system to be investigated. Results of the computer simulation shows that within a large region of configuration parameters of the rotor bearing system, the self exciting vibration can be eliminated or greatly reduced during operation by properly controlled deformation of the flexible sleeve and optimal choice of the hydraulic damper parameters.

Keywords: Oil Bearings, Active Control, Rotor System, Stability, Modelling, Analysis

INTRODUCTION

A multi-bearing rotor system is statically indeterminate. Its dynamic behaviour depends on the relative positions of the bearings as well as properties of its subsystems. The relative positions of the bearings are usually referred to as the system configuration, or bearing

alignment (Parszewski Krodkiewski [1986]). Therefore, the dynamic properties of the multi-bearing rotor system are a function of the rotating speed and the system configuration. Consequently, the investigation of the dynamic behaviour of the rotor system can be undertaken in two domains. One is in the rotating speed domain, the other is in the configuration domain.

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Self-excited vibration (oil whip), is one of the main causes of severe rotor vibration. Great efforts have been made on the analysis and improvement of the stability of rotor systems. Lund [1974] calculated the threshold speed and the damped critical speeds of a flexible rotor in journal bearings using a linearised model of the bearing. Adams [1980] used a non-linear model to simulate the response of a multi-bearing rotor system to study the instability threshold speed. Holmes *et al.* [1978] studied the self-excited vibration of a two-rotor four-bearing system due to vertical misalignment of the bearings. The effect of bearing alignment on the stability boundary has been studied by Hori and Uematsu [1980]. Parszewski and Li [1989] presented a method to obtain the optimal configuration of a multi-bearing rotor system with respect to the forced vibration.

The system configuration may change during operation due to the thermal expansion, deformation of the supporting structures, etc. The change affects the dynamic response of the rotor system and in many cases its improvement is necessary during operation. Therefore, the active vibration control has been paid growing attention in the dynamic design of the rotating machinery. The design of active oil bearings could be divided into two categories. The first design is based on actively changing the properties of the oil film. It can be done either by pumping the oil or air with different pressure into the oil film (Goodwin *et al.* [1989]) or by altering the thickness and geometry of the oil film. Another design includes movable elements (ring or pads). The motion of these elements is activated either by hydraulic systems (Ulbrich and

Althaus [1989]) or by piezoelectric pushers (Adams and McCloskey [1990]). The change of the characteristics of the rotor system is achieved by applying control forces (normally by a feed back loop) to the rotor via the oil film. The active tilting pad bearing can be used in both of the above designs.

In this paper the influence of modification of the active bearing described by Krodiewski and Sun [1995] is investigated. To improve attenuation of vibrations of a rotor bearing system the active bearing was supplied with a hydraulic damper.

DESCRIPTION OF THE ACTIVE BEARING

The flexible sleeve (see Fig. 1) can be considered as a new feature of the proposed bearing. The oil film and the pressure chamber are separated by the flexible sealing. Equilibrium position of the flexible sleeve and the bearing journal is determined by the bearing load and the pressure p_{in} which can be controlled during operation. Vibration of the flexible sleeve around its equilibrium position causes flow of the chamber oil through the hydraulic damper. Parameters of this damper can be also varied during operation. If the inlet pressure p_{in} and parameters of the damper are function of the system configuration or its rotating speed, the bearing can be considered as adaptive one. Search for their optimal values is the main purpose of this paper.

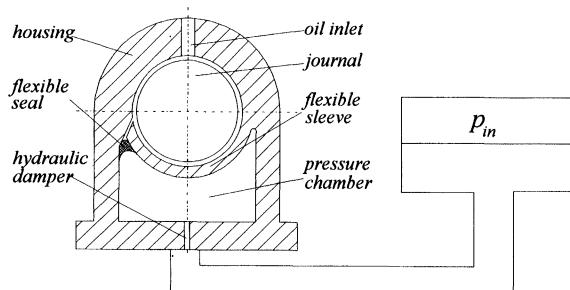


FIGURE 1 Schematic of the active bearing.

MODELLING OF THE ROTOR-ACTIVE BEARING SYSTEMS

Modelling of The Rotor And The Flexible Sleeve

The physical model of a three-bearing rotor system is shown in Fig. 2. The active bearing is located in the middle of the rotor. The rotor was approximated by the two degrees of freedom system with concentrated mass at the location of the active bearing. Its instantaneous position O , with respect to the absolute sys-

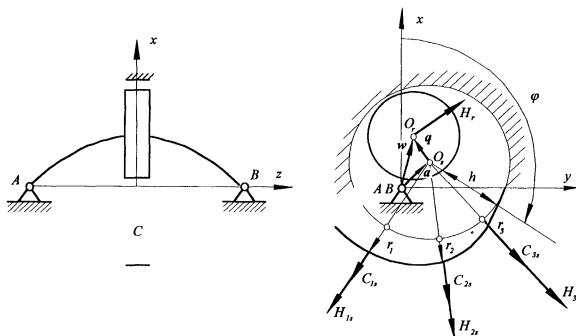


FIGURE 2 The physical model of the rotor.

tem of coordinates xyz is determined by vector w . Configuration of the rotor bearing system (the relative position of its bearings) is defined by vector a . The mathematical model of the rotor was adopted in the following form:

$$\mathbf{M}_r \ddot{\mathbf{w}} + \mathbf{K}_r \mathbf{w} = \mathbf{H}_r(\mathbf{q}, \dot{\mathbf{q}}, \mathbf{r}, \dot{\mathbf{r}}) + \mathbf{Q}_r + \mathbf{F}_r. \quad (1)$$

The Guyan condensation technique was used to produce the inertia \mathbf{M}_r and stiffness \mathbf{K}_r matrices. The hydrodynamic force that acts on the journal is denoted by \mathbf{H}_r , while \mathbf{Q}_r and \mathbf{F}_r stand for vector of the gravity forces and the external excitation respectively. Since:

$$\mathbf{w} = \mathbf{a} + \mathbf{q} \quad (2)$$

where \mathbf{q} represents the relative motion of the journal with respect to the bearing, the equation of motion of the rotor takes the following form:

$$\mathbf{M}_r \ddot{\mathbf{q}} + \mathbf{K}_r \mathbf{q} = \mathbf{H}_r(\mathbf{q}, \dot{\mathbf{q}}, \mathbf{r}, \dot{\mathbf{r}}) - \mathbf{K}_r \mathbf{a} + \mathbf{Q}_r + \mathbf{F}_r. \quad (3)$$

The flexible sleeve was modelled as a curved beam and its dynamic properties (the mass matrix \mathbf{M}_s and the stiffness matrix \mathbf{K}_s) were determined along coordinates \mathbf{r} (see Fig. 2.) by means of FEM and the Guyan condensation technique. Its motion can be approximated by the equation below.

$$\mathbf{M}_s \ddot{\mathbf{r}} + \mathbf{K}_s \mathbf{r} = \mathbf{H}_s(\mathbf{q}, \dot{\mathbf{q}}, \mathbf{r}, \dot{\mathbf{r}}) + \mathbf{C}_s. \quad (4)$$

In the equation (4) \mathbf{H}_s stands for vector of the hydrodynamic force caused by oil film pressure p and \mathbf{C}_s stands for vector of the hydrodynamic forces caused by pressure in the control chamber p_c . Both of them were computed from the instantaneous pressure distribution by means of the virtual work principle. Distribution of the oil film pressure p was modelled by means of the Reynolds equations.

$$\frac{1}{R^2} \frac{\partial}{\partial \varphi} \left(\frac{h^3}{\eta} \frac{\partial p}{\partial \varphi} \right) + \frac{\partial}{\partial z} \left(\frac{h^3}{\eta} \frac{\partial p}{\partial z} \right) = 6\Omega \frac{\partial h}{\partial \varphi} + 12 \frac{\partial h}{\partial t}. \quad (5)$$

More details on modelling of the rotor and the flexible sleeve can be found in work by Krodkiewski and Sun [1995]. The chamber pressure p_c depends upon the inlet pressure p_{in} (see Fig. 1), parameters of the hydraulic damper and motion of the flexible sleeve \mathbf{r} .

Modelling of The Chamber Pressure

If the initial volume of the chamber is denoted by V_o (see Fig. 3) and the fluctuation of this volume due to motion of the flexible sleeve by ΔV , the instantaneous volume can be expressed as follows:

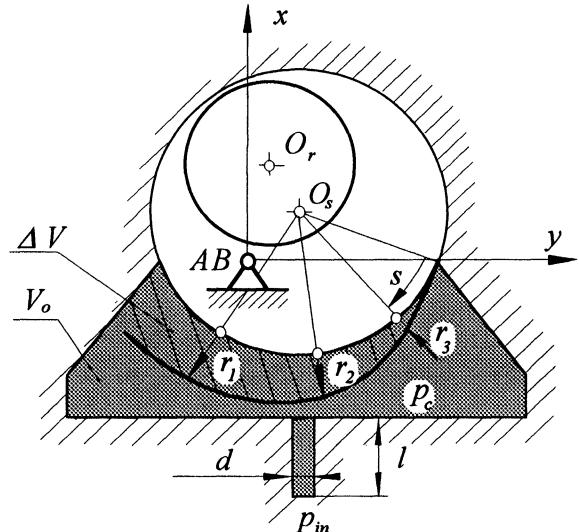


FIGURE 3 The physical model of the chamber.

$$V = V_o - \Delta V, \quad (6)$$

where

$$\Delta V = W_s \int_0^{L_s} r(s) ds. \quad (7)$$

Length and width of the flexible sleeve in the last equation are denoted by L_s and W_s respectively.

Rate of change of volume of oil V_c , due to its compressibility, can be determined by the following equation:

$$\frac{dV_c}{dt} = -\frac{V}{B} \frac{dp_c}{dt}, \quad (8)$$

where B is the bulk modulus of oil.

To fulfil the continuity of flow requirements, the rate of change of the oil flow through the damper has to obey the following formula:

$$Q = \frac{d\Delta V}{dt} + \frac{dV_c}{dt}. \quad (9)$$

On the other hand, the rate of flow through the damper is determined by its diameter d , length l and the inlet pressure p_{in} and chamber p_c pressure

$$Q = \frac{\pi d^4}{128\eta l} (p_c - p_{in}). \quad (10)$$

Equations (10), (9) and (8) allow to develop the relationship between motion of the flexible sleeve and the chamber pressure.

$$\frac{dp_c}{dt} = k_1 \left(k_2 (p_{in} - p_c) + \frac{d\Delta V}{dt} \right), \quad (11)$$

where:

$$k_1 = \frac{B}{V} \quad k_2 = \frac{\pi d^4}{128\eta l}. \quad (12)$$

The equation (11) governs correctly the chamber pressure if its magnitude is greater than the cavitation threshold p_{cav} . If during integration of the system equations of motion the chamber pressure dropped below this threshold value, the adopted for computation pressure was equal to p_{cav} (see Fig. 4). The chamber pressure was kept equal to p_{cav} as long as the expression (13) was greater than zero.

$$V_B = (W_s \int_0^{L_s} r(s, t) ds - W_s \int_0^{L_s} r(s, t_i) ds) - \\ k_2 (p_{in} + p_{cav}) (t - t_i) > 0. \quad (13)$$

The above expression represents part of volume of chamber filled by vapour of oil due to the cavitation process.

NUMERICAL SOLUTION OF THE MATHEMATICAL MODEL AND ITS PARAMETERS

The equations (3) and (4) form simultaneous set of five differential equations

$$\begin{cases} M_r \ddot{q} + K_r q = H_r(q, \dot{q}, r, \dot{r}) - K_r a + Q_r + F_r \\ M_s \ddot{r} + K_s r = H_s(q, \dot{q}, r, \dot{r}) + C_s \end{cases} \quad (14)$$

$$q = [x, y]^T \quad Hr = [r_1, r_2, r_3]^T$$

These equations were solved by means of the Runge-Kutta method. The necessary pressure distributions

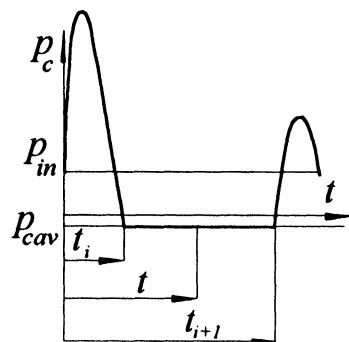


FIGURE 4 Chamber pressure as a function of time due to the cavitation process.

TABLE I Parameters of the Rotor Bearing System.

m	ω_I	S	D	L/D	c
11.24 kg	8.95 Hz	1.7 m	0.05 m	0.8	0.0003 m
β	ψ	η	Ω	$a_x a_y$	p_{in}
105°	160°	0.04 Pa s	3000 RPM	-1.85, 0 mm	0.1 MPa
p_{cav}	b	k_I	k_2		
-0.03 MPa	$4 \cdot 10^{-3}$ m	10^{11} Pa m ⁻³	$3 \cdot 10^{11}$ N ⁻¹ m ⁵ s ⁻¹		

inside the oil film as well as pressure in the chamber, at each step of integration, were computed from equations (5) and (11). Integration of the pressure along surface of the flexible sleeve and journal yields the hydrodynamic forces H_r , H_s and C_s .

Parameters of the physical model of the rotor bearing system adopted for the numerical analysis are collected in Table I.

The listed in Table I parameters correspond to an existing rotor—bearing system. The following numerical simulation of motion of the rotor system was aimed to obtain optimal parameters of the hydraulic damper.

RESULTS OF THE NUMERICAL COMPUTATION

Influence of the following parameters on improvement of stability of the rotor system response was investigated:

1. bulk modulus B (k_I)
2. length l and diameter d of the damper capillary (coefficient k_2)
3. thickness of the flexible sleeve b .

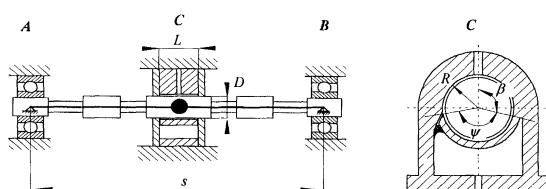


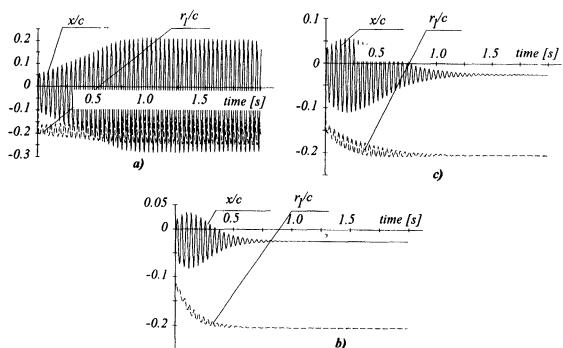
FIGURE 5 Schematic of the rotor system.

The following diagrams present displacement of the journal along x direction and displacement of the flexible sleeve along coordinate r_I (see Fig. 2) in dimensionless form as a function of time.

Influence of Bulk Modulus (k_I)

The bulk modulus of the oil depends on quantity of the air mixed with the oil and its magnitude is usually less than $2 \cdot 10^9$ Pa. With increasing amount of air, magnitude of the bulk becomes smaller and smaller amount of oil flows through the damper capillary. Therefore smaller value of the bulk modulus results in poorer improvement of stability of the system equilibrium position. Some results of the numerical simulation, for three different magnitudes of the bulk modulus ($B = 10^6, 10^7, 10^8$ Pa) and the chamber volume $V = 10^{-4}$ m³, are shown in Fig. 6.

For small magnitude of the bulk modulus B (Fig. 6.a), the journal and the flexible sleeve move along large limit cycles. Increment in B (see Fig. 6.b) re-

FIGURE 6 Influence of the bulk modulus; a) $k_I = 10^{10}$ [Nm⁻⁵], b) $k_I = 10^{11}$ [Nm⁻⁵], c) $k_I = 10^{12}$ [Nm⁻⁵].

sults in stabilisation of the equilibrium position. Further increment of value of the bulk modulus (see Fig. 6.c) insignificantly improves stability of the equilibrium position.

Influence of Damper Parameters

The hydraulic damper properties are determined by its diameter and length (see formula (12)). The computer simulation indicates that for any working parameters of the rotor system there exists an optimal set of the damper parameters determined by the coefficient k_2 . Sample of this computation is shown in Fig. 7.

If the bearing is not supplied with the damper (Fig. 7a), its journal as well as its sleeve perform large self exciting vibrations along a limit cycle around the unstable equilibrium position. For $k_2 = 1 \cdot 10^{-12} [N^{-1} m^5 s^{-1}]$ (Fig. 7b) the equilibrium position is still unstable but the dimensionless peak to peak amplitude of the journal limit cycle is as small as 0.09. Case $k_2 = 3 \cdot 10^{-11} [N^{-1} m^5 s^{-1}]$ (Fig. 7c) seems to be optimal. As can be seen, in this case the equilibrium position is stable. Further increment in the coefficient k_2 to $3 \cdot 10^{-10} [N^{-1} m^5 s^{-1}]$ (see Fig. 7d) results in the unstable equilibrium position.

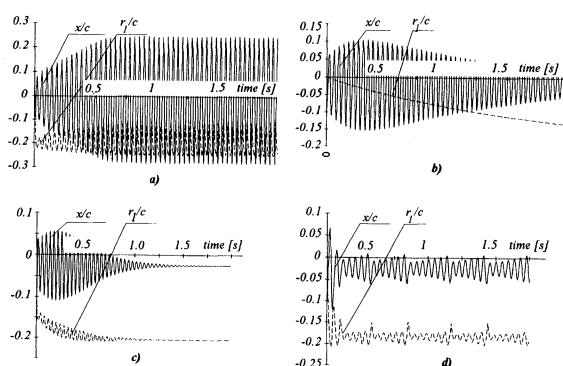


FIGURE 7 Influence of the throttling parameter k_2 ; a) $k_2 = \infty$ (p_c = constant), b) $k_2 = 1 \cdot 10^{-12} [N^{-1} m^5 s^{-1}]$, c) $k_2 = 3 \cdot 10^{-11} [N^{-1} m^5 s^{-1}]$, d) $k_2 = 3 \cdot 10^{-10} [N^{-1} m^5 s^{-1}]$.

Influence of Thickness of The Flexible Sleeve

The described above simulation was carried out for different thickness of the flexible sleeve b and for different magnitudes of coefficient k_2 . For any combination of these parameters the peak to peak amplitude of limit cycle A was recorded. Its dimensionless magnitudes A/c as a function of the dimensionless thickness b/R and the coefficient k_2 are presented in Fig. 8.

As one can see from Figure 8, thinner sleeve results in wider range of stability.

CONCLUSIONS

The presented numerical computation indicates that stability of the equilibrium position of rotor—oil bearing system can be significantly improved by means of the proposed active bearing with the hydraulic damper.

Since the optimal parameters depend on the system current configuration and rotating speed, the diameter of the capillary as well as the inlet pressure has to be controlled during operation.

To reduce influence of the ever existing compressibility of oil which is reflect by the bulk modulus, volume of the chamber should be as small as possible.

To reduce the influence of cavitation, the inlet pressure should be possibly high. Hence in case of heavily loaded bearings the proposed hydraulic damper is more effective. In case of lightly loaded bearings a compromise should be sought between the thickness of the sleeve and the inlet pressure.

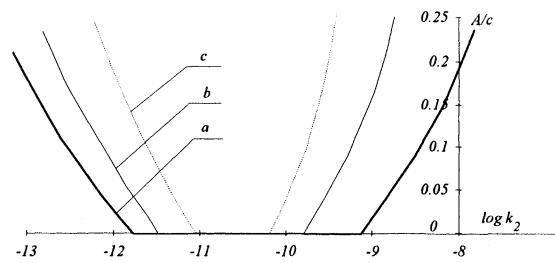


FIGURE 8 The peak to peak amplitude of limit cycle as a function of the throttling coefficient; a) for $b/R = 0.12$, b) for $b/R = 0.16$, c) for $b/R = 0.2$.

Optimisation of the forced vibrations (in case the equilibrium position is stable) as well as the experimental investigations are considered as the next step towards implementation of this type of active bearing to control vibrations of rotor—bearing system.

NOMENCLATURE

A	- the peak to peak amplitude of limit cycle	L_s	- length of the flexible sleeve, m
$\mathbf{a} = [a_x, a_y]$	- vector of configuration	l	- length of the damper's capillary, m
a_x, a_y	- configuration parameters, m	\mathbf{M}_r	- matrix of inertia of the rotor, kg
B	- bulk modulus, Pa	m	- weight of rotor, kg
b	- thickness of the flexible sleeve	p_{ni}	- input pressure, Pa
$\mathbf{C}_s = [C_{1s}, C_{2s}, C_{3s}]$	- vector of the hydrodynamic forces acting on the flexible sleeve due to the chamber pressure, N	p	- the oil film pressure, Pa
c	- bearing radial clearance, m	P_{cav}	- cavitation threshold, Pa
D	- diameter of journal, m	p_c	- the chamber pressure, Pa
d	- diameter of the damper's capillary, m	Q	- rate of flow, m^3/s
\mathbf{F}_r	- vector of the external excitation acting on journal along coordinates xy , N	\mathbf{Q}_r	- vector of the static load acting on journal along coordinates xy , N
$\mathbf{H}_s = [H_{1s}, H_{2s}, H_{3s}]$	- vector the hydrodynamic forces acting on the flexible sleeve due to the oil film pressure, N	$\mathbf{q} = [q_x, q_y]$	- vector of the relative coordinates of journal with respect to the centre of bearing
$\mathbf{H}_r = [H_{rx}, H_{ry}]$	- vector of the hydrodynamic forces acting on the journal due to the oil film pressure, N	$\dot{\mathbf{q}}$	- vector of the relative velocities of journal with respect to bearing, m/s
$h(\varphi, t)$	- instantaneous thickness of the oil film, m	$\mathbf{r} = [r_1, r_2, r_3]$	- vector of the absolute position of the flexible sleeve, m
\mathbf{K}_r	- matrix of stiffness of the rotor, N/m	R	- radius of journal, m
$k_I = B/V$	- coefficient of the oil compressibility, Nm^{-5}	s	- length of the flexible sleeve associated with coordinate φ , m
$k_2 = \pi d^4/128\eta l$	- throttling coefficient, $\text{N}^{-1}\text{m}^5\text{s}^{-1}$	S	- span of rotor, m
L	- length of the journal bearing, m	t	- time, s
		V	- instantaneous volume of the chamber, m^3
		V_o	- initial volume of the chamber, m^3
		ΔV	- fluctuation of the chamber volume due to motion of the flexible sleeve, m^3
		dV/dt	- rate of change of volume due to the oil compressibility, m^3/s
		W_s	- width of the flexible sleeve, m
		$\mathbf{w} = [x, y]$	- vector of the absolute coordinates of journal

xyz	- the absolute system of coordinates	Analyses Of Multi-Bearing Rotor System, <i>Journal of Sound and Vibration</i> , Vol. 164(2), pp. 267–280.
x, y	- coordinates of the journal, m	Goodwin M. J., Boroomand T. and Hooke C. J., 1989. Variable Impedance Hydrodynamic Journal Bearings for Controlling Flexible Rotor Vibrations, <i>The 1989 ASME Design Technology Conferences, Mechanical Vibration and Noise</i> , Montreal, pp. 261–267.
β, ψ	- angles defined in Fig. 5, rad	Guyan, R. J., 1965. Reduction of Stiffness and Mass Matrices, <i>American Institute of Aeronautics and Astronautics Journal</i> , Vol. 3, No. 2, pp. 380–385.
φ	- angular coordinates defined in Fig. 2, rad	Holmes A. G., Ettles C. M. McC. and Mayes I. W., 1978. The Dynamics of Multi-Rotor Systems Supported on Oil Film Bearings, <i>Journal of Mechanical Design, Transaction of the ASME</i> , Vol. 100, pp. 156–164.
ω_1	- first natural frequency of the rotor supported on the ball bearings only.	Hori Y. and Uematsu R, 1980. Influence of Misalignment of Support Journal Bearings on Stability of a Multi-Rotor System, <i>Tribology International</i> , pp. 249–252.
Ω	- angular velocity of the rotor, rad/s	Krodkiewski J. M., Sun L., 1995. Stability Control of Rotor-Bearing System by An Active Journal Bearing, <i>Proc. of the International Conference on Vibration and Noise</i> , Venice, pp. 217–225.

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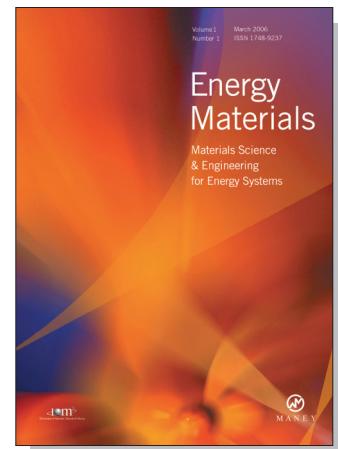
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