

## Research Article

# Numerical Investigation of Pressure Fluctuation Characteristics in a Centrifugal Pump with Variable Axial Clearance

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Clearance flows in the sidewall gaps of centrifugal pumps are unsteady as well as main flows in the volute casing and impeller, which may cause vibration and noise, and the corresponding pressure fluctuations are related to the axial clearance size. In this paper, unsteady numerical simulations were conducted to predict the unsteady flows within the entire flow passage of a centrifugal pump operating in the design condition. Pressure fluctuation characteristics in the volute casing, impeller, and sidewall gaps were investigated with three axial clearance sizes. Results show that an axial clearance variation affects the pressure fluctuation characteristics in each flow domain by different degree. The greatest pressure fluctuation occurs at the blade pressure surface and is almost not influenced by the axial clearance variation which has a certainly effect on the pressure fluctuation characteristics around the tongue. The maximum pressure fluctuation amplitude in the sidewall gaps is larger than that in the volute casing, and different spectrum characteristics show up in the three models due to the interaction between the clearance flow and the main flow as well as the rotor-stator interaction. Therefore, clearance flow should be taken into consideration in the hydraulic design of centrifugal pumps.

## 1. Introduction

Unsteady flows in centrifugal pumps are always paid attention to in engineering, which can cause vibration, noise, and even fatigue failure. Pressure fluctuation is one of the most common phenomena of unsteady flows. Considerable investigations [1–12] have been performed on pressure fluctuation characteristics in volute casings and impellers of centrifugal pumps, concluding that blade-tongue interaction plays a dominant role in causing fluctuating main flows, and the influencing factors including flow rate and impeller-tongue gap size have been analyzed. But flows in sidewall gaps, which refer to the space between impeller shrouds and casing covers, are usually ignored. However, Uy and Brennen [13] found that the leakage path had considerable influence on the rotor dynamic force when they did some experiments for centrifugal pumps whose shrouds were in different shapes. Increasing scholars are paying attention to the clearance flows in centrifugal pumps in recent years. Spence and Amaral-Teixeira [14] carried out an unsteady numerical simulation

for a double suction pump and thought that the deviation in pressure fluctuations between the numerical and experimental results was due to the simplified front shroud in the simulation. Park and Morrison [15] probed into the interaction between the back clearance flow and impeller passage flow by calculating the pressure fluctuations in a centrifugal pump with a semiopen impeller. Pei et al. [16–19] claimed the unstable flow in the sidewall gaps could not be ignored for the sake of an accurate calculation and analysis of the pressure fluctuation intensity and turbulence intensity distribution in the entire flow passage of a single-blade centrifugal pump. In addition, the size of the leakage path has certain influence on the main flow field. Radial clearance sizes, usually referring to wear ring sizes, were studied in a centrifugal pump by Liu et al. [20] and they found that the velocity and pressure distribution in the front wear ring was related to its radial dimension. Many scholars [21–24] also experimentally and theoretically discovered that changing the distance between a rotating wall and a stationary wall would make the flow structure between the two walls show different characters. However,



TABLE 1: Basic parameters of the pump.

Parameter	Value
Impeller inlet diameter	$D_0 = 85 \text{ mm}$
Impeller outlet diameter	$D_2 = 310 \text{ mm}$
Blade number	$Z_1 = 3$
Blade inlet angle	$\beta_1 = 44^\circ$
Blade outlet angle	$\beta_2 = 29^\circ$
Blade outlet width	$B = 79.2 \text{ mm}$
Back-blade number	$Z_2 = 9$
Back-blade width (axial direction)	$b = 1.9 \text{ mm}$
Back-blade inlet angle	$\beta_{b1} = 19^\circ$
Back-blade outlet angle	$\beta_{b2} = 26^\circ$
Volute casing width at impeller outlet	$B_v = 110 \text{ mm}$
Blade-tongue gap	$\Delta d = 0.29 D_2$
Specific speed*	$n_s = 165$
Design flow rate	$Q_0 = 0.08 \text{ m}^3/\text{s}$
Design head	$H_0 = 11.3 \text{ m}$

\*Specific speed is defined as  $n_s = 3.65 n Q_0^{1/2} / H_0^{3/4}$ , where  $n$  is the rated rotational speed, r/min. The unit of  $n_s$  is usually ignored.

TABLE 2: Research models.

Model	$c_f$ (mm)	$c_b$ (mm)
M1	0.22	0.78
M2	0.42	0.58
M3	0.62	0.38
M4	None	None

1.63 million or more elements had little difference. Thus, the grid scheme with 1.63 million elements was chosen to balance the calculation accuracy and cost. Grids for the four models have similar topological structure, with almost the same amount of elements in the volute casing, impeller, and suction pipe; the element amount of sidewall gaps proportionally changes with the variation of axial clearance sizes. Figure 2 shows the grid of back sidewall gap, back-blade domain, and front gap of Model M2, with partial enlarged views.

**2.2. Numerical Method and Boundary Condition.** With commercial CFD software package ANSYS CFX14.5, we used Reynolds-Averaged Navier-Stokes (RANS) equations to calculate the incompressible turbulent flow in the centrifugal pump. For better prediction of flow separation and rotational flows, SST  $k-\omega$  turbulence model with automatic wall function was applied to close the governing equations. After checking timestep independence, the timestep was set as 1/150 of one revolution; that is, the impeller rotates  $2.4^\circ$  after each timestep. During each timestep, the code keeps running until the iteration residual drops to 0.0001.

A total pressure was specified at the suction pipe inlet and a mass flow rate was specified at the volute casing outlet. The impeller and back-blade domain had the same rotational speed, 1000 r/min, and other domains were all set to be stationary. Steady simulations were first carried out

and then the results were used as the initial flow fields for unsteady simulations. In steady simulations, the interface between rotor and stator adopted Frozen Rotor Model, while in unsteady simulations it used Transient Frozen Rotor Model. The Frozen Rotor Model treats the flow from one component to the next by changing the frame of reference and maintaining the relative position of the components, while the Transient Rotor-Stator Model, which considers all interaction effects between components that are in relative motion to each other, updates the transient relative motion between the components on each side of the interface every timestep [25]. Walls were all set as no-slip walls. In front and back sidewall gaps, the walls contacting front or back shrouds were set as rotational walls that had the same speed of impeller.

Unsteady numerical simulations were processed with design flow rate  $Q_0$  for models M1, M2, and M3. When the calculations became stable, we continued the calculation for another six periods. The data of the last six periods are used to obtain the pump performance. Pressure variations were also recorded at several locations within the entire flow passage during the last six periods. Pressure fluctuations in the volute casing were studied via Points VC0~VC8 and in the impeller via Points BP1~BP5 located on the pressure surface and Points BS1~BS5 located on the suction surface. These points are all on the middle section of the volute casing as shown in Figure 3(a). In the front and back sidewall gaps, pressure fluctuations at Points FC1~FC4 and BC1~BC4 as shown in Figure 3(b) were analyzed, respectively. The axial and radial positions of these points are illustrated in Figure 1 (0.15 mm away from the casing cover).

**2.3. Validation of Simulation Results.** Validation of the numerical method is conducted by comparing with the corresponding model test results from Zhang's master thesis [26]. The model test was conducted in a closed hydraulic machinery test rig whose essential measuring equipment has been introduced in our previous work [27]. Flow rate was measured by an intelligent electromagnetic flowmeter with an uncertainty of  $\pm 0.2\%$ . Head was measured by an intelligent differential pressure transmitter with an uncertainty of  $\pm 0.1\%$ . Torque was measured by intelligent torque and rotational speed transducer with an uncertainty of  $\pm 0.1\%$ . The rotational speed of the test pump was 1000 r/min as in the simulation. The efficiency was calculated with the time-averaged head and axial power. Meanwhile, pressure data were recorded through CGYL-201 high-precision hydraulic pressure fluctuation transducers which were flush-mounted on the casing wall at Point VC5. The uncertainty of the pressure transducer is  $\pm 0.5\%$  and the dynamic response range is from 0 to 1500 Hz. The experimental sampling frequency was 4096 Hz and the sampling duration was 20 s for each condition.

Experimental and numerical results of the efficiency and head in the design condition are given in Table 3. When considering sidewall gaps, that is, for M1, M2, and M3, the relative errors are all within 5%, and the error rises as  $c_f$  increases. This may be related to the complicated flow states in the front sidewall gaps. Experimental and numerical results

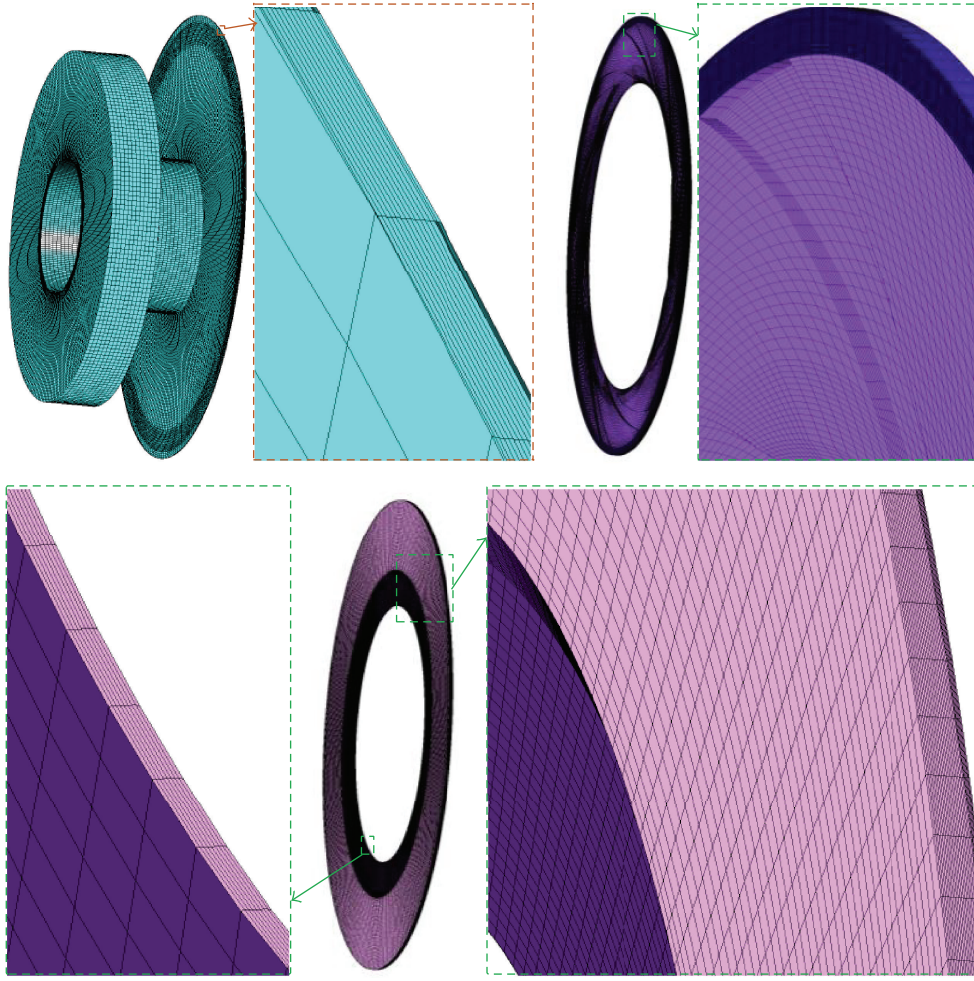


FIGURE 2: Grid of the sidewall gaps.

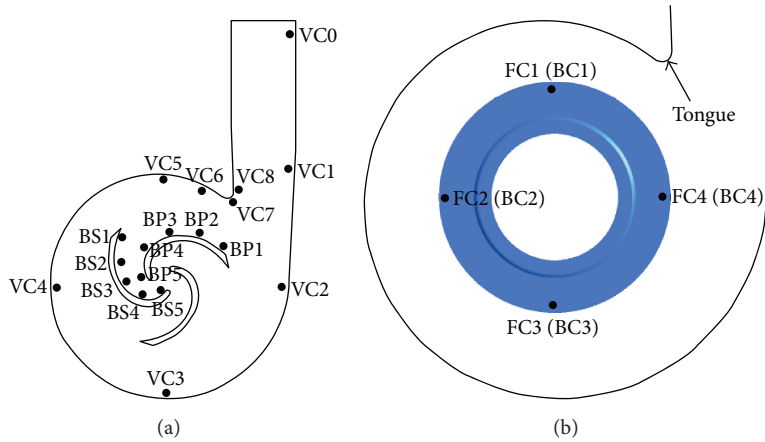


FIGURE 3: Distribution of pressure recording points.

exhibit a similar trend: the efficiency and head both drop as  $c_f$  increases. In our previous study [27], it has been found that as  $c_f$  increases, the interaction between the clearance flow and the main flow becomes stronger, leading to a decline of the simulation accuracy. The error would be larger if the sidewall

gaps were ignored in the simulation. The relative error of the efficiency of M4 is over 10% since it does not take volumetric losses and mechanical losses into account and moreover it is hard to reflect the hydraulic losses caused by the interaction between the clearance flow and the main flow.

TABLE 3: Comparison of numerical results with experiment results of the pump performance.

	Efficiency			Head		
	Experiment (%)	Simulation (%)	Error* (%)	Experiment (m)	Simulation (m)	Error* (%)
M1	82.8	83.0	0.24	11.45	11.36	-0.79
M2	81.0	82.6	1.98	11.30	11.33	0.44
M3	78.3	81.8	4.47	10.88	11.23	3.22
M4	78.3	86.7	10.73	10.88	10.34	-4.96

\*Error = |simulation result – experiment result|/(experiment result) 100%.

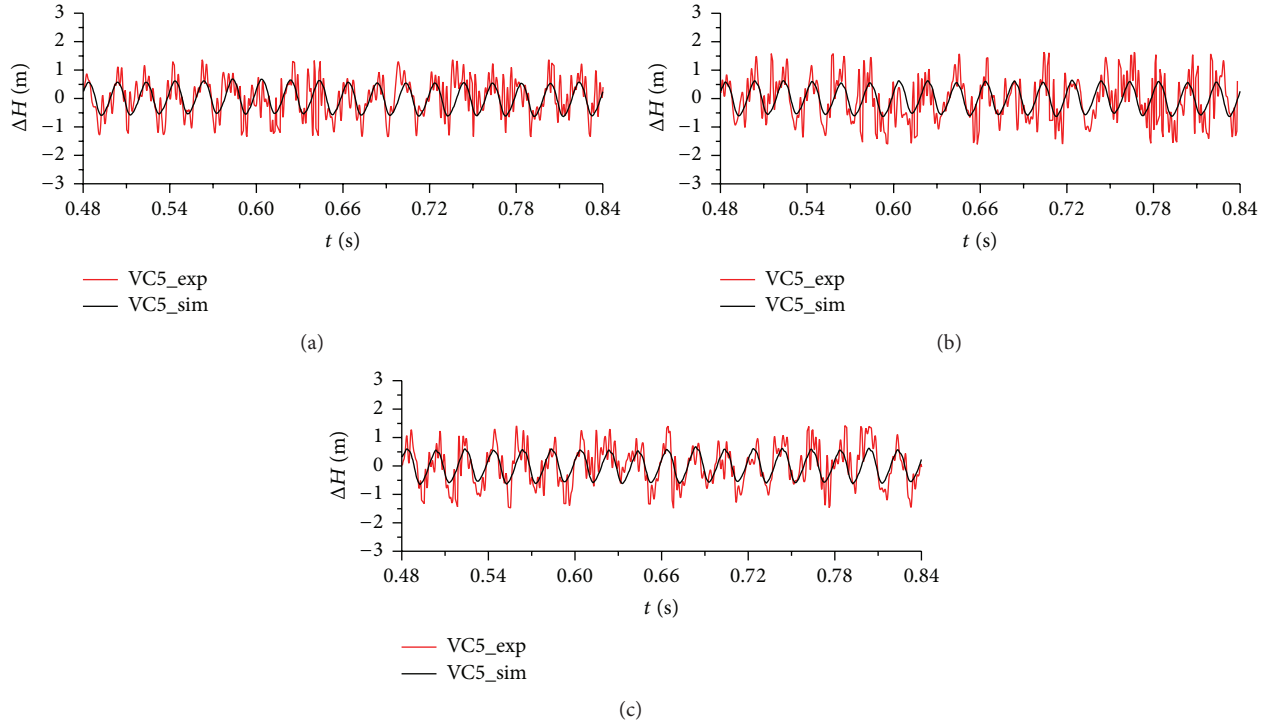


FIGURE 4: Comparison of the experimental and numerical pressure variations in three models: (a) M1, (b) M2, and (c) M3.

Pressure fluctuation data obtained by the numerical simulations and model test are compared for M1, M2 and M3. Relative pressure was got by subtracting the average pressure of six periods. We take 97% confidence interval for both numerical and experimental data to avoid unreliable data. Pressure variations with time are displayed in Figure 4. The numerical results show similar patterns as the experimental results except that the predicted peak-peak amplitudes are slightly smaller than the experiment results and the numerical simulations are not able to predict the small high-frequency undulations. Fast Fourier Transform (FFT) was conducted for the pressure data and the frequency spectrums are presented in Figure 5. Numerical results are in high agreement with experimental results at dominant frequency, but the numerical secondary-frequency amplitudes are smaller than the experimental ones especially at low frequencies.

The comparison shown in Figures 4 and 5 in both external performance and pressure fluctuation data demonstrates that the numerical method used in this paper meets the

accuracy requirement. Then this numerical method is applied to investigate the pressure fluctuation characteristics in the centrifugal pump with different axial clearance sizes.

### 3. Results and Discussion

#### 3.1. Pressure Fluctuations in the Main Flow Region

**3.1.1. Volute Casing.** Unsteady flow in volute casings is always a research hotspot for centrifugal pumps. For this pump, the strongest pressure fluctuation in the volute casing occurs around the tongue, which is consistent with the conclusions of many previous studies. From Figure 6 (the vertical coordinate denotes the ratio of the pressure fluctuation peak-peak amplitude and the corresponding pump head), the peak-peak amplitude first decreases along the flow direction (counter-clockwise from Point VC7 to Point VC3) and then rises in the diffusing section (from Point VC2 to Point VC0). The blade-tongue interaction affects the flow field in the tongue region

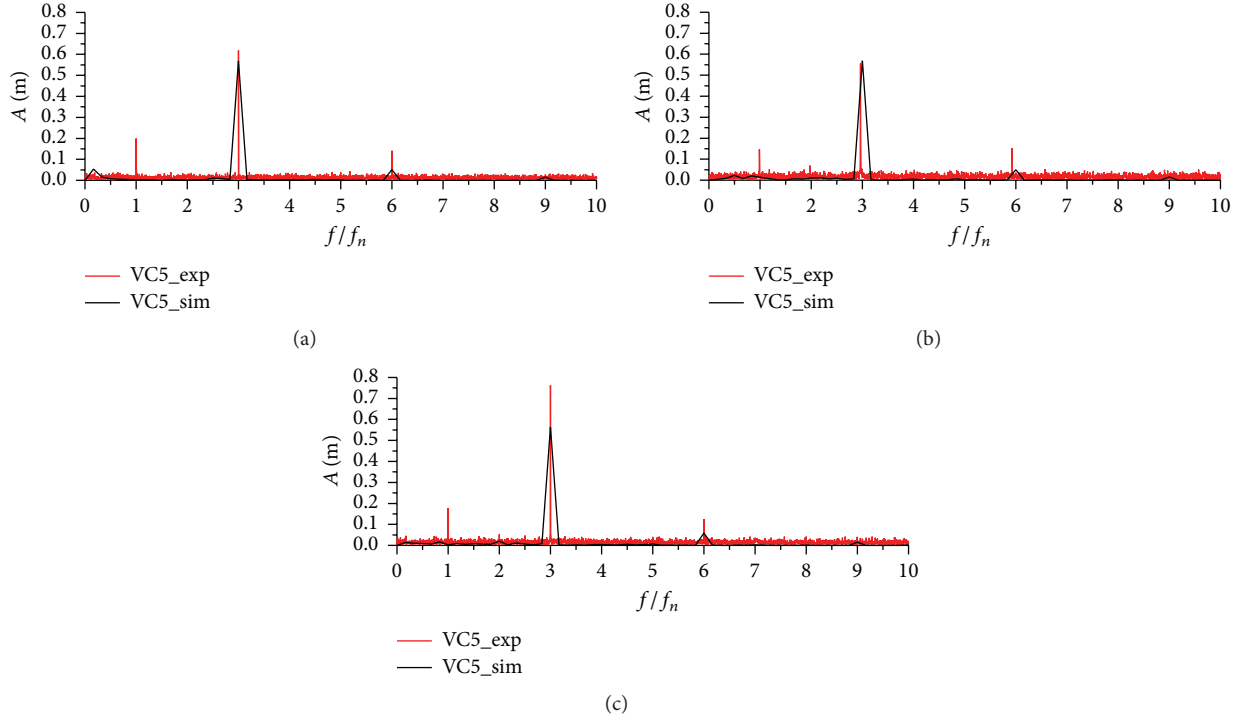


FIGURE 5: Comparison of the experimental and numerical frequency spectrum in three models: (a) M1, (b) M2, and (c) M3.  $f_n$  denotes rotating frequency.

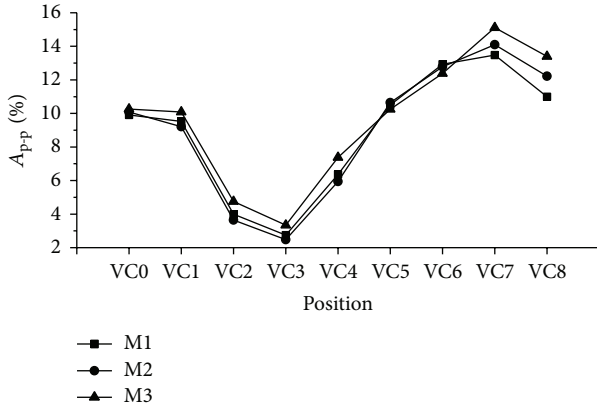


FIGURE 6: Peak-peak amplitudes at the points in the volute casing.

and causes stronger pressure fluctuations around the tongue; as the pressure wave propagates away from the tongue, the pressure fluctuation amplitude decreases. An axial clearance variation has little influence for most points in the volute casing; as  $c_f$  increases, the peak-peak amplitude only has an evident rise at the points near the tongue, that is, Point VC7 and Point VC8.

Via FFT algorithm, we find that the dominant frequency at every point in the volute casing is all triple rotating frequency, that is, blade passing frequency, which again proves the effect of the blade-tongue interaction. Figure 7 compares the frequency spectrum at Point VC7 for the three models. The amplitude at the dominant frequency is nearly not

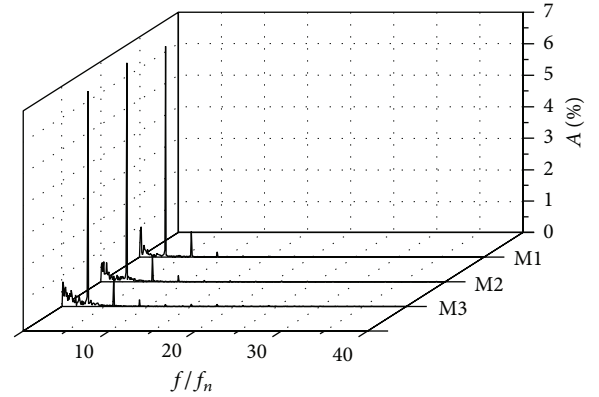


FIGURE 7: Frequency spectrum of the pressure fluctuations at Point VC7.

influenced by the axial clearance variation, only with a slightly higher value in M2. There are rich low-frequency components at Point VC7 which may be related to the complicated vortex structures around the tongue, and a larger  $c_f$  leads to richer low-frequency components. The vorticity distribution on the middle section of the pump exhibited in Figure 8 confirms the above conclusion. The increasing  $c_f$  results in more nonuniform vorticity distribution. It is related to the interference of the front clearance flow to the main flow, which is particularly severe near the tongue due to the rotor-stator interaction. A larger  $c_f$  brings more leakage fluid and accordingly induces severer disturbance around the tongue [27].

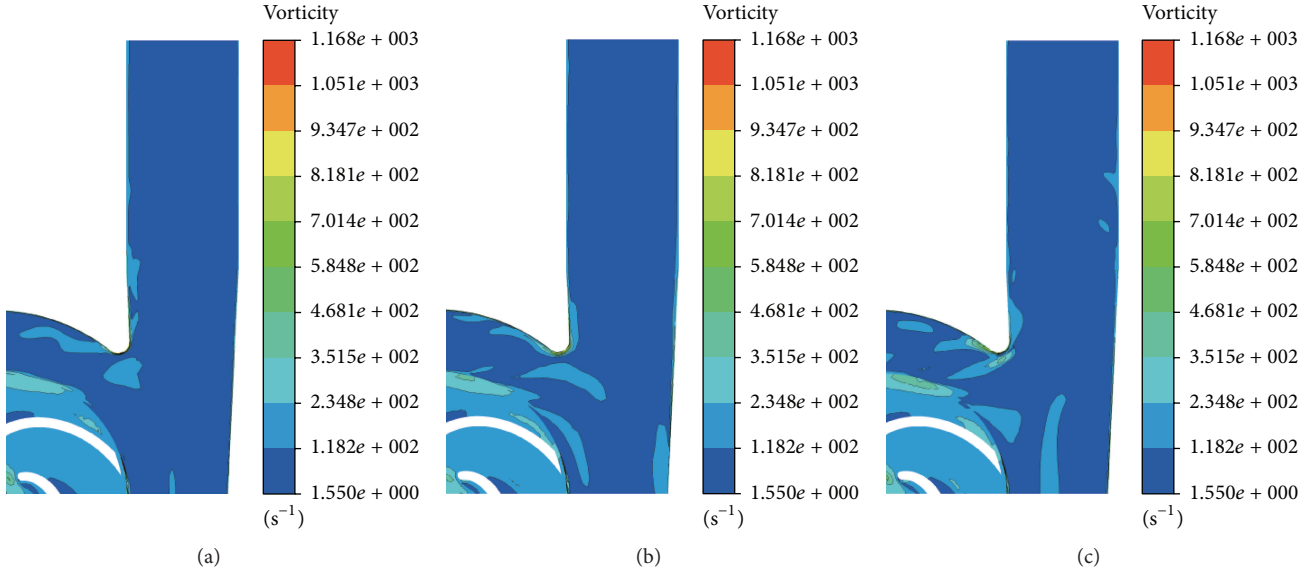


FIGURE 8: Vorticity distribution on the middle section: (a) M1, (b) M2, and (c) M3.

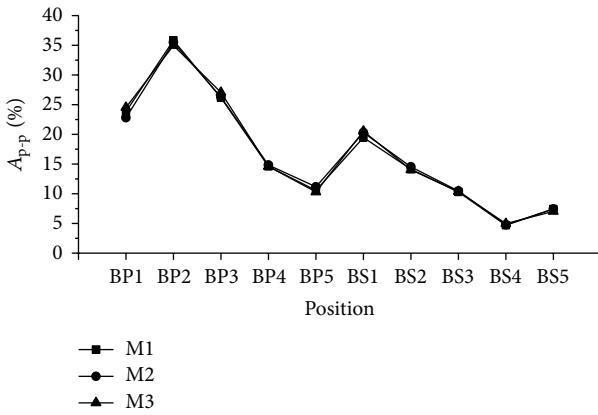


FIGURE 9: Peak-peak amplitudes at the points in impeller.

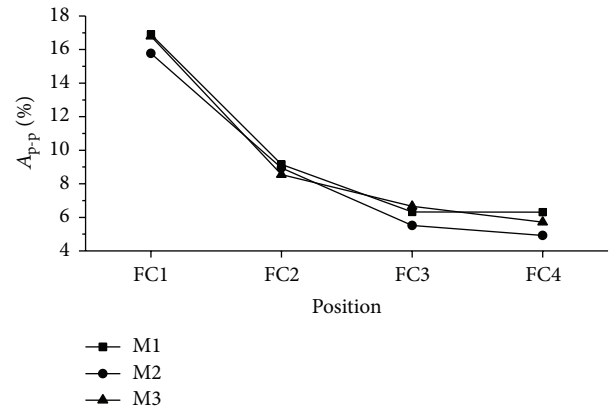


FIGURE 10: Peak-peak amplitude at the points in the front gap.

**3.1.2. Impeller.** Flow states in blade channels have significant influence on the pump performance. It is not easy to experimentally measure the flow field in rotating components, while CFD can be a good tool. Figure 9 shows the peak-peak amplitudes of the pressure fluctuations on the blade surfaces, which are larger on the pressure surface than on the suction surface. The effect of rotor-stator interaction is greater at the impeller outlet; thus the pressure fluctuation peak-peak amplitude rises from the impeller inlet to outlet. In addition, axial clearance variation has nearly no influence on the pressure fluctuations in the impeller. Based on our previous research [27], the clearance flows directly interact with the flows in the volute casing (near the impeller outlet) and in the suction pipe (just in front of the impeller inlet) rather than with the flows inside the impeller. So the clearance variation scarcely affects the pressure fluctuation characteristics in the blade channels. We also process a spectral analysis and find that the dominant frequency in the impeller is the rotating frequency induced by the blade-tongue interaction.

Additionally, it demonstrates that the pressure fluctuations on the pressure surface of the blade are severer than those in the volute casing; thus the unsteady flow in the rotating domain is still a research focus in terms of the centrifugal pump flow instability.

**3.2. Pressure Fluctuations in the Side Chambers.** Clearance flows in sidewall gaps are often ignored. However, through numerical simulations, we find that noteworthy pressure fluctuations exist in the sidewall gaps.

**3.2.1. Front Sidewall Gap.** The pressure fluctuation characteristics at Points FC1~FC4 in the front gap were analyzed. The pressure fluctuation peak-peak amplitudes gradually decrease along the flow direction (counterclockwise from Point FC1 to Point FC4) and at Point FC1 it even exceeds the maximum value in the volute casing, as displayed in Figure 10. The results of M1 and M3 have little difference, but they are

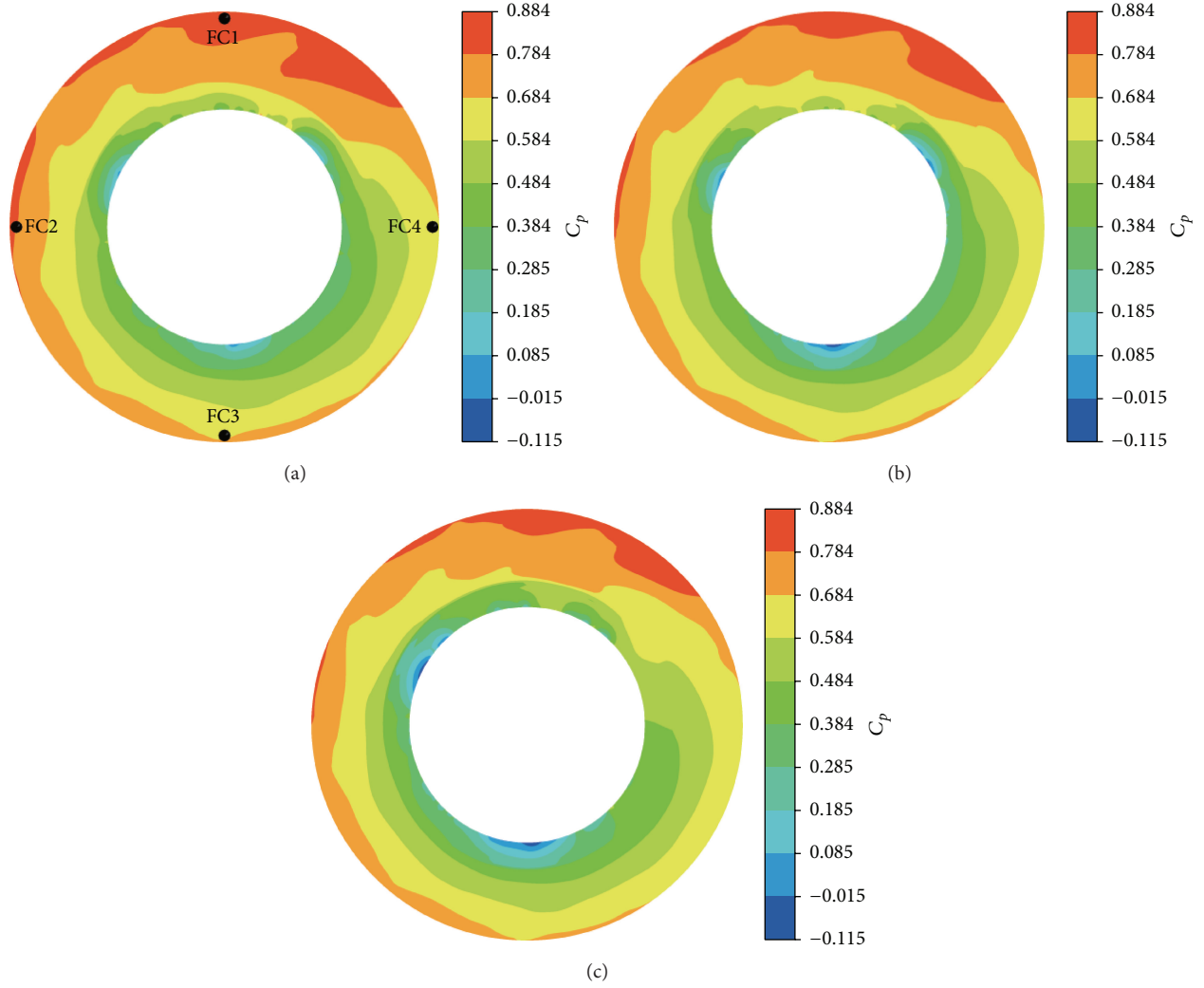


FIGURE 11: Pressure coefficient distribution on one section of the front gap at the last timestep: (a) M1, (b) M2, and (c) M3.

TABLE 4: Comparison of the dominant frequencies at the points in the front gap.

%	M1	M2	M3
	$f_d/f_n$	$f_d/f_n$	$f_d/f_n$
FC1	3	3	3
FC2	3	3	3
FC3	9	9	9
FC4	1/6	9	3

$f_d$  is the dominant frequency.

slightly larger than that of M2. Thus the pressure fluctuation intensity does not monotonously go up or down with the increase of  $c_f$ .

The dominant frequencies of the pressure fluctuations in the front gap are listed in Table 4. In each model, the dominate frequencies are not totally the same for the four points; the dominant frequencies at Point FC4 are also different for the three models. The clearance flow in the front gap is disturbed by the rotating component. The blade passing

frequency ( $3f_n$ ) dominates the pressure fluctuation in most region, while the back-blade passing frequency ( $9f_n$ ) plays a more important role at some lower regions. At Point FC4 in M1, there even exists a frequency of  $1/6f_n$  which may be related to the combined effect of the impeller blades and back blades. Consider the complexity of the flow states in the front gap.

Figure 11 displays the pressure coefficient distribution on the section where Points FC1~FC4 locate. The pressure coefficient is defined as  $C_p = 2(p - p_{\text{inlet}})/(\rho u_2^2)$ , where  $p_{\text{inlet}}$  is the area-averaged pressure at the pump inlet and  $u_2$  is the velocity at the impeller outlet. The pressure patterns look globally similar for the three models: the pressure distribution is not circumferentially uniform with evidently higher values in the top, which confirms the effect of the uneven flows in the volute casing. In addition, as  $c_f$  increases, the pressure level slightly drops. In each model,  $C_p$  decreases from FC1 to FC4 which is consistent with the corresponding pressure fluctuation peak-peak amplitude distribution feature. It indicates that the pressure at the top region is not only with higher values but also fluctuating more greatly.

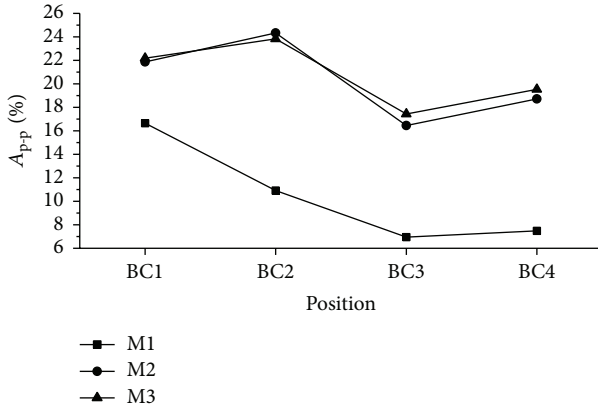


FIGURE 12: Peak-peak amplitudes at the points in the back sidewall gap.

**3.2.2. Back Sidewall Gap.** The pressure fluctuations in the back sidewall gap evidently differ from those in the front gap. As shown in Figure 12, in M1 the peak-peak amplitudes of the pressure fluctuations at the four points in the back sidewall gap are similar to those in the front gap, while in M2 and M3 the peak-peak amplitudes are much higher and slightly undulate along the flow direction (from Point BC1 to Point BC4).

A frequency spectrum analysis indicates that the four points in the back sidewall gap of each model have the same frequency components. We take Point BC1 as an example to illustrate the effect of the axial clearance variation. Figure 13 shows the corresponding frequency spectrum. In M1 the dominant frequency is the blade passing frequency, while in M2 and M3 the dominant frequency is the rotating frequency. Due to the viscous force, the fluid in the back sidewall gap flows with the rotating back shroud. The decreasing back clearance strengthens the viscous effect; thus the pressure fluctuation is weaker in M1. For M2 and M3, the flow fields are relatively symmetric owing to the smaller back clearance, and the rotating back shroud contributes to the dominant frequency (i.e., the rotating frequency). However, in M1 there is more mass transfer between the back sidewall gap and the main flow domain; thus the pressure wave dominated by the blade passing frequency propagates to the back sidewall gap.

Figure 14 displays the pressure coefficient distribution on the section where Points BC1~BC4 locate. Compared to the pressure field in the front gap, the pressure in the back sidewall gap distributes more evenly along the circumferential direction, and it gradually increases with the radius. Agreeing with the explanation in the last paragraph, the pressure presents more uneven pattern in M1 due to a bigger distance from the back shroud.

To sum up, the unsteady flow in the entire flow passage of the centrifugal pump is strongly associated with the rotor-stator interaction. The blade-tongue interaction is the main source of the pressure fluctuations in the main domains. Meanwhile, the front shroud, back shroud, and back blades that rotate with the impeller make the fluid in the sidewall gaps rotate at some degree, which may interact with the main flow. Flow states in the sidewall gaps are sensitive to the axial

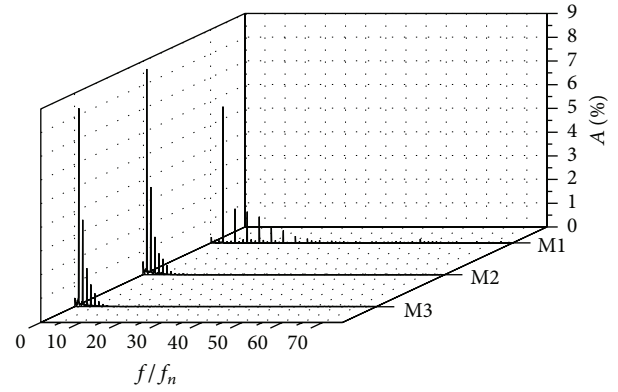


FIGURE 13: Frequency spectrum of the pressure fluctuations at Point BC1.

clearance size. A variation of the axial clearance affects the unsteady flow characteristics in the sidewall gaps and changes the level of the interaction between the clearance flow and the main flow, which leads to different pressure fluctuation frequencies. Flows in the impeller have no direct contact with clearance flows; thus the pressure fluctuation characteristics in the blade channels are not obviously affected by the axial clearance variation. It is meaningful to investigate the unsteady flow characteristics and the origin of the pressure fluctuations by doing numerical simulations for the entire flow passage of a centrifugal pump.

## 4. Conclusions

In this paper, unsteady numerical simulations were conducted in the design condition for a centrifugal pump with three axial clearance sizes, and the influence of axial clearance variation on the pressure fluctuation characteristics was analyzed for each flow domain. The numerical results of pump performance and pressure fluctuations with different axial clearance sizes all agree well with the experimental results. Through spectral analysis for several key points, we find that the axial clearance variation affects the pressure fluctuation characteristics in different flow domains by different degrees. Within the entire flow passage, the most intensive pressure fluctuation occurs at the blade pressure surface, which comes from the blade-tongue interaction, and it is almost not influenced by the axial clearance variation. The range of pressure fluctuation peak-peak amplitudes in the volute casing and in the sidewall gaps is mostly similar, but the maximum amplitude in the sidewall gaps is larger than that in the volute casing. Axial clearance sizes certainly affect the pressure fluctuation characteristics around the tongue and its effect is embodied in the low frequencies. Pressure fluctuation peak-peak amplitudes in the sidewall gaps do not change monotonously with the increase of the axial clearance size. In the front gap, the dominant frequency is not the same value for the four points in each model and also for the same point in different models. In the back sidewall gap, the largest back axial clearance leads to a much smaller pressure fluctuation which is related to viscous force from the rotating

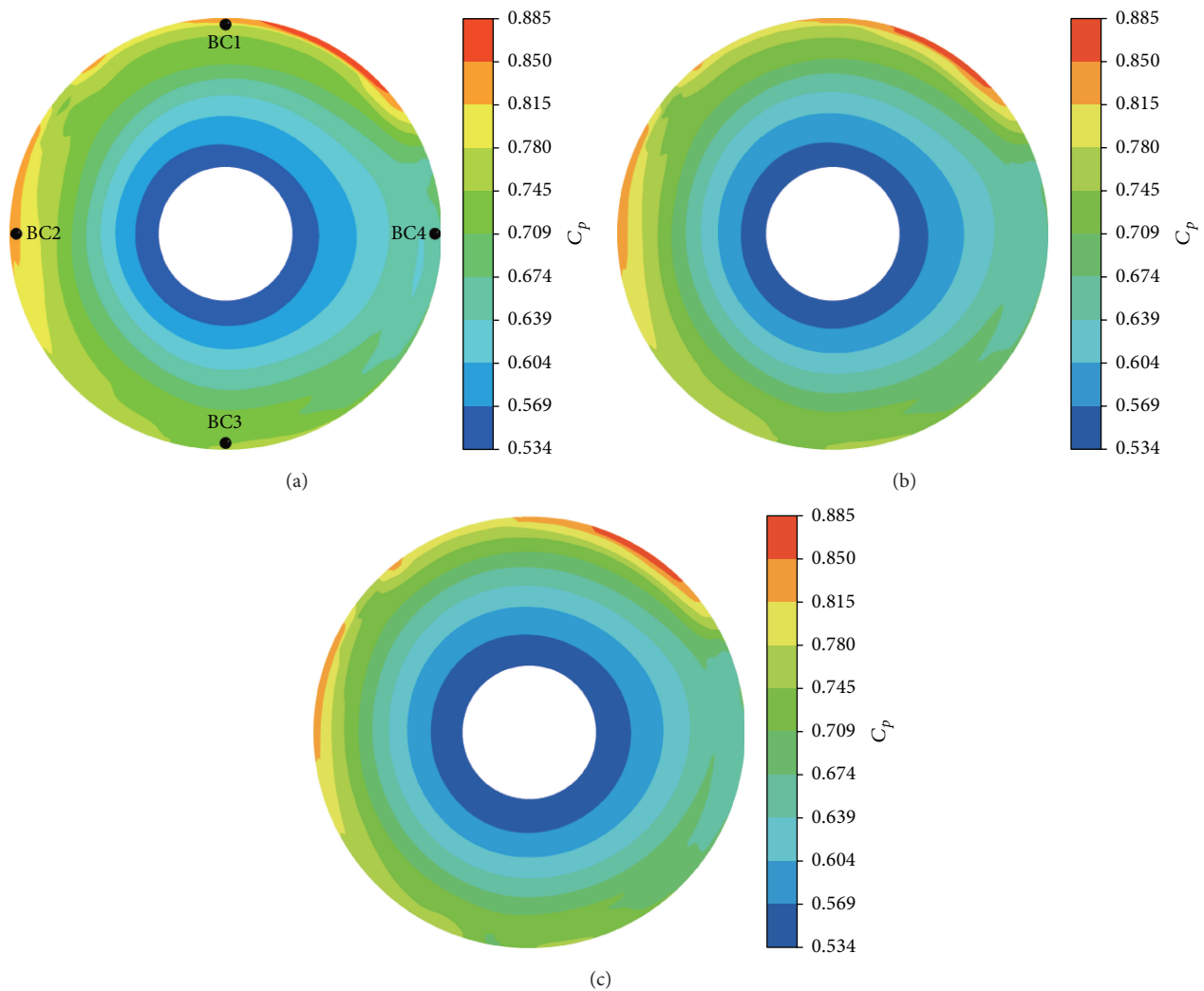


FIGURE 14: Pressure coefficient distribution on one section of the back sidewall gap at the last timestep: (a) M1, (b) M2, and (c) M3.

back shroud. Overall, clearance flows cannot be ignored and the axial clearance size has a certain influence on the pressure fluctuation characteristics in main flow domains as well as in sidewall gaps. The influence of the axial clearance on both pump performance and unsteady flow characteristics should be comprehensively considered in the hydraulic design of centrifugal pumps.

### Competing Interests

The authors declared that there are no competing interests related to this paper.

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