

Research Article

Spatial Central Configurations with Two Twisted Regular 4-Gons

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We study the configuration formed by two squares in two parallel layers separated by a distance. We picture the two layers horizontally with the z -axis passing through the centers of the two squares. The masses located on the vertices of each square are equal, but we do not assume that the masses of the top square are equal to the masses of the bottom square. We prove that the above configuration of two squares forms a central configuration if and only if the twist angle is equal to $k\pi/2$ or $(\pi/4 + k\pi/2)$ ($k = 1, 2, 3, 4$).

1. Introduction and Main Results

This paper uses the same notations as [1]. The Newtonian N -body problems [2, 3] concern the motions of N particles with masses $m_j \in \mathbb{R}^+$ and positions $q_j \in \mathbb{R}^3$ ($j = 1, 2, \dots, N$). The motion is governed by Newton's second law and the Universal law:

$$m_j \ddot{q}_j = -\frac{\partial U(q)}{\partial q_j}, \quad (1)$$

where $q = (q_1, q_2, \dots, q_N)$ with the Newtonian potential:

$$U(q) = \sum_{1 \leq j < k \leq N} \frac{m_j m_k}{|q_j - q_k|}. \quad (2)$$

Consider the space

$$X = \left\{ q = (q_1, q_2, \dots, q_N) \in \mathbb{R}^{3N} : \sum_{j=1}^N m_j q_j = 0 \right\}; \quad (3)$$

that is, suppose that the center of mass is fixed at the origin of the origin of the coordinate axis, because the potential is singular when two particles have the same position. It is

natural to assume that the configuration avoids the collision set $\Delta = \{q : q_j = q_k \text{ for some } k \neq j\}$. The set $X \setminus \Delta$ is called the configuration space.

Definition 1 (see [2, 3]). A configuration $q = (q_1, q_2, \dots, q_N) \in X \setminus \Delta$ is called a central configuration if there exists a constant λ such that

$$\sum_{j=1, j \neq k}^N \frac{m_j m_k}{|q_j - q_k|^3} (q_j - q_k) = -\lambda m_k q_k, \quad 1 \leq k \leq N. \quad (4)$$

The value of constant λ in (4) is uniquely determined by

$$\lambda = \frac{U}{I}, \quad (5)$$

where

$$I = \sum_{k=1}^N m_k |q_k|^2. \quad (6)$$

Consider the configuration in \mathbb{R}^3 consisting of two layer regular N -gons ($N \geq 2$) with distance $h \geq 0$. It is assumed that the lower layer regular N -gons lie in horizontal plane, and the upper regular N -gons parallel the lower

one, and z -axis passes through both centers of two regular N -gons. Suppose that the lower layer particles have masses m_1, m_2, \dots, m_N and the upper layer particles have masses $\tilde{m}_1, \tilde{m}_2, \dots, \tilde{m}_N$, respectively; then these assumptions can be interpreted more precisely by the following. Let ρ_k denote for all N complex roots of unity, that is,

$$\rho_k = \exp\left(\frac{2\pi k}{N}i\right). \quad (7)$$

Let

$$\tilde{\rho}_k = a\rho_k e^{i\varphi}, \quad (8)$$

where $a > 0$, $i = \sqrt{-1}$, $0 \leq \varphi \leq 2\pi$, and φ is called twist angle. It is assumed that m_k ($1 \leq k \leq N$) locates at the vertex q_k of the lower layer regular N -gons; \tilde{m}_k ($1 \leq k \leq N$) locates on the vertex \tilde{q}_k of the upper layer regular N -gons:

$$\begin{aligned} q_k &= (\rho_k, 0), \\ \tilde{q}_k &= (\tilde{\rho}_k, h), \end{aligned} \quad (9)$$

where $h \geq 0$ is the distance between the two layers. Then the center of masses is

$$z_0 = \frac{\sum_{j=1}^N (m_j q_j + \tilde{m}_j \tilde{q}_j)}{M}, \quad (10)$$

where

$$M = \sum_{j=1}^N (m_j + \tilde{m}_j). \quad (11)$$

Let

$$\begin{aligned} P_k &= q_k - z_0, \\ \tilde{P}_k &= \tilde{q}_k - z_0. \end{aligned} \quad (12)$$

If m_1, m_2, \dots, m_N and $\tilde{m}_1, \tilde{m}_2, \dots, \tilde{m}_N$ form a central configuration, then there is $\lambda \in \mathbb{R}^+$ such that

$$\begin{aligned} \sum_{j \neq k}^N m_j \frac{P_j - P_k}{|P_j - P_k|^3} + \sum_{j=1}^N \tilde{m}_j \frac{\tilde{P}_j - P_k}{|\tilde{P}_j - P_k|^3} &= -\lambda P_k, \quad 1 \leq k \leq N, \\ \sum_{j=1}^N m_j \frac{P_j - \tilde{P}_k}{|P_j - \tilde{P}_k|^3} + \sum_{j \neq k}^N \tilde{m}_j \frac{\tilde{P}_j - \tilde{P}_k}{|\tilde{P}_j - \tilde{P}_k|^3} &= -\lambda \tilde{P}_k, \quad 1 \leq k \leq N. \end{aligned} \quad (13)$$

Under the case that twist angle $\varphi = 0$, Moeckel-Simo proved.

Theorem 2 (see [4]). If $N < 473$, there is a unique pair of spatial central configurations of parallel regular N -gons. If $N \geq 473$, there is no such central configuration for $b < b_0(N)$, where b is the mass ratio. At a unique pair bifurcates from the planar central configuration with the smaller masses on the inner polygon. This remains the unique pair of spatial central configurations until $b = 1/b_0$, where a similar bifurcation occurs in reverse, so that $b > 1/b_0$, and only the planar central configurations remain.

Xie et al. [5] studied the necessary conditions for the masses of two layer regular polygon central configurations in \mathbb{R}^3 , and they proved the following theorems.

Theorem 3. Under the assumptions of (9), if m_1, m_2, \dots, m_N and $\tilde{m}_1, \tilde{m}_2, \dots, \tilde{m}_N$ form a central configuration, then

(i)

$$\lambda = \frac{M}{N} \sum_{j=1}^N \frac{1}{|q_j - \tilde{q}_N|^3}, \quad (14)$$

(ii) for $N \geq 2$ ($\varphi \neq \pi/2$, when $N = 2$),

$$m_1 = m_2 = \dots = m_N, \quad \tilde{m}_1 = \tilde{m}_2 = \dots = \tilde{m}_N. \quad (15)$$

Without loss of generality, suppose that $\tilde{m}_j = bm_j$ and $N \geq 2$ ($\varphi \neq \pi/2$, when $N = 2$), under the above assumptions, there are four parameters, and ratio b of masses ratio, a of radius of two regular N -gons, the distance h between two layers, and the phase difference φ . As for these parameters, Xie et al. [5] proved.

Theorem 4. Under the assumptions of (9) and $\tilde{m}_j = bm_j$, then m_1, m_2, \dots, m_N and $\tilde{m}_1, \tilde{m}_2, \dots, \tilde{m}_N$ form a central configuration if and only if the parameters b, a, h , and φ satisfy the following relationships:

$$\begin{aligned} \frac{\lambda}{M} &= \frac{1}{N(1+b)} \left(\sum_{j \neq N} \frac{1 - \rho_j}{|q_j - q_N|^3} + \sum_j \frac{b(1 - a\rho_j e^{i\varphi})}{|\tilde{q}_j - q_N|^3} \right), \\ \frac{\lambda}{M} &= \frac{1}{N} \sum_{j=1}^N \frac{1}{|q_j - \tilde{q}_N|^3}, \\ \frac{\lambda}{M} &= \frac{e^{-i\varphi}}{Na(1+b)} \left(\sum_j \frac{ae^{i\varphi} - \rho_j}{|q_j - \tilde{q}_N|^3} + \sum_{j \neq N} \frac{b(a - a\rho_j e^{i\varphi})}{|\tilde{q}_j - \tilde{q}_N|^3} \right). \end{aligned} \quad (16)$$

Zhang and Zhu [1] proved the sufficient conditions for special cases $a = 1$, $b = 1$, and $\varphi = \pi/N$.

Theorem 5. Under the assumptions of (9) and (15) and $\tilde{m}_j = bm_j$ if $a = 1$, $b = 1$, and $\varphi = \pi/N$, then for every N there exists a unique central configuration.

When $h = 0$, Bang and Elmabsout [6] study the twist angle.

Theorem 6. Let Π_1 (resp., Π_2) be a regular N -gon centred around a mass m_0 at O , m_1 being at each of its vertices (resp., m_2). Then m_0, Π_1 , and Π_2 are relative equilibriums if and only if they are homothetic or cursed with an angle equal to (π/N) (and suitable ratio of radii).

A natural interesting problem is that whether (9) form a central configuration for $h > 0$ if and only if twist angle $\varphi = 0$ or π/N . In this paper, one will prove the necessary condition of twisted angle φ for a special case $N = 4$.

Theorem 7. Under the assumptions of (9) and (15) and $\tilde{m}_j = bm_j$ if $N = 4$, then $\varphi = \pi/4 + (\pi/2)k$ or $\varphi = (\pi/2)k$, ($k = 0, 1, 2, 3$).

2. The Proof of Theorem 7

2.1. *Some Lemmas.* We need three lemmas.

Lemma 1 (see [7]). For $N \geq 3$ and $\tilde{m}_1 = \tilde{m}_2 = \dots = \tilde{m}_N$, if $\tilde{m}_1, \tilde{m}_2, \dots, \tilde{m}_N$ locate on vertices of a regular polygon, then they form a central configuration.

One has $m_1^* = m_2^* = m_3^* = m_4^*$. Also m_1^*, m_2^*, m_3^* , and m_4^* are located on vertices of a square S ; q_i^* is the position m_i^* ($i = 1, 2, 3, 4$) and q^* is the position of m^* . Plane P contains the square S . One has $F_* = \sum_{i=1}^4 m_i^* m^* (q_i^* - q^*) / |q_i^* - q^*|^3$.

Lemma 2. F_* 's projection on P is directed toward the center of the square if and only if the point m^* is on a vertical plane of symmetry.

2.2. *The Proof of Lemma 2.* Without loss of generality, suppose that $m_1^* = m_2^* = m_3^* = m_4^* = 1$, $q_1^* = (1, 0, 0)$, $q_2^* = (0, 1, 0)$, $q_3^* = (-1, 0, 0)$, $q_4^* = (0, -1, 0)$, and $q^* = (\rho \cos \theta, \rho \sin \theta, h)$, where $\rho > 0$. Consider the following:

$$\begin{aligned} F_* = & \frac{m^* (1 - \rho \cos \theta, -\rho \sin \theta, -h)}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} \\ & + \frac{m^* (-\rho \cos \theta, 1 - \rho \sin \theta, -h)}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \\ & + \frac{m^* (-1 - \rho \cos \theta, -\rho \sin \theta, -h)}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \\ & + \frac{m^* (-\rho \cos \theta, -1 - \rho \sin \theta, -h)}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}}. \end{aligned} \quad (17)$$

If F_* 's projection on P is directed toward the center of the square, there exists $\lambda^* > 0$ such that

$$\begin{aligned} & \frac{m^* (1 - \rho \cos \theta)}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} + \frac{m^* (-\rho \cos \theta)}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \\ & + \frac{m^* (-1 - \rho \cos \theta)}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \\ & + \frac{m^* (-\rho \cos \theta)}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}} \\ & = -\lambda^* \rho \cos \theta, \end{aligned} \quad (18)$$

$$\begin{aligned} & = \frac{m^* (-\rho \sin \theta)}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} \\ & + \frac{m^* (1 - \rho \sin \theta)}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \\ & + \frac{m^* (-\rho \sin \theta)}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \\ & + \frac{m^* (-1 - \rho \sin \theta)}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}} \\ & = -\lambda^* \rho \sin \theta. \end{aligned} \quad (19)$$

If $\theta = (\pi/2)k$, m^* is on a vertical plane of symmetry.

When $\theta \neq (\pi/2)k$, multiplying both sides of (18) and (19), respectively, by $\sin \theta$ and $\cos \theta$, we have

$$\begin{aligned} & \frac{\sin \theta - \rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} + \frac{-\rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \\ & + \frac{-\sin \theta - \rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \\ & + \frac{-\rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}} \\ & = -\lambda^* \rho \cos \theta \sin \theta, \\ & \frac{-\rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} + \frac{\cos \theta - \rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \\ & + \frac{-\rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \\ & + \frac{-\cos \theta - \rho \cos \theta \sin \theta}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}} \\ & = -\lambda^* \rho \cos \theta \sin \theta. \end{aligned} \quad (20)$$

By (20) we have

$$\begin{aligned} & \frac{\sin \theta}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} + \frac{-\cos \theta}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \\ & + \frac{-\sin \theta}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \\ & + \frac{\cos \theta}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}} = 0. \end{aligned} \quad (21)$$

Then

$$\begin{aligned} & \frac{1}{\sin \theta} \left(\frac{1}{|1 + \rho^2 + h^2 - 2\rho \sin \theta|^{3/2}} \right. \\ & \quad \left. - \frac{1}{|1 + \rho^2 + h^2 + 2\rho \sin \theta|^{3/2}} \right) \\ & = \frac{1}{\cos \theta} \left(\frac{1}{|1 + \rho^2 + h^2 - 2\rho \cos \theta|^{3/2}} \right. \\ & \quad \left. - \frac{1}{|1 + \rho^2 + h^2 + 2\rho \cos \theta|^{3/2}} \right). \end{aligned} \quad (22)$$

Let

$$f(x) = \frac{1}{x} \left(\frac{1}{|1 + \rho^2 + h^2 - 2\rho x|^{3/2}} - \frac{1}{|1 + \rho^2 + h^2 + 2\rho x|^{3/2}} \right); \quad (23)$$

then the system (22) is equivalent to $f(\sin \theta) = f(\cos \theta)$. It is obvious that $f(x)$ is even and $f(x) > 0$ when $x > 0$. We will prove that $f(x)$ is a strictly increasing function when $x > 0$.

When $0 < x \leq 1$ we compute

$$\begin{aligned} \frac{df(x)}{dx} &= -\frac{1}{x^2} \left(\frac{1}{|1 + \rho^2 + h^2 - 2\rho x|^{3/2}} - \frac{1}{|1 + \rho^2 + h^2 + 2\rho x|^{3/2}} \right) \\ &\quad + \frac{1}{x} \left(\frac{3\rho}{|1 + \rho^2 + h^2 - 2\rho x|^{5/2}} + \frac{3\rho}{|1 + \rho^2 + h^2 + 2\rho x|^{5/2}} \right) \\ &= \frac{(1 + \rho^2 + h^2 + 2\rho x)(1 + \rho^2 + h^2 - 2\rho x)^{5/2}}{x^2|1 + \rho^2 + h^2 + 2\rho x|^{5/2}|1 + \rho^2 + h^2 - 2\rho x|^{5/2}} \\ &\quad - \frac{(1 + \rho^2 + h^2 - 2\rho x)(1 + \rho^2 + h^2 + 2\rho x)^{5/2}}{x^2|1 + \rho^2 + h^2 + 2\rho x|^{5/2}|1 + \rho^2 + h^2 - 2\rho x|^{5/2}} \\ &\quad + \frac{3\rho x(1 + \rho^2 + h^2 + 2\rho x)^{5/2}}{x^2|1 + \rho^2 + h^2 + 2\rho x|^{5/2}|1 + \rho^2 + h^2 - 2\rho x|^{5/2}} \\ &\quad + \frac{3\rho x(1 + \rho^2 + h^2 - 2\rho x)^{5/2}}{x^2|1 + \rho^2 + h^2 + 2\rho x|^{5/2}|1 + \rho^2 + h^2 - 2\rho x|^{5/2}}. \end{aligned} \quad (24)$$

Let

$$\begin{aligned} g(x) &= (1 + \rho^2 + h^2 + 2\rho x) \\ &\quad \times (1 + \rho^2 + h^2 - 2\rho x)^{5/2} \\ &\quad - (1 + \rho^2 + h^2 - 2\rho x) \\ &\quad \times (1 + \rho^2 + h^2 + 2\rho x)^{5/2} \\ &\quad + 3\rho x(1 + \rho^2 + h^2 + 2\rho x)^{5/2} \\ &\quad + 3\rho x(1 + \rho^2 + h^2 - 2\rho x)^{5/2}. \end{aligned} \quad (25)$$

Then

$$\frac{df(x)}{dx} = \frac{g(x)}{x^2|1 + \rho^2 + h^2 + 2\rho x|^{5/2}|1 + \rho^2 + h^2 - 2\rho x|^{5/2}}, \quad (26)$$

$$g(0) = (1 + \rho^2 + h^2)^{7/2} - (1 + \rho^2 + h^2)^{7/2} = 0, \quad (27)$$

$$\begin{aligned} \frac{dg(x)}{dx} &= 2\rho(1 + \rho^2 + h^2 - 2\rho x)^{5/2} \\ &\quad - 5\rho(1 + \rho^2 + h^2 + 2\rho x) \\ &\quad \times (1 + \rho^2 + h^2 - 2\rho x)^{3/2} \\ &\quad + 2\rho(1 + \rho^2 + h^2 + 2\rho x)^{5/2} \\ &\quad - 5\rho(1 + \rho^2 + h^2 - 2\rho x) \\ &\quad \times (1 + \rho^2 + h^2 + 2\rho x)^{3/2} \\ &\quad + 3\rho(1 + \rho^2 + h^2 + 2\rho x)^{5/2} \\ &\quad + 15\rho^2 x(1 + \rho^2 + h^2 + 2\rho x)^{3/2} \\ &\quad + 3\rho(1 + \rho^2 + h^2 - 2\rho x)^{5/2} \\ &\quad - 15\rho^2 x(1 + \rho^2 + h^2 - 2\rho x)^{3/2}, \end{aligned} \quad (28)$$

$$\begin{aligned} \left. \frac{dg(x)}{dx} \right|_{x=0} &= 10\rho(1 + \rho^2 + h^2)^{5/2} \\ &\quad - 10\rho(1 + \rho^2 + h^2)^{5/2} = 0, \end{aligned} \quad (29)$$

$$\begin{aligned} \frac{d^2g(x)}{dx^2} &= -10\rho^2(1 + \rho^2 + h^2 - 2\rho x)^{3/2} \\ &\quad - 10\rho^2(1 + \rho^2 + h^2 + 2\rho x)^{3/2} \\ &\quad + 15\rho^2(1 + \rho^2 + h^2 + 2\rho x) \\ &\quad \times (1 + \rho^2 + h^2 - 2\rho x)^{1/2} \\ &\quad + 10\rho^2(1 + \rho^2 + h^2 + 2\rho x)^{3/2} \\ &\quad + 10\rho^2(1 + \rho^2 + h^2 - 2\rho x)^{3/2} \\ &\quad - 15\rho^2(1 + \rho^2 + h^2 - 2\rho x) \\ &\quad \times (1 + \rho^2 + h^2 + 2\rho x)^{1/2} \\ &\quad + 15\rho^2(1 + \rho^2 + h^2 + 2\rho x)^{3/2} \\ &\quad + 15\rho^2(1 + \rho^2 + h^2 - 2\rho x)^{3/2} \\ &\quad + 45\rho^3 x(1 + \rho^2 + h^2 + 2\rho x)^{1/2} \end{aligned}$$

$$\begin{aligned}
& -15\rho^2(1+\rho^2+h^2-2\rho x)^{3/2} \\
& -15\rho^2(1+\rho^2+h^2-2\rho x)^{3/2} \\
& +45\rho^3x(1+\rho^2+h^2-2\rho x)^{1/2} \\
& \geq 45\rho^3x(1+\rho^2+h^2+2\rho x)^{1/2} \\
& +45\rho^3x(1+\rho^2+h^2-2\rho x)^{1/2} \\
& > 0, \quad \forall 0 < x \leq 1.
\end{aligned} \tag{30}$$

Hence by (26), (27), (29), and (30) we have

$$g(x) > 0, \quad \forall 0 < x \leq 1. \tag{31}$$

Hence $df(x)/dx > 0$, when $0 < x \leq 1$. From $f(\sin \theta) = f(\cos \theta)$ we obtain $|\sin \theta| = |\cos \theta|$; that is, $\theta = (\pi/4) + (\pi/2)k$, so m^* is on a vertical plane of symmetry.

It is obvious that if the point m^* is on a vertical plane of symmetry, F_* 's projection on P is directed toward the center of the square.

The proof of Lemma 2 is completed.

2.3. The Proof of Theorem 7. Without loss of generality, suppose that $\tilde{m}_1 = \tilde{m}_2 = \tilde{m}_3 = \tilde{m}_4 = m$, $m_1 = m_2 = m_3 = m_4 = 1$. We can let

$$\begin{aligned}
q_1 &= (1, 0, 0), & q_2 &= (0, 1, 0), \\
q_3 &= (-1, 0, 0), & q_4 &= (0, -1, 0), \\
\tilde{q}_1 &= (\rho \cos \varphi, \rho \sin \varphi, h), \\
\tilde{q}_2 &= \left(\rho \cos \left(\frac{2\pi}{4} + \varphi \right), \rho \sin \left(\frac{2\pi}{4} + \varphi \right), h \right), \\
\tilde{q}_3 &= \left(\rho \cos \left(\frac{2\pi}{4} \times 2 + \varphi \right), \rho \sin \left(\frac{2\pi}{4} \times 2 + \varphi \right), h \right), \\
\tilde{q}_4 &= \left(\rho \cos \left(\frac{2\pi}{4} \times 3 + \varphi \right), \rho \sin \left(\frac{2\pi}{4} \times 3 + \varphi \right), h \right),
\end{aligned} \tag{32}$$

where $\rho h > 0$. It is obvious that $z_0 = \sum_{j=1}^4 (m_j q_j + \tilde{m}_j + \tilde{q}_j)/M = (0, 0, 4mh/(4+4m)) = (0, 0, mh/(1+m))$. By Lemma 1, since $\tilde{m}_1, \tilde{m}_2, \tilde{m}_3$, and \tilde{m}_4 locate on vertices of a regular polygon, so they form a central configuration; then there exists a constant $\tilde{\lambda}$ (notice that it can be different from λ in (16)), such that

$$\sum_{j \neq k}^4 \tilde{m}_j \frac{\tilde{q}_j - \tilde{q}_k}{|\tilde{q}_j - \tilde{q}_k|} = -\tilde{\lambda} (\tilde{q}_k - \tilde{z}_0), \tag{33}$$

where $\tilde{z}_0 = (\sum_{j=1}^4 \tilde{q}_j \tilde{m}_j)/4m = (0, 0, h)$.

By (1) and (33) we have

$$\begin{aligned}
-\lambda m \tilde{P}_k &= \sum_{j=1}^4 m \frac{q_j - \tilde{q}_k}{|q_j - \tilde{q}_k|^3} + \sum_{j \neq k}^4 m m \frac{\tilde{q}_j - \tilde{q}_k}{|\tilde{q}_j - \tilde{q}_k|^3} \\
&= \sum_{j=1}^4 m \frac{P_j - \tilde{P}_k}{|P_j - \tilde{P}_k|^3} + \sum_{j \neq k}^4 m m \frac{\tilde{P}_j - \tilde{P}_k}{|\tilde{P}_j - \tilde{P}_k|^3} \\
&= \sum_{j=1}^4 m \frac{q_j - \tilde{q}_k}{|q_j - \tilde{q}_k|^3} - \tilde{\lambda} m (\tilde{q}_k - \tilde{z}_0) \quad (k = 1, 2, 3, 4).
\end{aligned} \tag{34}$$

Then

$$\begin{aligned}
& \sum_{j=1}^4 m \frac{(q_j - \tilde{q}_k)}{|q_j - \tilde{q}_k|^3} \\
&= -\lambda m (\tilde{q}_k - z_0) + \tilde{\lambda} m (\tilde{q}_k - \tilde{z}_0) \\
&= m (\tilde{\lambda} - \lambda) \tilde{q}_k + m \left(0, 0, \frac{\lambda m h}{1+m} \right) - m (0, 0, \tilde{\lambda} h).
\end{aligned} \tag{35}$$

Letting $\tilde{F}_1 = \sum_{j=1}^4 m(q_j - \tilde{q}_1)/|q_j - \tilde{q}_1|^3$, \tilde{F}_1 is the force generated by the m_1, m_2, m_3 , and m_4 .

By (19), \tilde{F}_1 's projection on the plane P_1 is $(\tilde{\lambda} - \lambda)(\rho \cos \varphi, \rho \sin \varphi, 0)$, where P_1 is the plane containing m_1, m_2, m_3 , and m_4 . It is obvious that the projection is directed toward the center of the lower layer regular 4-gons. By Lemma 2, we have $\varphi = \pi/4 + (\pi/2)k$ or $\varphi = (\pi/2)k$ ($k = 0, 1, 2, 3$).

The proof of Theorem 7 is completed.

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References

- [1] S. Zhang and C. Zhu, "Central configurations consist of two layer twisted regular polygons," *Science in China A*, vol. 45, no. 11, pp. 1428–1438, 2002.
- [2] D. G. Saari, *Collisions, Rings and Other Newtonian N-body Problems*, vol. 104, AMS, Providence, RI, USA, 2005.
- [3] A. Wintner, *The Analytical Foundations of Celestial Mechanics*, Princeton University Press, Princeton, NJ, USA, 1941.
- [4] R. Moeckel and C. Simó, "Bifurcation of spatial central configurations from planar ones," *SIAM Journal on Mathematical Analysis*, vol. 26, no. 4, pp. 978–998, 1995.
- [5] Z. F. Xie, S. Q. Zhang, and Q. Zhou, "Two layer regular polygonal central configurations in R^3 ," Preprint, 2000.
- [6] D. Bang and B. Elmabsout, "Representations of complex functions, means on the regular n -gon and applications to gravitational potential," *Journal of Physics A*, vol. 36, no. 45, pp. 11435–11450, 2003.

- [7] L. M. Perko and E. L. Walter, "Regular polygon solutions of the N -body problem," *Proceedings of the American Mathematical Society*, vol. 94, no. 2, pp. 301–309, 1985.

