

Research Article

Ground State Solutions of Schrödinger-Kirchhoff Equations with Potentials Vanishing at Infinity

Dongdong Sun 🕑

School of Mathematics, Qilu Normal University, Jinan 250013, China

Correspondence should be addressed to Dongdong Sun; sundd@amss.ac.cn

Received 27 June 2023; Revised 7 August 2023; Accepted 17 August 2023; Published 31 August 2023

Academic Editor: Baowei Feng

Copyright © 2023 Dongdong Sun. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

In this paper, we deal with the following Schrödinger-Kirchhoff equation with potentials vanishing at infinity: $-(\varepsilon^2 a + \varepsilon b \int_{\mathbb{R}^3} |\nabla u|^2) \Delta u + V(x)u = K(x)|u|^{p-1}u$ in \mathbb{R}^3 and u > 0, $u \in H^1(\mathbb{R}^3)$, where $V(x) \sim |x|^{-\alpha}$ and $K(x) \sim |x|^{-\beta}$ with $0 < \alpha < 2$, and $\beta > 0$. We first prove the existence of positive ground state solutions $u_{\varepsilon} \in H^1(\mathbb{R}^3)$ under the assumption that $\sigma for some <math>\sigma = \sigma_{\alpha,\beta}$, then we show that u_{ε} concentrates at a global minimum point of $\mathcal{A}(x) = V^{2/(p-1)-1/2}(x)/K^{2/(p-1)}(x)$.

1. Introduction

In this paper, the following Schrödinger-Kirchhoff equations with potentials vanishing at infinity are studied:

$$\begin{cases} -\left(\varepsilon^{2}a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla u|^{2}\right) \Delta u + V(x)u = K(x)|u|^{p-1}u \text{ in } \mathbb{R}^{3}, \\ u > 0, u \in H^{1}(\mathbb{R}^{3}), \end{cases}$$
(1)

where a, b > 0 are constants, and $\varepsilon > 0$ is a small parameter. The potentials V(x) and K(x) satisfy

(V) $V : \mathbb{R}^3 \longrightarrow \mathbb{R}$ is Hölder continuous and $\bar{a}/(1 + |x|^{\alpha}) \le V(x) \le A$ for some $\bar{a}, A > 0$, and $0 < \alpha < 2$

(K) $K : \mathbb{R}^3 \longrightarrow \mathbb{R}$ is Hölder continuous and $0 < K(x) \le k/(1 + |x|^{\beta})$ for some $\beta, k > 0$

Problem (1) is related to the stationary analogue of the equation

$$u_{\rm tt} - \left(a + b \int_{\Omega} |\nabla u|^2 \right) \Delta u = g(x, t), \tag{2}$$

proposed by Kirchhoff [1] as an extension of the classical D'Alembert's wave equation for free vibrations of elastic

strings. Early studies on Kirchhoff equations (2) were Bernstein [2] and Pohozaev [3]. In recent years, lots of interesting results on the elliptic Kirchhoff equations have been obtained. Here, we only refer to [4–12] and references therein.

Let b = 0 in (1), then equation (1) becomes the wellknown Schrödinger equation

$$-\varepsilon^2 \Delta u + V(x)u = f(u) \text{ in } \mathbb{R}^N.$$
(3)

There have been enormous results on (3). Since we cannot give a complete list of references here, we only refer to [13-22].

Recently, many authors considered the existence and concentration of positive solutions for Schrödinger-Kirchhoff equations. In [23], He and Zou studied the Schrödinger-Kirchhoff equation with subcritical nonlinearity. In [24], Wang et al. studied the Schrödinger-Kirchhoff equation with critical nonlinearity. In [25], Figueiredo et al. treated the Schrödinger-Kirchhoff equation with the almost optimal Berestycki-Lions type nonlinearity. In [26], Sun and Zhang investigated the existence and concentration of ground state solutions for Schrödinger-Kirchhoff equations with competing potentials. In [27], Sun and Zhang studied the ground state solutions for Schrödinger-Kirchhoff equations with a critical frequency. Further results can be seen in [28–36] etc.

In [37], Ambrosetti et al. studied a class of nonlinear Schrödinger equations with potentials vanishing at infinity. Inspired by [37], we consider the Schrödinger-Kirchhoff equation (1) which contains a nonlocal term $\int_{\mathbb{R}^3} |\nabla v|^2$ in it. Because of the nonlocal term, several special difficulties will occur in our arguments. For instance, the estimate (34) in our paper would not be obtained if we follow the proof of Lemma 17 in [37] directly. As a result, we can only prove Lemma 6 for small b > 0 in (1). To overcome this difficulty, new definition of $R_{n,\varepsilon}$ is given, and the desired inequality (34) is finally proved by delicate analysis.

Let

$$\sigma \coloneqq \sigma_{\alpha,\beta} = \begin{cases} 5 - \frac{4\beta}{\alpha}, & \text{if } 0 < \beta < \frac{\alpha}{2}, \\ 3, & \text{otherwise.} \end{cases}$$
(4)

Our main results are as follows.

Theorem 1. Assume that (V) and (K) are satisfied. Let

$$\sigma$$

Then, for every $\varepsilon > 0$, there exists a positive ground state solution $u_{\varepsilon} \in H^1(\mathbb{R}^3)$ of (1).

Concerning the concentration behaviour of the ground state solutions u_{ε} obtained in Theorem 1, we have the following theorem.

Theorem 2. Under the same assumptions as in Theorem 1, the positive ground state solutions $\{u_{\varepsilon}\}$ concentrate at a global minimum point x^* of $\mathscr{A}(x) = V^{2/(p-1)-1/2}(x)/K^{2/(p-1)}(x)$. That is, u_{ε} has a unique maximum point x_{ε} and $x_{\varepsilon} \longrightarrow x^*$ as $\varepsilon \longrightarrow 0^+$, and

$$u_{\varepsilon}(x) = U^* \left(\frac{x - x_{\varepsilon}}{\varepsilon}\right) + \omega_{\varepsilon}(x) \text{ as } \varepsilon \longrightarrow 0^+, \tag{6}$$

where $\omega_{\varepsilon} \longrightarrow 0$ in $L^{\infty}(\mathbb{R}^3)$ and in $C^2_{loc}(\mathbb{R}^3)$ as $\varepsilon \longrightarrow 0^+$ and U^* is the unique positive ground state solution of

$$-\left(a+b\int_{\mathbb{R}^{3}}|\nabla U^{*}|^{2}\right)\Delta U^{*}+V(x^{*})U^{*}=K(x^{*})(U^{*})^{p}.$$
 (7)

Remark 3. The uniqueness of ground state solution of (7) can be seen in [26].

2. Existence of Positive Ground State Solutions

2.1. Preliminaries. To prove our results, we work in the following weighted Sobolev spaces:

$$\mathscr{H}_{\varepsilon} \coloneqq \left\{ v \in \mathscr{D}^{1,2}(\mathbb{R}^3) \colon \int_{\mathbb{R}^3} \left(\varepsilon^2 a |\nabla v|^2 + V(x) v^2 \right) < +\infty \right\}.$$
(8)

 $\mathcal{H}_{\varepsilon}$ is a Hilbert space, and the scalar product and norm are as follows:

$$\|v\|_{\varepsilon}^{2} = \int_{\mathbb{R}^{3}} \left(\varepsilon^{2} a |\nabla v|^{2} + V(x)v^{2}\right),$$

$$(u, v)_{\varepsilon} = \int_{\mathbb{R}^{3}} \left(\varepsilon^{2} a \nabla u \nabla v + V(x)uv\right).$$
(9)

Set $\mathcal{H} = \mathcal{H}_1$ with norm $\|\cdot\|_{\mathcal{H}}$. Let L_K^q be the weighted space of measurable functions $u : \mathbb{R}^3 \longrightarrow \mathbb{R}$ such that

$$|u|_{q,K} = \left(\int_{\mathbb{R}^3} K(x)|u|^q dx\right)^{1/q} < \infty.$$
 (10)

The following result can be seen in [38].

Theorem 4. Assume that (V) and (K) are satisfied. Then $\mathcal{H}_{\varepsilon} \subset L_{K}^{p+1}$ provided $\sigma \leq p \leq 5$, and

$$\|\nu\|_{p+1,K} \le C_{\varepsilon} \|\nu\|_{\varepsilon}, \forall \nu \in \mathscr{H}_{\varepsilon}, \tag{11}$$

where $C_{\varepsilon} > 0$ is a constant. Furthermore, the embedding is compact if (5) holds.

2.2. Proof of Theorem 1. In this section, we will prove Theorem 1.

Define $I_{\varepsilon}: \mathscr{H}_{\varepsilon} \longrightarrow \mathbb{R}$ by

$$I_{\varepsilon}(u) \coloneqq \frac{1}{2} \int_{\mathbb{R}^3} \left(\varepsilon^2 a |\nabla u|^2 + V(x) u^2 \right) + \frac{b}{4} \varepsilon \left(\int_{\mathbb{R}^3} |\nabla u|^2 \right)^2 - \frac{1}{p+1} \int_{\mathbb{R}^3} K(x) |u|^{p+1}.$$
(12)

Then, $I_{\varepsilon} \in C^{1}(\mathscr{H}_{\varepsilon}, \mathbb{R})$ and the critical point of the energy functional I_{ε} are just a weak solution of problem (1).

We have the following theorem.

Theorem 5. *Assume that (V) and (K) are satisfied, and let (5) hold. Then*

$$c_{\varepsilon} = \inf_{\nu \in \mathscr{H}_{\varepsilon} \setminus \{0\}} \max_{t \ge 0} I_{\varepsilon}(t\nu)$$
(13)

is the ground energy level of I_{ε} , and equation (1) has a positive ground state solution $u_{\varepsilon} \in \mathscr{H}_{\varepsilon}$. Furthermore,

$$\|u_{\varepsilon}\|_{\mathfrak{s}}^{2} \leq Cc_{\varepsilon},\tag{14}$$

where C > 0 is a constant.

Proof. As Lemma 3.2 in [26] (or Lemma 2.2 in [27]), we know that c_{ε} is the ground energy level of I_{ε} , i.e.,

$$c_{\varepsilon} = \inf_{v \in \mathcal{N}_{\varepsilon}} I_{\varepsilon}(v), \tag{15}$$

where $\mathscr{N}_{\varepsilon}$ is the Nehari manifold of $I_{\varepsilon}: \mathscr{H}_{\varepsilon} \longrightarrow \mathbb{R}$ defined

by

$$\mathcal{N}_{\varepsilon} = \left\{ v \in \mathcal{H}_{\varepsilon} \setminus \{0\} : \int_{\mathbb{R}^{3}} \left(\varepsilon^{2} a |\nabla v|^{2} + V(x)v^{2} \right) + \varepsilon b \left(\int_{\mathbb{R}^{3}} |\nabla v|^{2} \right)^{2} = \int_{\mathbb{R}^{3}} K(x) |v|^{p+1} \right\}.$$
(16)

It is easy to verify the Mountain-Pass geometry of $I_{\varepsilon} : \mathscr{H}_{\varepsilon} \longrightarrow \mathbb{R}$ (see [26] or [27] for example). Now, we prove that I_{ε} satisfies the Palais-Smale condition.

Let $\{v_n\} \in \mathcal{H}_{\varepsilon}$ be such that $I_{\varepsilon}(v_n) \longrightarrow c$ for some c > 0 and $I_{\varepsilon}'(v_n) \longrightarrow 0$. First standard arguments show that $\{v_n\}$ is bounded in $\mathcal{H}_{\varepsilon}$. Thus, there exists $v \in \mathcal{H}_{\varepsilon}$ and if necessary a subsequence of $\{v_n\}$ such that $v_n \rightarrow v$ in $\mathcal{H}_{\varepsilon}$. Since the embedding of $\mathcal{H}_{\varepsilon}$ into L_K^{p+1} is compact (see Theorem 4), we can get a subsequence of $\{v_n\}$ (also denoted by $\{v_n\}$) such that $v_n \longrightarrow v$ in L_K^{p+1} . Note that

$$\begin{split} \left\langle I_{\varepsilon}'(v_{n}) - I_{\varepsilon}'(v), v_{n} - v \right\rangle \\ &= \left(\varepsilon^{2}a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla v_{n}|^{2} \right) \int_{\mathbb{R}^{3}} \nabla v_{n} \cdot \nabla (v_{n} - v) + \int_{\mathbb{R}^{3}} V(x)(v_{n} - v)^{2} \\ &- \left(\varepsilon^{2}a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla v|^{2} \right) \int_{\mathbb{R}^{3}} \nabla v \cdot \nabla (v_{n} - v) - \int_{\mathbb{R}^{3}} K(x) \left(|v_{n}|^{p-1}v_{n} - |v|^{p-1}v \right) \\ &\cdot (v_{n} - v) = \left(\varepsilon^{2}a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla v_{n}|^{2} \right) \cdot \int_{\mathbb{R}^{3}} |\nabla (v_{n} - v)|^{2} + \int_{\mathbb{R}^{3}} V(x)(v_{n} - v)^{2} \\ &+ \varepsilon b \left(\int_{\mathbb{R}^{3}} |\nabla v_{n}|^{2} - \int_{\mathbb{R}^{3}} |\nabla v|^{2} \right) \cdot \int_{\mathbb{R}^{3}} \nabla v \cdot \nabla (v_{n} - v) \\ &- \int_{\mathbb{R}^{3}} K(x) \left(|v_{n}|^{p-1}v_{n} - |v|^{p-1}v \right) (v_{n} - v) \geq ||v_{n} - v||_{\varepsilon}^{2} \\ &- \varepsilon b \left(\int_{\mathbb{R}^{3}} |\nabla v|^{2} - \int_{\mathbb{R}^{3}} |\nabla v_{n}|^{2} \right) \int_{\mathbb{R}^{3}} \nabla v \cdot \nabla (v_{n} - v) \\ &- \int_{\mathbb{R}^{3}} K(x) \left(|v_{n}|^{p-1}v_{n} - |v|^{p-1}v \right) (v_{n} - v), \end{split}$$
(17)

which implies that

$$\|\boldsymbol{v}_{n} - \boldsymbol{v}\|_{\varepsilon}^{2} \leq \left\langle I_{\varepsilon}'(\boldsymbol{v}_{n}) - I_{\varepsilon}'(\boldsymbol{v}), \boldsymbol{v}_{n} - \boldsymbol{v} \right\rangle + \varepsilon b \left(\int_{\mathbb{R}^{3}} |\nabla \boldsymbol{v}|^{2} - \int_{\mathbb{R}^{3}} |\nabla \boldsymbol{v}_{n}|^{2} \right) \cdot \int_{\mathbb{R}^{3}} \nabla \boldsymbol{v} \cdot \nabla (\boldsymbol{v}_{n} - \boldsymbol{v}) + \int_{\mathbb{R}^{3}} K(\boldsymbol{x}) \left(|\boldsymbol{v}_{n}|^{p-1} \boldsymbol{v}_{n} - |\boldsymbol{v}|^{p-1} \boldsymbol{v} \right) (\boldsymbol{v}_{n} - \boldsymbol{v}).$$
(18)

Since $\{v_n\}$ is bounded in $\mathscr{H}_{\varepsilon}, I'_{\varepsilon}(v_n) \longrightarrow 0$ and $v_n \longrightarrow v$ in $\mathscr{H}_{\varepsilon}$, we have that $\langle I_{\varepsilon}'(v_n) - I_{\varepsilon}'(v), v_n - v \rangle \longrightarrow 0$ and $\varepsilon b(\int_{\mathbb{R}^3} |\nabla v_n|^2 - \int_{\mathbb{R}^3} |\nabla v_n|^2) \int_{\mathbb{R}^3} \nabla v \cdot \nabla (v_n - v) \longrightarrow 0$ as $n \longrightarrow \infty$. Furthermore,

$$\begin{split} \left| \int_{\mathbb{R}^{3}} K(x) \left(|v_{n}|^{p-1} v_{n} - |v|^{p-1} v \right) (v_{n} - v) \right| \\ &= \left| \int_{\mathbb{R}^{3}} K^{p/(p+1)}(x) \left(|v_{n}|^{p-1} v_{n} - |v|^{p-1} v \right) K^{1/(p+1)}(x) (v_{n} - v) \right| \\ &\leq \left(\int_{\mathbb{R}^{3}} K(x) |v_{n}|^{p-1} v_{n} - |v|^{p-1} v |^{(p+1)/p} \right)^{p/(p+1)} \\ &\cdot \left(\int_{\mathbb{R}^{3}} K(x) |v_{n} - v|^{p+1} \right)^{1/(p+1)}. \end{split}$$
(19)

Since $\{v_n\}$ is bounded in $L_K^{p+1}(\mathbb{R}^3)$ and $v_n \longrightarrow v$ in $L_K^{p+1}(\mathbb{R}^3)$, we have

$$\int_{\mathbb{R}^3} K(x) \left(|\nu_n|^{p-1} \nu_n - |\nu|^{p-1} \nu \right) (\nu_n - \nu) \longrightarrow 0, \qquad (20)$$

as $n \longrightarrow \infty$. Thus, we have $||v_n - v||_{\varepsilon} \longrightarrow 0$ as $n \longrightarrow \infty$, i.e., $v_n \longrightarrow v$ in $\mathcal{H}_{\varepsilon}$.

From the above arguments, we know that $I_{\varepsilon}: \mathscr{H}_{\varepsilon} \longrightarrow \mathbb{R}$ satisfies the Mountain-Pass geometry and the Palais-Smale condition; hence, by the Mountain-Pass theorem, we can get a critical point $u_{\varepsilon} \in \mathscr{H}_{\varepsilon}$ of I_{ε} with $c_{\varepsilon} = I_{\varepsilon}(u_{\varepsilon})$. As in [26], we can also know that $u_{\varepsilon} > 0$. From

$$-\left(\varepsilon^{2}a+\varepsilon b\int_{\mathbb{R}^{3}}|\nabla u_{\varepsilon}|^{2}\right)\Delta u_{\varepsilon}+V(x)u_{\varepsilon}=K(x)u_{\varepsilon}^{p},\qquad(21)$$

we have

$$\int_{\mathbb{R}^3} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) + \varepsilon b \left(\int_{\mathbb{R}^3} |\nabla u_{\varepsilon}|^2 \right)^2 = \int_{\mathbb{R}^3} K(x) |u_{\varepsilon}|^{p+1}.$$
(22)

Thus,

$$\begin{aligned} c_{\varepsilon} &= I_{\varepsilon}(u_{\varepsilon}) = \left(\frac{1}{2} - \frac{1}{p+1}\right) \int_{\mathbb{R}^3} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2\right) \\ &+ \left(\frac{1}{4} - \frac{1}{p+1}\right) \varepsilon b \left(\int_{\mathbb{R}^3} |\nabla u_{\varepsilon}|^2\right)^2 \ge \left(\frac{1}{2} - \frac{1}{p+1}\right) \|u_{\varepsilon}\|_{\varepsilon}^2, \end{aligned}$$

$$(23)$$

which implies (14).

Next, we will show that the positive ground state solution u_{ε} of (1) found in Theorem 5 belongs indeed to H^1 (\mathbb{R}^3). We first prove

Lemma 6. Let u_{ε} be solutions of (1) found in Theorem 5 and suppose there exists $\Gamma > 0$ such that

$$\|u_{\varepsilon}\|_{\varepsilon}^{2} \leq \Gamma \varepsilon^{3}.$$
 (24)

Then, there exist $K_{\Gamma} > 0$ and $R_{\Gamma} > 0$ such that for $R \ge R_{\Gamma}$ and $\Omega_{n,\varepsilon} \subset \mathbb{R}^3 \setminus B_R$,

$$\int_{\Omega_{n+1,\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) \le \frac{3}{4} \int_{\Omega_{n,\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right),$$
(25)

where $\Omega_{n,\varepsilon} = \mathbb{R}^3 \setminus B_{R_{n,\varepsilon}}$ and $R_{n,\varepsilon} = \varepsilon K_{\Gamma} n^{2/(2-\alpha)}$.where $\overline{C}_{\Gamma} > 0$ is a constant.

Proof. By (24), we have

$$\int_{\mathbb{R}^3} a |\nabla u_{\varepsilon}|^2 \le \Gamma \varepsilon.$$
 (26)

Then

$$\varepsilon^2 a + \varepsilon b \int_{\mathbb{R}^3} |\nabla u_{\varepsilon}|^2 \le \varepsilon^2 \left(a + b \frac{\Gamma}{a} \right). \tag{27}$$

Let $M_1 \coloneqq a + b(\Gamma/a)$, and choose $K_{\Gamma} > 0$ large enough such that

$$\frac{M_1^2}{a} K_{\Gamma}^{\alpha-2} \le \frac{\bar{a}}{2},\tag{28}$$

where \bar{a} is defined in (V). Let $R_{n,\varepsilon} \coloneqq \varepsilon K_{\Gamma} n^{2/(2-\alpha)}$, and let $\chi_{n,\varepsilon}(r)$ be piecewise affine functions such that

$$\chi_{n,\varepsilon}(r) = 0 \text{ for } r \le R_{n,\varepsilon},$$

$$\chi_{n,\varepsilon}(r) = 1 \text{ for } r \ge R_{n+1,\varepsilon}.$$
(29)

Then

$$\begin{split} \left| R_{n+1,\varepsilon} - R_{n,\varepsilon} \right| &= \varepsilon K_{\Gamma} \left| (n+1)^{2/(2-\alpha)} - n^{2/(2-\alpha)} \right| \\ &= \varepsilon K_{\Gamma} (n+1)^{\alpha/(2-\alpha)} \left| (n+1) - n \left(\frac{n}{n+1}\right)^{\alpha/(2-\alpha)} \right| \\ &= \varepsilon^{(2-\alpha)/2} K_{\Gamma}^{(2-\alpha)/2} \left(\varepsilon K_{\Gamma} (n+1)^{2/(2-\alpha)} \right)^{\alpha/2} \\ &\cdot \left| (n+1) - n \left(\frac{n}{n+1}\right)^{\alpha/(2-\alpha)} \right| \\ &\geq \varepsilon^{(2-\alpha)/2} K_{\Gamma}^{(2-\alpha)/2} R_{n+1,\varepsilon}^{\alpha/2}, \end{split}$$
(30)

which yields

$$\left|R_{n+1,\varepsilon} - R_{n,\varepsilon}\right|^{-2} \le \varepsilon^{\alpha-2} K_{\Gamma}^{\alpha-2} R_{n+1,\varepsilon}^{-\alpha}.$$
 (31)

By (28) and (31), we have

$$\frac{M_1^2}{a}\varepsilon^2 \left| \nabla \chi_{n,\varepsilon}(x) \right|^2 = \frac{M_1^2}{a}\varepsilon^2 \left| R_{n+1,\varepsilon} - R_{n,\varepsilon} \right|^{-2} \le \frac{M_1^2}{a}\varepsilon^{\alpha} K_{\Gamma}^{\alpha-2} R_{n+1,\varepsilon}^{-\alpha}$$
$$\le \frac{M_1^2}{a} K_{\Gamma}^{\alpha-2} R_{n+1,\varepsilon}^{-\alpha} \le \frac{\bar{a}}{2} R_{n+1,\varepsilon}^{-\alpha} \le V(x), x \in \mathbb{R}^3.$$
(32)

Now, we test (1) on $\chi_{n,\varepsilon}u_{\varepsilon}$. We get

$$\left(\varepsilon^{2} a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla u_{\varepsilon}|^{2} \right) \int_{\Omega_{n,\varepsilon}} \chi_{n,\varepsilon} |\nabla u_{\varepsilon}|^{2} + \int_{\Omega_{n,\varepsilon}} \chi_{n,\varepsilon} V u_{\varepsilon}^{2}$$

$$= \int_{\Omega_{n,\varepsilon}} \chi_{n,\varepsilon} K u_{\varepsilon}^{p+1} - \left(\varepsilon^{2} a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla u_{\varepsilon}|^{2} \right) \int_{\Omega_{n,\varepsilon}} \nabla u_{\varepsilon} \nabla \chi_{n,\varepsilon} u_{\varepsilon}.$$

$$(33)$$

Now, by (32), we have

$$-\left(\varepsilon^{2}a + \varepsilon b \int_{\mathbb{R}^{3}} |\nabla u_{\varepsilon}|^{2}\right) \int_{\Omega_{n\varepsilon}} \nabla u_{\varepsilon} \nabla \chi_{n,\varepsilon} u_{\varepsilon}$$

$$\leq \varepsilon^{2} M_{1} \int_{\Omega_{n\varepsilon}} |\nabla u_{\varepsilon}| |\nabla \chi_{n,\varepsilon}| |u_{\varepsilon}|$$

$$= \int_{\Omega_{n\varepsilon}} \left(\varepsilon \sqrt{a} |\nabla u_{\varepsilon}|\right) \left(\varepsilon \frac{1}{\sqrt{a}} M_{1} |\nabla \chi_{n,\varepsilon}| |u_{\varepsilon}|\right) \qquad (34)$$

$$\leq \frac{1}{2} \int_{\Omega_{n\varepsilon}} \varepsilon^{2} a |\nabla u_{\varepsilon}|^{2} + \frac{M_{1}^{2}}{a} \varepsilon^{2} |\nabla \chi_{n,\varepsilon}(x)|^{2} u_{\varepsilon}^{2}$$

$$\leq \frac{1}{2} \int_{\Omega_{n\varepsilon}} \varepsilon^{2} a |\nabla u_{\varepsilon}|^{2} + V(x) u_{\varepsilon}^{2}.$$

Then (33) and (34) imply that

$$\begin{split} \int_{\Omega_{n+1,\varepsilon}} (\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2) &\leq \int_{\Omega_{n,\varepsilon}} \chi_{n,\varepsilon} (\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2) \\ &\leq \left(\varepsilon^2 a + \varepsilon b \int_{\mathbb{R}^3} |\nabla u_{\varepsilon}|^2 \right) \int_{\Omega_{n,\varepsilon}} \chi_{n,\varepsilon} |\nabla u_{\varepsilon}|^2 \\ &+ \int_{\Omega_{n,\varepsilon}} \chi_{n,\varepsilon} V u_{\varepsilon}^2 \leq \int_{\Omega_{n,\varepsilon}} K u_{\varepsilon}^{p+1} \\ &+ \frac{1}{2} \int_{\Omega_{n,\varepsilon}} (\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2). \end{split}$$
(35)

Then, by Proposition 11 in [37], let $\delta > 0$ be fixed, and for sufficiently large R > 0, we have

$$\int_{\Omega_{n+1,\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right)
\leq \frac{1}{2} \int_{\Omega_{n\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right)
+ \delta \varepsilon^{-3(p-1)/2} \left(\int_{\Omega_{n\varepsilon}} \varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right)^{(p+1)/2}.$$
(36)

By (24), we have

$$\left(\int_{\Omega_{n,\varepsilon}} \varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2\right)^{(p+1)/2}$$

$$\leq \Gamma^{(p-1)/2} \varepsilon^{3(p-1)/2} \int_{\Omega_{n,\varepsilon}} \varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2.$$
(37)

Thus, by the above two estimates, we know

$$\int_{\Omega_{n+1,\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) \le \left(\frac{1}{2} + \delta \Gamma^{(p-1)/2} \right) \int_{\Omega_{n\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right).$$
(38)

Choosing $\delta > 0$ sufficiently small, we can get that

$$\int_{\Omega_{n+1,\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) \le \frac{3}{4} \int_{\Omega_{n,\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right).$$
(39)

Lemma 7. Let u_{ε} be solutions of (1), and let Γ , R_{Γ} , and K_{Γ} be as in Lemma 6. Then, for all $\rho \ge 2R_{\Gamma}$,

$$\begin{aligned} \int_{|x|>\rho} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) \\ &\leq \bar{C}_{\Gamma} \varepsilon^3 \exp\left\{ -\frac{1}{2} \left| \log \frac{3}{4} \right| (K_{\Gamma} \varepsilon)^{-(2-\alpha)/2} \left(\rho^{(2-\alpha)/2} - R_{\Gamma}^{(2-\alpha)/2} \right) \right\}, \end{aligned}$$

$$\tag{40}$$

Proof. The proof is just as the proof of Lemma 18 in [37], we give it here for the sake of completeness.

Let $\rho > 2R_{\Gamma}$ and choosing two positive integers $\tilde{n} > \bar{n}$ with

$$\begin{aligned} &R_{\bar{n},\varepsilon} \le \rho \le R_{\bar{n}+1,\varepsilon}, \\ &R_{\bar{n},\varepsilon} \le R_{\Gamma} \le R_{\bar{n}+1,\varepsilon}, \end{aligned} \tag{41}$$

where $R_{n,\varepsilon}$ is defined in Lemma 6. From Lemma 6, we can know that

$$\begin{split} \int_{|x|>\rho} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) &\leq \int_{|x|>R_{\bar{n},\varepsilon}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right) \\ &\leq \left(\frac{3}{4} \right)^{\bar{n}-\bar{n}} \int_{|x|>R_{\Gamma}} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2 \right). \end{split}$$

$$(42)$$

Then, from (24), we have

$$\int_{|x|>\rho} \left(\varepsilon^2 a |\nabla u_{\varepsilon}|^2 + V(x) u_{\varepsilon}^2\right) \le \left(\frac{3}{4}\right)^{\bar{n}-\bar{n}} \Gamma \varepsilon^3.$$
(43)

By the choices of \tilde{n}, \bar{n} ,

$$\rho \sim \varepsilon K_{\Gamma} \tilde{n}^{2/(2-\alpha)},$$

$$R_{\Gamma} \sim \varepsilon K_{\Gamma} \bar{n}^{2/(2-\alpha)},$$
(44)

which implies

$$\tilde{n} - \bar{n} \ge \frac{1}{2} \left(K_{\Gamma} \varepsilon \right)^{-(2-\alpha)/2} \left(\rho^{(2-\alpha)/2} - R_{\Gamma}^{(2-\alpha)/2} \right).$$
(45)

Then, from (43) and the above formula, we can get the estimate (40).

Now, we are in position to prove Theorem 1. To prove $u_{\varepsilon} \in H^1(\mathbb{R}^3)$, we actually need to show that $u_{\varepsilon} \in L^2(\mathbb{R}^3)$. In the following, we follow [37] to give a proof here.

First, for the simplicity of the notation, we can take ε to 1. Let $u \in \mathcal{H}$ be a solution of (1)(with $\varepsilon = 1$), and we choose $y \in \mathbb{R}^3$ such that |y| > 2. Then

$$\int_{B_1(y)} u^2 = \int_{B_1(y)} V(x) u^2 \cdot \frac{1}{V(x)} \le c_1 |y|^{\alpha} \int_{B_1(y)} V(x) u^2.$$
(46)

For R = (1/2)|y| we know that

$$\int_{B_1(y)} V(x)u^2 \le \int_{\mathbb{R}^3 \setminus B_R} V(x)u^2.$$
(47)

From the above two estimates and Lemma 2.8, we have

$$\int_{B_1(y)} u^2 \le C_1 |y|^{\alpha} \exp\left\{-C_2 |y|^{1-\alpha/2}\right\}, \forall |y| \gg 1.$$
(48)

Let $B_5 \setminus B_2 \subset \bigcup_{i=1}^m B_1(y_i)$ where $m \in \mathbb{N}$ and $y_i \in \mathbb{R}^3$. Define $y_{i,k} \coloneqq 2^k y_i$. Then

$$\int_{\mathbb{R}^{3}\setminus B_{2}} u^{2} \leq \sum_{k=0}^{\infty} \int_{2^{k}(B_{5}\setminus B_{2})} u^{2} \leq \sum_{i,k} \int_{B_{2^{k}}(y_{i,k})} u^{2}.$$
 (49)

By (48) and $0 < \alpha < 2$, we have

$$\int_{\mathbb{R}^{3} \setminus B_{2}} u^{2} \leq C_{1} \sum_{i,k} |y_{i,k}|^{\alpha} \exp\left\{-C_{2} |y_{i,k}|^{1-\alpha/2}\right\} < \infty.$$
(50)

Thus, we have that $u \in L^2(\mathbb{R}^3)$, whence $u \in H^1(\mathbb{R}^3)$. Finally, standard elliptic estimate (see Theorem 4.1 in [39] for example) shows that $\lim_{|x| \to \infty} u_{\varepsilon}(x) = 0$.

3. Concentration Behaviour of Ground State Solutions

In this section, we will study the concentration behaviour of the positive ground state solutions $\{u_{\varepsilon}\}$ obtained in Theorem 1 as $\varepsilon \longrightarrow 0^+$. Assume that (V) and (K) are satisfied, and let (5) hold. By (5), we can know that $\mathscr{A}(x) \longrightarrow \infty$ as $|x| \longrightarrow \infty$, where $\mathscr{A}(x) = V^{2/(p-1)-1/2}(x)/K^{2/(p-1)}(x)$ is defined in Theorem 2, and therefore, $\mathscr{A}(x)$ has a global minimum point x_0 in \mathbb{R}^3 . Now, let $u_0 > 0$ be the unique positive ground state solution (see [26]) of

$$-\left(a+b\int_{\mathbb{R}^3}|\nabla u|^2\right)\Delta u+V(x_0)u=K(x_0)u^p, x\in\mathbb{R}^3, \quad (51)$$

and let c_0 be the ground energy level associated to (51),

i.e., $c_0 = I_0(u_o)$ where

$$I_{0}(u) \coloneqq \frac{1}{2} \int_{\mathbb{R}^{3}} \left(a |\nabla u|^{2} + V(x_{0})u^{2} \right) + \frac{b}{4} \left(\int_{\mathbb{R}^{3}} |\nabla u|^{2} \right)^{2} - \frac{1}{p+1} \int_{\mathbb{R}^{3}} K(x_{0}) |u|^{p+1}.$$
(52)

Let c_{ε} be as in Theorem 5, then we have the following lemma.

Lemma 8.

$$\limsup_{\varepsilon \longrightarrow 0^+} \varepsilon^{-3} c_{\varepsilon} \le c_0.$$
(53)

Proof. Since the proof is just the same as the proof of Lemma 11 in [26], we omit it here. \Box

Remark 9. In particular, there exists $C^* > 0$ such that $c_{\varepsilon} \leq C^* \varepsilon^3$ for sufficiently small $\varepsilon > 0$.

Corollary 10. Let u_{ε} be the ground state solutions obtained in Theorem 1. Then by Lemma 8 and (14), we have that there exists $\Gamma > 0$ such that for sufficiently small $\varepsilon > 0$,

$$\|u_{\varepsilon}\|_{\varepsilon}^{2} \leq \Gamma \varepsilon^{3}.$$
(54)

Now, we give a uniform pointwise decay estimate for the solutions u_{ε} :

Lemma 11. Let Γ , R_{Γ} , K_{Γ} and u_{ε} be as in Lemma 6. Then, there exist constants C > 0 and d > 0, such that for $|x| \ge 2R_{\Gamma} + C$,

$$|u_{\varepsilon}(x)| \le C|x|^{d} \varepsilon^{-d} \exp\left\{-\frac{1}{4} \left|\log \frac{3}{4}\right| (K_{\Gamma} \varepsilon)^{-(2-\alpha)/2} \left(|x|^{(2-\alpha)/2} - R_{\Gamma}^{(2-\alpha)/2}\right)\right\},$$
(55)

where C depends only on Γ and p and d depends on p, α , and β .

Proof. Let $x_0 \in \mathbb{R}^3$ be such that $|x_0| \ge 2R_{\Gamma} + 2$, and let η be a smooth cut-off function satisfying $\eta = 1$ for $x \in B_1(x_0)$, $\eta = 0$ for $x \in \mathbb{R}^3 \setminus B_2(x_0)$, and $|\nabla \eta| \le 2$. For simplicity of notation, we let $v = u_{\varepsilon}$. For L > 0 and $s \ge 0$, define $\phi = \phi_{s,L} := v$ min $\{|v|^{2s}, L^2\}\eta^2$. The function v satisfies

$$-\left(\varepsilon^{2}a+\varepsilon b\int_{\mathbb{R}^{3}}|\nabla v|^{2}\right)\Delta v+V(x)v=K(x)v^{p},$$
(56)

and by (24), we have $\int_{\mathbb{R}^3} \varepsilon^2 a |\nabla v|^2 \leq \Gamma \varepsilon^3$, which implies that $\int_{\mathbb{R}^3} |\nabla v|^2 \leq (\Gamma/a)\varepsilon$. Now, test (56) on ϕ we can obtain

that

$$\begin{aligned} \varepsilon^{2} \int a |\nabla v|^{2} \min \left\{ |v|^{2s}, L^{2} \right\} \eta^{2} \\ &+ \frac{s}{2} \varepsilon^{2} \int_{\{ |v|^{s} \leq L \}} a |\nabla (|v|^{2})|^{2} v^{2s-2} \eta^{2} + \int V(x) v^{2} \eta^{2} \min \left\{ |v|^{2s}, L^{2} \right\} \\ &\leq -2 \varepsilon^{2} \int av\eta \min \left\{ |v|^{2s}, L^{2} \right\} \nabla v \cdot \nabla \eta - 2 \varepsilon b \int_{\mathbb{R}^{3}} |\nabla v|^{2} \\ &\cdot \int v\eta \min \left\{ |v|^{2s}, L^{2} \right\} \nabla v \cdot \nabla \eta + \int K v^{p+1} \eta^{2} \min \left\{ |v|^{2s}, L^{2} \right\} \\ &\leq 2 \varepsilon^{2} \int av\eta \min \left\{ |v|^{2s}, L^{2} \right\} |\nabla v \cdot \nabla \eta| \\ &+ \frac{2 \varepsilon^{2} b \Gamma}{a} \int v\eta \min \left\{ |v|^{2s}, L^{2} \right\} |\nabla v \cdot \nabla \eta| \\ &+ \frac{2 \varepsilon^{2} b \Gamma}{a} \int v\eta \min \left\{ |v|^{2s}, L^{2} \right\} |\nabla v \cdot \nabla \eta| \\ &+ \frac{2 \varepsilon^{2} b \Gamma}{a} \int v\eta \min \left\{ |v|^{2s}, L^{2} \right\} |\nabla v \cdot \nabla \eta| \\ &+ C \varepsilon^{2} \int av^{2} \min \left\{ |v|^{2s}, L^{2} \right\} |\nabla \eta|^{2} + \int K v^{p+1} \eta^{2} \min \left\{ |v|^{2s}, L^{2} \right\}. \end{aligned}$$

$$(57)$$

Next, we can follow the proof of Lemma 22 in [37] directly, so we omit the details here.

By Lemma 11, we know that $\lim_{|x|\to\infty} u_{\varepsilon} = 0$ for any small and fixed $\varepsilon > 0$. Thus, there exists a maximum point x_{ε} in \mathbb{R}^3 for u_{ε} . We have the following lemma.

Lemma 12. Let u_{ε} be the solution of (1) satisfying (24), and let x_{ε} be any maximum point of u_{ε} . Then, for sufficiently small $\varepsilon > 0$, $|x_{\varepsilon}| \le C$, where $C = C(\Gamma)$.

Proof. As x_{ε} is a maximum point of u_{ε} , we know that $\Delta u_{\varepsilon}(x_{\varepsilon}) \leq 0$, and furthermore, $-(\varepsilon^2 a + \varepsilon b \int_{\mathbb{R}^3} |\nabla u_{\varepsilon}|^2) \Delta u_{\varepsilon}(x_{\varepsilon}) \geq 0$. Therefore, from (1), we have that

$$V(x_{\varepsilon})K^{-1}(x_{\varepsilon}) \le u_{\varepsilon}^{p-1}(x_{\varepsilon}).$$
(58)

Now, by (V) and (K), we know that

$$c|x_{\varepsilon}|^{\beta-\alpha} \le V(x_{\varepsilon})K^{-1}(x_{\varepsilon}), \tag{59}$$

where c > 0 is a constant. Combining (58), (59), and (55), we have that for $|x_{\varepsilon}| \ge 2R_{\Gamma}$,

$$\begin{aligned} c|x_{\varepsilon}|^{\beta-\alpha} &\leq |x_{\varepsilon}|^{d(p-1)} \varepsilon^{-d(p-1)} \\ &\cdot \exp\left\{-\frac{1}{4}(p-1)\left|\log\frac{3}{4}\right| (K_{\Gamma}\varepsilon)^{-(2-\alpha)/2} \left(|x_{\varepsilon}|^{(2-\alpha)/2} - R_{\Gamma}^{(2-\alpha)/2}\right)\right\}. \end{aligned}$$

$$\tag{60}$$

This shows that $|x_{\varepsilon}|$ must be bounded as $\varepsilon \longrightarrow 0^+$, and Lemma 12 has been proved.

Lemma 13. Let u_{ε} be as in Lemma 12. Then, for sufficiently small $\varepsilon > 0$, $||u_{\varepsilon}||_{L^{\infty}} \ge C$ where C > 0 is a constant.

Proof. From (1) we know that

$$\|u_{\varepsilon}\|_{\varepsilon}^{2} = \int_{\mathbb{R}^{3}} \left(\varepsilon^{2} a |\nabla u_{\varepsilon}|^{2} + V(x) u_{\varepsilon}^{2}\right) \leq \int_{\mathbb{R}^{3}} K(x) u_{\varepsilon}^{p+1}.$$
(61)

Now, let $\delta < \Gamma^{-(p-1)/2}$ be fixed. By Proposition 11 in [37], we know that there exists R > 0 such that

$$\int_{|x|>R} K(x) u_{\varepsilon}^{p+1} \le \delta \varepsilon^{-3(p-1)/2} \|u_{\varepsilon}\|_{\varepsilon}^{p+1}.$$
(62)

Then, from the proof of Lemma 25 in [37], we have

$$\int_{|x|\leq R} K(x) u_{\varepsilon}^{p+1} \leq \frac{k}{\bar{a}} (1+R^{\alpha}) \left\| u_{\varepsilon}^{p-1} \right\|_{L^{\infty}} \| u_{\varepsilon} \|_{\varepsilon}^{2}.$$
(63)

Combining (61), (62), and (63), we get

$$\begin{aligned} \|u_{\varepsilon}\|_{\varepsilon}^{2} &\leq \int_{\mathbb{R}^{3}} K(x) u_{\varepsilon}^{p+1} \leq \delta \varepsilon^{-3(p-1)/2} \|u_{\varepsilon}\|_{\varepsilon}^{p+1} \\ &+ \frac{k}{\bar{a}} (1+R^{\alpha}) \|u_{\varepsilon}^{p-1}\|_{L^{\infty}} \|u_{\varepsilon}\|_{\varepsilon}^{2}, \end{aligned}$$
(64)

which implies that

$$1 \le \delta \varepsilon^{-3(p-1)/2} \| u_{\varepsilon} \|_{\varepsilon}^{p-1} + \frac{k}{\bar{a}} (1+R^{\alpha}) \| u_{\varepsilon}^{p-1} \|_{L^{\infty}}.$$
 (65)

By Corollary 10, we know that $||u_{\varepsilon}||_{\varepsilon}^{p-1} \leq \Gamma^{(p-1)/2} \varepsilon^{3(p-1)/2}$. Then, from (65), we get

$$1 \le \delta \Gamma^{(p-1)/2} + \frac{k}{\bar{a}} (1+R^{\alpha}) \left\| u_{\varepsilon}^{p-1} \right\|_{L^{\infty}}.$$
 (66)

Now, by $\delta < \Gamma^{-(p-1)/2}$, we have that

$$\|u_{\varepsilon}\|_{\varepsilon}^{p-1} \ge \left(1 - \delta\Gamma^{(p-1)/2}\right) \frac{\bar{a}}{k(1+R^{\alpha})} > 0.$$
(67)

Thus, we have proved the lemma.

Proof of Theorem 14. Let x_{ε} be a global maximum point of u_{ε} . From Lemma 12, we know that going to a subsequence if necessary, there exists $x^* \in \mathbb{R}^3$ such that $x_{\varepsilon} \longrightarrow x^*$. Since x_0 is a global minimum point of $\mathscr{A}(x)$ in \mathbb{R}^3 , then by Lemma 13 of [26], we have

$$c(x^*) \ge c_0, \tag{68}$$

where $c(x^*)$ is the ground energy level corresponding to equation (51) with x_0 replaced by x^* in it. Set $\phi_{\varepsilon}(x) \coloneqq u_{\varepsilon}$ $(\varepsilon x + x_{\varepsilon})$, then ϕ_{ε} satisfies

$$-\left(a+b\int_{\mathbb{R}^{3}}\left|\nabla\phi_{\varepsilon}\right|^{2}\right)\Delta\phi_{\varepsilon}+V(\varepsilon x+x_{\varepsilon})\phi_{\varepsilon}=K(\varepsilon x+x_{\varepsilon})\phi_{\varepsilon}^{p}.$$
(69)

From Corollary 10 and (V), we know that

$$\Gamma \geq \varepsilon^{-3} \|u_{\varepsilon}\|_{\varepsilon}^{2} \geq \varepsilon^{-3} \int_{\mathbb{R}^{3}} \varepsilon^{2} a |\nabla \phi_{\varepsilon}|^{2} + \frac{\bar{a}}{1+|x|^{\alpha}} u_{\varepsilon}^{2}$$

$$= \int_{\mathbb{R}^{3}} a |\nabla \phi_{\varepsilon}|^{2} + \frac{\bar{a}}{1+|\varepsilon x+x_{\varepsilon}|^{\alpha}} \phi_{\varepsilon}^{2}.$$
(70)

From Lemma 12, we know that $|\varepsilon x + x_{\varepsilon}| \le C(1 + |x|)$, and therefore,

$$\int_{\mathbb{R}^3} a |\nabla \phi_{\varepsilon}|^2 + \frac{\bar{a}}{1+|x|^{\alpha}} \phi_{\varepsilon}^2 \le C',$$
(71)

where C' > 0 is a constant. Then $\{\phi_{\varepsilon}\}$ is bounded in $\overline{H} := \{v \in \mathcal{D}^{1,2}(\mathbb{R}^3): \int_{\mathbb{R}^3} \overline{a}/1 + |x|^{\alpha} v^2 < +\infty\}$, and therefore, going to a subsequence if necessary there exists $U^* \in \overline{H}$ such that $\phi_{\varepsilon} \to U^*$ in $\overline{H}, \phi_{\varepsilon} \longrightarrow U^*$ in L_K^{p+1} , and $\phi_{\varepsilon} \longrightarrow U^*$ a.e. in \mathbb{R}^3 . By standard arguments (see [37]), we have that ϕ_{ε} converges in $C_{loc}^2(\mathbb{R}^3)$ to U^* and $\phi_{\varepsilon} \longrightarrow U^*$ in $L^{\infty}(\mathbb{R}^3)$. Now, we assume that up to a subsequence, $\int_{\mathbb{R}^3} |\nabla \phi_{\varepsilon}|^2 \longrightarrow A$, then $A \ge \int_{\mathbb{R}^3} |\nabla U^*|^2$ and

$$-(a+bA)\Delta U^* + V(x^*)U^* = K(x^*)(U^*)^p, x \in \mathbb{R}^3,$$
 (72)

since 0 is the maximum point of ϕ_{ε} , so does U^* . Moreover, Lemma 13 shows that $\phi_{\varepsilon}(0) = u_{\varepsilon}(x_{\varepsilon}) = ||u_{\varepsilon}||_{L^{\infty}} \ge C$ for some C > 0, thus $U^*(0) = \max U^* \ge C > 0$. Therefore, $U^* \ne 0$ and by the maximum principle we have that $U^* > 0$. Now, for $R_n \longrightarrow \infty$, by Corollary 10, we have

$$\int_{\bar{B}_{R_n}} a |\nabla \phi_{\varepsilon}|^2 + V(\varepsilon x + x_{\varepsilon}) \phi_{\varepsilon}^2 \le \varepsilon^{-3} ||v_{\varepsilon}||_{\varepsilon}^2 \le \Gamma.$$
(73)

Thus, by the Dominated Convergence Theorem, $\phi_{\varepsilon} \longrightarrow U^*$ in $C^1(\bar{B}_{R_n})$, and (73) we have

$$\int_{\bar{B}_{R_n}} a |\nabla U^*|^2 + V(x^*) (U^*)^2 \le \Gamma.$$
(74)

Now, let $R_n \longrightarrow \infty$, we get that $U^* \in H^1(\mathbb{R}^3)$.

Next, we prove actually $A = \int_{\mathbb{R}^3} |\nabla U^*|^2$, where *A* is in (72). Otherwise, we assume $A > \int_{\mathbb{R}^3} |\nabla U^*|^2$. Then, from (72), we have

$$\int_{\mathbb{R}^{3}} \left(a |\nabla U^{*}|^{2} + V(x^{*})(U^{*})^{2} \right) + Ab \int_{\mathbb{R}^{3}} |\nabla U^{*}|^{2} - \int_{\mathbb{R}^{3}} K(x^{*}) |U^{*}|^{p+1} = 0,$$
(75)

which implies that $\langle I_{x^*}'(U^*), U^* \rangle < 0$, where $\langle I_{x^*}'(U^*), U^* \rangle$ is defined by

$$\left\langle I_{x^*}'(U^*), U^* \right\rangle = \int_{\mathbb{R}^3} \left(a |\nabla U^*|^2 + V(x^*)(U^*)^2 \right)$$

+ $b \left(\int_{\mathbb{R}^3} |\nabla U^*|^2 \right)^2 - \int_{\mathbb{R}^3} K(x^*) |U^*|^{p+1}.$ (76)

Then, there is a unique 0 < t < 1 such that

$$\left\langle I_{x^{*}}'(tU^{*}), tU^{*} \right\rangle = 0.$$
 (77)

Now,

$$\begin{split} c(x^{*}) &\leq I_{x^{*}}(tU^{*}) - \frac{1}{4} \left\langle I_{x^{*}}'(tU^{*}), tU^{*} \right\rangle \\ &= \frac{t^{2}}{4} \int_{\mathbb{R}^{3}} \left(a |\nabla U^{*}|^{2} + V(x^{*})U^{*2} \right) \\ &+ \left(\frac{t^{p+1}}{4} - \frac{t^{p+1}}{p+1} \right) \int_{\mathbb{R}^{3}} K(x^{*}) |U^{*}|^{p+1} \\ &\leq \frac{1}{4} \int_{\mathbb{R}^{3}} \left(a |\nabla U^{*}|^{2} + V(x^{*})U^{*2} \right) \\ &+ \left(\frac{1}{4} - \frac{1}{p+1} \right) \int_{\mathbb{R}^{3}} K(x^{*}) |U^{*}|^{p+1} \\ &\leq \underbrace{\lim_{\varepsilon \longrightarrow 0}} \left(\frac{1}{4} \int_{\mathbb{R}^{3}} \left(a |\nabla \phi_{\varepsilon}|^{2} + V(\varepsilon x + x_{\varepsilon}) \phi_{\varepsilon}^{2} \right) \\ &+ \left(\frac{1}{4} - \frac{1}{p+1} \right) \int_{\mathbb{R}^{3}} K(\varepsilon x + x_{\varepsilon}) |\phi_{\varepsilon}|^{p+1} \right) \\ &= \underbrace{\lim_{\varepsilon \longrightarrow 0}} \left(\overline{I}_{\varepsilon}(\phi_{\varepsilon}) - \frac{1}{4} \left\langle \overline{I}'_{\varepsilon}(\phi_{\varepsilon}), \phi_{\varepsilon} \right\rangle \right), \end{split}$$
(78)

where $\bar{I}_{\varepsilon}(v)$: $\mathscr{H}_{\varepsilon} \longrightarrow \mathbb{R}$ is defined by

$$\bar{I}_{\varepsilon}(\nu) \coloneqq \frac{1}{2} \int_{\mathbb{R}^{3}} \left(a |\nabla \nu|^{2} + V(\varepsilon x + x_{\varepsilon})\nu^{2} \right) + \frac{b}{4} \left(\int_{\mathbb{R}^{3}} |\nabla \nu|^{2} \right)^{2} - \frac{1}{p+1} \int_{\mathbb{R}^{3}} K(\varepsilon x + x_{\varepsilon}) |\nu|^{p+1}.$$
(79)

Then, by Lemma 8,

$$\underline{\lim_{\varepsilon \to 0}} \bar{I}_{\varepsilon}(\phi_{\varepsilon}) = \underline{\lim_{\varepsilon \to 0}} \varepsilon^{-3} I_{\varepsilon}(v_{\varepsilon}) = \underline{\lim_{\varepsilon \to 0}} \varepsilon^{-3} c_{\varepsilon} \leq \underline{\lim_{\varepsilon \to 0}} \varepsilon^{-3} c_{\varepsilon} \leq c_{0}.$$
(80)

Thus, by (78), we have $c(x^*) < c_0$ which is a contradiction to (68), and hence, we have proved $A = \int_{\mathbb{R}^3} |\nabla U^*|^2$. From the above arguments, we know that U^* is the unique positive ground state solution of

$$-\left(a+b\int_{\mathbb{R}^{3}}|\nabla U^{*}|^{2}\right)\Delta U^{*}+V(x^{*})U^{*}=K(x^{*})(U^{*})^{p}.$$
 (81)

Furthermore, from (78), we know that $c(x^*) = c_0$, which implies that x^* is a minimum point of the function

 $\mathscr{A}(x) = V^{2/(p-1)-1/2}(x)/K^{2/(p-1)}(x)$. Since the proof of the uniqueness of maximum point is just similar as in [37], we omit it here. Now, we have proved Theorem 2.

Data Availability

The author confirm that the data supporting the findings of this study are available within the article.

Conflicts of Interest

The author declares that they have no conflicts of interest.

Acknowledgments

The paper is supported by the Natural Science Foundation of Shandong Province (ZR2021MA087).

References

- [1] G. Kirchhoff, Mechanik, Teubner, Leipzig, 1883.
- [2] S. Bernstein, "Sur une classe d'équations fonctionnelles aux dérivées partielles," *Izvestiya Akademii Nauk SSSR. Seriya Matematicheskaya*, vol. 4, pp. 17–26, 1940.
- [3] S. I. Pohožaev, "A certain class of quasilinear hyperbolic equations," *Matematicheskii Sbornik*, vol. 96, no. 138, pp. 152–166, 1975.
- [4] A. Azzollini, "The elliptic Kirchhoff equation in R^N perturbed by a local nonlinearity," *Differential and Integral Equations*, vol. 25, no. 5-6, pp. 543–554, 2012.
- [5] G. B. Li and H. Y. Ye, "Existence of positive ground state solutions for the nonlinear Kirchhoff type equations in R³," *Journal of Differential Equations*, vol. 257, no. 2, pp. 566–600, 2014.
- [6] Y. H. Li, F. Y. Li, and J. P. Shi, "Existence of a positive solution to Kirchhoff type problems without compactness conditions," *Journal of Differential Equations*, vol. 253, no. 7, pp. 2285– 2294, 2012.
- [7] Z. P. Liang, F. Y. Li, and J. P. Shi, "Positive solutions to Kirchhoff type equations with nonlinearity having prescribed asymptotic behavior," *Annales de l'IHP Analyse non linéaire*, vol. 31, no. 1, pp. 155–167, 2014.
- [8] D. Naimen, "The critical problem of Kirchhoff type elliptic equations in dimension four," *Journal of Differential Equations*, vol. 257, no. 4, pp. 1168–1193, 2014.
- [9] K. Perera and Z. T. Zhang, "Nontrivial solutions of Kirchhofftype problems via the Yang index," *Journal of Differential Equations*, vol. 221, no. 1, pp. 246–255, 2006.
- [10] X. Wu, "Existence of nontrivial solutions and high energy solutions for Schrödinger-Kirchhoff-type equations in R^N," *Nonlinear Analysis: Real World Applications*, vol. 12, no. 2, pp. 1278–1287, 2011.
- [11] Y. Baoqiang, D. O'Regan, and R. P. Agarwal, "Existence of solutions for Kirchhoff-type problems via the method of lower and upper solutions," *Electronic Journal of Differential Equations*, vol. 54, pp. 1–19, 2019.
- [12] Z. T. Zhang and K. Perera, "Sign changing solutions of Kirchhoff type problems via invariant sets of descent flow," *Journal* of Mathematical Analysis and Applications, vol. 317, no. 2, pp. 456–463, 2006.

- [13] A. Ambrosetti, M. Badiale, and S. Cingolani, "Semiclassical states of nonlinear Schrödinger equations," *Archive for Rational Mechanics and Analysis*, vol. 140, no. 3, pp. 285–300, 1997.
- [14] J. Byeon and L. Jeanjean, "Standing waves for nonlinear Schrödinger equations with a general nonlinearity," Archive for Rational Mechanics and Analysis, vol. 185, no. 2, pp. 185– 200, 2007.
- [15] S. Cingolani and N. Lazzo, "Multiple positive solutions to nonlinear Schrödinger equations with competing potential functions," *Journal of Differential Equations*, vol. 160, no. 1, pp. 118–138, 2000.
- [16] A. Floer and A. Weinstein, "Nonspreading wave packets for the cubic Schrödinger equation with a bounded potential," *Journal of Functional Analysis*, vol. 69, no. 3, pp. 397–408, 1986.
- [17] M. Del Pino and P. L. Felmer, "Local mountain passes for semilinear elliptic problems in unbounded domains," *Calculus* of Variations and Partial Differential Equations, vol. 4, no. 2, pp. 121–137, 1996.
- [18] P. H. Rabinowitz, "On a class of nonlinear Schrödinger equations," Zeitschrift für Angewandte Mathematik und Physik, vol. 43, no. 2, pp. 270–291, 1992.
- [19] X. F. Wang, "On concentration of positive bound states of nonlinear Schrödinger equations," *Communications in Mathematical Physics*, vol. 53, no. 2, pp. 224–229, 1993.
- [20] X. F. Wang and B. Zeng, "On concentration of positive bound states of nonlinear Schrödinger equations with competing potential functions," *SIAM Journal on Mathematical Analysis*, vol. 28, no. 3, pp. 633–655, 1997.
- [21] J. JZhang, W. Zhang, and V. D. Rădulescu, "Double phase problems with competing potentials: concentration and multiplication of ground states," *Mathematische Zeitschrift*, vol. 301, no. 3, pp. 4037–4078, 2022.
- [22] J. Zhang and W. Zhang, "Semiclassical states for coupled nonlinear Schrödinger system with competing potentials," *Journal* of Geometric Analysis, vol. 32, no. 4, p. 114, 2022.
- [23] X. M. He and W. M. Zou, "Existence and concentration behavior of positive solutions for a Kirchhoff equation in ℝ³," *Journal of Differential Equations*, vol. 252, no. 2, pp. 1813–1834, 2012.
- [24] J. Wang, L. X. Tian, J. X. Xu, and F. B. Zhang, "Multiplicity and concentration of positive solutions for a Kirchhoff type problem with critical growth," *Journal of Differential Equations*, vol. 253, no. 7, pp. 2314–2351, 2012.
- [25] G. M. Figueiredo, N. Ikoma, and J. R. Santos Júnior, "Existence and concentration result for the Kirchhoff type equations with general nonlinearities," *Archive for Rational Mechanics and Analysis*, vol. 213, no. 3, pp. 931–979, 2014.
- [26] D. D. Sun and Z. T. Zhang, "Uniqueness, existence and concentration of positive ground state solutions for Kirchhoff type problems in ℝ³," *Journal of Mathematical Analysis and Applications*, vol. 461, no. 1, pp. 128–149, 2018.
- [27] D. D. Sun and Z. T. Zhang, "Existence and asymptotic behaviour of ground state solutions for Kirchhoff-type equations with vanishing potentials," *Zeitschrift für Angewandte Mathematik und Physik*, vol. 70, no. 1, 2019.
- [28] H. N. Fan and X. C. Liu, "On the multiplicity and concentration of positive solutions to a Kirchhoff-type problem with competing potentials," *Journal of Mathematical Physics*, vol. 63, no. 1, article 011512, 2022.

- [29] Y. He, G. B. Li, and S. J. Peng, "Concentrating bound states for Kirchhoff type problems in R³ involving critical Sobolev exponents," *Advanced Nonlinear Studies*, vol. 14, no. 2, pp. 483– 510, 2014.
- [30] T. Hu and W. Shuai, "Multi-peak solutions to Kirchhoff equations in ℝ³ with general nonlinearity," *Journal of Differential Equations*, vol. 265, pp. 3587–3617, 2018.
- [31] G. B. Li, P. Luo, S. J. Peng, C. H. Chun, and C. L. Xiang, "A singularly perturbed Kirchhoff problem revisited," *Journal of Differential Equations*, vol. 268, pp. 541–589, 2020.
- [32] D. D. Sun, "Multiple positive solutions to Kirchhoff equations with competing potential functions in R³," *Boundary Value Problems*, vol. 85, 2019.
- [33] D. D. Sun and Z. T. Zhang, "Existence of positive solutions to Kirchhoff equations with vanishing potentials and general nonlinearity," SN Partial Differential Equations and Applications, vol. 1, no. 2, p. 8, 2020.
- [34] Q. Xie and X. Zhang, "Semi-classical solutions for Kirchhoff type problem with a critical frequency," *Proceedings of the Royal Society of Edinburgh Section A*, vol. 151, no. 2, pp. 761–798, 2021.
- [35] W. Zhang, J. Zhang, and V. D. Rădulescu, "Concentrating solutions for singularly perturbed double phase problems with nonlocal reaction," *Journal of Differential Equations*, vol. 347, pp. 56–103, 2023.
- [36] H. Zhang, J. X. Xu, and F. B. Zhang, "Semiclassical ground states for a class of nonlinear Kirchhoff-type problems," *Applicable Analysis*, vol. 96, no. 13, pp. 2267–2284, 2017.
- [37] A. Ambrosetti, V. Felli, and A. Malchiodi, "Ground states of nonlinear Schrödinger equations with potentials vanishing at infinity," *Journal of the European Mathematical Society*, vol. 7, pp. 117–144, 2005.
- [38] B. Opic and A. Kufner, "Hardy-type inequalities," *Pitman Research Notes in Mathematics Series*, vol. 219, 1990Zbl 0698.26007 MR92b: 26028.
- [39] Q. Han and F. Lin, *Elliptic Partial Differential Equations, Lecture Notes*, American Mathematical Society, Providence, 2000.