

Research Article

Impact of Multiplicative Noise on the Exact Solutions of the Fractional-Stochastic Boussinesq-Burger System

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In this paper, we consider the fractional-stochastic Boussinesq-Burger system (FSBBS) generated by the multiplicative Brownian motion. The Jacobi elliptic function techniques are used to create creative elliptic, hyperbolic, and rational fractional-stochastic solutions for FSBBS. Furthermore, we draw 2D and 3D graphs by using the MATLAB Package for some obtained solutions of FSBBS to discuss the influence of the Brownian motion on these solutions. Finally, we indicate that the Brownian motion stabilizes the solutions of FSBBS around zero.

1. Introduction

Nonlinear partial differential equations (NLPDEs) have grown in popularity in the area of nonlinear science, owing to their large variety of uses in economics [1], engineering [2], civil engineering [3], soil mechanics [4], physics [5], quantum mechanics [6], statistical mechanics [7], solid-state physics [8], population ecology [9], etc. Solitons are among the most common in the setting of NLPDE solutions, and they are essential for understanding nonlinear physical phenomena. Solitons are utilized to understand the properties of nonlinear media in various areas including quantum electronics, plasma physics, nonlinear optics, and fluid dynamics [10–13]. Recently, the searching of exact soliton solutions to NLPDEs has become an enthralling research topic in engineering and applied sciences. Many techniques have been used to determine exact solutions for NLPDE including tanh-sech [14, 15], Darboux transformation [16], sine-cosine [17, 18], $\exp(-\phi(\zeta))$ -expansion [19], (G'/G) -expansion

[20–22], Lie symmetry analysis method [23], improved F-expansion method [24, 25], Hirota's function [26], the Jacobi elliptic function [27, 28], and perturbation [29, 30].

The fractional differential equation is extensively used in fluid mechanics, solid state physics, optical fibers, neural physics, quantum field theory, mathematical biology, plasma physics, and other areas [31–34]. Researchers recommend fractional-order derivative over ordinary order derivative because integer-order derivative is essentially a local operator, but fractional-order derivative is so much more. Also, they explain physical phenomena such as quantum mechanics, diffusion, gravity, heat, elasticity, fluid dynamics, electrodynamics, electrostatics, and sound. Recently, the exact solutions with conformable derivative have been obtained in many papers for instance [35–40].

On the other hand, a wide variety of complex nonlinear physical phenomena can be represented using stochastic partial differential equations (SPDEs). These kind of equations can be found in many fields, such as physics and finance.

On the other side, stochastic partial differential equations (SPDEs) can be used to represent a wide range of complicated nonlinear physical processes. These kind of equations appear in a variety of areas including engineering, geophysical, biology, climate dynamics, finance, and physics [41–43].

To realize a higher level of qualitative accord, we take the following fractional-stochastic Boussinesq-Burger system (FSBBS) perturbed in the itô sense by multiplicative noise:

$$d\Phi + \left[2\Phi \mathbb{D}_x^\alpha \Phi - \frac{1}{2} \mathbb{D}_x^{3\alpha} \Phi \right] dt = \sigma \Phi d\mathbb{B}, \quad (1)$$

$$d\Psi + \left[2\mathbb{D}_x^\alpha (\Phi \Psi) - \frac{1}{2} \mathbb{D}_x^{3\alpha} \Phi \right] dt = \sigma \Psi d\mathbb{B}, \quad (2)$$

where $\Phi(x, t)$ denotes the horizontal velocity field. \mathbb{D}^α is the conformable derivative (CD) [44]. $\Psi(x, t)$ is the height of the water surface above the bottom horizontal level. $\mathbb{B}(t)$ is a Brownian motion (BM) and σ is the noise strength.

The Boussinesq-Burgers system (BBS), with $\alpha = 1$ and $\sigma = 0$, appears in fluid flow research and explains how shallow water waves spread. Due to the importance of BBS, many researchers have created its exact solutions by using various methods such as Hirota method [45], Lie symmetry method [46], sine-Gordon expansion method [47], Jacobi elliptic function method [48], extended homogeneous balance [49], Darboux transformation [50], The modified $\exp(-\phi(\zeta))$ -expansion function method [51], and Exp-function method [52]. On the other side, many techniques have been documented for fractional BBS, including a domain decomposition method [53] and generalized Kudryashov method [54]. The exact solutions of the FSBBS (1-2) have not yet been studied.

Our novelty of this paper is to find the exact fractional stochastic solutions of FSBBS (1-2). In the presence of a stochastic term and the fractional space, this study is the first to obtain analytical solutions to the FSBBS (1-2). Numerous solutions, including those involving elliptic, trigonometric, rational, and hyperbolic functions, can be obtained using the Jacobi elliptic function technique. Moreover, we utilize MATLAB to build 2D and 3D figures for some of the obtained solutions in this study to display how the BM influences on the solutions of FSBBS (1-2).

The layout of the document is as follows: in Sec. 2, we define and give some properties of the CD and BM. In Sec. 3, we use an effective wave transformation to establish the FSBBS (1-2) wave equation. In Sec.4, we use the Jacobi elliptic function method to generate the analytic of FSBBS (1-2). While, in Sec.5, the effect of the BM on the solutions obtained is studied. In Sec. 6, the document's conclusion is shown.

2. Preliminaries

In this section, we define and clarify some characteristics of the BM and CD. In the following, we define BM $\mathbb{B}(t)$ as:

Definition 1 (see [55]). Stochastic process $\{\mathbb{B}(t)\}_{t \geq 0}$ is said a BM if the following conditions satisfy: $B(t)$ is continuous function of $t \geq 0$; $B(0) = 0$; for $\tau_1 < \tau_2$, $B(\tau_2) - B(\tau_1)$ is independent; and $B(\tau_2) - B(\tau_1)$ has a Gaussian distribution $\mathcal{N}(0, \tau_2 - \tau_1)$.

Lemma 2 (see [55]). $\mathbb{E}(e^{\sigma \mathbb{B}(t)}) = e^{((1/2)\sigma^2 t)}$ for $\sigma \geq 0$.

Definition 3 (see [44]). Let $\phi : (0, \infty) \rightarrow \mathbb{R}$, then the CD of ϕ of order $\alpha \in (0, 1]$ is defined as

$$\mathbb{D}_x^\alpha \phi(x) = \lim_{\kappa \rightarrow 0} \frac{\phi(x + \kappa x^{1-\alpha}) - \phi(x)}{\kappa}. \quad (3)$$

Let us go through some of the CD's features. If a, b are constant, then

- (1) $\mathbb{D}_x^\alpha [a] = 0$,
- (2) $\mathbb{D}_x^\alpha [x^b] = b x^{b-\alpha}$,
- (3) $\mathbb{D}_x^\alpha [a\Theta_1(x) + b\Theta_2(x)] = a\mathbb{D}_x^\alpha \Theta_1(x) + b\mathbb{D}_x^\alpha \Theta_2(x)$,
- (4) $\mathbb{D}_x^\alpha \Theta(x) = x^{1-\alpha} (d\Theta/dx)$,
- (5) $\mathbb{D}_x^\alpha (\Theta_1 \circ \Theta_2)(x) = x^{1-\alpha} \Theta_2'(x) \Theta_1'(\Theta_2(x))$.

3. Wave Equation of FSBBS

The next wave transformation is used

$$\begin{aligned} \Phi(x, t) &= \varphi(\xi) e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \Psi(x, t) \\ &= \psi(\xi) e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \xi = \frac{1}{\alpha} x^\alpha + \omega t, \end{aligned} \quad (4)$$

in order to attain the wave equation of FSBBS (1-2). Where ω is a constant, φ and ψ are deterministic functions. Plugging Equation (4) into Equations (1) and (2) and utilizing

$$\begin{aligned} d\Phi &= \left[\omega \varphi' dt + \sigma \varphi d\mathbb{B} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \\ d\Psi &= \left[\omega \psi' dt + \sigma \psi d\mathbb{B} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \\ \mathbb{D}_x^\alpha \Phi &= \varphi' e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \mathbb{D}_x^\alpha \Psi = \psi' e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \\ \mathbb{D}_x^{3\alpha} \Phi &= \varphi''' e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \mathbb{D}_x^{3\alpha} \Psi = \psi''' e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \end{aligned} \quad (5)$$

we get

$$\omega \varphi' + 2\omega \varphi' e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))} - \frac{1}{2} \psi' = 0, \quad (6)$$

$$\omega \psi' + 2(\varphi \psi)' e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))} - \frac{1}{2} \varphi''' = 0. \quad (7)$$

Taking expectation $\mathbb{E}(\cdot)$ for Equations (6) and (7), we

attain

$$\omega\varphi' + 2\varphi\varphi' e^{-((1/2)\sigma^2 t)} \mathbb{E}\left(e^{\sigma\mathbb{B}(t)}\right) - \frac{1}{2}\psi' = 0, \quad (8)$$

$$\omega\psi' + 2(\varphi\psi)' e^{-((1/2)\sigma^2 t)} \mathbb{E}\left(e^{\sigma\mathbb{B}(t)}\right) - \frac{1}{2}\varphi''' = 0. \quad (9)$$

Since $\mathbb{B}(t)$ is a Gaussian distribution, then $\mathbb{E}(e^{\sigma\mathbb{B}(t)}) = e^{((\sigma^2/2)t)}$. Now Equations (8) and (9) become

$$\omega\varphi' + 2\varphi\varphi' - \frac{1}{2}\psi' = 0, \quad (10)$$

$$\omega\psi' + (\varphi\psi)' - \frac{1}{2}\varphi''' = 0. \quad (11)$$

Integrating Equations (10) and (11) and putting the constants of integration equal zero, we have

$$\psi = 2\omega\varphi + 2\varphi^2, \quad (12)$$

$$\omega\psi + 2(\varphi\psi) - \frac{1}{2}\varphi'' = 0. \quad (13)$$

Plugging Equation (12) into (13), we attain

$$\varphi'' - 8\varphi^3 - 12\omega\varphi^2 - 4\omega^2\varphi = 0. \quad (14)$$

4. Exact Solutions of FSBBS

We use the Jacobi elliptic functions approach described by Peng [56] to find the solutions of Equation (14). Consequently, we can therefore derive the exact solutions of FSBBS (1-2).

4.1. *Jacobi Elliptic Functions Method.* First, we suppose the solutions of Equation (14) are

$$\varphi(\xi) = \sum_{i=1}^N a_i \chi^i, \quad (15)$$

where χ is the solution of

$$\chi' = \sqrt{\frac{1}{2}p\chi^4 + q\chi^2 + r}, \quad (16)$$

where r, q and p are real parameters.

We note from the next Table 1 that Equation (16) has different types of solutions relying on r, q and p :

Where $dn(\xi, m) = dn(\xi, m)$, $cn(\xi) = cn(\xi, m)$, $sn(\xi) = sn(\xi, m)$, for $0 < m < 1$ are the Jacobi elliptic functions (JEFs). If $m \rightarrow 1$, then JEFs are converted into the hyperbolic functions as follows:

$$\begin{aligned} dn(\xi) &\longrightarrow \operatorname{sech}(\xi)sn(\xi) \longrightarrow \tanh(\xi), \quad cn(\xi) \longrightarrow \operatorname{sech}(\xi), \\ cs(\xi) &\longrightarrow \operatorname{csch}(\xi), \quad ds \longrightarrow \operatorname{csch}(\xi). \end{aligned} \quad (17)$$

TABLE 1: All possible solutions for Equation (16).

Case	p	q	r	$\chi(\xi)$
1	$2m^2$	$-(1+m^2)$	1	$sn(\xi)$
2	2	$2m^2-1$	$-m^2(1-m^2)$	$ds(\xi)$
3	2	$2-m^2$	$(1-m^2)$	$cs(\xi)$
4	$-2m^2$	$2m^2-1$	$(1-m^2)$	$cn(\xi)$
5	-2	$2-m^2$	(m^2-1)	$dn(\xi)$
6	$\frac{m^2}{2}$	$\frac{(m^2-2)}{2}$	$\frac{1}{4}$	$\frac{sn(\xi)}{1 \pm dn(\xi)}$
7	$\frac{m^2}{2}$	$\frac{(m^2-2)}{2}$	$\frac{m^2}{4}$	$\frac{sn(\xi)}{1 \pm dn(\xi)}$
8	$\frac{-1}{2}$	$\frac{(m^2+1)}{2}$	$\frac{-(1-m^2)^2}{4}$	$mcn(\xi) \pm dn(\xi)$
9	$\frac{m^2-1}{2}$	$\frac{(m^2+1)}{2}$	$\frac{(m^2-1)}{4}$	$\frac{dn(\xi)}{1 \pm sn(\xi)}$
10	$\frac{1-m^2}{2}$	$\frac{(1-m^2)}{2}$	$\frac{(1-m^2)}{4}$	$\frac{cn(\xi)}{1 \pm sn(\xi)}$
11	$\frac{(1-m^2)^2}{2}$	$\frac{(1-m^2)^2}{2}$	$\frac{1}{4}$	$\frac{sn(\xi)}{dn \pm cn(\xi)}$
12	2	0	0	$\frac{c}{\xi}$
13	0	1	0	ce^{ξ}

4.2. *Solutions of FSBBS.* By balancing φ'' with φ^3 in Equation (14), we can calculate the parameter N as

$$N + 2 = 3N \Rightarrow N = 1. \quad (18)$$

Hence, Equation (15) with $N = 1$ becomes

$$\varphi = a_0 + a_1\chi. \quad (19)$$

Differentiating Equation (19) twice, we have, by using (16),

$$\varphi'' = a_1q\chi + a_1p\chi^3. \quad (20)$$

Substituting Equation (19) and Equation (20) into Equation (14) we obtain

$$\begin{aligned} (a_1p - 8a_1^3)\chi^3 - (24a_0a_1^2 + 12\omega a_1^2)\chi^2 \\ + (a_1q - 24a_0^2a_1 - 24\omega a_0a_1 - 4\omega^2a_1)\chi \\ - (8a_0^3 + 12\omega a_0^2 + 4\omega^2a_0) = 0. \end{aligned} \quad (21)$$

Equating each coefficient of χ^k , for $k = 0, 1, 2, 3$, to zero,

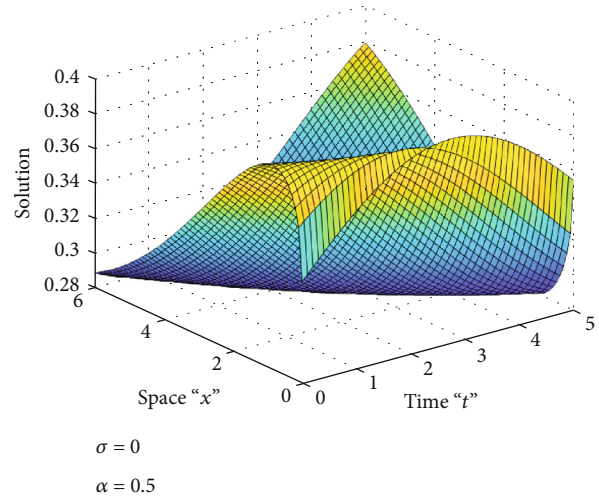
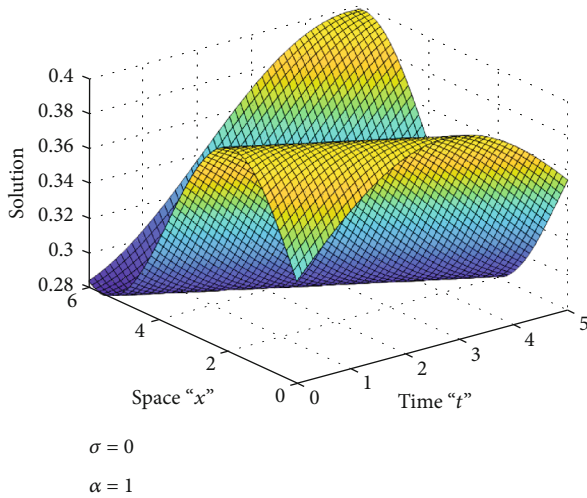


FIGURE 1: 3D-graph of Equation (33) with $\sigma = 0$ and various values of $\alpha = 1, 0.5$.

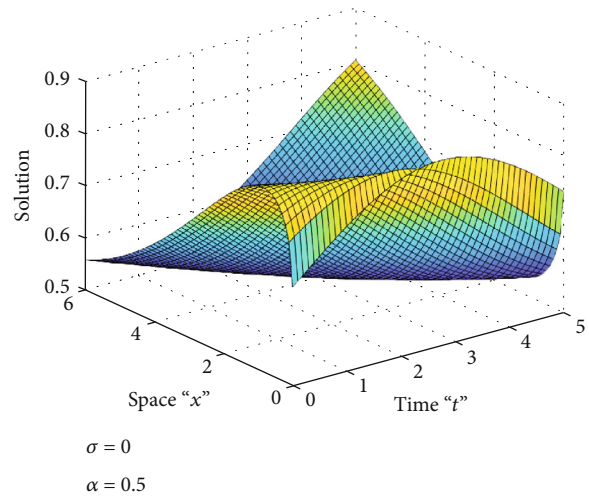
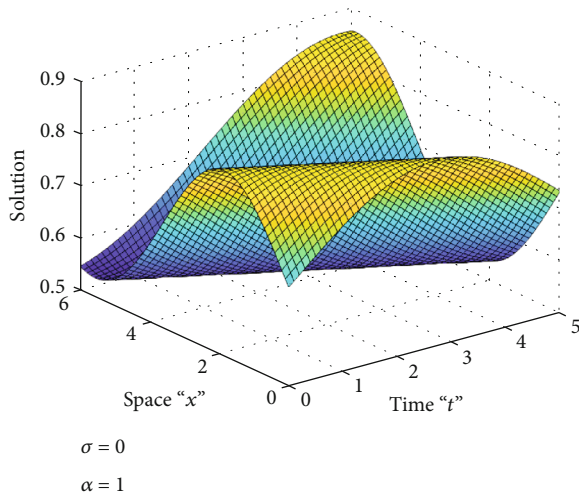


FIGURE 2: 3D-graph of Equation (34) with $\sigma = 0$ and various values of $\alpha = 1, 0.5$.

we have

$$\begin{aligned} a_1 p - 8a_1^3 &= 0, \\ 24a_0 a_1^2 + 12\omega a_1^2 &= 0, \\ a_1 q - 24a_0^2 a_1 - 24\omega a_0 a_1 - 4\omega^2 a_1 &= 0, \end{aligned} \tag{22}$$

and

$$8a_0^3 + 12\omega a_0^2 + 4\omega^2 a_0 = 0. \tag{23}$$

We get by solving these equations:

$$a_0 = \frac{1}{2} \sqrt{-\frac{1}{2}q}, a_1 = \pm \sqrt{\frac{p}{8}}, \omega = \pm \sqrt{-\frac{1}{2}q}, \tag{24}$$

for $p > 0$ and $q \leq 0$. Then, the Equation (14) has the

solutions:

$$\varphi(\xi) = \sqrt{-\frac{1}{8}q} \pm \sqrt{\frac{p}{8}\chi(\xi)}. \tag{25}$$

Therefore, by utilizing (4) and (12), the solution of the FSBBS (2-1) are

$$\begin{aligned} \Phi(x, t) &= \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}\chi(\xi)} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \\ \Psi(x, t) &= \left[-\frac{3q}{4} + \sqrt{-pq}\chi(\xi) + \frac{p}{4}\chi^2 \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}. \end{aligned} \tag{26}$$

There are numerous cases, by using the previous Table 1, for $q \leq 0, p > 0$ and r as follows:

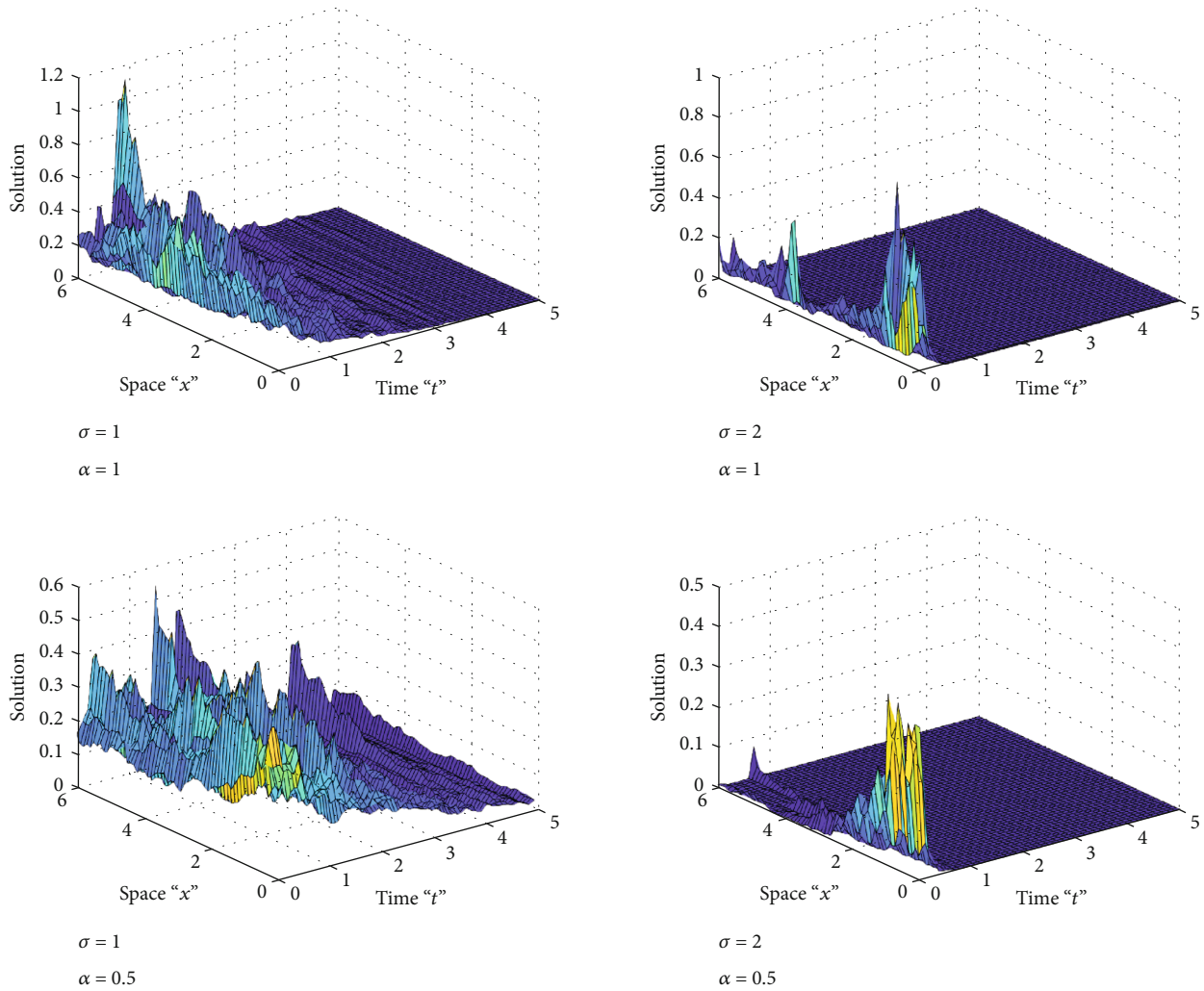


FIGURE 3: 3D-graph of Equation (33) with $\sigma = 1, 2$ and $\alpha = 1, 0.5$.

Case 1. If $q = -(1 + m^2), p = 2m^2$ and $r = 1$, then $\chi(\xi) = sn(\xi)$. Therefore, the FSBBS (1-2) has the solution

$$\Phi(x, t) = \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}} sn \left(\frac{1}{\alpha} x^\alpha + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}, \tag{27}$$

$$\Psi(x, t) = \left[-\frac{3q}{4} + \sqrt{-pq} sn \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) + \frac{p}{4} sn^2 \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}. \tag{28}$$

If $m \rightarrow 1$, then Equations (27) and (28) degenerates to

$$\begin{aligned} \Phi(x, t) &= \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}} \tanh \left(\frac{1}{\alpha} x^\alpha + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}, \\ \Psi(x, t) &= \left[-\frac{3q}{4} + \sqrt{-pq} \tanh \left(\frac{1}{\alpha} x^\alpha + \sqrt{\frac{-q}{2}} t \right) + \frac{p}{4} \tanh^2 \left(\frac{1}{\alpha} x^\alpha + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}. \end{aligned} \tag{29}$$

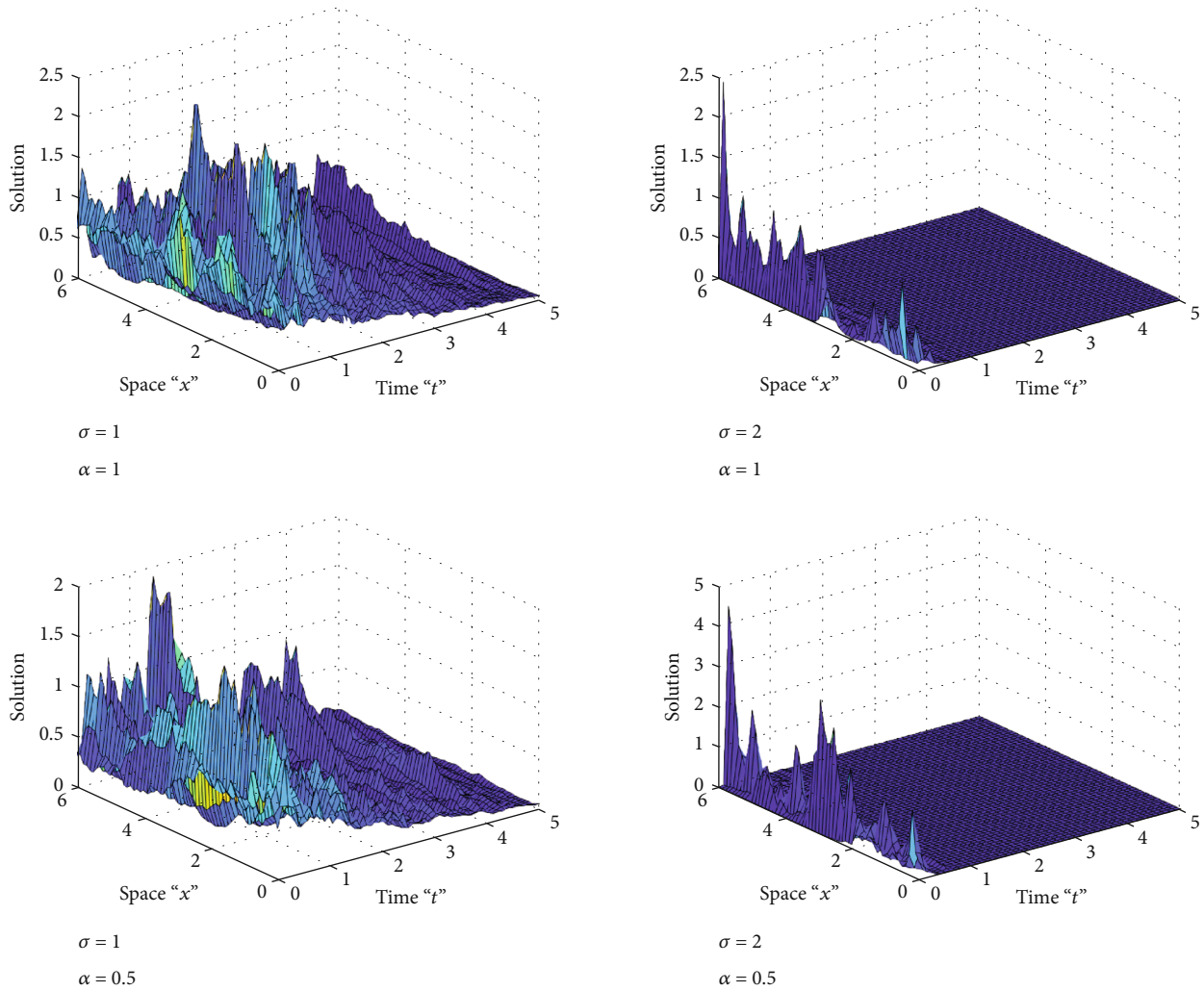


FIGURE 4: 3D-graph of Equation (34) with $\sigma = 1, 2$ and $\alpha = 1, 0.5$.

Case 2. If $q = 2m^2 - 1$ for $m \leq (1/\sqrt{2})$, $p = 2$ and $r = -m^2(1 - m^2)$, then $\chi(\xi) = ds(\xi)$. So, the FSBBS (1-2) has the solution:

$$\Phi(x, t) = \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}} ds \left(\frac{1}{\alpha} x^\alpha + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}, \tag{30}$$

$$\Psi(x, t) = \left[-\frac{3q}{4} + \sqrt{-pq} ds \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) + \frac{p}{4} ds^2 \cdot \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}. \tag{31}$$

If $m \rightarrow 1$, then Equations (30) and (31) degenerates to

$$\begin{aligned} \Phi(x, t) &= \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}} \operatorname{csch} \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}, \\ \Psi(x, t) &= \left[-\frac{3q}{4} + \sqrt{-pq} \operatorname{csch} \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) + \frac{p}{4} \operatorname{csch}^2 \left(\frac{x^\alpha}{\alpha} + \sqrt{\frac{-q}{2}} t \right) \right] e^{(\sigma B(t) - ((1/2)\sigma^2 t))}. \end{aligned} \tag{32}$$

Case 3. If $q = ((m^2 - 2)/2), p = (m^2/2)$ and $r = (1/4)$ (or $r = (m^2/4)$), then $\chi(\xi) = ((sn(\xi))/(1 \pm dn(\xi)))$. Therefore, the FSBBS (1-2) has the solution:

$$\Phi(x, t) = \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}} \frac{\left(\operatorname{sn} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)}{\left(1 \pm \operatorname{dn} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \tag{33}$$

$$\Psi(x, t) = \left[-\frac{3q}{4} + \sqrt{-pq} \frac{\left(\operatorname{sn} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)}{\left(1 \pm \operatorname{dn} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)} + \frac{p}{4} \frac{\left(\operatorname{sn}^2 \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)}{\left(1 \pm \operatorname{dn} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)^2} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}. \tag{34}$$

If $m \rightarrow 1$, then Equations (33) and (34) degenerates to

$$\Phi(x, t) = \left[\sqrt{-\frac{1}{8}q} + \sqrt{\frac{p}{8}} \frac{\left(\tanh \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)}{\left(1 \pm \operatorname{sech} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \tag{35}$$

$$\Psi(x, t) = \left[-\frac{3q}{4} + \sqrt{-pq} \frac{\left(\tanh \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)}{\left(1 \pm \operatorname{sech} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)} + \frac{p}{4} \frac{\left(\tanh^2 \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)}{\left(1 \pm \operatorname{sech} \left(\left(\frac{1}{\alpha} \right) x^\alpha + \sqrt{-q/2} t \right) \right)^2} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}. \tag{36}$$

Case 4. If $q = 0, p = 2$ and $r = 0$, then $\chi(\xi) = (C/\xi)$. Hence, the FSBBS (1-2) has the solution:

$$\begin{aligned} \Phi(x, t) &= \left[\frac{\alpha C}{2} x^{-\alpha} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}, \\ \Psi(x, t) &= \left[\frac{\alpha C}{2} x^{-2\alpha} \right] e^{(\sigma \mathbb{B}(t) - ((1/2)\sigma^2 t))}. \end{aligned} \tag{37}$$

Remark 4. If we set $\sigma = 0$ and $\alpha = 1$ in Equations (27) and (36), then we get the same results as reported in [48].

5. The Influence of Fractional Derivative and Noise

Here, the influence of noise on the achieved solutions of FSBBS (1-2) is explained. For various values of α (the fractional derivative order) and σ (noise strength), some graphs are provided using the MATLAB tools.

Firstly the Fractional Derivative Influence. In Figures 1 and 2, if $\sigma = 0$ and $m = 0.4$, we can observe that the surface shrinks when α is decreasing:

Secondly the Noise Influence. In Figures 3 and 4, when noise is introduced, the surface flattens significantly if its strength is increased $\sigma = 1, 2$

In Figure 5, we introduce 2D-graph of the $\Phi(x, t)$ in (33) with $\alpha = 1$ and with $\sigma = 0, 0.5, 1, 2$, which highlights the previous outcomes:

We may deduct from Figures 1–5 that:

- (1) When the fractional-order α increases, the surface expands,
- (2) The multiplicative noise stabilizes the solutions of FSBBS at zero.

This results show that it is important to add the stochastic term into the Boussinesq-Burger equation in order to obtain accurate solutions.

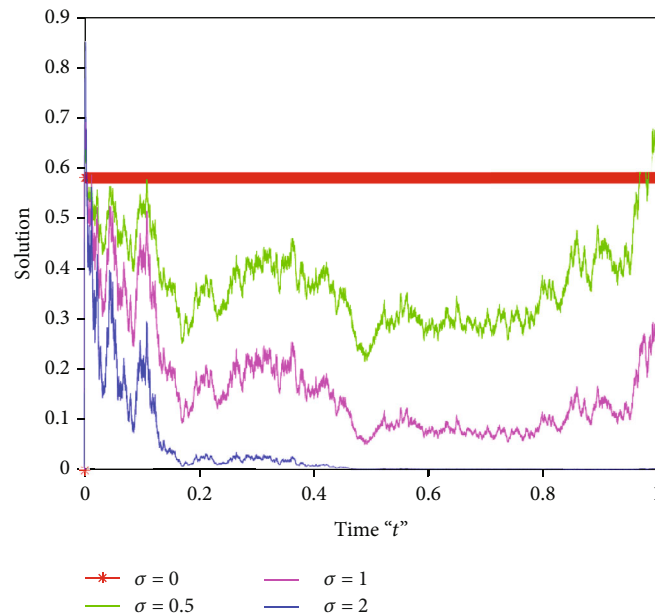


FIGURE 5: 2D graph of Equation (33)

6. Conclusions

In this article, the exact fractional-stochastic solutions of the fractional-stochastic Boussinesq-Burger system (1-2) driven by multiplicative noise were successfully obtained by using the Jacobi elliptic function method. Numerous analytical solutions for FSBBS (1-2) including elliptic, trigonometric, rational, and hyperbolic functions can be determined using the Jacobi elliptic function method. Because of the importance of FSBBS in fluid flow research and in explaining the propagation of shallow water waves, the acquired solutions are far more beneficial and efficient in understanding several critical complicated physical phenomena. In addition, we utilized the MATLAB package to demonstrate how multiplicative noise and fractional derivative influenced the solutions of FSBBS. As a result, we concluded that the stabilization of the solutions of the FSBBS (1-2) is affected by the multiplicative noise.

Data Availability

All data are available in this paper

Conflicts of Interest

The authors declare that they have no competing interests.

Authors' Contributions

All authors contributed equally to the writing of this paper. All authors read and approved the final manuscript.

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