# On the Existence and Stability of Positive Solutions of Eigenvalue Problems for a Class of P-Laplacian $\psi$-Caputo Fractional Integro-Differential Equations 

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#### Abstract

This article introduces a generalized approach for analyzing stability and establishing the existence of positive solutions in a specific type of differential equations known as p-Laplacian $\psi$-Caputo fractional differential equations with fractional integral boundary conditions. The study utilizes various techniques, including the analysis of Green's function properties and the application of Guo-Krasnovelsky's fixed point theorem on cones. By employing these methods, the research establishes novel findings concerning the existence and nonexistence of positive solutions. The investigation relies on fractional integrals, differential operators, and fundamental lemmas as fundamental tools. To assess solution stability, the Hyers-Ulam concept is employed, which extends prior research and introduces a specific definition. The article also provides numerical examples that support the obtained results, thereby demonstrating the practical applicability and accuracy of the proposed methods. Moreover, the study contributes to a deeper understanding of this subject matter and highlights real-life applications for these types of problems. Overall, this study offers a comprehensive analysis of stability and solution existence in a specific class of differential equations, with implications that extend to real-world scenarios such as engineering systems, financial modeling, population dynamics, epidemiology, and ecological studies. These types of problems arise in various fields where modeling and analyzing complex phenomena are necessary.


## 1. Introduction

Fractional calculus, a field of mathematical analysis, expands upon the principles of differentiation and integration to encompass noninteger orders. The roots of this branch can be traced back to the 18th century, where notable mathematicians such as Leibniz and Euler made early contributions. However, it was not until the mid-19th century that fractional calculus gained increased attention from the scientific and mathematical communities. In recent years, fractional calculus has gained increasing attention and has become an active area of research. Many mathematical techniques and numerical methods have been developed to solve fractional differential equations and perform fractional calculus operations. This field has diverse applications in physics, engineering, signal
processing, finance, and control theory [1-4]. By providing an extension to conventional calculus, fractional calculus enables a more comprehensive comprehension of intricate systems and phenomena, and this is due to its remarkable properties and wide-ranging applications of this subject that render it a captivating area of exploration for mathematicians, scientists, and engineers alike. Several studies have presented proof that fractional models provide a more accurate and organized depiction of natural phenomena when compared to conventional models based on integerorder and ordinary time-derivatives. This has been extensively discussed in scholarly works such as [5-7], and the relevant literature cited therein.

Here, we present various examples showcasing the applications of fractional calculus. In a study by Mohammadi et al. [7], a novel fractional system was utilized for modeling
hearing loss associated with the mumps virus. This system employed a nonsingular kernel fractional operator called the Caputo-Fabrizio derivative. The fractional-order derivative, unlike its integer-order counterpart, effectively accounted for memory effects and nonlocal behavior. This research marked the first instance of employing the Caputo-Fabrizio derivative for modeling hearing loss with the mumps virus. Another significant contribution was made by Rezapour et al. [8], where the authors introduced a generalized fractional system for analyzing the dynamics of anthrax disease transmission. They employed the Caputo-Fabrizio derivative, a fractional operator without a singular kernel, to develop this novel approach. The authors observed that the approximate solutions obtained from the fractional Capu-to-Fabrizio model gradually approached those derived from the classical integer-order system as time progressed. Moreover, several researchers have successfully implemented fractional extensions of mathematical models originally formulated with integer orders, demonstrating their natural representation of real-world phenomena. Notable examples include the works of Aydogan et al. [9], Baleanu et al. [5, 10-12], Mohammadi et al. [13], and George et al. [14], among others. These studies systematically incorporated fractional calculus into existing models, further enriching their accuracy and applicability.

Fractional boundary value problems (FBVPs) have gained considerable attention due to their broad applications in various fields. FBVPs offer a robust framework for modeling complex phenomena and provide valuable insights into the behavior of fractional systems. These problems have found applications in diverse areas such as heat conduction, control systems, and population dynamics. Numerous investigations have been devoted to studying the existence of positive solutions for fractional differential equations with integral boundary conditions. These investigations employ techniques such as the fixed point theory and upper and lower solution methods. Relevant references for further information include ( $[15-31]$ ) and references therein. Furthermore, significant progress has been made recently in the exploration of p -Laplacian operators and the treatment of eigenvalue problems associated with fractional differential equations. Notable contributions in this field have been made by researchers such as Bai et al. Noteworthy references for this topic include ( $[17,29,32-40]$ ) and references therein.

Researchers have also made significant advancements in understanding the spectral properties and solution behavior of eigenvalue problems in fractional differential equations. Several studies have investigated various aspects, including the existence, uniqueness, and stability of eigenvalues and eigenfunctions. Contributions by researchers such as Bai et al. have greatly enhanced our understanding of eigenvalue problems in the context of fractional differential equations. Key references for further exploration include ( $[10,17,20,21,33,38,42-44]$ ) among others. These works have greatly enriched our understanding of eigenvalue problems in the context of fractional differential equations.

Regarding the function $\psi$, fractional derivatives serve as generalizations of the Riemann-Liouville derivatives. The $\psi$-Caputo fractional derivative differs from the classical derivative due to the presence of kernel terms. Recently, Almeida re-evaluated this derivative and provided a Caputo-type regularization of the existing definition with intriguing properties. Additional properties and applications of the $\psi$-Caputo fractional derivatives can be found in references such as ([45-51]), and references therein.

In [40], the authors investigated the existence of positive solutions for an eigenvalue problem utilizing the method of upper and lower solutions and the Schauder fixed point theorem. The problem can be formulated as follows:

$$
\left\{\begin{array}{l}
-D_{t}^{\beta}\left(\phi_{p}\left(D_{t}^{\alpha} u\right)\right)(t)=\lambda f(t, u(t)), \quad t \in(0,1)  \tag{1}\\
u(0)=D_{t}^{\alpha} u(t)=0, u(1)=\int_{0}^{1} u(s) d A(s)
\end{array}\right.
$$

In this equation, $D_{t}^{\beta}$ and $D_{t}^{\alpha}$ represent the standard Riemann-Liouville derivatives, with $1<\alpha \leq 2$ and $0<\beta \leq 1$. The function $A$ is of bounded variation, and $\int_{0}^{1} u(s) d A(s)$ denotes the Riemann-Stieltjes integral of $u$ with respect to $A$. The $p$-Laplacian operator $\phi_{p}$ is defined as $\phi_{p}=|s|^{p-2} s$, where $p>1$. The function $f(t, x):(0,1) \times(0,+\infty) \longrightarrow[0$, $+\infty)$ is continuous and may exhibit singularity at $t=0$, $t=1$, and $x=0$.

The authors in [21] conducted a study on the presence of positive solutions for an eigenvalue problem related to a nonlinear fractional differential equation. The equation involves a generalized p-Laplacian operator and is described by the following system:

$$
\left\{\begin{array}{l}
D_{0^{+}}^{\beta}\left(\phi_{p}\left(D_{0^{+}}^{\alpha} u(t)\right)\right)=\lambda f(u(t)), \quad t \in(0,1)  \tag{2}\\
u(0)=u^{\prime}(0)=u^{\prime \prime}(0)=0, \phi_{p}\left(D_{0^{+}}^{\alpha} u(0)\right)=\left(\phi_{p}\left(D_{0^{+}}^{\alpha} u(1)\right)\right)^{\prime}=0
\end{array}\right.
$$

In this system, the values of $2<\alpha \leq 3$ and $1<\beta \leq 2$ are real numbers. The operator $\phi_{p}$ represents a generalized $p$ Laplacian operator, $\lambda>0$ is a parameter, and $f(t):(0,+\infty) \longrightarrow(0,+\infty)$ is a continuous function. The study utilized the properties of Green functions and the

Guo-Krasnoselskii's fixed-point theorem on cones to establish several results regarding the existence of at least one or two positive solutions within different eigenvalue intervals. In addition, the nonexistence of positive solutions was also examined in relation to the parameter $\lambda$.

In [29], the authors investigated the eigenvalue problem concerning a boundary fractional differential equation
involving the p-Laplacian operator. The equation is given as follows:

$$
\left\{\begin{array}{l}
-{ }^{c} D_{0^{+}}^{\beta}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{\alpha} u\right)\right)(t)=\lambda f(t, u(t)), \quad t \in(0,1)  \tag{3}\\
u(0)=\int_{0}^{1} g_{1}(s) u(s) \mathrm{d} s, u(1)=\int_{0}^{1} g_{2}(s) u(s) \mathrm{d} s, u^{\prime \prime}(0)=\int_{0}^{1} g_{3}(s) u(s) \mathrm{d} s,\left.{ }^{c} D_{0^{+}}^{\alpha} u(t)\right|_{t=0}=0
\end{array}\right.
$$

where ${ }^{c} D_{0^{+}}^{\beta}$ and ${ }^{c} D_{0^{+}}^{\alpha}$ are the standard Caputo derivatives with $2<\alpha \leq 3$, and $0<\beta \leq 1$. The p -Laplacian operator $\phi_{p}$ is defined as $\phi_{p}=|s|^{p-2} s$ with $p>1$. The functions $f(t, x):[0,1] \times[0,+\infty) \longrightarrow[0,+\infty)$ and $g_{i}(s) \in C[0,1]$ (where $i=1,2,3$ ) are continuous. The interval $J$ is defined as $J=[0,1]$, and $\lambda$ is a positive parameter.

In the work by Matar et al. [52], a newly proposed pLaplacian nonperiodic boundary value problem is examined, which involves generalized Caputo fractional derivatives. The authors extensively investigate the existence and uniqueness of solutions for this problem:

$$
\left\{\begin{array}{l}
\frac{d}{d t}\left(\phi_{p}\left({ }^{c} D^{\alpha, \rho} u(t)\right)\right)=q\left(t, u(t),{ }^{c} D^{\gamma, \rho} u(t)\right), \quad t \in[0,1]  \tag{4}\\
u(0)+\mu u(1)=\theta(u(0), u(1)), u^{\prime}(1)=v(u(0), u(1))
\end{array}\right.
$$

where ${ }^{c} D^{\alpha, \rho}$ and ${ }^{c} D^{\gamma, \rho}$ are the $G C_{p} F r$ derivatives with $\alpha \in(1,2), \gamma \in(0,1), \phi_{p}$ is the p -Laplacian operator with $p>1, \mu \neq-1$, and the nonlinear nonlinear functions $q:[0,1] \times R \times R \longrightarrow R$ and $\theta: R \times R \longrightarrow R$ are given continuous functions.

However, as far as we know, there are few papers studying the eigenvalue problem for the p-Laplacian fractional differential equations involving the integral boundary condition. Inspired by the abovementioned works, in this paper, we will explore the following p-Laplacian $\psi$-Caputo fractional integro-differential equation:

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\beta, \psi}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{\alpha, \psi} u\right)\right)(t)+\lambda f(t, u(t))=0, \quad t \in[0, T]  \tag{5}\\
u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
u(T)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{2}(s) u(s) d s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
\left.{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)\right|_{t=0}=0
\end{array}\right.
$$

where ${ }^{c} D_{0^{+}}^{\alpha, \psi}$ and ${ }^{c} D_{0^{+}}^{\beta, \psi}$ are the $\psi$-Caputo fractional derivatives of respective fractional orders $\alpha \in(2,3]$ and
$\beta \in(0,1]$. The IBVP 1 also involves the $p$-Laplacian operator $\phi_{p}$, which is a nonlinear operator defined as $\phi_{p}(s)=|s|^{p-2} s$, where $p>1$. The operator $\phi_{p}$ is used to model nonlinear phenomena such as turbulence and phase transitions. The boundary conditions of the IBVP involve integrals of the form $\int_{0}^{T} \psi^{\prime} \gamma-1 g_{i}(s) u(s) d s$, where $\gamma$ is a parameter between 0 and 1 , and $g_{i}$ are continuous functions on $[0, T]$ for $i=1,2,3$. These boundary conditions involve memory effects and nonlocality, which are typical of fractional order problems. The functions $f(t, u): I \times R \longrightarrow R$, and $g_{i}: I \longrightarrow R ; \quad(i=1,2,3)$, and $u: I \longrightarrow R$ are continuous functions on $I=[0, T]$.

It is clear that the IBVP (5) is an equation that combines fractional derivatives and integrals. It comprises a fractional differential equation governed by the p-Laplacian operator and a nonlinearity function $f(t, u)$, along with three boundary conditions. These boundary conditions involve fractional integrals that incorporate the function $u$ and three distinct functions $g_{1}(t), g_{2}(t)$, and $g_{3}(t)$. The differential equation and boundary conditions employ $\psi$-Caputo derivatives, which are an extension of the classical Caputo derivative. The $\psi$-Caputo derivatives are defined using the Riemann-Liouville fractional derivative and a scaling function $\psi(t)$ that satisfies specific conditions. The pLaplacian operator, a nonlinear differential operator, is utilized to model various physical phenomena such as fluid flow, diffusion, and image processing. Within this problem, there exists a positive constant parameter $\lambda$ referred to as the eigenvalue, which characterizes the system. The objective of the research is to determine the values of $\lambda$ that allow for nontrivial solutions, known as eigenfunctions, to exist for the problem. By analyzing the eigenvalues and eigenfunctions, it is possible to study the stability and dynamics of the system represented by the problem.

The motivation behind studying such type of IBVPs is to understand the behavior of the solutions of this problem, particularly in the presence of nonlinearity and nonlocality. This is an important research area with applications in various fields, including physics, engineering, and biology. In particular, the Caputo-type fractional IBVP in equation (5) has numerous applications in modeling systems with memory effects, nonlocality, and long-term dependencies. Its specific applications depend on the choice of the functions $f(t, u)$ and $g_{i}(t)$ used in each application depend on the properties of the system being modeled, indicating the flexibility and adaptability of the IBVP to different contexts. Thus, the Caputo-type fractional IBVP (5) is a challenging and interesting problem with significant potential for a range
of real-world applications in various fields of science and engineering.

The utilization of eigenvalue problems for a specific class of $p$-Laplacian $\psi$-Caputo fractional integro-differential equations has significant implications for various fields. For instance, in the context of mumps virus transmission, these eigenvalue problems offer valuable insights into the system's stability, critical thresholds, and spatial patterns, facilitating the development of effective strategies for controlling the spread of the virus (see [7] and related references). Moreover, eigenvalues play a crucial role in other domains as well. In control theory, they determine the stability and performance of quadcopter systems, allowing for stable flight and improved control [53, 54]. In signal processing, eigenvalues are employed for signal classification, compression, and denoising, with applications in the analysis of electroencephalogram signals related to Alzheimer's disease [10, 55, 56]. Fractional IBVPs using these eigenvalue problems are also used to model non-Newtonian fluid dynamics in porous media and heat conduction in biological tissues, offering insights into memory effects and thermal responses [22, 57-60]. These diverse applications underscore the broad scope and significance of employing eigenvalue problems for understanding complex phenomena and informing practical solutions.

According to our knowledge, numerous research studies have been conducted on the eigenvalue problems related to fractional differential equations that incorporate the $p$-Laplacian and integral boundary conditions. But the initial boundary value problem (IBVP) (5), presented in our work, serves as a basis for obtaining various investigations outlined in the monograph by utilizing necessary and appropriate parameters as follows:
(i) If $\psi(t)=t, \beta=0,1<\alpha \leq 2, p=2, \lambda=1, T=1$, and $g_{1}=g_{2}=g_{3}=0$, then we obtain the following
boundary value problem of nonlinear fractional differential equations:

$$
\left\{\begin{array}{l}
D_{0^{+}}^{\alpha} u(t)+f(t, u(t))=0, \quad t \in(0,1)  \tag{6}\\
u(0)=u(1)=0
\end{array}\right.
$$

which is similar to those outcomes obtained in [17].
(ii) If $\psi(t)=t, \beta=0,2<\alpha \leq 3, p=2, \lambda=1, T=1$, and $g_{1}=g_{2}=g_{3}=0$, then we obtain the following boundary value problem of nonlinear fractional differential equations:

$$
\left\{\begin{array}{l}
D_{0^{+}}^{\alpha} u(t)+f(t, u(t))=0, \quad t \in(0,1)  \tag{7}\\
u(0)=u(1)=u^{\prime}(0)=0
\end{array}\right.
$$

which is similar to those outcomes obtained by Zhao et al. [31].
(iii) If $\psi(t)=t, \quad 1<\alpha \leq 2, \quad 0<\beta \leq 1, \quad \lambda=1, \quad T=1$, $g_{1}=g_{2}=g_{3}=0$, then we obtain the following boundary value problem of nonlinear fractional differential equations:

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\beta}\left(\phi_{p}^{c} D_{0^{+}}^{\alpha} u\right)(t)+f(t, u(t))=0, \quad t \in(0,1),  \tag{8}\\
u(0)=u(1)=0, \text { and }^{c} D_{0^{+}}^{\alpha} u(0)=0,
\end{array}\right.
$$

which is similar to the results obtained by Chai [18] for $\sigma=0$.
(iv) If $\psi(t)=t, T=1,0<\beta \leq 1,2<\alpha \leq 3, \gamma=1, g_{2}=1$, and $g_{1}=g_{3}=0$, then obtained outcomes in the current paper incorporate the investigation of Zhang et al. [40] for $A(s)=s$

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\beta}\left(\phi_{p}^{c} D_{0^{+}}^{\alpha} u\right)(t)+\lambda f(t, u(t))=0, \quad t \in(0,1)  \tag{9}\\
u(0)=0, u(1)=0, u(1)=\int_{0}^{1} u(s) d A(s), \text { and }\left.{ }^{c} D_{0^{+}}^{\alpha} u(t)\right|_{t=0}=0
\end{array}\right.
$$

(v) If $\psi(t)=t, p=2, \beta=0,2<\alpha \leq 3, \gamma=1, \lambda=q(t)$, $T=1$, and $g_{1}=g_{3}=0$, then we obtain the following fractional differential equation with the following integral boundary conditions:

$$
\left\{\begin{array}{l}
D_{0^{+}}^{\alpha} u(t)+q(t) f(t, u(t))=0, \quad t \in(0,1)  \tag{10}\\
u(0)=u^{\prime}(0)=0, \text { and } u(1)=\int_{0}^{1} g(s) u(s) d s
\end{array}\right.
$$

which is similar to that studied by Sun and Zhao [29].
(vi) Also, if $\psi(t)=t, T=1$, and $\gamma=1$ then we have the $p$-Laplacian fractional differential equations involving the integral boundary condition.
$\left\{\begin{array}{l}{ }^{c} D_{0^{+}}^{\beta}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{\alpha} u\right)\right)(t)+\lambda f(t, u(t))=0, \quad t \in[0,1], \\ u(0)=\int_{0}^{1} g_{1}(s) u(s) d s, u(1)=\int_{0}^{1} g_{2}(s) u(s) d s, \\ u^{\prime \prime}(0)=\int_{0}^{1} g_{3}(s) u(s) d s, \text { and }\left.^{c} D_{0^{+}}^{\alpha} u(t)\right|_{t=0}=0,\end{array}\right.$
which is the same result obtained in Su et al. [36].

## 2. Main Results

2.1. Mathematical Background. In this section, we introduce some notations, definitions, lemmas, and theorems that are considered prerequisites for our work, such as fractional order integrals and derivatives, Green's functions, completely continuous operators, and others.

Definition 1 (see [47]). For any real number $\alpha>0$, the leftsided $\psi$-Riemann-Liouville fractional integral of order $\alpha$ for an integrable function $u: I \longrightarrow R$ with respect to another function $\psi: I \longrightarrow R$, which is an increasing differentiable function such that $\psi^{\prime}(t) \neq 0$ for all $t \in I$ is defined by the following equation:

$$
\begin{equation*}
I^{\alpha, \psi} u(t)=\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \psi^{\prime}(s)(\psi(t)-\psi(s))^{\alpha-1} u(s) \mathrm{d} s \tag{12}
\end{equation*}
$$

where $\Gamma$ is the classical Euler Gamma function.

Definition 2 (see [47]). If $n \in N$ and $\psi, u \in C^{\mathrm{n}}(I, R)$ are two functions such that $\psi$ is increasing and $\psi^{\prime}(t) \neq 0$ for all $t \in I$, then the left-sided $\psi$-Caputo fractional derivative of a function u of order $\alpha$ is defined by the following equation:

$$
\begin{align*}
{ }^{c} D^{\alpha, \psi} u(t) & =I^{n-\alpha, \psi}\left(\frac{1}{\psi^{\prime}(t)} \frac{d}{d t}\right)^{n} u(t) \\
& =\frac{1}{\Gamma(n-\alpha)} \int_{0}^{t} \psi^{\prime}(s)(\psi(t)-\psi(s))^{n-\alpha-1} u_{\psi}^{[n]}(s) d s \tag{13}
\end{align*}
$$

where $u_{\psi}^{[n]}(t)=\left(1 / \psi^{\prime}(t) d / d t\right)^{n} u(t)$ and $n=[\alpha]+1$ for $\alpha \notin N$, and $n=\alpha$ for $\alpha \in N$.

Lemma 3. Let $\alpha>0$, then the differential equation $\left({ }^{c} D_{a+}^{\alpha, \psi} h\right)(t)=0$ has the following solution:

$$
\begin{align*}
h(t)= & c_{0}+c_{1}(\psi(t)-\psi(0))+c_{2}(\psi(t)-\psi(0))^{2}  \tag{14}\\
& +\cdots+c_{n-1}(\psi(t)-\psi(0))^{n-1},
\end{align*}
$$

where $c_{i} \in R, i=0,1,2, \ldots, n-1$, and $n=[\alpha]+1$.

Lemma 4 (see [48]). Let $\alpha, \beta \in R^{+}$, and $f(t) \in L_{1}(I)$. Then, $I_{a+}^{\alpha, \psi} I_{a+}^{\beta, \psi} f(t)=I_{a+}^{\beta, \psi} I_{a+}^{\alpha, \psi} f(t)=I_{a+}^{\alpha+\beta, \psi} f(t)$, and $\left(I_{a+}^{\alpha, \psi}\right)^{n} f(t)$ $=I_{a+}^{n \alpha, \psi} f(t)$, where $n \in N$.

Definition 5 (see [15]). Let $X$ be any space and let $f: X \longrightarrow X$. A point $x \in X$ is called a fixed point for mapping $f$ if $x=f(x)$.

Theorem 6 (see [48]). Arzela-Ascoli Theorem.

Let $X$ be the Banach space of real or complex valued continuous functions normed by $\|f\|=\sup _{t \in X}|f(t)|$. If $A=$ $\left\{f_{n}\right\}$ is a sequence in $X$ such that $f_{n}$ is uniformly bounded and equi-continuous, then $\bar{A}$ is compact.

Theorem 7 (see [27]). Fixed point theorem on a cone.
Let $X$ be a Banach space, and let $S \subset X$ be a cone in $X$. Assume that $\Omega_{1}$ and $\Omega_{2}$ are bounded open subsets of $X$ with $0 \in \Omega_{1}, \bar{\Omega}_{1} \subset \Omega_{2}$, and let $f: S \cap\left(\bar{\Omega}_{2} \backslash \Omega_{1}\right) \longrightarrow S$ be a completely continuous operator such that either
(1) $\|f u\| \leq\|u\|, u \in S \cap\left(\partial \Omega_{1}\right)$, and $\|f u\| \geq\|u\|, u \in S \cap\left(\partial \Omega_{2}\right)$,
(2) $\|f u\| \geq\|u\|, u \in S \cap\left(\partial \Omega_{1}\right)$, and $\|f u\| \leq\|u\|, u \in S \cap\left(\partial \Omega_{2}\right)$,
Then, $f$ has a fixed point in $S \cap\left(\bar{\Omega}_{2} \backslash \Omega_{1}\right)$.
Lemma 8 (see [35]). Let $\phi_{p}$ be a p-Laplacian operator. Then, we have
(1) If $1<p \leq 2, \quad \zeta_{1} \zeta_{2}>0$, and $\left|\zeta_{1}\right|,\left|\zeta_{2}\right| \geq \varrho$, then $\left|\phi_{p}\left(\zeta_{1}\right)-\phi_{p}\left(\zeta_{2}\right)\right| \leq(p-1) \varrho^{p-2}\left|\zeta_{1}-\zeta_{2}\right|$
(2) If $p>2$, and $\left|\zeta_{1}\right|,\left|\zeta_{2}\right| \leq \varrho$, then

$$
\left|\phi_{p}\left(\zeta_{1}\right)-\phi_{p}\left(\zeta_{2}\right)\right| \leq(p-1) \varrho^{p-2}\left|\zeta_{1}-\zeta_{2}\right|
$$

2.2. Preliminary Results. In the following, we present the required results that will be used in our later discussion of the existence of positive solutions for the IBVP (5).

Lemma 9. The following IBVP

$$
\left\{\begin{array}{l}
-{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)=h(t), \quad t \in I=[0, T]  \tag{15}\\
u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) \mathrm{d} s \\
u(T)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{2}(s) u(s) \mathrm{d} s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) \mathrm{d} s
\end{array}\right.
$$

has a unique solution for any continuous function $h(t): I \longrightarrow R$, and $g_{i} \in C[0, T](i=1,2,3)$, such that

$$
\begin{align*}
u(t)= & \int_{0}^{T} \psi^{\prime}(s) G(t, s) h(s) d s \\
& +\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u(s) \mathrm{d} s \tag{16}
\end{align*}
$$

where

$$
G(t, s)=\frac{1}{\Gamma(\alpha)} \begin{cases}\frac{\varphi(t)}{\varphi(T)}(\psi(T)-\psi(s))^{\alpha-1}-(\psi(t)-\psi(s))^{\alpha-1}, & \text { if } 0 \leq s \leq t \leq T  \tag{17}\\ \frac{\varphi(t)}{\varphi(T)}(\psi(T)-\psi(s))^{\alpha-1}, & \text { if } 0 \leq t \leq s \leq T\end{cases}
$$

where $\quad \varphi(t)=(\psi(t)-\psi(0))\left(2\left(\psi^{\prime}(0)\right)^{2}-(\psi(t)-\right.$ $\left.\psi(0)) \psi^{\prime \prime}(0)\right)$, which is an increasing function such that $\varphi(t) \neq 0$ for all $t \in[0, T]$, and

$$
\Phi(t, s)=\frac{1}{\Gamma(\gamma)}\left[\left(1-\frac{\varphi(t)}{\varphi(T)}\right) g_{1}(s)+\frac{\varphi(t)}{\varphi(T)} g_{2}(s)+\frac{\chi(t)}{\varphi(T)} g_{3}(s)\right]
$$

where $\chi(t)=(\psi(T)-\psi(0))(\psi(t)-\psi(0))(\psi(t)-\psi(T))$.

Proof. From Lemma 3, we have

$$
\begin{align*}
u(t) & =-I_{0^{+}}^{\alpha, \psi} h(t)+c_{0}+c_{1}(\psi(t)-\psi(0))+c_{2}(\psi(t)-\psi(0))^{2} \\
& =-\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \psi^{\prime}(s)[\psi(t)-\psi(s)]^{\alpha-1} h(s) d s+c_{0}+c_{1}(\psi(t)-\psi(0))+c_{2}(\psi(t)-\psi(0))^{2} \tag{19}
\end{align*}
$$

Using the boundary conditions in (15), we obtain

$$
\begin{align*}
c_{0} & =u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s  \tag{20}\\
u(T) & =-\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \psi^{\prime}(s)[\psi(T)-\psi(s)]^{\alpha-1} h(s) d s+c_{0}+c_{1}(\psi(T)-\psi(0))+c_{2}(\psi(T)-\psi(0))^{2},
\end{align*}
$$

and
Solving (20) and (21) for $c_{1}$ and $c_{2}$, we obtain

$$
\begin{equation*}
u^{\prime \prime}(0)=c_{1} \psi^{\prime \prime}(0)+2 c_{2}\left(\psi^{\prime}(0)\right)^{2} \tag{21}
\end{equation*}
$$

$$
\begin{equation*}
c_{1}=\frac{-u^{\prime \prime}(0)[\psi(T)-\psi(0)]^{2}+2\left(\psi^{\prime}(0)\right)^{2}\left(1 / \Gamma(\alpha) \int_{0}^{T} \psi^{\prime}(s)[\psi(T)-\psi(s)]^{\alpha-1} h(s) d s+u(T)-u(0)\right)}{(\psi(T)-\psi(0))\left[2\left(\psi^{\prime}(0)\right)^{2}-(\psi(T)-\psi(0)) \psi^{\prime \prime}(0)\right]} \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
c_{2}=\frac{u^{\prime \prime}(0)[\psi(T)-\psi(0)]-\psi^{\prime \prime}(0)\left(1 / \Gamma(\alpha) \int_{0}^{T} \psi^{\prime}(s)[\psi(T)-\psi(s)]^{\alpha-1} h(s) d s+u(T)-u(0)\right)}{(\psi(T)-\psi(0))\left[2\left(\psi^{\prime}(0)\right)^{2}-(\psi(T)-\psi(0)) \psi^{\prime \prime}(0)\right]} \tag{23}
\end{equation*}
$$

Hence, if we take $\varphi(t)=(\psi(t)-\psi(0))\left(2\left(\psi^{\prime}(0)\right)^{2}-\right.$ $\left.(\psi(t)-\psi(0)) \psi^{\prime \prime}(0)\right) \quad$ where $\quad 2\left(\psi^{\prime}(0)\right)^{2}-(\psi(t)-$ $\psi(0)) \psi^{\prime \prime}(0)>1$, then

$$
\begin{align*}
u(t)= & -\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \psi^{\prime}(s)[\psi(t)-\psi(s)]^{\alpha-1} h(s) d s+c_{0}+c_{1}(\psi(t)-\psi(0))+c_{2}(\psi(t)-\psi(0))^{2} \\
= & -\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \psi^{\prime}(s)[\psi(t)-\psi(s)]^{\alpha-1} h(s) d s \\
& +\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
& -\frac{(\psi(T)-\psi(0))^{2}(\psi(t)-\psi(0))}{\Gamma(\gamma) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
& +2 \frac{(\psi(t)-\psi(0))\left(\psi^{\prime}(0)\right)^{2}}{\Gamma(\alpha) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)[\psi(T)-\psi(s)]^{\alpha-1} h(s) d s  \tag{24}\\
& +2 \frac{(\psi(t)-\psi(0))\left(\psi^{\prime}(0)\right)^{2}}{\varphi(T) \Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\left(g_{2}(s)-g_{1}(s)\right) u(s) d s \\
& +\frac{(\psi(T)-\psi(0))(\psi(t)-\psi(0))^{2}}{\Gamma(\gamma) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
& -\frac{(\psi(t)-\psi(0))^{2} \psi^{\prime \prime}(0)}{\Gamma(\alpha) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)[\psi(T)-\psi(s)]^{\alpha-1} h(s) d s \\
& -\frac{(\psi(t)-\psi(0))^{2} \psi^{\prime \prime}(0)}{\Gamma(\gamma) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\left(g_{2}(s)-g_{1}(s)\right) u(s) d s \\
&
\end{align*}
$$

Consequently from the fact that $\int_{0}^{T}=\int_{0}^{t}+\int_{t}^{T}$, we obtain

$$
\begin{align*}
u(t)= & -\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \psi^{\prime}(s)[\psi(t)-\psi(s)]^{\alpha-1} h(s) d s+\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
& -\frac{(\psi(T)-\psi(0))^{2}(\psi(t)-\psi(0))}{\Gamma(\gamma) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
& +2 \frac{(\psi(t)-\psi(0))\left(\psi^{\prime}(0)\right)^{2}}{\Gamma(\alpha) \varphi(T)} \int_{0}^{t} \psi^{\prime}(s)[\psi(t)-\psi(s)]^{\alpha-1} h(s) d s \\
& +2 \frac{(\psi(t)-\psi(0))\left(\psi^{\prime}(0)\right)^{\prime}}{\left.\varphi(T) \Gamma(\gamma)[\psi(T)-\psi(s)]^{\alpha-1} h(s) d s\right]} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\left(g_{2}(s)-g_{1}(s)\right) u(s) d s  \tag{25}\\
& +\frac{(\psi(T)-\psi(0))(\psi(t)-\psi(0))^{2} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s}{\Gamma(\gamma) \varphi(T)} \\
& -\frac{(\psi(t)-\psi(0))^{2} \psi^{\prime \prime}(0)}{\Gamma(\alpha) \varphi(T)}\left[\int_{0}^{t} \psi^{\prime}(s)[\psi(t)-\psi(s)]^{\alpha-1} h(s) d s\right] \\
& -\frac{(\psi(t)-\psi(0))^{2} \psi^{\prime \prime}(0)}{\Gamma(\gamma) \varphi(T)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\left(g_{2}(s)-g_{1}(s)\right) u(s) d s
\end{align*}
$$

which implies that $u(t)=\int_{0}^{T} \psi^{\prime}(s) G(t$, $s) h(s) d s+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u(s) d s$.

Therefore, the proof is completed.
Remark 10. Throughout our work, we take the following remarks in consideration:
(1) If $\quad m=\min \{\Phi(t, s):(t, s) \in I \times I\} \quad$ and $M=\max \{\Phi(t, s):(t, s) \in I \times I\}$, then we assume that

$$
\begin{equation*}
0<m \leq|\Phi(t, s)| \leq M<T, \quad\left(A_{0}\right) \tag{26}
\end{equation*}
$$

(2) Since $\psi(t)$ is an increasing function with $\psi^{\prime}(t) \neq 0$ for all $t \in I=[0, T]$, then $G(t, s)$ is continuous and positive for all $t, s \in I$.
(3) According to the definition of $G(t, s)$ and by $\left(A_{0}\right)$, we deduce that $\Phi(t, s) \in C\left(I \times I, R^{+}\right)$is continuous.

Lemma 11. For any $s \in I$, we have

$$
\begin{equation*}
\max _{t \in I} G(t, s) \leq \frac{\varphi(t)}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1} \tag{27}
\end{equation*}
$$

and for any $s \in[0, T]$, there exist $\varepsilon \in(0, T / 2), q(t)=\psi(t)-$ $\psi(t)^{\alpha-1}$ which is a concave function, and $\rho=\min _{t \in[\varepsilon, T-\varepsilon]} q(t)$ $=\min \left\{(\psi(\varepsilon)-\psi(0))-(\psi(\varepsilon)-\psi(0))^{\alpha-1},(\psi(T)-\psi\right.$
( $\left.\varepsilon))-(\psi(T)-\psi(\varepsilon))^{\alpha-1}\right\}$ such that

$$
\begin{equation*}
\min _{t \in[\varepsilon, T-\varepsilon]} G(t, s) \geq \frac{\rho}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1} \tag{28}
\end{equation*}
$$

Proof. From the expression for $G(t, s)$, it could be easily obtained that

$$
\begin{equation*}
\max _{t \in I} G(t, s) \leq \frac{\varphi(t)}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1} \tag{29}
\end{equation*}
$$

Now, if $0 \leq t \leq s \leq T$, then for $t \in[\varepsilon, T-\varepsilon]$ and since $\varphi$ is an increasing function, then we have:

$$
\begin{align*}
G(t, s) & =\frac{1}{\Gamma(\alpha) \varphi(T)} \varphi(t)(\psi(T)-\psi(s))^{\alpha-1} \\
& \geq \frac{1}{\Gamma(\alpha) \varphi(T)} \varphi(\varepsilon)(\psi(T)-\psi(s))^{\alpha-1} \tag{30}
\end{align*}
$$

Since $\psi$ and $\varphi$ are increasing functions, then

$$
\begin{align*}
G(t, s) & \geq \frac{1}{\Gamma(\alpha) \varphi(T)}[(\psi(\varepsilon)-\psi(0)) \underbrace{\left(2\left(\psi^{\prime}(0)\right)^{2}-(\psi(t)-\psi(0)) \psi^{\prime \prime}(0)\right)}_{>1}](\psi(T)-\psi(s))^{\alpha-1} \\
& \geq \frac{1}{\Gamma(\alpha) \varphi(T)}[\psi(\varepsilon)-\psi(0)](\psi(T)-\psi(s))^{\alpha-1}  \tag{31}\\
& \geq \frac{1}{\Gamma(\alpha) \varphi(T)}\left[(\psi(\varepsilon)-\psi(0))-(\psi(\varepsilon)-\psi(0))^{\alpha-1}\right](\psi(T)-\psi(s))^{\alpha-1} \\
& \geq \frac{\rho}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1}
\end{align*}
$$

If $0 \leq s \leq t \leq T$, then for $t \in[\varepsilon, T-\varepsilon]$ and $s \in[0, T-\varepsilon]$, we have the following equation:

$$
\begin{aligned}
G(t, s) & =\frac{1}{\Gamma(\alpha)}\left[\frac{\varphi(t)}{\varphi(T)}(\psi(T)-\psi(s))^{\alpha-1}-(\psi(t)-\psi(s))^{\alpha-1}\right] \\
& =\frac{1}{\Gamma(\alpha)}\left[\frac{\varphi(t)}{\varphi(T)}-\left(\frac{\psi(t)-\psi(s)}{\psi(T)-\psi(s)}\right)^{\alpha-1}\right](\psi(T)-\psi(s))^{\alpha-1} \\
& =\frac{1}{\Gamma(\alpha)}\left[\frac{\varphi(t)}{\varphi(T)}-\left(1-\frac{\psi(T)-\psi(t)}{\psi(T)-\psi(s)}\right)^{\alpha-1}\right](\psi(T)-\psi(s))^{\alpha-1} \\
& \geq \frac{1}{\Gamma(\alpha)}\left[\frac{\psi(t)-\psi(0)}{\psi(T)-\psi(0)}-\left(1-\frac{1-\psi(t)-\psi(0) / \psi(T)-\psi(0)}{1-\psi(s)-\psi(0) / \psi(T)-\psi(0)}\right)^{\alpha-1}\right](\psi(T)-\psi(s))^{\alpha-1} \\
& \geq \frac{\rho}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1} .
\end{aligned}
$$

Hence, we deduce that $\min _{t \in[\varepsilon, T-\varepsilon]} G(t, s) \geq \rho \quad$ Therefore, we obtain for all $t, s \in[0, T]$ that $(1 / \Gamma(\alpha) \varphi(T))(\psi(T)-\psi(s))^{\alpha-1}$ for any $s \in[0, T]$.

$$
\begin{equation*}
\rho \frac{1}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1} \leq G(t, s) \leq \frac{\varphi(t)}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1} \tag{33}
\end{equation*}
$$

In the following, let $X$ be a Banach space of all continuous functions from $I=[0, T]$ to $R$ with norm $\|u\|=\max _{t \in I}|u(t)|$, and let $P=\{u \in X: u(t) \geq 0, t \in I\}$ be a nonempty, bounded, closed, and convex cone in $X$.

Definition 12. Let $\mathscr{A}: C(I, R) \longrightarrow C(I, R)$ be an operator defined by the following equation:

$$
\begin{equation*}
\mathscr{A}(u(t))=\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u(s) \mathrm{d} s . \tag{34}
\end{equation*}
$$

Lemma 13. If assumption $\left(A_{0}\right)$ holds, then $\mathscr{A}$ satisfies the Proof following properties:
(1) $\mathscr{A}$ is continuous linear operator
(2) $\mathscr{A}$ is a bounded
(3) $\mathscr{A}(P) \subset P$
(4) $\mathscr{A}$ is reversible
(5) $\left\|(I-A)^{-1}\right\| \leq 1+N /(1-N)(\psi(T)-\psi(0))^{\gamma} / \gamma$, where $N=\left((\psi(T)-\psi(0))^{\gamma} / \gamma\right) M$
(1) $\mathscr{A}$ is continuous

Consider a sequence $\left\{u_{n}\right\} \subset P$ that converges to $u \in p$, i.e., $u_{n} \longrightarrow u$ in $P$ as $n \longrightarrow \infty$. Then,

$$
\begin{equation*}
\left|\mathscr{A}\left(u_{n}(t)\right)-\mathscr{A}(u(t))\right| \leq \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}|\Phi(t, s)|\left|u_{n}(s)-u(s)\right| \mathrm{d} s \tag{35}
\end{equation*}
$$

where $u_{n}, u \in C(I, R)$. Taking supremum for all $t \in I$, we obtain

$$
\begin{equation*}
\left\|\mathscr{A}\left(y_{n}\right)-\mathscr{A}(y)\right\| \leq M\left\|u_{n}-u\right\| \frac{(\psi(T)-\psi(0))^{\gamma}}{\gamma} \tag{36}
\end{equation*}
$$

Thus, $\left\|\mathscr{A}\left(y_{n}\right)-\mathscr{A}(y)\right\| \longrightarrow 0$ as $n \longrightarrow \infty$ and consequently $\mathscr{A}$ is continuous.
(2) $\mathscr{A}$ is bounded.

$$
\begin{align*}
\|\mathscr{A}(u(t))\| & \leq \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}|\Phi(t, s)||u(s)| d s \\
& \leq M\|u\| \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} d s \\
& \leq M\|u\| \frac{(\psi(T)-\psi(0))^{\gamma}}{\gamma} . \tag{37}
\end{align*}
$$

Thus, $\|\mathscr{A}\| \leq M(\psi(T)-\psi(0))^{\gamma} / \gamma$ which is bounded.
(3) $\mathscr{A}$ is compact

If $u(t) \in P$, then since for all $t, s \in I$, we have $\Phi(t, s)>0, u(t)>0$, and $\psi(t)$ is an increasing function defined on $R^{+}$, which implies that:

$$
\begin{equation*}
\mathscr{A}(u(t))=\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u(s) d s \geq 0 \tag{38}
\end{equation*}
$$

Thus, $\mathscr{A}(P) \subset P$.
(4) $\mathscr{A}$ is reversible since $\mathscr{A}$ is continuous, bounded, and compact.
(5) Let $v \in C[0, T]$ be defined as $v(t)=u(t)-\mathscr{A} u(t)$. Then,
$u(t)=(I-\mathscr{A})^{-1} v(t) \quad$ for all $t \in[0, T]$,
$u(t)=v(t)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u(s) d s$.

By using the iteration method, let
$\square$

$$
\begin{equation*}
u_{i}(t)=v(t)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u_{i-1}(s) d s ; \quad i=1,2,3, \ldots, \tag{41}
\end{equation*}
$$

with $u_{0}(t)=u(t)$, and $\Phi_{1}(t, s)=\Phi(t, s)$, we obtain

$$
\begin{equation*}
u(t)=v(t)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, s) v(s) d s \tag{42}
\end{equation*}
$$

$$
\begin{align*}
\mathscr{R}(t, s) & =\sum_{j=i}^{\infty} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, \tau) \Phi_{j-1}(\tau, s) d \tau \\
& =\sum_{j=i}^{\infty} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi_{j}(t, s) \tag{43}
\end{align*}
$$

where
where

$$
\Phi_{j}(t, s)=\int_{0}^{T} \psi^{\prime}(\tau)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, \tau) \Phi_{j-1}(\tau, s) \mathrm{d} \tau
$$

But, $0 \leq \Phi(t, s) \leq M \leq T$. Then, $\mathscr{R}(t, s) \geq 0$, and

$$
\begin{align*}
\mathscr{R}(t, s) & =\sum_{j=i}^{\infty} \psi^{\prime}(s)(\psi(T)-\psi(0))^{\gamma-1} \Phi_{j}(t, s) \\
& <\frac{(\psi(T)-\psi(0))^{\gamma}}{\gamma} M+\frac{(\psi(T)-\psi(0))^{2 \gamma}}{\gamma^{2}} M^{2}+\cdots+\frac{(\psi(T)-\psi(0))^{n \gamma}}{\gamma^{n}} M^{n}+\cdots  \tag{45}\\
& <N+N^{2}+\cdots+N^{n}+\cdots=\frac{N}{1-N}
\end{align*}
$$

Thus,

$$
\begin{align*}
\left|(I-\mathscr{A})^{-1} v(t)\right| & \leq|v(t)|+\frac{N}{1-N}\left|\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\right| v(s)|d s|  \tag{46}\\
& \leq\|v\|+\frac{N}{1-N} \frac{(\psi(T)-\psi(0))^{\gamma}}{\gamma}\|v\|,
\end{align*}
$$

which implies that

$$
\begin{equation*}
\left\|(I-\mathscr{A})^{-1}\right\| \leq 1+\frac{N}{1-N} \frac{(\psi(T)-\psi(0))^{\gamma}}{\gamma} . \tag{47}
\end{equation*}
$$

The proof is completed.
Lemma 14. Suppose that assumption $\left(A_{0}\right)$ holds, $u,{ }^{c} D^{\alpha, \psi} u \in C(I)$, then the integral boundary value problem
is equivalent to the following integral equation

$$
\begin{equation*}
u(t)=\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) f(s, u(s)) \mathrm{d} s \tag{49}
\end{equation*}
$$

where

$$
\mathscr{H}(t, s)=G(t, s)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, \tau) G(\tau, s) d \tau .
$$

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)+f(t, u(t))=0, \quad t \in I=[0, T]  \tag{50}\\
u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
u(T)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{2}(s) u(s) d s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s
\end{array}\right.
$$

Proof. Define the nonlinear operator $\mathscr{B}: C(I) \longrightarrow C(I)$ as

$$
\begin{equation*}
\mathscr{B} u(t)=\int_{0}^{T} \psi^{\prime}(s) G(t, s) f(s, u(s)) d s \tag{51}
\end{equation*}
$$

From (34) and (51), the IBVP (48) is equivalent to the following integral equation:
(48)

$$
\begin{align*}
u(t) & =\int_{0}^{T} \psi^{\prime}(s) G(t, s) f(s, u(s)) d s+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \Phi(t, s) u(s) d s  \tag{52}\\
& =\mathscr{B} u(t)+\mathscr{A} u(t)
\end{align*}
$$

By Lemma 13, we have the following equation:

$$
\begin{equation*}
(I-\mathscr{A})^{-1} \mathscr{B} u(t)=u(t) \tag{53}
\end{equation*}
$$

$$
\begin{aligned}
u(t)= & \int_{0}^{T} \psi^{\prime}(s) G(t, s) f(s, u(s)) d s \\
& +\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, s) \int_{0}^{T} \psi^{\prime}(s) G(s, \tau) f(\tau, u(\tau)) d \tau d s \\
= & \int_{0}^{T} \psi^{\prime}(s) G(t, s) f(s, u(s)) d s \\
& +\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\left(\int_{0}^{T} \mathscr{R}(t, \tau) \psi^{\prime}(s) G(\tau, s) f(s, u(s)) d \tau\right) d s \\
= & \int_{0}^{T} \psi^{\prime}(s) G(t, s) f(s, u(s)) d s \\
& +\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1}\left(\int_{0}^{T} \mathscr{R}(t, \tau) \psi^{\prime}(s) G(\tau, s) f(s, u(s)) d \tau\right) d s \\
= & \int_{0}^{T} \psi^{\prime}(s)\left(G(t, s)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, \tau) G(\tau, s) d \tau\right) f(s, u(s)) d s \\
= & \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) f(s, u(s)) d s,
\end{aligned}
$$

where $\quad \mathscr{H}(t, s)=G(t, s)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\quad$ (1) $\mathscr{H}(t, s)$ is continuous and positive $\psi(s))^{\gamma-1} \mathscr{R}(t, \tau) G(\tau, s) d \tau$.
(2) $n /(1-n) \leq \mathscr{R}(t, s) \leq N /(1-N)$

Lemma 15. Assumption that $\left(A_{0}\right)$ holds, then for any $t, s \in[0, T]$, functions $\mathscr{H}$ and $\mathscr{R}$ satisfy the following properties:

$$
\begin{aligned}
& \left(\rho n \psi^{\prime}(s) / \Gamma(\alpha)(1-n)\right)\left((\psi(T)-\psi(s))^{\alpha+\gamma-2} / \varphi(T)\right)(T-2 \varepsilon) \leq \mathscr{H}(t, s) \\
& \leq\left((\psi(T)-\psi(s))^{\alpha-1} / \Gamma(\alpha) \varphi(T)\right)\left[\varphi(t)+T N /(1-N)(\psi(T)-\psi(s))^{\gamma-1}\right]
\end{aligned}
$$

Proof. Let $t, s \in[0, T]$, then
(1) It is clear from expression (11) that $\mathscr{H}(t, s)$ is continuous. In addition, by Remark 10 and expression (9), we have $\mathscr{H}(t, s) \geq 0$.
(2) From expression (9), we have

$$
\begin{equation*}
\frac{n}{1-n} \leq \mathscr{R}(t, s) \leq \frac{N}{1-N} \tag{55}
\end{equation*}
$$

where $\quad n=\left((\psi(T)-\psi(s))^{(\gamma-1)} / \gamma\right) m \quad$ and $\quad N=$ $\left((\psi(T)-\psi(s))^{(\gamma-1)} / \gamma\right) M$.
(3) Suppose that $\left(A_{0}\right)$ holds, then from (45) and (50), and Lemma 11 we have the following equation:

$$
\begin{align*}
\mathscr{H}(t, s) & =G(t, s)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, \tau) G(\tau, s) d \tau \\
& \geq \int_{\varepsilon}^{T-\varepsilon} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, \tau) G(\tau, s) d \tau \\
& \geq \frac{\rho n}{\Gamma(\alpha)(1-n)} \int_{\varepsilon}^{T-\varepsilon} \psi^{\prime}(s) \frac{(\psi(T)-\psi(s))^{\alpha-1}}{\varphi(T)}(\psi(T)-\psi(s))^{\gamma-1} d \tau  \tag{56}\\
& \geq \frac{\rho n \psi^{\prime}(s)}{\Gamma(\alpha)(1-n)} \frac{(\psi(T)-\psi(s))^{\alpha+\gamma-2}}{\varphi(T)}(T-2 \varepsilon) .
\end{align*}
$$

Similarly,

$$
\begin{align*}
\mathscr{H}(t, s) & =G(t, s)+\int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} \mathscr{R}(t, \tau) G(\tau, s) d \tau \\
& \leq \frac{\varphi(t)}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1}+\frac{N}{1-N} \int_{0}^{T} \frac{(\psi(T)-\psi(s))^{\alpha+\gamma-2}}{\Gamma(\alpha) \varphi(T)} d \tau \\
& \leq \frac{\varphi(t)}{\Gamma(\alpha) \varphi(T)}(\psi(T)-\psi(s))^{\alpha-1}+\frac{T N}{1-N} \frac{(\psi(T)-\psi(s))^{\alpha+\gamma-2}}{\Gamma(\alpha) \varphi(T)}  \tag{57}\\
& \leq \frac{(\psi(T)-\psi(s))^{\alpha-1}}{\Gamma(\alpha) \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right] .
\end{align*}
$$

Consider the $p$-Laplacian operator $\phi_{p}$ defined as $\phi_{p}=$ $|s|^{p-2} s$ with $p>1$, and let $q>1$ be a real number satisfying the relation $1 / p+1 / p=1$, and $\phi_{p}^{-1}=\phi_{q}$. In the following, we
consider the associated linear $p$-Laplacian fractional differential equations involving integral boundary conditions

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\beta, \psi}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)\right)\right)+h(t)=0, \quad t \in I=[0, T]  \tag{58}\\
u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
u(T)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{2}(s) u(s) d s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
\left.{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)\right|_{t=0}=0
\end{array}\right.
$$

where $h \in C[0, T]$ and $h \geq 0$.
Lemma 16. The associated linear p-Laplacian fractional differential equations (58) has the unique solution

$$
\begin{equation*}
u(t)=\frac{1}{\Gamma(\beta)} \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} h(\tau) d \tau\right) \mathrm{d} s \tag{59}
\end{equation*}
$$

Proof. Let $w={ }^{c} D_{0^{+}}^{\alpha, \psi} u$, and $v=\phi_{p}(w)$. Then, the solution of the initial value problem

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\beta, \psi} v(t)+h(t)=0, \quad t \in I=[0, T]  \tag{60}\\
v(0)=0
\end{array}\right.
$$

is given by $v(t)=-{ }^{c} I_{0^{+}}^{\beta, \psi} h(t)+c_{0}$ for any $t \in[0, T]$. Since $v(0)=0,0<\beta \leq 1$, then we obtain $c_{0}=0$, which implies that

$$
\begin{equation*}
v(t)=-{ }^{c} I_{0^{+}}^{\beta, \psi} h(t), \quad t \in[0, T] . \tag{61}
\end{equation*}
$$

But $w=\phi_{q}(v)$ which implies that the solution of the IBVP (58) satisfies the following IBVP:

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)=\phi_{q}\left(-{ }^{c} I_{0^{+}}^{\beta, \psi} h(t)\right), \quad t \in I=[0, T]  \tag{62}\\
u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
u(T)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{2}(s) u(s) d s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
\left.{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)\right|_{t=0}=0
\end{array}\right.
$$

Since $h(t) \geq 0$, we have $\phi_{q}\left(-^{c} I_{0^{+}}^{\beta, \psi} h(t)\right)=-\phi_{q}\left({ }^{c} I_{0^{+}}^{\beta, \psi} h(t)\right)$.
Hence, by Lemma 11, we deduce that the solution of the IBVP (62) can be written as follows:

$$
\begin{equation*}
u(t)=\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s)\left(\phi_{q}\left({ }^{c} I_{0^{+}}^{\beta, \psi} h(s)\right)\right) \mathrm{d} s \tag{63}
\end{equation*}
$$

which implies that

$$
\begin{equation*}
u(t)=\frac{1}{\Gamma(\beta)} \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} h(\tau) d \tau\right) \mathrm{d} s \tag{64}
\end{equation*}
$$

Lemma 17. Suppose that assumption $\left(A_{0}\right)$ holds, and $u,{ }^{c} D_{0^{+}}^{\alpha, \psi} u \in C[0, T]$, the the following integral boundary value problem:

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{\beta, \psi}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{\alpha, \psi} u\right)\right)(t)+\lambda f(t, u(t))=0, \quad t \in[0, T]  \tag{65}\\
u(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{1}(s) u(s) d s \\
u(T)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{2}(s) u(s) d s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(\gamma)} \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\gamma-1} g_{3}(s) u(s) d s \\
\left.{ }^{c} D_{0^{+}}^{\alpha, \psi} u(t)\right|_{t=0}=0
\end{array}\right.
$$

is equivalent to the following integral equation

$$
\begin{equation*}
u(t)=\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) \mathrm{d} s \tag{66}
\end{equation*}
$$

Definition 18. Let $X$ be a Banach space with norm $\|u\|=\max _{t \in[0, T]}|u(t)|$. Define the reproductory cone $\mathscr{P}_{0} \subset X$ to be the ${ }^{t \in[0, T]}$

$$
\begin{equation*}
\mathscr{P}_{0}=\{u \in P: u(t) \geq \sigma\|u\|, \quad t \in I\} \tag{67}
\end{equation*}
$$

where $\quad \sigma=\rho n(T-2 \varepsilon)(\psi(T)-\psi(s))^{\gamma-1} /(1-n)[\varphi(t)+T$ $\left.N / 1-N(\psi(T)-\psi(s))^{\gamma-1}\right]$ such that $0<\sigma<T$.

Definition 19. Let $X$ be a Banach space, define the operator $\mathscr{C}_{\lambda}: X \longrightarrow X$ as follows:

$$
\begin{equation*}
\mathscr{C}_{\lambda} u(t)=\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) \mathrm{d} \tau\right) \mathrm{d} s \tag{68}
\end{equation*}
$$

Consider the following assumptions:
$\left(A_{1}\right)$ : The nonlinear function $f: I \times R^{2} \longrightarrow R$ is continuous and there exist $\mu \in C\left(I, R^{+}\right)$with norm $\|\mu\|$ such that:

$$
\begin{equation*}
|f(t, u(t))-f(t, v(t))| \leq \mu(t)|u(t)-v(t)|, \quad \text { for all } t \in I \tag{69}
\end{equation*}
$$

Lemma 20. Suppose that assumption $\left(A_{0}\right)$ and $\left(A_{1}\right)$ hold, then $\mathscr{C}_{\lambda}: \mathscr{P}_{0} \longrightarrow \mathscr{P}_{0}$ is completely continuous

Proof. First, the operator $\mathscr{C}_{\lambda}$ is continuous.
Consider a sequence $\left\{u_{n}\right\} \subset \mathscr{P}_{0}$ that converges to $y \in \mathscr{P}_{0}$, i.e., $u_{n} \longrightarrow u$ in $\mathscr{P}_{0}$ as $n \longrightarrow \infty$. Then, by $\left(A_{1}\right)$ we have

$$
\begin{align*}
& \left|\mathscr{C}_{\lambda}\left(u_{n}(t)\right)-\mathscr{C}_{\lambda}(u(t))\right| \\
& \quad \leq \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}\left|f\left(\tau, u_{n}(\tau)\right)-f(\tau, u(\tau))\right| d \tau\right) d s \\
& \leq \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|\mu(t)|\left|u_{n}(\tau)-u(\tau)\right| d \tau\right) d s  \tag{70}\\
& \leq \frac{1}{\Gamma(\alpha) \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right] \\
& \quad \times \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\alpha-1} \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|\mu(t)|\left|u_{n}(\tau)-u(\tau)\right| d \tau\right) d s .
\end{align*}
$$

Taking supremum for all $t \in I$, then applying Lebesgue dominated convergence theorem, we get

$$
\begin{align*}
& \left|\mathscr{C}_{\lambda}\left(u_{n}(t)\right)-\mathscr{C}_{\lambda}(u(t))\right| \\
& \quad \leq \frac{\lambda\|\mu\|\left\|u_{n}-u\right\|}{\Gamma(\alpha) \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right]  \tag{71}\\
& \quad \times \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\alpha-1} \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} d \tau\right) d s .
\end{align*}
$$

Since $u_{n} \longrightarrow u$, then we obtain $\mathscr{C}_{\lambda}\left(u_{n}(t)\right) \longrightarrow \mathscr{C}_{\lambda}(u(t))$ as $n \longrightarrow \infty$ for each $t \in I$, and hence $\left\|\mathscr{C}_{\lambda}\left(u_{n}(t)\right)-\mathscr{C}_{\lambda}(u(t))\right\| \longrightarrow 0$ as $n \longrightarrow \infty$ and consequently $\mathscr{A}$ is continuous.

In other words, since $H(t, s)$ and $f(s, u(s))$ are nonnegative and continuous, then it is clear that operator $\mathscr{C}_{\lambda}$ is continuous.

Second, the operator $\mathscr{C}_{\lambda}$ maps bounded sets in $\mathscr{P}_{0}$ into bounded sets in $\mathscr{P}_{0}$.

From Lemma 14, we have

$$
\begin{align*}
\mathscr{C}_{\lambda} u(t)= & \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s \\
\geq & \frac{\rho n(T-2 \varepsilon)}{\Gamma(\alpha)(1-n)} \frac{(\psi(T)-\psi(s))^{\gamma-1}}{\varphi(T)}  \tag{72}\\
& \quad \times \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\alpha-1} \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s
\end{align*}
$$

and

$$
\begin{align*}
\mathscr{C}_{\lambda} u(t) \leq & \frac{1}{\Gamma(\alpha) \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right]  \tag{73}\\
& \times \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\alpha-1} \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s
\end{align*}
$$

which implies that

$$
\begin{align*}
\mathscr{C}_{\lambda} u(t) & \geq \frac{\rho n(T-2 \varepsilon)(\psi(T)-\psi(s))^{\gamma-1}}{(1-n)\left[\varphi(t)+(T N /(1-N))(\psi(T)-\psi(s))^{\gamma-1}\right]}\left\|\mathscr{C}_{\lambda} u(t)\right\|  \tag{74}\\
& \geq \sigma\left\|\mathscr{C}_{\lambda} u(t)\right\|
\end{align*}
$$

Hence, $\mathscr{C}_{\lambda}\left(\mathscr{P}_{0}\right) \subset \mathscr{P}_{0}$.

Third, $\mathscr{C}_{\lambda}(\Omega)$ is uniformly bounded for every $\Omega$ is bounded in $\mathscr{P}_{0}$.

Let $\Omega \subset \mathscr{P}_{0}$ be bounded. Then, there exists a real number $\ell>0$ such that for any $u \in \Omega,\|u\| \leq \ell$. Thus,

$$
\begin{align*}
\left|\mathscr{C}_{\lambda} u(t)\right|= & \left|\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s\right| \\
\leq & \int_{0}^{T} \psi^{\prime}(s)|\mathscr{H}(t, s)| \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|f(\tau, u(\tau))| d \tau\right) d s \\
\leq & \frac{\lambda^{q-1}}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right]  \tag{75}\\
& \times \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|\mu(\tau)||u(\tau)| d \tau\right) d s .
\end{align*}
$$

Taking supremum for all $t \in I$, we get

$$
\begin{align*}
\left\|\mathscr{C}_{\lambda} u(t)\right\| & \leq \frac{\lambda^{q-1} \ell\|\mu\|}{\Gamma(\alpha) \Gamma(\beta) \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right] \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\frac{(\psi(s)-\psi(0))^{\beta}}{\beta}\right) d s \\
& \leq \frac{\lambda^{q-1} \ell\|\mu\|}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right] \int_{0}^{T} \psi^{\prime}(s)\left(\frac{(\psi(s)-\psi(0))^{\beta}}{\beta}\right)^{q-1} d s  \tag{76}\\
& \leq \frac{\lambda^{q-1} \ell\|\mu\|}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right] \frac{1}{q}\left(\frac{(\psi(T)-\psi(0))^{\beta}}{\beta}\right)^{q} \\
& \leq \frac{\lambda^{q-1} \ell\|\mu\|}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \frac{(\psi(T)-\psi(0))^{q \beta}}{q \beta^{q}}<+\infty
\end{align*}
$$

Finally, $\mathscr{C}_{\lambda}$ is equi-continuous.

Since $H(t, s)$ is continuous and hence uniformly continuous on $[0,1] \times[0,1]$, then for every $\varepsilon>0$, there exist $\delta>0$ and $t_{1}, t_{2} \in I$, with $t_{1}<t_{2}$ and $\left|t_{2}-t_{1}\right|<\delta$
such that $\left|\mathscr{H}\left(t_{2}, s\right)-\mathscr{H}\left(t_{1}, s\right)\right|<\varepsilon$. Then, for all $u \in \mathscr{P}_{0}$ we have

$$
\begin{align*}
& \left|\mathscr{C}_{\lambda} u\left(t_{2}\right)-\mathscr{C}_{\lambda} u\left(t_{1}\right)\right| \\
& \quad \leq\left|\int_{0}^{T} \psi^{\prime}(s)\left(\mathscr{H}\left(t_{2}, s\right)-\mathscr{H}\left(t_{1}, s\right)\right) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s\right| \\
& \quad \leq \int_{0}^{T} \psi^{\prime}(s)\left|\mathscr{H}\left(t_{2}, s\right)-\mathscr{H}\left(t_{1}, s\right)\right| \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|f(\tau, u(\tau))| d \tau\right) d s \\
& \quad \leq \frac{\lambda^{q-1} \varepsilon}{\Gamma(\beta)^{q-1}} \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|\mu(\tau)||u(\tau)| d \tau\right) d s  \tag{77}\\
& \quad \leq \frac{\lambda^{q-1}\|\mu\|\|u\| \varepsilon}{\Gamma(\beta)^{q-1}} \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} d \tau\right) d s \\
& \quad \leq \frac{\lambda^{q-1}\|\mu\|\|u\| \varepsilon}{\Gamma(\beta)^{q-1}} \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\frac{(\psi(s)-\psi(0))^{\beta}}{\beta}\right) d s \\
& \quad \leq\left(\frac{\lambda^{q-1}\|\mu\|\|u\| \|}{\Gamma(\beta)^{q-1}} \frac{(\psi(T)-\psi(0))^{q \beta}}{q \beta^{q}}\right) \varepsilon .
\end{align*}
$$

As $t_{1} \longrightarrow t_{2}$, the right-hand side of the above inequality is not dependent on $u$ and tends to zero. Consequently,

$$
\begin{equation*}
\left|\mathscr{A}\left(y\left(t_{2}\right)\right)-\mathscr{A}\left(y\left(t_{1}\right)\right)\right| \longrightarrow 0, \forall\left|t_{2}-t_{1}\right| \longrightarrow 0 \tag{78}
\end{equation*}
$$

This show that $\left\{\mathscr{C}_{\lambda} u\right\}$ is equi-continuous on $\mathscr{P}_{0}$, and $\mathscr{C}_{\lambda}$ is a compact and completely continuous operator by the Arzela-Ascoli Theorem 6.
2.3. Existence of Positive Solutions. In this subsection, we study the existence of positive solutions for the IBVP (5) by and sufficient conditions for having at least one positive solution are presented.

Notation 21. Consider the following notations:
(1) $f^{0}=\lim _{u \rightarrow 0^{+}} \sup \sup _{t \in[0, T]} f(t, u) / \phi_{p}(u)$
(2) $f^{\infty}=\lim _{u \rightarrow+\infty} \sup \sup _{t \in[0, T]} f(t, u) / \phi_{p}(u)$
(3) $A_{1}=\left(\|\mu\| / \Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)\right)[\varphi(t)+T N /(1-N)$ $\left.(\psi(T)-\psi(0))^{\gamma-1}\right](\psi(T)-\psi(0))^{q \beta} / q \beta^{q q \beta}$
(4) applying the fixed point Theorem 7 on a cone. The required

$$
\begin{gathered}
A_{2}=\left(\rho n(T-2 \varepsilon) /\left(\left(\Gamma(\alpha) \Gamma(\beta)^{(q-1)}(1-n)\right)\right)(\psi(T)-\psi(0))^{(\gamma-1)}\right) /(\varphi(T)) \\
(1 / \beta(\beta+1)(\beta+2))\left[\begin{array}{c}
(\psi(T)-\psi(0))^{\beta+1}((\beta+1)(\psi(T)-\psi(0))+\psi(0)) \\
-(\beta+2)(\psi(T)-\psi(0))
\end{array}\right]
\end{gathered}
$$

Theorem 22. Assume that $\left(A_{0}\right)$ holds. If we have $\zeta, \aleph>0$ such that $f^{0}<\zeta$ and $f^{\infty}>\aleph$, then for each real number $\lambda$ satisfying

$$
\begin{equation*}
\frac{1}{A_{2}^{p-1} \aleph} \leq \lambda \leq \frac{1}{A_{1}^{p-1} \zeta} \tag{79}
\end{equation*}
$$

The IBVP (5) has at least one positive solution.

Proof. By Notation 21, there exist two positive real numbers $\widetilde{\rho}, \rho_{1}$, and $\rho_{2}$ such that $\rho_{2}>(1 / \sigma) \widetilde{\rho}$ and $\rho_{2}=\max \left\{2 \rho_{1},(1 / \sigma)\right.$ $\widetilde{\rho}\}$. Then, for any $t \in[0, T]$, we have

$$
\begin{equation*}
f(t, u) \leq \zeta \phi_{p}(u), \quad \text { for } u \in\left(0, \rho_{1}\right] \tag{80}
\end{equation*}
$$

and

$$
\begin{equation*}
f(t, u) \geq \aleph \phi_{p}(u), \quad \text { for } u \in\left[\rho_{2},+\infty\right) \tag{81}
\end{equation*}
$$

## Define

$$
\begin{align*}
\left\|\mathscr{C}_{\lambda} u(t)\right\| \leq & \sup _{t \in[0, T]} \int_{0}^{T} \psi^{\prime}(s)|\mathscr{H}(t, s)| \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|f(\tau, u(\tau))| d \tau\right) d s \\
\leq & \frac{\lambda^{q-1}}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(s))^{\gamma-1}\right] \\
& \times \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} \zeta \phi_{p}\left(\rho_{1}\right) d \tau\right) d s  \tag{83}\\
\leq & \frac{\lambda^{q-1} \zeta^{q-1}\|\mu\|}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \frac{(\psi(T)-\psi(0))^{q \beta}}{q \beta^{q}} \rho_{1} \\
\leq & \lambda^{q-1} \zeta^{q-1} A_{1} \rho_{1}
\end{align*}
$$

where

$$
\begin{equation*}
A_{1}=\frac{\|\mu\|}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \frac{(\psi(T)-\psi(0))^{q \beta}}{q \beta^{q}} \tag{84}
\end{equation*}
$$

Hence,

$$
\begin{equation*}
\left\|\mathscr{C}_{\lambda} u\right\| \leq\|u\|, \quad \text { for any } u \in \mathscr{P}_{0} \cap \partial \Omega_{1} \tag{85}
\end{equation*}
$$

$$
\begin{align*}
u(t) & \geq \frac{\rho n(T-2 \varepsilon)(\psi(T)-\psi(s))^{\gamma-1}}{(1-n)\left[\varphi(t)+(T N /(1-N))(\psi(T)-\psi(s))^{\gamma-1}\right]}\|u(t)\| \\
& \geq \sigma\|u\| \\
& \geq \sigma \rho_{2}>\tilde{\rho} . \tag{86}
\end{align*}
$$

Then, by (79) and (81), we have

$$
\begin{align*}
\left\|\mathscr{C}_{\lambda} u(t)\right\| \geq & \sup _{t \in[0, T]} \int_{0}^{T} \psi^{\prime}(s)|\mathscr{H}(t, s)| \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|f(\tau, u(\tau))| d \tau\right) d s \\
\geq & \frac{\rho n(T-2 \varepsilon)}{\Gamma(\alpha)(1-n)} \frac{(\psi(T)-\psi(s))^{\gamma-1}}{\varphi(T)} \\
& \times \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\alpha-1} \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} \aleph \phi_{p}\left(\rho_{2}\right) d \tau\right) d s \\
\geq & \frac{\rho n(T-2 \varepsilon)}{\Gamma(\alpha)(1-n)} \frac{(\psi(T)-\psi(s))^{\gamma-1}}{\varphi(T)}\left(\frac{\lambda \aleph}{\Gamma(\beta)}\right)^{q-1}  \tag{87}\\
& \times \int_{0}^{T} \psi^{\prime}(s)(\psi(T)-\psi(s))^{\alpha-1} \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} \phi_{p}\left(\rho_{2}\right) d \tau\right) d s \\
\geq & (\lambda \aleph)^{q-1} \frac{\rho n(T-2 \varepsilon)}{\Gamma(\alpha) \Gamma(\beta)^{q-1}(1-n)} \frac{(\psi(T)-\psi(0))^{\gamma-1}}{\varphi(T)} \frac{1}{(\beta+2)(\beta+1) \beta} \\
& \times\left[(\psi(T)-\psi(0))^{\beta+1}((\beta+1)(\psi(T)-\psi(0))+\psi(0))-(\beta+2)(\psi(T)-\psi(0))\right] \rho_{2} \\
\geq & \lambda^{q-1} \aleph^{q-1} A_{2} \rho_{2} \\
\geq & \rho_{2}=\|u\|,
\end{align*}
$$

where

$$
\begin{align*}
A_{2}= & \frac{\rho n(T-2 \varepsilon)}{\Gamma(\alpha) \Gamma(\beta)^{q-1}(1-n)} \frac{(\psi(T)-\psi(0))^{\gamma-1}}{\varphi(T)} \\
& \times \frac{1}{(\beta+2)(\beta+1) \beta}\left[\begin{array}{c}
(\psi(T)-\psi(0))^{\beta+1}((\beta+1)(\psi(T)-\psi(0))+\psi(0)) \\
-(\beta+2)(\psi(T)-\psi(0))
\end{array}\right] \tag{88}
\end{align*}
$$

This implies that

$$
\begin{equation*}
\left\|\mathscr{C}_{\lambda} u\right\| \geq\|u\|, \quad \text { for any } u \in \mathscr{P}_{0} \cap \partial \Omega_{2} . \tag{89}
\end{equation*}
$$

Finally, by 20, 21, and Theorem 7, we conclude that By (81) and (89) and Theorem 22, we conclude $\mathscr{C}_{\lambda}$ has a fixed point $u \in \mathscr{P}_{0} \cap\left(\bar{\Omega}_{2} \backslash \Omega_{1}\right)$ with $\rho_{1} \leq\|u\| \leq \rho_{2}$, and obviously the IBVP (5) has at least one positive solution $u$.

Now, if we use other conditions for $f^{0}$ and $f^{\infty}$, we obtain the following theorem.

Theorem 23. Assume that $\left(A_{0}\right)$ holds. If we have $\zeta, \aleph>0$ such that $f^{0}>\aleph$, and $f^{\infty}<\zeta$, then for each real number $\lambda$ satisfying

$$
\begin{equation*}
\frac{1}{A_{2}^{p-1} \aleph} \leq \lambda \leq \frac{1}{A_{1}^{p-1} \zeta} \tag{90}
\end{equation*}
$$

The IBVP (5) has at least one positive solution.

Proof. Suppose that $\lambda$ satisfies (90), then there exists $\rho_{1}>0$ for any $t \in[0, T]$ so that

$$
\begin{equation*}
f(t, u) \geq \aleph \phi_{p}(u), \quad \text { for } u \in\left(0, \rho_{2}\right] \tag{91}
\end{equation*}
$$

Hence, if $u \in \mathscr{P}_{0}$ with norm $\|u\|=\rho_{2}$, then we can choose $\Omega_{1}$ and $\Omega_{2}$ as carried out in the previous theorem such that

$$
\begin{equation*}
\left\|\mathscr{C}_{\lambda} u\right\| \geq\|u\|, \quad \text { for all } u \in \mathscr{P}_{0} \cap \partial \Omega_{1} \tag{92}
\end{equation*}
$$

In addition, for any $t \in[0, T]$, there exists $\rho_{2}>0$ so that

$$
\begin{equation*}
f(t, u) \geq \zeta \phi_{p}(u), \quad \text { for } u \in\left[\rho_{2},+\infty\right) . \tag{93}
\end{equation*}
$$

Hence, if $u \in \mathscr{P}_{0}$ with norm $\|u\|=\rho_{2}$, then we have

$$
\begin{equation*}
\left\|\mathscr{C}_{\lambda} u\right\| \leq\|u\|, \quad \text { for all } u \in \mathscr{P}_{0} \cap \partial \Omega_{2} \tag{94}
\end{equation*}
$$

By 20, 21, and Theorem 7, we deduce that $\mathscr{C}_{\lambda}$ has a fixed point $u \in \mathscr{P}_{0} \cap\left(\bar{\Omega}_{2} / \Omega_{1}\right)$ with $\rho_{1} \leq\|u\| \leq \rho_{2}$, and obviously the IBVP (5) has at least one positive solution $u$.
2.4. Nonexistence of Positive Solutions. In this part, we present the conditions for the nonexistence of positive solutions for the IBVP (5).

Theorem 24. If $f^{0}<+\infty$ and $f^{\infty}<+\infty$, then for any $t \in[0, T]$ there exist positive constants $K_{1}, K_{2}, \rho_{1}$, and $\rho_{2}$ such that $\rho_{1}<\rho_{2}$ and

$$
\begin{equation*}
f(t, u) \leq K_{1} \phi_{p}(u), \quad \text { for } u \in\left[0, \rho_{1}\right] \tag{95}
\end{equation*}
$$

and

$$
\begin{equation*}
f(t, u) \geq K_{2} \phi_{p}(u), \quad \text { for } u \in\left[\rho_{2},+\infty\right) \tag{96}
\end{equation*}
$$

Denote by $\bar{K}=\max \left\{K_{1}, K_{2}, \max _{\rho_{1} \leq u \leq \rho_{2}, 0 \leq t \leq T} f(t\right.$, $\left.u) / \phi_{p}(u)\right\}$, then

$$
\begin{equation*}
f(t, u) \leq \bar{K} \phi_{p}(u), \quad \text { for any } t \in[0, T], u \in[0,+\infty] . \tag{97}
\end{equation*}
$$

Assume conversely that $v(t)$ is a positive solution of the IBVP (5). Then, since $C_{\lambda} v(t)=v(t)$ for any $t \in[0, T]$, we have

$$
\begin{align*}
\|v\|= & \left\|C_{\lambda} v\right\| \leq \frac{\lambda^{q-1}}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \\
& \times \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|f(\tau, u(\tau))| d \tau\right) d s \\
\leq & \frac{\lambda^{q-1}}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right]  \tag{98}\\
& \times \int_{0}^{T} \psi^{\prime}(s) \phi_{q}\left(\int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} \bar{M} \phi_{p}(v(\tau)) d \tau\right) d s \\
\leq & (\lambda \bar{K})^{q-1}\|v\| A_{1} \\
< & \|v\|,
\end{align*}
$$

which is a contradiction when $0<\lambda<\bar{K}^{-1}\left(A_{1}^{-1}\right)^{p-1}$. Hence, the IBVP (5) has no positive solutions in this case.
2.5. Hyers-Ulam Stability of Positive Solutions. In the following, we investigate the Hyers-Ullam stability of our proposed IBVP (5). We give a similar definition of the Hyers-Ullam stability as that given in [61].

Definition 25. The integral equation (66) is Hyers-Ulam stable if there exists a positive constant $\delta$ such that for every $\epsilon>0$, if

$$
\begin{equation*}
\left|u(t)-\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s\right| \leq \epsilon \tag{99}
\end{equation*}
$$

then, there exists

$$
\begin{equation*}
u^{*}(t)=\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f\left(\tau, u^{*}(\tau)\right) d \tau\right) d s \tag{100}
\end{equation*}
$$

where

$$
\begin{equation*}
\left|u(t)-u^{*}(t)\right| \leq \delta \epsilon . \tag{101}
\end{equation*}
$$

Proof. Suppose that $u(t)$ is a positive solution of (66) and $u^{*}(t)$ is an approximation satisfying (100). Then, by Lemmas 3 and 8, we obtain

Theorem 26. With the assumption $\left(A_{0}\right)$, the IBVP (5) is Hyers-Ulam stable.

$$
\begin{align*}
\left|u(t)-u^{*}(t)\right|= & \left\lvert\, \int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f(\tau, u(\tau)) d \tau\right) d s\right. \\
& \left.-\int_{0}^{T} \psi^{\prime}(s) \mathscr{H}(t, s) \phi_{q}\left(\frac{\lambda}{\Gamma(\beta)} \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1} f\left(\tau, u^{*}(\tau)\right) d \tau\right) d s \right\rvert\, \\
\leq & \frac{\lambda^{q-1}(p-1) \varrho^{p-2}}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \\
& \times \int_{0}^{T} \psi^{\prime}(s) \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}\left|f(\tau, u(\tau))-f\left(\tau, u^{*}(\tau)\right)\right| d \tau d s  \tag{102}\\
\leq & \frac{\lambda^{q-1}(p-1) \varrho^{p-2}}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \\
& \times \int_{0}^{T} \psi^{\prime}(s) \int_{0}^{s} \psi^{\prime}(\tau)(\psi(s)-\psi(\tau))^{\beta-1}|\mu(t)|\left|u_{n}(\tau)-u(\tau)\right| d \tau d s \\
\leq & \frac{\lambda^{q-1}(p-1) \varrho^{p-2}\|\mu\|}{\Gamma(\alpha) \Gamma(\beta)^{q-1} \varphi(T)}\left[\varphi(t)+\frac{T N}{1-N}(\psi(T)-\psi(0))^{\gamma-1}\right] \frac{(\psi(T)-\psi(0))^{\beta+1}}{\beta^{2}+\beta}\left\|u_{n}-u\right\| .
\end{align*}
$$

Hence, if $\quad \delta=\lambda^{q-1}(p-1) \varrho^{p-2}\|\mu\| / \Gamma(\alpha) \Gamma(\beta)^{q-1}$ $\varphi(T)\left[\varphi(t)+T N / 1-N(\psi(T)-\psi(0))^{\gamma-1}\right](\psi(T)-\psi(0))^{\beta+1} /$ $\left(\beta^{2}+\beta\right)$, then we deduce that the integral (66) is Hyers-Ulam stable. Consequently, the IBVP (5) is Hyers-Ulam stable.

## 3. Numerical Examples

To demonstrate the practical implications of the key findings derived above, we provide the following numerical examples.

Example 1. Consider the following p-Laplacian $\psi$-Caputo fractional differential equations involving integral boundary conditions.

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{1 / 3, t}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{7 / 3, t} u\right)\right)(t)+\lambda\left(t^{3}+2\right)\left(65 u-\frac{8449}{130} \sin u\right)=0, \quad t \in[0,1]  \tag{103}\\
u(0)=\int_{0}^{1} \frac{199}{205} u(s) d s, u(T)=\int_{0}^{1} \frac{199}{205} u(s) d s, u^{\prime \prime}(0)=\int_{0}^{1} \frac{3}{205} u(s) d s \\
\left.{ }^{c} D_{0^{+}}^{7 / 3, \psi} u(t)\right|_{t=0}=0
\end{array}\right.
$$

where $\psi(t)=t, T=1, \gamma=1, f(t, u(t))=\left(t^{3}+2\right)(65 u-$ $8449 / 130 \sin u), \quad \alpha=7 / 3, \quad \beta=1 / 3, \quad p=2, \quad \varepsilon=1 / 6$,
$g_{1}(s)=g_{2}(s)=199 / 205, \quad g_{3}(s)=3 / 205, \quad \varphi(t)=2 t, \quad$ and
$\chi(t)=t(t-1)$. This implies that

$$
\begin{align*}
\Phi(t, s) & =(1-t) g_{1}(s)+t g_{2}(s)+\frac{t(t-1)}{2} g_{3}(s)  \tag{104}\\
& =\frac{199}{205}(1-t)+\frac{199}{205} t+\frac{3}{410}\left(-t+t^{2}\right) .
\end{align*}
$$

Then, through straightforward computations, we can determine that $m=1589 / 1640, \quad M=199 / 205, \quad \rho=$ $5(6-\sqrt[3]{180}) / 36, \quad \sigma=95(6-\sqrt[3]{180}) / 1806, \quad A_{1}^{p-1}=205 /$ $6 \Gamma(11 / 3)=8.51569, \quad A_{2}^{p-1}=12624605(6-\sqrt[3]{180})^{2} /$ $51827526 \Gamma(11 / 3)=0.0075989, \quad f^{0}=3 / 130<0.03=\zeta$, $f_{\infty}=130>\mathcal{\aleph}=129, f_{0}=1 / 65$, and $f^{\infty}=195$. By using Theorem 23, it is clear that the inequality
$1 / A_{2}^{p-1} \aleph<\lambda<1 / A_{1}^{p-1} \zeta$ implies that the IBVP (103) has at least one positive solution for each $\lambda \in(1.02014,3.91434)$. Moreover, since $3 / 130 u<f(t, u)<195 u$, which implies by using Theorem 24 that the IBVP (103) has no positive solution for each $\lambda \in(0,0.000602207)$. Finally, since assumption $\left(A_{0}\right)$ holds for $m=1589 / 1640$ and $M=199 / 205$, i.e., $0<m \leq|\Phi(t, s)| \leq M<1$ then the IBVP (103) is Hyers-Ulam stable.

Example 2. Consider the following p-Laplacian $\psi$-Caputo fractional differential equations involving integral boundary conditions.

$$
\left\{\begin{array}{l}
{ }^{c} D_{0^{+}}^{1 / 2, t^{2}+1}\left(\phi_{p}\left({ }^{c} D_{0^{+}}^{5 / 2, t^{2}+1} u\right)\right)(t)+\lambda\left(t^{2}+1\right)\left(65 u-\frac{8449}{130} \sin u\right)=0, \quad t \in[0,1]  \tag{105}\\
u(0)=\frac{1}{\Gamma(1 / 2)} \int_{0}^{1} \frac{200}{201} 2 s\left(1-s^{2}\right) u(s) d s, u(T)=\frac{1}{\Gamma(1 / 2)} \int_{0}^{1} \frac{200}{201} 2 s\left(1-s^{2}\right) u(s) d s \\
u^{\prime \prime}(0)=\frac{1}{\Gamma(1 / 2)} \int_{0}^{1} \frac{4}{201} 2 s\left(1-s^{2}\right) u(s) d s,\left.{ }^{c} D_{0^{+}}^{5 / 2, t^{2}+1} u(t)\right|_{t=0}=0
\end{array}\right.
$$

where $\quad \psi(t)=t^{2}+1, \quad T=1, \quad \gamma=1 / 2, \quad f(t, u(t))=$ $\left(t^{3}+2\right)(65 u-8449 / 130 \sin u), \quad \alpha=5 / 2, \quad \beta=1 / 2, \quad p=2$, $\varepsilon=1 / 4, g_{1}(s)=g_{2}(s)=200 / 201, g_{3}(s)=4 / 201$,

$$
\begin{align*}
& \varphi(t)=(\psi(t)-\psi(0))\left(2\left(\psi^{\prime}(0)\right)^{2}-(\psi(t)-\psi(0)) \psi^{\prime \prime}(0)\right)=2 t^{4}  \tag{106}\\
& \chi(t)=(\psi(1)-\psi(0))(\psi(t)-\psi(0))(\psi(t)-\psi(1))=t^{2}\left(t^{2}-1\right)
\end{align*}
$$

This implies that

$$
\begin{align*}
\Phi(t, s) & =\frac{1}{\Gamma(\gamma)}\left[\left(1-\frac{\varphi(t)}{\varphi(T)}\right) g_{1}(s)+\frac{\varphi(t)}{\varphi(T)} g_{2}(s)+\frac{\chi(t)}{\varphi(T)} g_{3}(s)\right] \\
& =\frac{1}{\Gamma(1 / 2)}\left[\left(1-t^{4}\right) \frac{200}{201}+t^{4} \frac{200}{201}+\frac{t^{2}\left(t^{2}-1\right)}{2} \frac{200}{201}\right] \\
& =\frac{1}{\sqrt{\pi}}\left(\frac{200}{201} t^{4}+\frac{100}{201} t^{2}\left(t^{2}-1\right)+\frac{200}{201}\left(1-t^{4}\right)\right) . \tag{107}
\end{align*}
$$

Then, through straightforward computations, we can determine that $m=175 /(201 \sqrt{\pi})=0.49121, M=200 /(201$ $\sqrt{\pi})=0.561383, \rho=(3 / 8)(\sqrt{3}-2)=0.100481, \sigma=175(3 /$ $4-3 \sqrt{3} / 8)(1-200 /(201 \sqrt{\pi}) / 402(1-175 /(201 \sqrt{\pi})) \quad \sqrt{\pi}=$ $0.0212749, \quad A_{1}^{p-1}=1 /(6(1-200 /(201 \sqrt{\pi})))=0.379982$, $A_{2}^{p-1}=30625(3 / 4-3 \sqrt{3} / 8)^{2} \quad(1-200 /(201 \sqrt{\pi})) / 969624(1-$ $175 /(201 \sqrt{\pi}))^{2} \quad \pi=0.000171988, \quad f^{0}=1 / 65<1 / 64=\zeta$, $f_{\infty}=65>\mathcal{\aleph}=64, f_{0}=(1 / 130)$, and $f^{\infty}=130$. By using Theorem 23, it is clear that the inequality $1 /\left(A_{2}^{(p-1)} \aleph\right)<\lambda<1 /\left(A_{1}^{(p-1)} \zeta\right)$ implies that the IBVP (105) has
at least one positive solution for each $\lambda \in(90.8492,168.429)$. Moreover, since $(1 / 130) u<f(t, u)<130 u$, which implies by using Theorem 24 that the IBVP (105) has no positive solution for each $\lambda \in(0,(3 / 98)(1-200 /(201 \sqrt{\pi}))=$ ( $0,0.0134271$ ). Finally, since assumption $\left(A_{0}\right)$ holds for $m=$ $175 /(201 \sqrt{\pi})$ and $M=200 /(201 \sqrt{\pi})$, i.e., $0<m \leq|\Phi(t, s)| \leq$ $M<1$ then the IBVP (105) is Hyers-Ulam stable.

## 4. Conclusion

In conclusion, our study focuses on investigating the stability and existence of positive solutions for a specific type of p-Laplacian $\psi$-Caputo fractional differential equations with fractional integral boundary conditions. Through the use of Green's function properties and the application of Guo-Krasnovelsky's fixed point theorem on cones, we have established novel existence results, ensuring the presence of at least one positive solution. Our analysis encompasses a range of parameter values, providing a comprehensive understanding of the problem. We have laid a solid foundation for our investigation by utilizing fractional integrals, differential operators, and fundamental lemmas. Our study
contributes to the existing knowledge in this field by exploring both the existence and nonexistence of positive solutions. We have also examined the Hyers-Ulam stability of solutions by introducing a defined notion and building upon previous research. To validate the effectiveness and accuracy of our approach, we have presented numerical examples. Overall, our study significantly advances the understanding of stability and existence of positive solutions for the considered class of fractional differential equations. Future research can focus on conducting more extensive numerical studies to gain further insights into the system's behavior and properties under different parameter regimes. In addition, exploring different choices of $\psi(t)$ and other parameter values can provide a deeper understanding of the existence and stability of positive solutions. In summary, our study makes substantial contributions to the understanding of eigenvalue problems with the $\psi$-Caputo integrodifferential operator, and its insights can be extended to other investigations in this field. By continuing to conduct further numerical studies and explore different parameter choices, researchers can advance our knowledge and uncover new aspects of this intriguing problem.

## Data Availability

The required data used to support the findings of this study are included within the article.

## Conflicts of Interest

The authors declare that they have no conflicts of interest.

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