


Research Article

Design and Prediction of a Novel Two-Dimensional Carbon Nanostructure with In-Plane Negative Poisson's Ratio

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The intrinsic negative Poisson's ratio effect in 2-dimensional nanomaterials have attracted a lot of research interests due to its superior mechanical properties, and new mechanisms have emerged in the nanoscale. In this paper, we designed a novel graphyne-like two-dimensional carbon nanostructure with a "butterfly" shape (GL-2D-1) and its configuration isomer with a "herring-bone" form (GL-2D-2) by means of density functional theoretical calculation and predicted their in-plane negative Poisson's ratio effect and other mechanical properties. Both GL-2D-1 and GL-2D-2 present a significant negative Poisson's ratio effect under different specific strains conditions. By contrast, GL-2D-2 presents a much stronger negative Poisson's ratio effect and mechanical stability than does GL-2D-1. It is hoped that this work could be a useful structural design strategy for the development of the 2D carbon nanostructure with the intrinsic negative Poisson's ratio.

1. Introduction

It is known that carbon can exist in several different structures. Besides the traditional allotropes [1], some other allotropes of carbon with nanostructures including fullerenes, carbon nanotubes, graphene, graphyne, T-carbon, and graphane were found one after another [2–12]. However, researchers in physical chemistry never stop to explore newer allotropes from zero to three dimensions [13, 14]. As Hobza pointed out [15], "the way of covalent interactions determine [*sic*] the primary structure of a molecule." The diversity of the carbon allotrope derives from its three kinds of hybridization ability, namely, sp , sp^2 , and sp^3 hybridizations. Two-dimensional (2D) carbon allotropes have received much attention in the last ten years due to their potential application in nanoelectronics, optoelectronics, biomaterials [16–24], energy storage [25], and gas separation [26]. But to date, the 2D carbon and analogous materials are not quite abundant. Therefore, it deserves exploration in theory and in experiment whether

there exists other 2D carbon allotropes which may be stable as a monolayer and exhibit novel properties, especially novel mechanical properties. If confirmative, it is likely to make important contributions to the multifunctional carbon nanomaterials.

Since the first observation of the auxetic phenomenon in polyurethane foam by Lakes [27] three decades ago, great efforts have been made to find or design new auxetic materials which are also known as negative Poisson's ratio materials. Most materials will be contracted in the lateral direction when they are stretched, generally resulting in a positive value for Poisson's ratio. Abnormally, negative Poisson's ratio materials expand laterally when stretched and contract laterally when compressed. This enables them to offer enormous potential important applications [28], such as biomedicine [29] and protective equipments [30].

In recent years, negative Poisson's ratio effect has also been reported in several monolayer 2D materials. It is known that the out-of-plane and in-plane negative Poisson's ratios

have been discovered in different black phosphorus [31–34]. The in-plane negative Poisson’s ratio was also theoretically predicted in borophene [35], prismanes [36], transition metal dichalcogenides [37], and some not-yet-synthesized materials (e.g., penta-graphene) [38]. Recently, Grima’s group [39] reported “how graphene can be modified to mimic the behavior of a highly and densely wrinkled paper model to the extent that it can exhibit negative Poisson’s [*sic*] ratio behavior under ambient conditions” for the first time. Significantly, Jiang and Park [40] theoretically explored the property of negative Poisson’s ratio in single-layer graphene ribbons, which results from the compressive edge stress-induced warping of the edges. They claimed that “the effect is robust, as the negative Poisson’s [*sic*] ratio is observed for graphene ribbons with widths smaller than about 10 nm, and for tensile strains smaller than about 0.5% with negative Poisson’s [*sic*] ratio values reaching as large as -1.51 .” Lately, they performed molecular dynamics (MD) simulations to investigate Poisson’s ratio of graphene oxide and found that “the Poisson’s ratio can be effectively tuned by increasing the degree of oxidation of graphene oxide.” More specifically, “the Poisson’s ratio decreases linearly from positive to negative with increasing of oxidation, and reaching a value of -0.567 for fully oxidized graphene” [41]. In addition, negative Poisson’s ratio in rippled graphene due to the dewrinkling effect has been simulated by MD calculation [42].

Negative Poisson’s ratio is theoretically possible within continuum elasticity for carbon nanostructures; it is often observed in engineered materials [43], as they are not intrinsic for many materials at the molecular level. Excitedly, also reported by Jiang’s group [44], the single-layer graphene may exhibit an intrinsic negative Poisson’s ratio, which is robust and independent of its size and temperature. They pointed out that “the negative Poisson’s [*sic*] ratio arises from the interplay between two intrinsic deformation pathways (one with positive Poisson’s ratio, the other with negative Poisson’s [*sic*] ratio).” It is known that one of the major causes for a negative Poisson’s ratio is the microstructure with a reentrant honeycomb geometric unit of the material at the molecular level [45, 46]. Moreover, using the two-dimensional classical example of reentrant honeycomb [47], “it is clear that the auxetic effect would initially be large and would gradually decay as the dent in the hexagon flattens out.” Accordingly, a pictorial representation of the system studied here and the auxetic structure designed to simulate are shown in Figure 1. This kind of topology geometry is expected to mimic the reentrant hexagonal honeycombs, which can be successfully created by the appropriate combination of sp^2 (phenyl) and sp (alkynyl) carbon atoms like in graphyne. This will be further illustrated thereafter. In this paper, we design a novel graphyne-like two-dimensional carbon nanostructure with a reentrant hexagonal structural unit optimized by means of density functional theory and then try to predict its in-plane negative Poisson’s ratio and other mechanical properties. We hope our work could offer a new strategy for the theoretical development of the 2D carbon nanostructure with the intrinsic negative Poisson’s ratio.

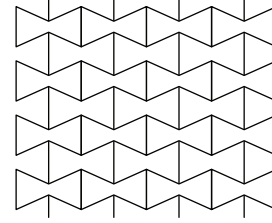


FIGURE 1: Pictorial representation of the reentrant hexagonal honeycombs.

2. Computational Method

All calculations, including the self-consistent energy calculations and structure optimization, were performed using the density functional theory of Perdew–Burke–Ernzerhof (PBE) exchange correlation with a generalized gradient approximation (GGA) functional [48] along with the projector-augmented wave (PAW) potentials [49, 50] as implemented in the Vienna Ab Initio Simulation Package (VASP) [51, 52]. An energy cutoff of 500 eV was used. The monolayer structure is modelled in an orthorhombic cell that contains one formula unit (Figure 2) and a 20 Å vacuum space inserted in the out-of-plane direction. A $3 \times 5 \times 1$ k-point grid was used to sample the Brillouin zone during structure relaxation. All atoms were fully relaxed until their atomic forces were $<0.005 \text{ eV \AA}^{-1}$. The effects of spin-orbit coupling on the structural deformation are considered to be minor and hence not included in our study. The vibration analyses for the relaxed structures were performed to confirm that these structures were thermodynamically stable.

A classical trajectory calculation using the atom-centered density matrix propagation ab initio molecular dynamics (MD) model [53–55], which can provide an equivalent functionality to Born-Oppenheimer molecular dynamics [56, 57] at considerably reduced computational cost, was performed for the preoptimized basic structural unit (neutral molecule) with different sizes so as to further infer and investigate the dynamic stabilities of the designed 2D structures.

Poisson’s ratio is calculated from the engineering strain (ϵ), which is defined as the change in length ΔL per unit of the original length L , that is, $\epsilon = \Delta L/L$. The applied uniaxial strain is realized in our calculations by fixing the lattice parameter to a value different from its equilibrium value during structural relaxation. The resulting strain in the transverse direction is extracted from the fully relaxed structure subjected to an applied strain.

3. Results and Discussions

The chemical structure of the designed graphyne-like 2D carbon nanomaterial is shown in Figure 3. It has a reentrant hexagon structural unit in which m , n , and l correspond to the number of alkynyl at three different positions between phenyl units. It is found that when $m = 1$, $n = 1$, and $l = 2$ (denoted as GL-2D-1, similar to Figure 2), the corresponding 2D structure is stable in thermodynamics; when $m = 1$, $n = 2$ and $l = 2$, the corresponding 2D structure is also stable in

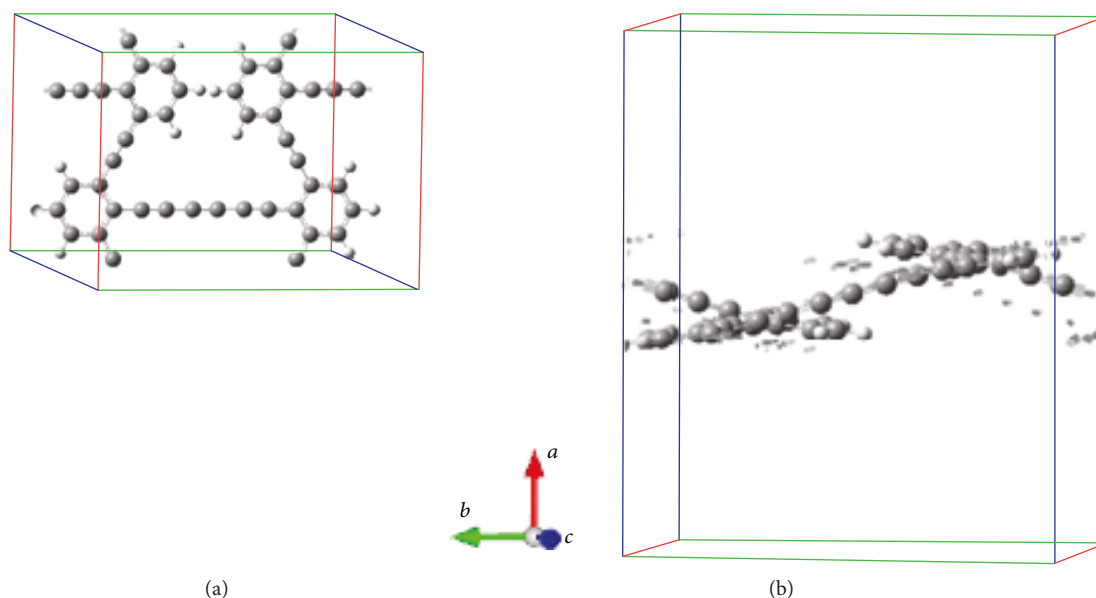


FIGURE 2: The designed periodical 2D carbon nanostructure (a) top view and (b) side view (carbon atoms in grey and hydrogen atoms in white).

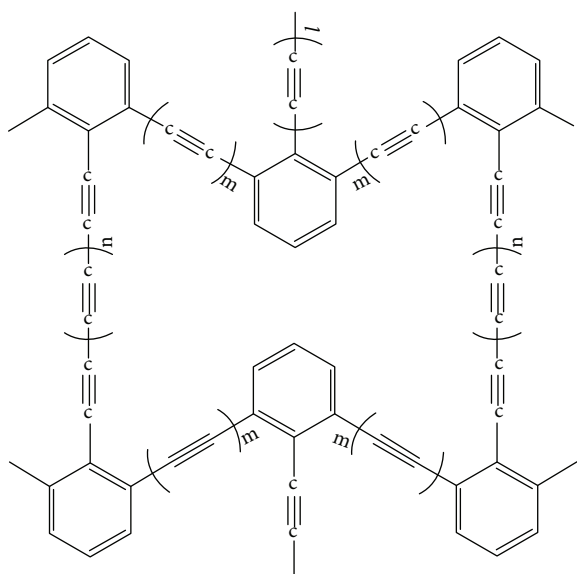


FIGURE 3: The chemical structure of the present designed graphdiyne-like 2D carbon nanomaterial ($m = 1$, $n = 1$, and $l = 2$).

thermodynamics, while when $m = 1$, $n = 2$ and $l = 3$, the thermodynamic stable 2D structure cannot be obtained theoretically. Therefore, it can be seen that the length of the alkyne-linking chain is very important for the design of a reasonable 2D structure. Namely, it is not that the alkyne-linking chain with reasonable geometric length can get the corresponding stable 2D structure. Presently, the following investigations and discussions about mechanic properties are based on the structure of GL-2D-1 and its configurational isomer.

Figure 4 plots the configuration of the supercell of GL-2D-1 at three orthogonal directions of a , b , c and standard

orientation together with the optimized lattice parameters of the unit cell. GL-2D-1 is a typical anisotropic 2D material, and it would have obvious differences in mechanical properties at a and b directions. Viewed from a direction, it is “butterfly”-shaped (Figure 4(a)).

In order to investigate the thermodynamic stabilities of the structures of GL-2D-1 and GL-2D-2, frequency analysis was performed at the same theoretical level to the structure optimizations. The results showed that there is no imaginary frequency for the two periodical 2D structures, indicating that these structures were thermodynamically stable on the potential energy surfaces. The data of frequency of the two 2D structures are listed in Table S1 and Table S2 of the Supporting Information.

Molecular dynamic stability is also as important as thermodynamic stability. Herein, *ab initio* MD was conducted for the preoptimized minimum basic structural unit (Figure S1) and hexamer of the basic structural unit (Figure S2) of GL-2D at room temperature, respectively, to evaluate the dynamic stabilities of the two designed 2D structures. Figure 5 shows the molecular dynamic simulations of the preoptimized minimum basic structural units of GL-2D-1 and GL-2D-2. It is clear that the largest changes in potential energies for the two structural units are only 5.78 and 4.98 $\text{kJ}\cdot\text{mol}^{-1}$ within 5000 fs, respectively, which are small enough to indicate the high molecular dynamic stabilities. Moreover, GL-2D-2 shows a slight higher molecular dynamic stability than does GL-2D-1 at 298 K. For completeness, the molecular dynamic simulations of the hexamer of the basic structural unit of GL-2D-1 and GL-2D-2 are shown in Figure S3 of the Supporting Information, and the largest changes in potential energies for the two hexamers are both smaller than 4.50 $\text{kJ}\cdot\text{mol}^{-1}$ (4.46 and 4.27 $\text{kJ}\cdot\text{mol}^{-1}$) within 1000 fs, respectively, also indicating the excellent dynamic stabilities of the two 2D structures.

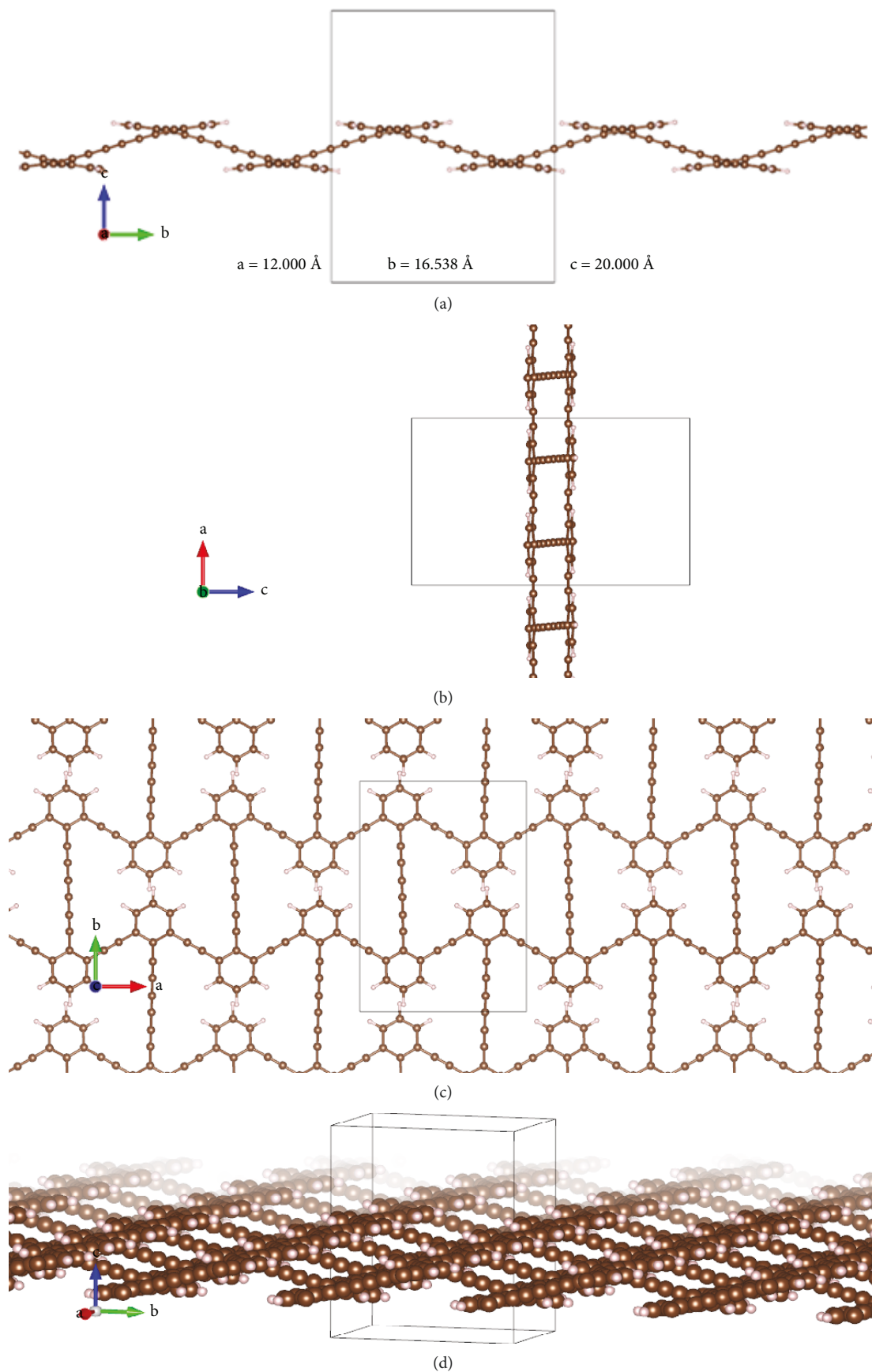


FIGURE 4: The configuration of the supercell of GL-2D-1 at three orthogonal directions of (a), (b), (c) and standard orientation (d).

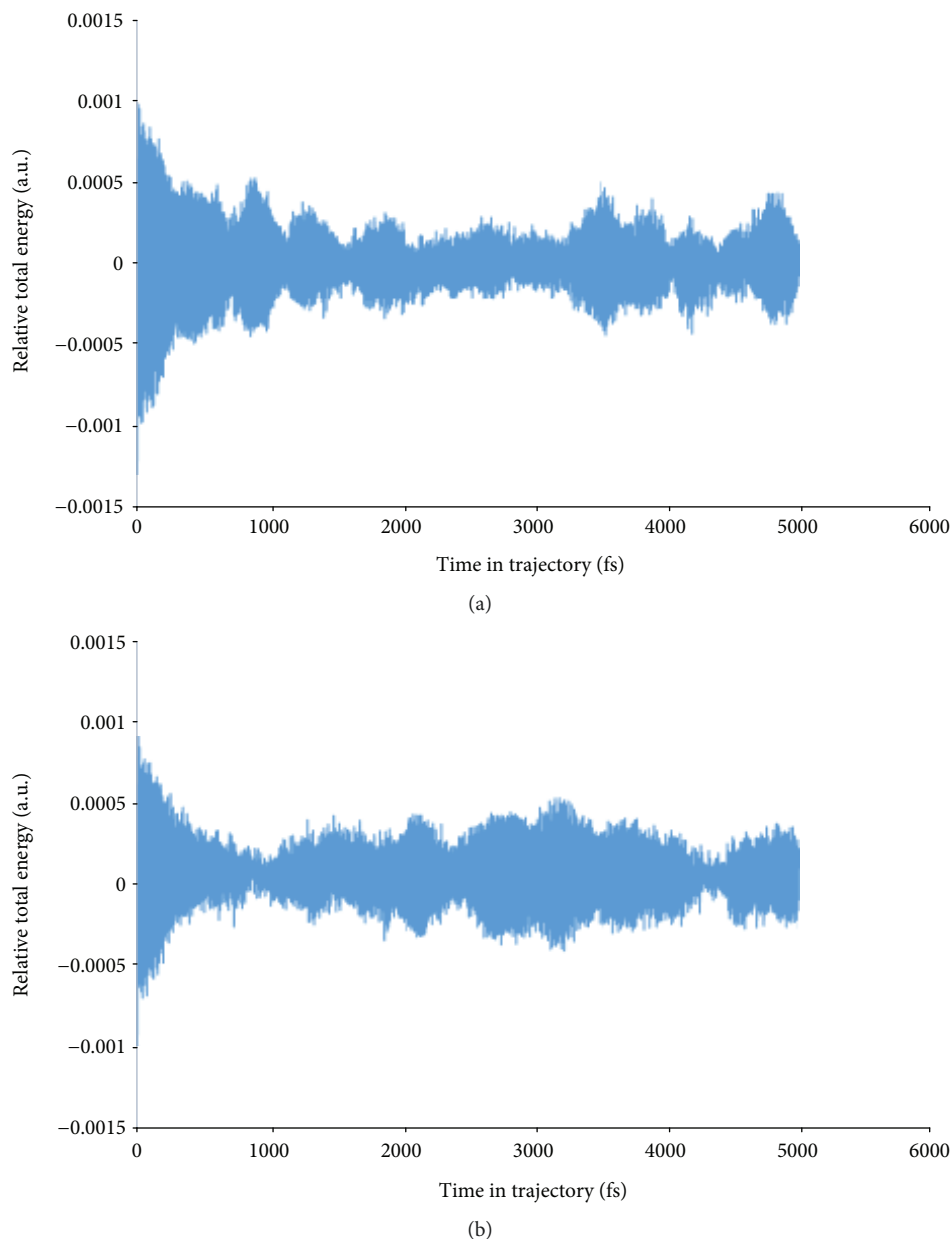


FIGURE 5: Molecular dynamic simulations of the preoptimized minimum basic structural units of (a) GL-2D-1 and (b) GL-2D-2.

Poisson's ratio is the most important mechanical property we concerned. Figures 6(a) and 6(b) are Poisson's ratio of GL-2D-1 *vs* strain in *a* and *b* directions, respectively. It is clear that there is no negative Poisson's ratio effect when compressing along the *a* direction. However, when stretching along the *a* direction, the negative Poisson's ratio effect of GL-2D-1 is significant (Figure 6(a)). At the initial stage of stretching, the material has a very large negative Poisson's ratio; with the increase in the tensile strain, the negative Poisson's ratio effect is weakened, and the Poisson's ratio ν converges to a certain value. As Figure 6(b) shows, the negative Poisson's ratio effect in the *b* direction is contrary to that of the *a* direction. There is no negative Poisson's ratio effect when stretching along the *b* direction, while when compressing along the *b* direction, the negative Poisson's ratio effect

of GL-2D-1 is obvious, although it is weaker than that in the case of stretching along the *a* direction. Similarly, with the increase in the compression strain in the *b* direction, the negative Poisson ratio effect is weakened, and then the Poisson's ratio ν converges to a certain value. In brief, GL-2D-1 can present the significant negative Poisson ratio effect both in *a* and *b* directions under different specific strains.

Additionally, other mechanical properties were further investigated. The stiffness matrix and its six eigenvalues are given in Table S3 and Table S4 of the Supporting Information. Seen from the stiffness matrix of GL-2D-1, it is not a positive definite matrix, manifesting that the mechanical stability of the designed GL-2D-1 is not as excellent as that of graphdiyne. However, it is noted that among the six eigenvalues ($\lambda_1 \sim \lambda_6$), $\lambda_1 \sim \lambda_4$ are negative

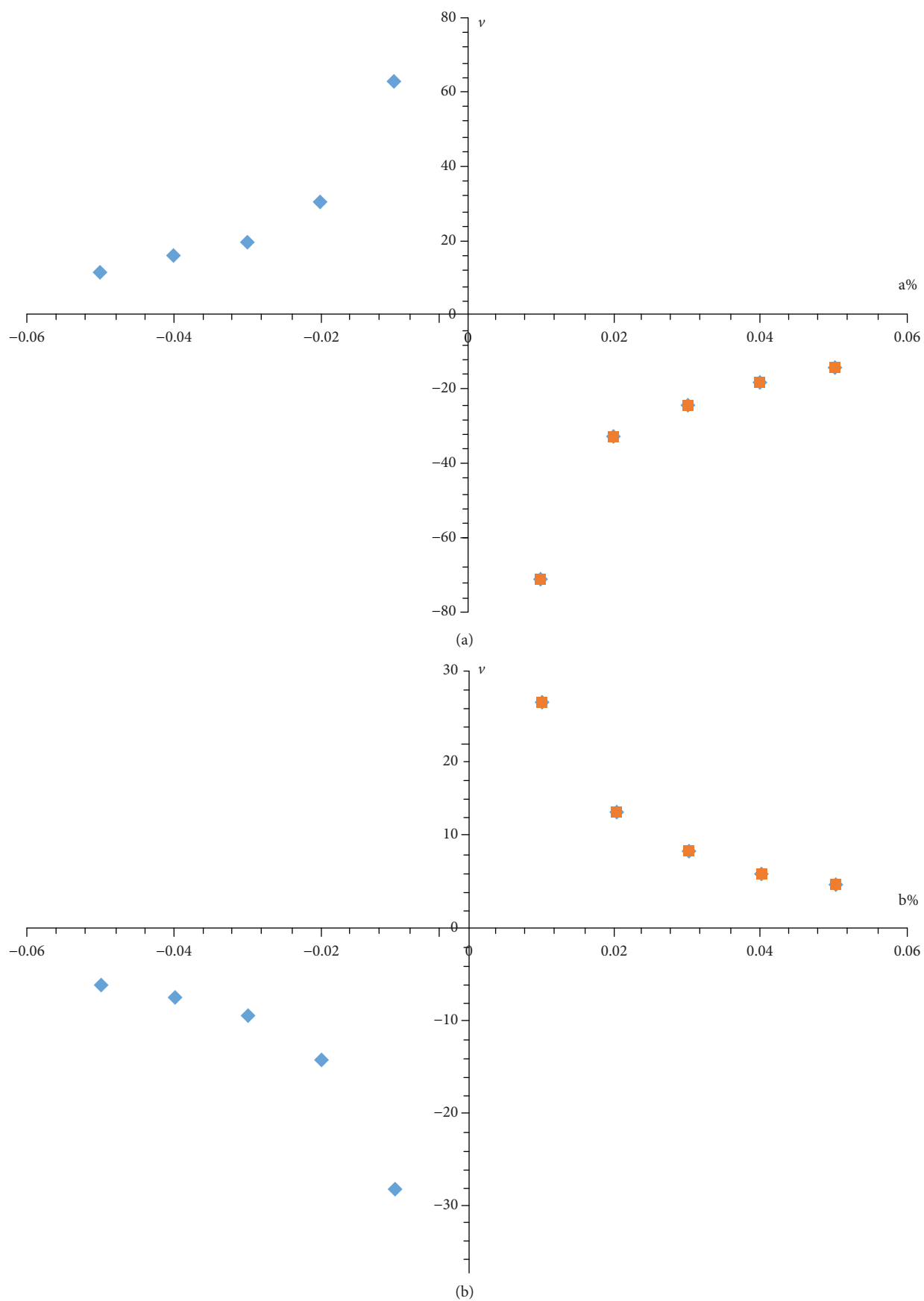


FIGURE 6: Poisson's ratio of GL-2D-1 vs strain in the a direction (a) and b direction (b).

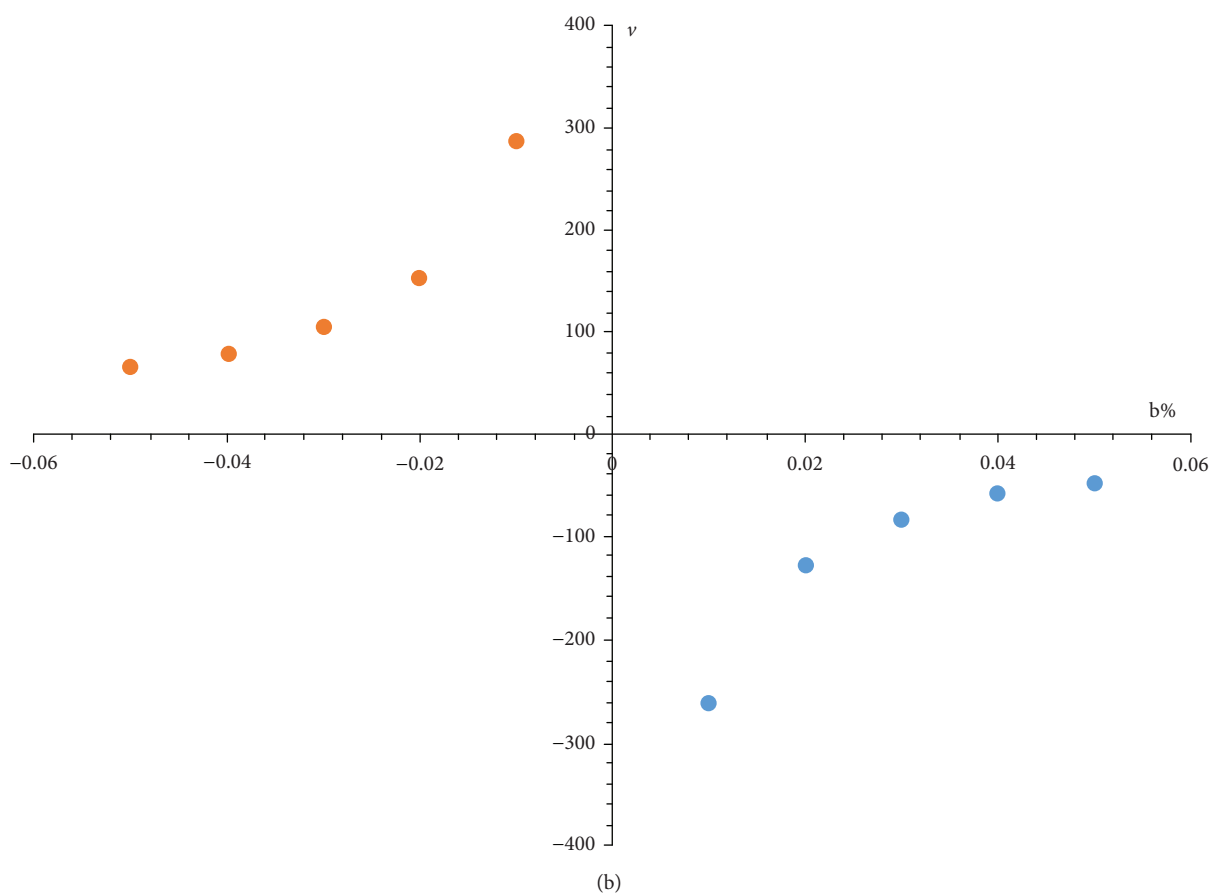
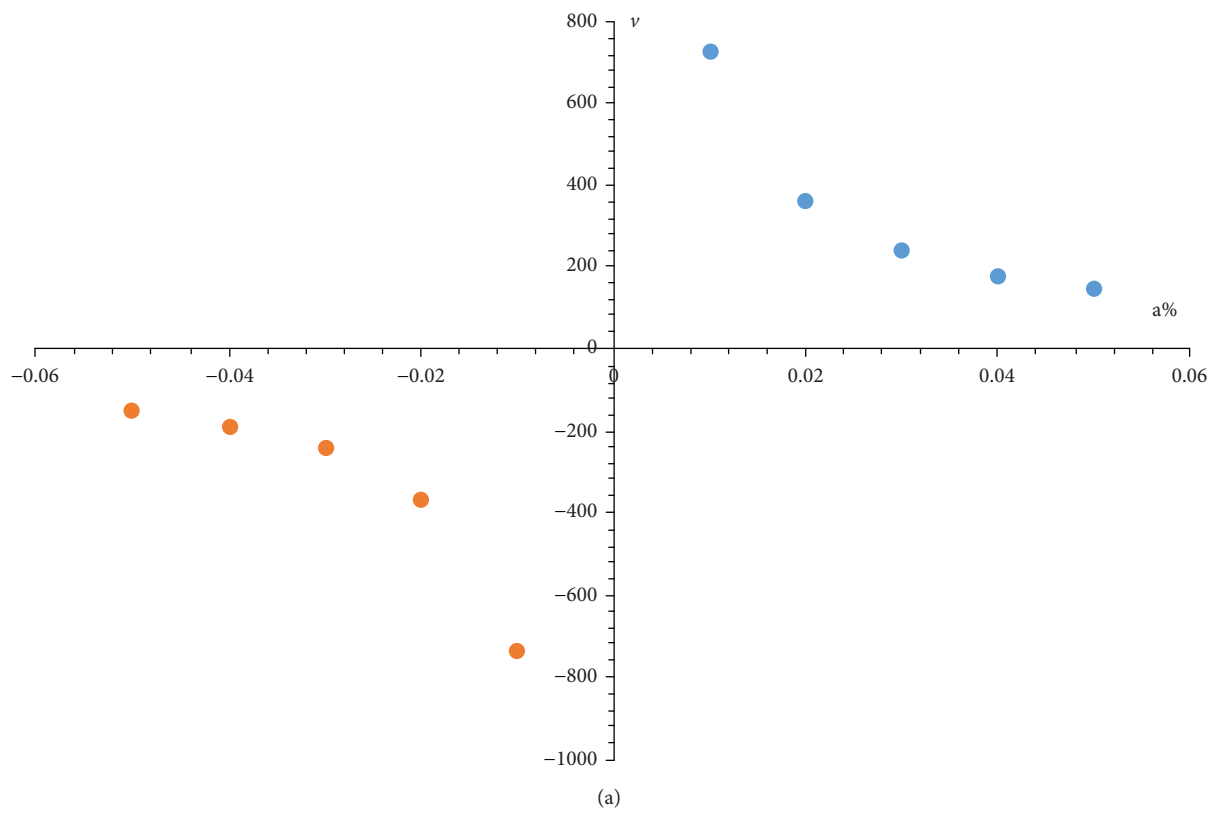


FIGURE 7: Poisson's ratio of GL-2D-2 vs strain in the *a* direction (a) and *b* direction (b).

with a smaller absolute value, and λ_5 and λ_6 are positive, indicating that GL-2D-1 has a certain mechanical stability. Moreover, the modulus of GL-2D-1 under different models including Voigt, Reuss, and Hill are also presented in Table S5 of the Supporting Information.

Interestingly, a configuration isomer of GL-2D-1 is found, denoted as GL-2D-2, which is also a thermodynamically stable 2D structure. Figure S4 of the Supporting Information shows the configuration of the supercell of GL-2D-2 at three orthogonal directions of a, b, c and standard orientation. Viewed from a direction, it is a “herring-bone” form (Figure S4(a)). Compared to GL-2D-1, the negative Poisson’s ratio effect of GL-2D-2 is exactly opposite. Figure 7 shows the Poisson’s ratio of GL-2D-2 *vs* strain in the a direction and b direction. It can be seen that GL-2D-2 presents the negative Poisson’s ratio effect when compressing along the a direction or stretching along the b direction and the normal Poisson’s ratio effect when stretching along the a direction or compressing along the b direction. Either in the initial stage of stretching or in that of compressing, the absolute values of Poisson’s ratio of GL-2D-2 are very large, but with the increasing of strains, they gradually decrease and converge to a certain value. Poisson’s ratio has a “hyperbolic” rule with the whole trend of strain variation.

It is particularly noteworthy that although the negative Poisson’s ratio effect of GL-2D-1 under certain strain conditions is significant, those of GL-2D-2 under the compression along the a direction or stretching along the b direction are distinctly larger than those of GL-2D-1 under stretching along the a direction or compression along the b direction. For example, Poisson’s ratio of GL-2D-2 is -730.37 when it is compressed by 0.01% along the a direction, while at the same stretching value along the same direction, that of GL-2D-1 is only -63.34. Thereby, it can be found that the small differences in configuration may bring very large differences in the negative Poisson’s ratio effect of the present 2D structures. This is of great significance for designing and recognizing 2D carbon nanostructures with negative Poisson’s ratio effect.

For comparison, the stiffness matrix and its six eigenvalues and the modulus of GL-2D-2 are also given in Table S6, Table S7, and Table S8 of the Supporting Information. Similar to that of GL-2D-1, the stiffness matrix of GL-2D-2 is not a positive definite matrix, meaning that the mechanical stability is not very ideal. However, we noted that among the six eigenvalues of the stiffness matrix, $\lambda_3 \sim \lambda_6$ are all positive values, and λ_1 and λ_2 are negative with very small absolute values, indicating that the mechanical stability of GL-2D-2 increased obviously compared to that of GL-2D-1.

4. Conclusion

In summary, a novel thermodynamic stable graphyne-like two-dimensional carbon nanostructure (GL-2D-1) with in-plane negative Poisson’s ratio is designed and predicted by means of density functional theory and PBE exchange correlation with a GGA functional along with the PAW potentials. Simultaneously, an unforeseen stable conformational

isomer of GL-2D-1, GL-2D-2, is also found theoretically. Excitedly, both GL-2D-1 and GL-2D-2 all present a significant negative Poisson’s ratio effect under different specific strain conditions. Additionally, it is found that the small differences in conformation may bring very large differences in the negative Poisson’s ratio effect of the 2D materials with the same configuration. We hope that our work could offer a new strategy for the theoretical development of the 2D carbon nanostructure with the intrinsic negative Poisson’s ratio.

Data Availability

The data used to support the findings of this study are included within the article and Supplementary materials.

Conflicts of Interest

The authors declare that there is no conflict of interest regarding the publication of this paper.

Acknowledgments

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Supplementary Materials

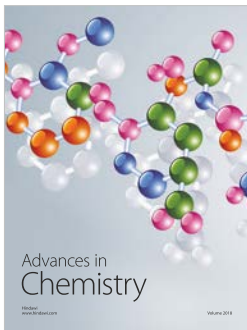
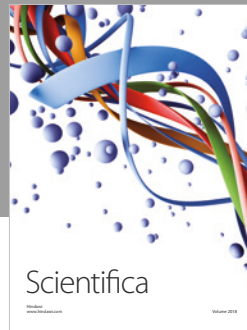
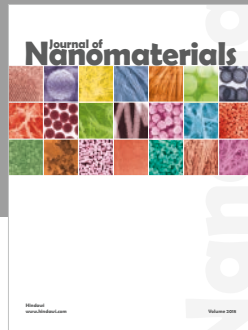
Figure S1: the preoptimized minimum basic structural unit. Figure S2: the preoptimized hexamer of the basic structural unit. Table S1: the frequency analysis result for the periodical structure of GL-2D-1. Table S2: the frequency analysis result for the periodical structure of GL-2D-2. Figure S3: the molecular dynamic simulations of the hexamer of the basic structural unit of GL-2D-1 and GL-2D-2. Table S3: stiffness matrix (coefficients in GPa) of GL-2D-1. Table S4: eigenvalue of the stiffness matrix of GL-2D-1. Table S5: modulus of GL-2D-1 under different models. Table S6: stiffness matrix (coefficients in GPa) of GL-2D-2. Table S7: eigenvalue of stiffness matrix of GL-2D-2. Table S8: modulus of GL-2D-2 under a different model. Figure S4: the configuration of the supercell of GL-2D-2 at three orthogonal directions of a, b, c and standard orientation. (*Supplementary Materials*)

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