

## Research Article

# Influence of Cu-Doping Concentration on the Structural and Optical Properties of SnO<sub>2</sub> Nanoparticles by Coprecipitation Route

Abebe G. Habte <sup>1</sup>, Fekadu Gashaw Hone <sup>2</sup>, and F. B. Dejene <sup>3</sup>

<sup>1</sup>Physics Department, Hawassa University, P.O. Box 5, Hawassa, Ethiopia

<sup>2</sup>Physics Department, Addis Ababa University, P.O. Box 1176, Addis Ababa, Ethiopia

<sup>3</sup>Department of Chemical and Physical Sciences, Walter Sisulu University, Private Bag XI, Code 5117, Mthatha, South Africa

Correspondence should be addressed to Fekadu Gashaw Hone; fekeye@gmail.com

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Tin dioxide (SnO<sub>2</sub>) nanoparticles doped with varying concentrations of copper were synthesized and characterized using various techniques. The X-ray diffraction analysis revealed that all doped and undoped SnO<sub>2</sub> samples had a rutile-type tetragonal structure. The average crystalline size of the doped samples estimated using Scherrer's formula and the Williamson–Hall plot decreased as dopant concentration increased. Images from scanning electron microscopy revealed spherical grains in the samples. The transmission electron microscope was used to examine the particle nature, and nearly spherical particles were discovered. The energy-dispersive X-ray spectroscopy analyses confirmed that the synthesized nanoparticles were nearly stoichiometrically composed of the expected elements copper, oxygen, and tin. The bandgap energy of doped and undoped SnO<sub>2</sub> nanoparticles was determined using UV–visible diffuse reflectance spectra, and it was found to decrease as Cu<sup>2+</sup> ion concentration increased. The photoluminescence study at the excitation wavelength of 300 nm revealed defect-oriented emissions between 350 and 500 nm. All the obtained results showed that the physical properties of SnO<sub>2</sub> can be easily engineered through Cu doping for various optoelectronic applications using a low-cost coprecipitation method.

## 1. Introduction

Tin dioxide (SnO<sub>2</sub>) is an important n-type, wide-bandgap semiconductor with excellent optical, electrochemical, and electrical properties [1]. SnO<sub>2</sub> with a tetragonal structure has an energy bandgap of about 3.6 eV at room temperature [2]. It has a wide range of applications, including lithium-ion storage batteries [3, 4], transparent conducting electrodes [5, 6], gas sensors [7, 8], optoelectronic devices [9, 10], photovoltaic cells, photo detectors, and so on. To obtain the required properties for specific applications, the morphology and size of SnO<sub>2</sub> nanoparticles must be controlled during synthesis [11–13]. Doping SnO<sub>2</sub> with specific metal ions is an efficient method for changing the optical, electrical, and magnetic properties of SnO<sub>2</sub> nanostructures [7]. Several doping elements, such as Cu, Zn, Ni, La, and Co, have been used to tune the properties of SnO<sub>2</sub> nanomaterials [14–17]. The practical performance of SnO<sub>2</sub> is determined

by crystallinity, surface morphology, and crystal size, all of which are determined by the synthesis method. To prepare SnO<sub>2</sub> nanostructures, various synthesis techniques such as solid-state reaction [6], sol–gel [18–20], coprecipitation [10], hydrothermal [21], and microwave assisted [22] have been developed. Among these methods, coprecipitation produces stoichiometric and homogeneous nanoparticles at a low cost. Nanocrystalline SnO<sub>2</sub> properties are highly dependent on preparation methods and dopant materials. Dopant elements can change the microstructure and defect chemistry of SnO<sub>2</sub>, resulting in changes in physical and chemical properties [23]. Many studies have been published on the influence of Cu doping on the physical properties of SnO<sub>2</sub> nanomaterials. Mishra et al. [24] studied the role of Cu doping on the formaldehyde sensing of SnO<sub>2</sub> nanoparticles using stannic tetrachloride hydrated (SnCl<sub>4</sub>·5H<sub>2</sub>O). Thermal evaporation method was employed to grow Cu-doped SnO<sub>2</sub> nanowires for H<sub>2</sub>S gas sensors by Kumar et al. [25]. Lavanya et al. [26]

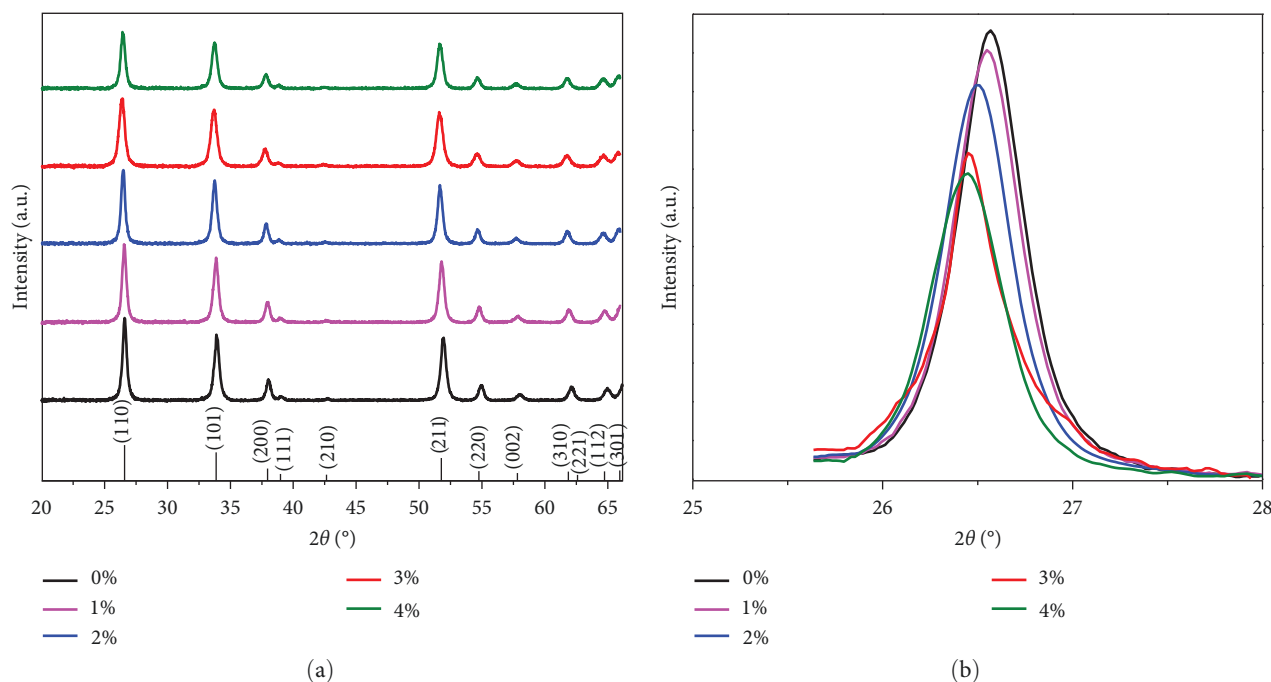


FIGURE 1: (a) The XRD pattern of pure and Cu-doped  $\text{SnO}_2$  nanoparticles; (b) magnified view of the shifting of (110) plane for pure and Cu-doped  $\text{SnO}_2$ .

reported the synthesis of Cu-doped  $\text{SnO}_2$  nanoparticles by microwave irradiation method for the precise determination of folic acid. Chetri et al. [27] investigated the structural and optical properties of Cu-doped  $\text{SnO}_2$  nanostructures using experiments and correlated the observations with density functional theory results. Thermal evaporation and condensation method was employed to synthesize Cu-doped  $\text{SnO}_2$  nanobelts by Li et al. [28] and the temperature dependence photoluminescence (PL) properties of the nanobelts and samples annealed in oxygen were studied. The main challenge in the addition of dopants is the synthesis of nanoparticles without secondary phases. In light of this, we chose cost-effective coprecipitation method, which allows the chemicals to be mixed at the atomic level, avoiding the formation of undetectable impurity phases. From the existing report, copper metal ions would be a perfect material to be used as a dopant. In this study, we reported the synthesis of pure and Cu-doped  $\text{SnO}_2$  via a simple co-precipitation method. The effects of Cu-doping on essential structural, optical, and morphological properties have been studied systematically. The obtained results revealed that Cu doping has a significant influence on the physical properties of the host nanomaterial.

## 2. Experiment

**2.1. Sample Preparation.** The synthesis of  $\text{SnO}_2$  nanoparticles was started by the dissolving of 0.132 mol of tin dichloride dihydrate ( $\text{SnCl}_2 \cdot 2\text{H}_2\text{O}$ ) in deionized water and allowed it to stirrer for 5 min. Copper (II) acetate monohydrate  $\text{Cu}(\text{CH}_3\text{COO})_2 \cdot \text{H}_2\text{O}$  was taken as the source of dopant. The dopant solution of the required amount (1%, 2%, 3%, and 4%) was slowly added drop by drop to the previous precursor solution with continuing magnetic stirring. The pH of the

solution was kept around 9.8 by adding aqueous ammonia ( $\text{NH}_3$ , 25%) solution drop wise. It was then heated at  $50^\circ\text{C}$  for 2 hr with vigorous stirring. The final solution was washed repeatedly with deionized water and ethanol to get chlorine ion-free  $\text{SnO}_2$  nanoparticles. The resulting product was kept for 24 hr at room temperature and then dried for 1 hr at  $70^\circ\text{C}$ . Finally, smaller and homogeneous powder particles were obtained. The prepared samples were annealed at  $600^\circ\text{C}$  for 2 hr to get crystallized nanoparticles.

**2.2. Characterization.** The structures of the prepared Cu-doped and pure  $\text{SnO}_2$  nanoparticles were measured by using a powder X-ray diffraction (XRD), model Bruker D8 advance, Germany equipped with  $\text{CuK}\alpha$  irradiation ( $\lambda = 1.5406 \text{ \AA}$ ) in the range of  $2\theta$  from  $20^\circ$  to  $66^\circ$ . The UV-Vis spectra of the samples were recorded in diffuse reflectance mode using a Perkin Elmer Lambda 950 UV-Vis spectrometer in the range of 250–800 nm. The room temperature PL measurements were carried out by an F-7000 FL spectrophotometer. The surface morphology and elemental composition of the samples were studied using a high-resolution JEOL JSM-7800F scanning electron microscope equipped with energy-dispersive X-ray (EDX) analysis. Transmission electron microscope (TEM) measurements were done using a Philips CM 100 microscope.

## 3. Results and Discussion

**3.1. Structure Analyses.** The XRD spectrum of undoped and Cu-doped  $\text{SnO}_2$  samples is displayed in Figure 1(a). The diffraction pattern corresponds to the rutile type phase of  $\text{SnO}_2$  with tetragonal structure (JCPDS 41-1445) showing a preferred orientation along the (110) plane. All the observed

TABLE 1: The structural parameters and optical bandgap of pure and Cu-doped SnO<sub>2</sub> nanoparticles.

Doping level	Scherrer's method		W-H method		Lattice constants (Å)		Unit cell volume (Å) <sup>3</sup>	Energy bandgap (eV)
	<i>D</i> (nm)	Density of dislocation (×10 <sup>15</sup> lines/m <sup>2</sup> )	<i>D</i> (nm)	Strain (×10 <sup>-3</sup> )	<i>a</i> = <i>b</i>	<i>c</i>		
0%	19.01	2.76	24.53	0.148	4.740	3.188	71.626	3.74
1%	18.23	3.00	23.56	0.70	4.744	3.193	71.860	3.69
2%	17.54	3.23	23.32	1.06	4.751	3.196	72.140	3.65
3%	14.87	4.52	16.56	1.22	4.761	3.196	72.444	3.61
4%	13.79	5.25	15.38	1.31	4.761	3.197	72.466	3.54

diffraction peaks were sharp and strong, which showed that the synthesized samples were fully crystallized and polycrystalline in nature. The addition of copper did not affect the preferred orientation along the (110) plane and crystal structure. No diffraction peaks due to copper or copper oxide were visible in the pattern, which revealed that the produced nanostructure had a tetragonal phase structure of SnO<sub>2</sub>. The Cu-doped samples indicate that SnO<sub>2</sub> is still in the rutile phase, suggesting a substitutional doping of Cu at Sn-site [26]. The diffraction peak intensity decreased with an increase in Cu<sup>2+</sup> ion concentrations, which revealed that the doping induces defects in the nanostructure, hindering the growth of the SnO<sub>2</sub> crystal grains [29]. The lattice parameters (*a* = *b*, and *c*) for the tetragonal phase structure were calculated from the XRD peaks using Equation (1):

$$\frac{1}{d_{hkl}^2} = \frac{h^2 + k^2}{a^2} + \frac{l^2}{c^2}, \quad (1)$$

where the plane spacing (*d*<sub>*hkl*</sub>) was calculated by Bragg's law,  $n\lambda = 2d_{hkl} \sin(\theta_{hkl})$ .

The diffraction peaks shifted toward the lower angle as the copper concentration increased. This indicated that Cu<sup>2+</sup> ions have been incorporated into the crystal lattice of SnO<sub>2</sub>, thereby developing lattice strain within the host crystal lattice [30, 31]. Figure 1(b) displays the (110) peak shift of SnO<sub>2</sub> prepared at different dopant concentrations. The diffraction peaks were affected by lattice strain and crystallite size, which increased the peak width and shifted the 2θ peak position [30]. The shifting of the peaks toward the lower angle increased the *d*-spacing, indicating the expansion of lattice parameters as well as the unit cell volume of the Cu-doped SnO<sub>2</sub> crystal. The average lattice constants determined from the diffraction planes were found to be *a* = *b* = 4.752 Å, *c* = 3.194 Å, and their ratio *c/a* = 0.672 (Table 1) which matches with the reference file data values *a* = 4.738 Å, *c* = 3.187 Å, and *c/a* = 0.672. The small difference comes from the microstructural tensile strain due to the incorporation of more Cu<sup>2+</sup> ions which have a larger ionic radius (0.073 nm) than the tin ions (0.069 nm) [27]. The *c/a* parameter has also been found to show a good match with the value for the ideally close-packed tetragonal structure (0.672). The lattice parameters variation with Cu-doping concentration is presented in Figure 2(a).

The increase of the full-width at half-maximum (FWHM) with Cu<sup>2+</sup> ion concentration is an indication of the decrease of crystallite size in the doped SnO<sub>2</sub> samples [32, 33] (Figure 2(b)). The average crystallite size was calculated by Scherrer's formula and Williamson–Hall (W–H) plot. The three highest diffraction peaks, (110), (101), and (211) were used to estimate the crystallite size (*D*) by Scherrer's formula (Equation (2)) [34]:

$$D_s = \frac{k\lambda}{\beta_s \cos \theta}, \quad (2)$$

where  $\lambda$  is X-ray wavelength ( $\lambda = 1.5406$  Å); *k* is the shape factor (0.9);  $\beta$  is full width at half maximum of the diffraction peak in radians and  $\theta$  is Bragg diffraction angle in radians.

The lattice strain ( $\epsilon$ ) and crystallite size were calculated from W–H plot,  $\beta \cos \theta$  versus  $4 \sin \theta$  (Figure 3(a)) [35]. We obtained the strain and crystal size from the slope of the linear-fitted data and the intercepts of the curve, respectively [36]. The reason for the development of the strain could be due to the presence of Cu<sup>2+</sup> ion dopant. The estimated strain increased with the addition of the dopant concentration (Table 1). The difference in the values of the crystal size determined from the Scherrer formula and W–H plot can be ascribed to the fact that W–H plot considers the broadening caused by the effects of tensile strain. The results from both methods indicated that as the doping concentration increased from 1% to 4%, the crystal size decreased. A similar trend of change in crystal size was reported by different literatures [37, 38] for Cu-doped SnO<sub>2</sub> nanoparticles. This decrease in crystal size of SnO<sub>2</sub> as the Cu-doping level increases suggesting that the growth was suppressed due to doping of Cu<sup>2+</sup> ion into Sn site. This can be attributed to the incorporation of Cu<sup>2+</sup> ion in the host lattice to create more nucleation sites which would increase the lattice strain by slowing down the growth of the crystals [37, 39]. The density of dislocation ( $\delta = 1/D^2$ ) [35, 40], where *D* is the crystallite size from Scherrer's formula, caused by the microstructural strain and other defects seen to increase as Cu-doping increased in the host lattice (Figure 3(b)).

**3.2. SEM, EDX, and TEM Analyses.** The study of the surface morphology of Cu-doped and undoped SnO<sub>2</sub> nanoparticle has been carried out using scanning electron microscopy

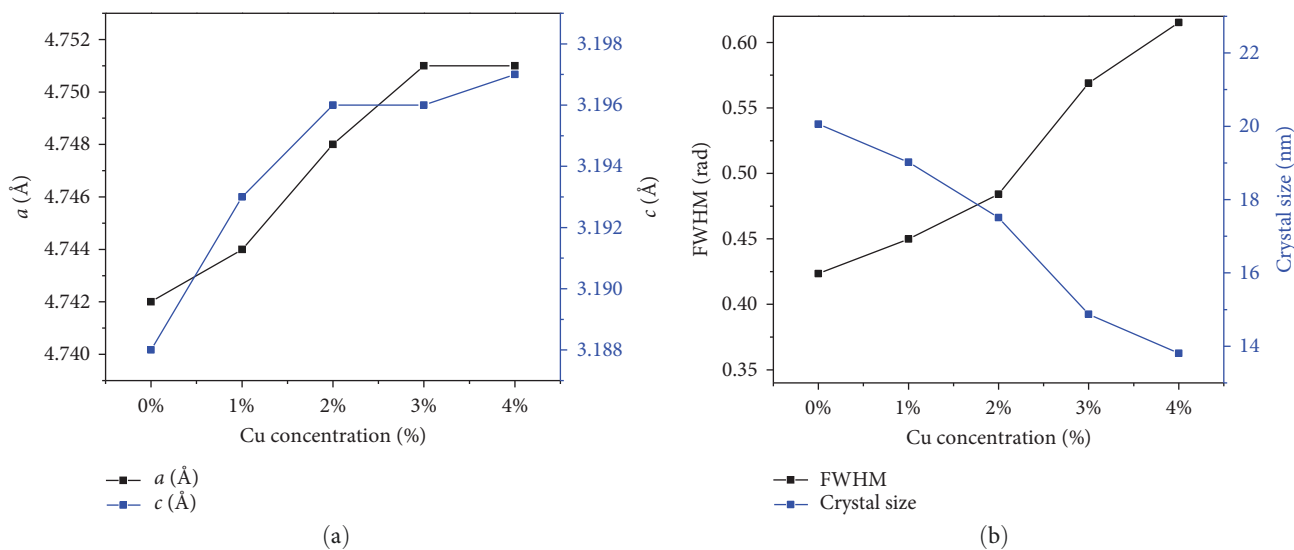


FIGURE 2: (a) Variation of lattice parameters  $a$  and  $c$  with Cu-doping concentration; (b) crystal size and FWHM analysis graph of SnO<sub>2</sub> with Cu-doping concentration.

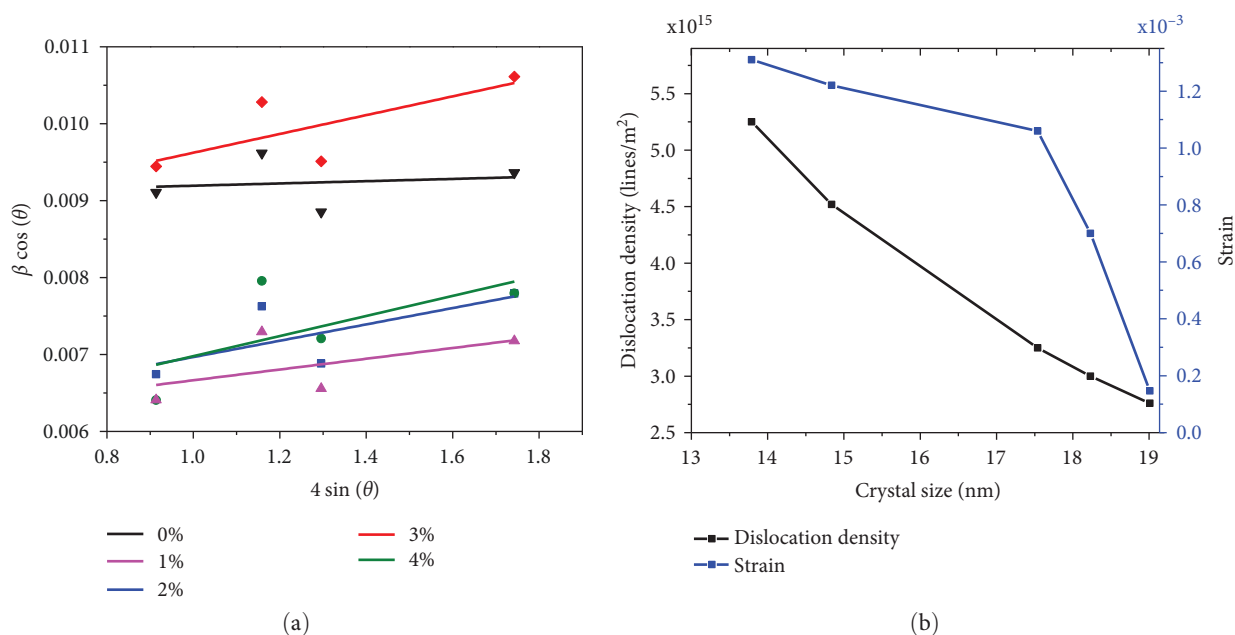


FIGURE 3: (a) Williamson-Hall plots of undoped and Cu-doped SnO<sub>2</sub> nanoparticle; (b) variation of strain and dislocation density with crystal size.

(SEM) coupled with EDX spectroscopy. The SEM images with 60 KX magnification are presented in Figure 4(a)–4(e). A closer inspection of the images showed that the Cu-doped and undoped samples contained spherical-shaped grains and the morphology of SnO<sub>2</sub> nanoparticles changed with increasing Cu-doping concentration. Depending on the dopant contents, the particles were agglomerated with different sizes. This could be an indication for the formation of larger grain sizes through clustering of the primary particles. The rate of nucleation and growth processes influences the size distribution of the grain and the extent of agglomeration [41]. Figure 4(f)–4(h) illustrates the TEM images of the synthesized

samples with 0%, 2%, and 4% of copper doping, respectively. They showed nearly spherical-shaped particles. The particle size distributions for Cu-doped SnO<sub>2</sub> nanoparticles were estimated from TEM images. It can be seen from the images that the particles are nonuniform in size and have an average grain size of 33.60, 29.42, and 26.71 nm for the samples with 0%, 2% and 4% Cu-doping, respectively.

The EDX study was also carried out for looking at the elemental composition of the synthesized samples. The EDX analysis was done at different places in the same sample which confirmed the fairly uniform composition of the nanocrystals. The EDX spectra and their percentage composition of 0%, 1%,

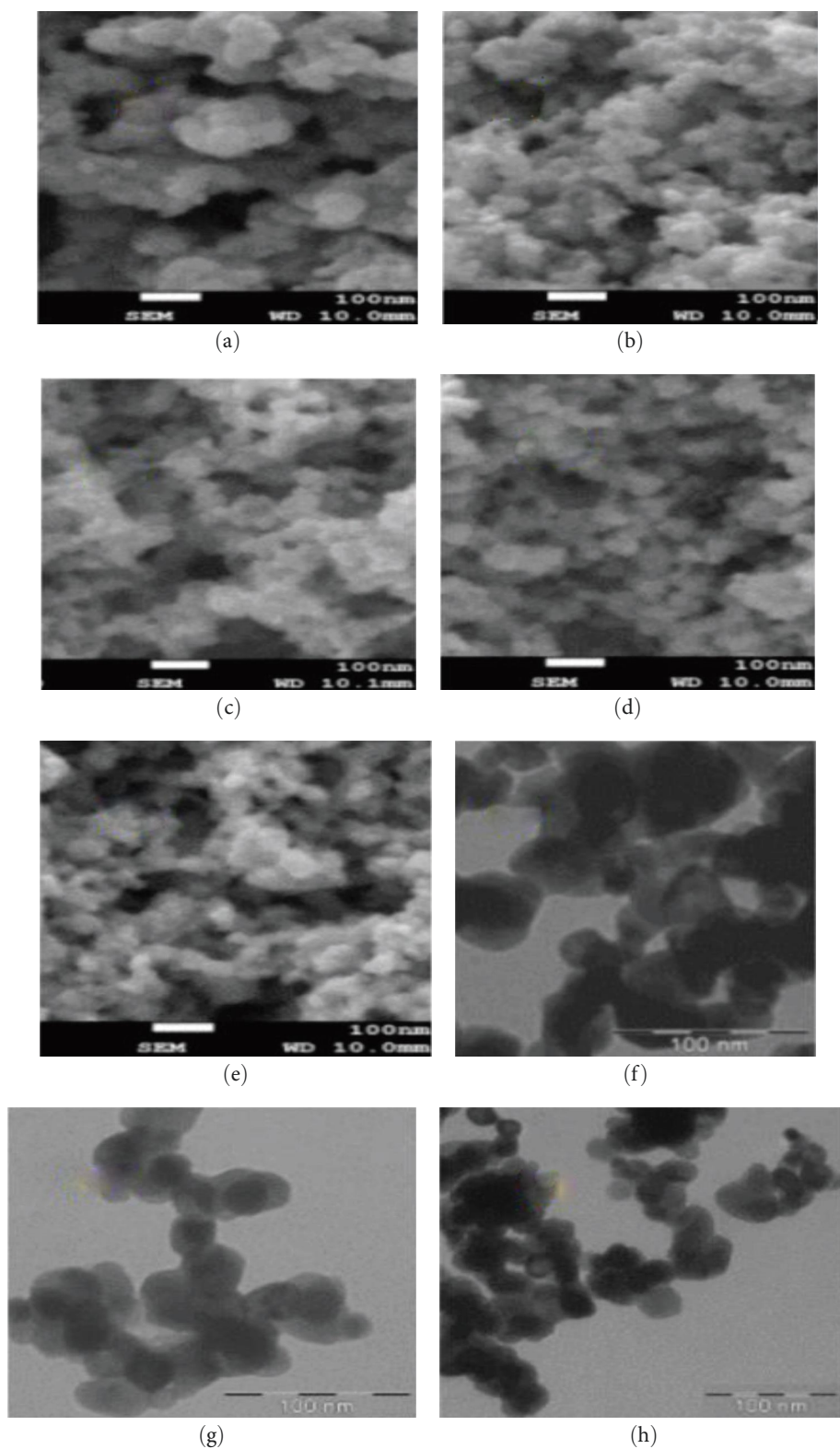


FIGURE 4: SEM images of (a) pure SnO<sub>2</sub>; (b) 1% Cu-doped; (c) 2% Cu-doped; (d) 3% Cu-doped; and (e) 4% Cu-doped SnO<sub>2</sub>. TEM micrographs of (f) pure SnO<sub>2</sub>; (g) 2% Cu-doped SnO<sub>2</sub>; (h) 4% Cu-doped SnO<sub>2</sub>.

3%, and 4% Cu-doped samples are shown in Figure 5(a)–5(d). The spectra for the doped samples confirm the incorporation of Cu<sup>2+</sup> ions into SnO<sub>2</sub> lattice. The atomic percentages composition is close to the stoichiometric ratio for all samples.

**3.3. Optical Properties Analyses.** The UV–Visible diffuse reflectance spectra of SnO<sub>2</sub> doped with copper were recorded in the 250–800 nm wavelength range. It provides useful information about the optical bandgap of the prepared Cu-doped

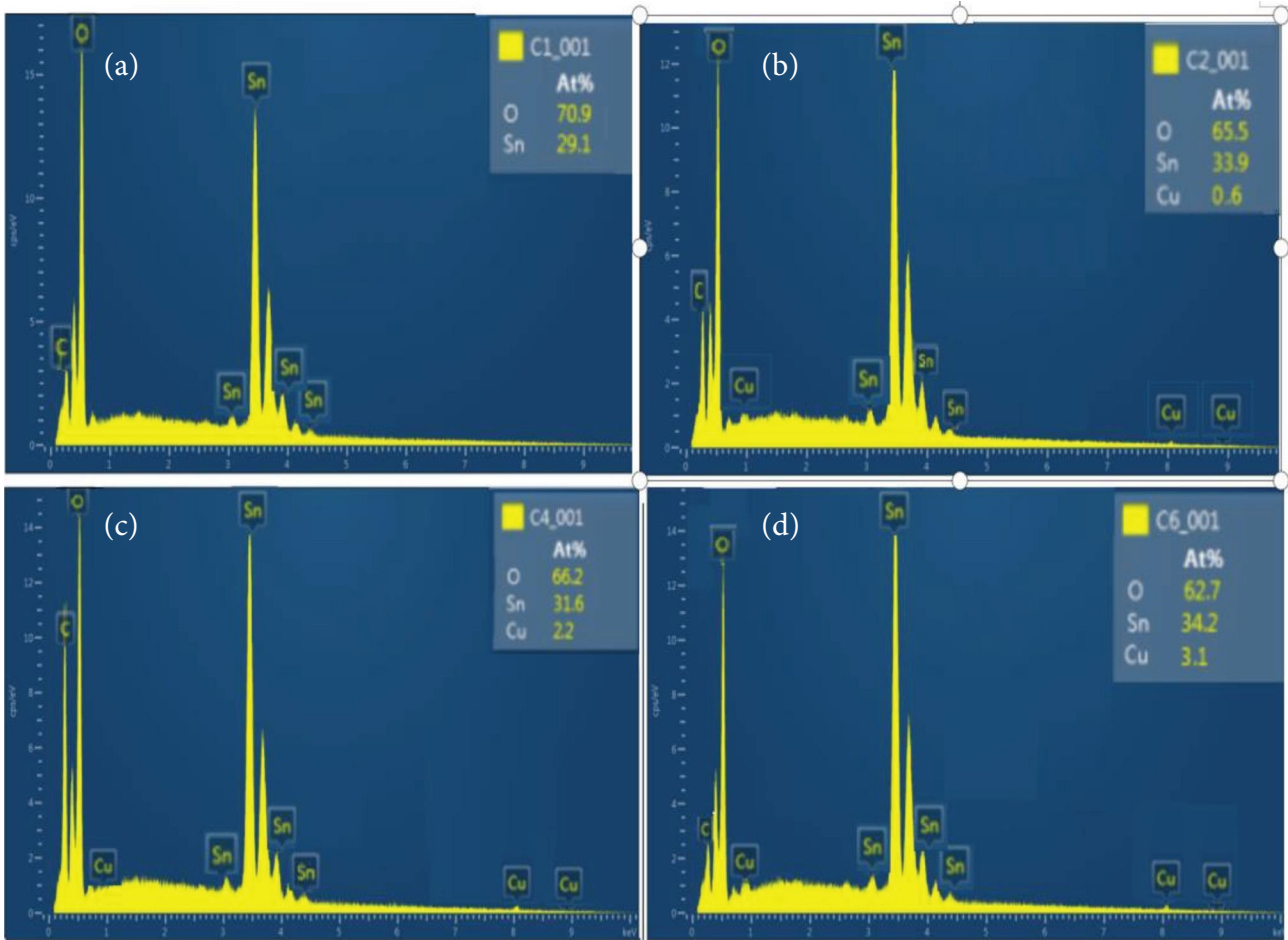


FIGURE 5: EDX spectra of (a) pure SnO<sub>2</sub>; (b) 1% Cu-doped; (c) 3% Cu-doped; (d) 4% Cu-doped SnO<sub>2</sub> nanoparticle.

and undoped samples. The percentage diffuse reflectance of the produced nanostructures is presented in Figure 6(a). The samples have high reflectance spectrum in the visible region but showed small reflectance in the UV region. This might be due to the transfer of charge from the valence to the conduction band of SnO<sub>2</sub> nanostructures [42]. Tauc's relation (Equation (3)) was used to calculate the bandgap ( $E_g$ ) of the prepared samples:

$$\alpha h\nu = A(h\nu - E_g)^n, \quad (3)$$

where  $A$  is a constant,  $h\nu$  is the photon energy,  $\alpha$  is the absorption coefficient and  $n$  is equal to 1/2 for the direct allowed transition. To determine the value of the optical bandgap of Cu-doped SnO<sub>2</sub>, the reflectance values were converted to absorbance by applying the Kubelka–Munk transformation function (Equation (4)):

$$F(R) = \frac{(1 - R)^2}{2R}, \quad (4)$$

where  $R$  is the percentage reflectance and  $F(R)$  is the Kubelka–Munk function which is proportional to the absorption coefficient ( $\alpha$ ).

In Figure 6(b), the absorption spectrum exhibits only one absorption edge with no visible shoulder on the lower energy side, indicating the absence of any secondary phase-related absorption edges. To estimate the bandgap, we extrapolate the linear portion of the  $(F(R)h\nu)^2$  versus  $h\nu$  curve to zero, where  $h\nu$  is the photon energy [43]. From Figure 6(b), the bandgap energy obtained from the curve was around 3.74 eV for undoped. The magnitude of bandgap was decreasing from 3.69 to 3.54 eV as Cu-dopant concentration increased from 1% to 4% (Table 1). A similar narrowing in the bandgap was reported in the literature for the doping of SnO<sub>2</sub> with copper [25, 37]. We find that the absorption edge is red shifted with increasing Cu<sup>2+</sup> ion content in SnO<sub>2</sub> lattice. Presence of dopant introduces impurity band levels into the bandgap of tin oxide. The narrowing bandgap energy is possibly due to the existence of copper impurities in the SnO<sub>2</sub> structure, which induce the formation of new recombination centers with lower emission energy [17]. Thus, with a greater amount of copper ions, the transition of electrons occurs from the filled valence band to the new recombination centers instead of the transition of electrons from the filled valence band to the empty conduction band which usually occurs [44].

**3.4. Photoluminescence Analyses.** The PL technique has been widely used to examine the defects and to understand the fate of photogenerated electrons and holes in semiconductors [45].

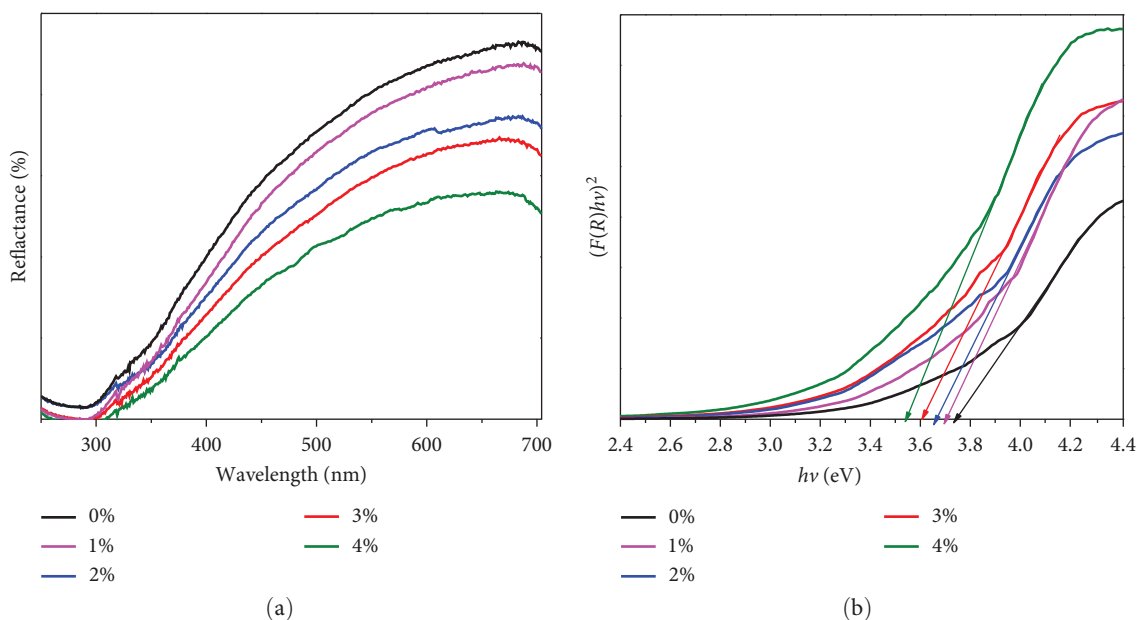


FIGURE 6: (a) UV-Vis diffuse reflectance spectra versus wavelength; (b) Tauc's plot of  $(F(R)hv)^2$  versus photon energy.

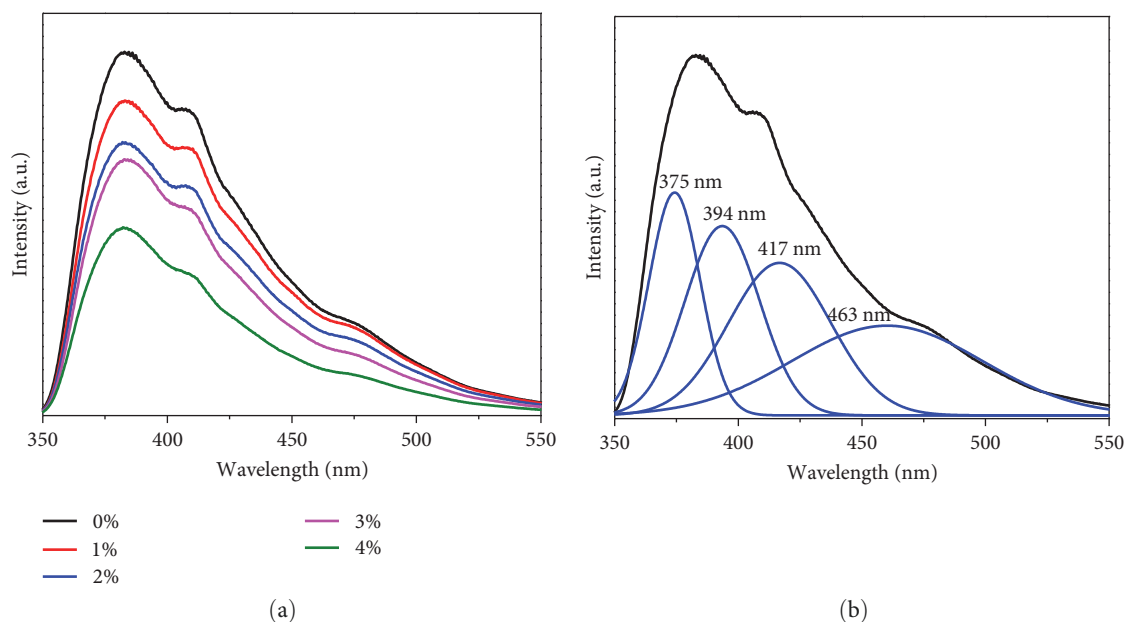


FIGURE 7: (a, b) PL emission spectra of undoped and Cu-doped SnO<sub>2</sub> nanoparticles excited at 280 nm.

Usually, metal-doped SnO<sub>2</sub> nanostructures display defect emissions. These defects might include oxygen vacancies and the luminescence center formed by tin interstitial or the presence of dangling bonds in SnO<sub>2</sub> nanoparticles [46]. The oxygen vacancies are the main defects, which form the donor levels in SnO<sub>2</sub> nanostructures [47]. The defects are mainly located on the surface of the nanostructures and could form a series of energy levels within the bandgap of SnO<sub>2</sub> by trapping electrons from the valence band. The room temperature PL spectrum was recorded at an excitation wavelength of 300 nm as shown in Figure 7(a). The broad spectrum measured from the undoped sample was deconvoluted by to four components centered at 375, 395, 417, and 463 nm using

Gaussian fitting, as shown in Figure 7(b). These emission energies are much lower than the energy bandgap of SnO<sub>2</sub> calculated from the diffuse reflectance spectra; therefore, the electron from the conduction band was not directly recombined to a hole in the valence band; instead, they are assigned to the electron transition, mediated by defect levels in the bandgap.

The UV emission around 375 nm may be attributed to the transition of electrons mediated by oxygen vacancies [1, 48]. Oxygen vacancies are the most common defects which form donor levels inside the bandgap [49]. The UV emission was due to the transition of electrons from the donor level to the valence band. The peak around 395 nm can be attributed to

the electron transition mediated by defect levels in the band. The defects are mainly located on the surface of the nanostructures and could form a series of energy levels within the bandgap of SnO<sub>2</sub> nanomaterials by trapping electrons from the valence band [50]. The peak around 417 nm is attributed to the tin interstitial present in the nanocrystals which forms a significant amount of trapped in the bandgap of SnO<sub>2</sub> nanocrystals [48]. It is reported that the SnO<sub>2</sub> nanostructures can exhibit emission in the visible region because of the presence of different defect energy levels inside the SnO<sub>2</sub> bandgap [10, 51]. The visible emission around 463 nm may be attributed to the formation of luminescent centers formed by defects [52, 53].

It can also be seen that the PL intensity of the undoped SnO<sub>2</sub> is higher than that in the spectra of Cu-doped samples, indicating that copper doping can inhibit the excited electron and hole recombination. An intensity reduction was observed for the emission peaks with an increase in Cu-doping concentration [37]. This could be attributed to a concentration quenching effect characterized by the energy exchange between a pair of Cu<sup>2+</sup> ions [24].

#### 4. Conclusion

Nanocrystalline Cu-doped SnO<sub>2</sub> was successfully synthesized via a simple coprecipitation technique at variable Cu-doping concentrations. The XRD analysis showed a tetragonal structure for all Cu-doped and undoped SnO<sub>2</sub> samples. The average crystallite size decreased with the increasing of Cu-dopant concentration. A shift in the (110) SnO<sub>2</sub> plane confirmed the incorporation of dopants in the crystal lattice. The doping of copper atoms, with greater ionic radius than Sn atoms in the host, created a tensile strain in the lattice and caused expanded lattice parameters of the produced nanoparticles. The SEM micrographs showed spherically shaped particles. The TEM images confirmed that the prepared undoped and Cu-doped SnO<sub>2</sub> were in nano size. The optical studies revealed that Cu doping in the SnO<sub>2</sub> lattice reduced the bandgap energy from 3.74 to 3.54 eV. The noticed red shift in the bandgap was due to the incorporation of Cu<sup>2+</sup> ions in the host lattice which created additional energy levels inside the bandgap. The PL spectra exhibited UV and visible emissions between 350 and 500 nm due to defects such as oxygen vacancies and the luminescence center formed by tin interstitial or the presence of dangling bonds in SnO<sub>2</sub> nanoparticles.

#### Data Availability

The data used to support the findings of this study are available from the corresponding author upon request.

#### Conflicts of Interest

The authors declare that they have no conflicts of interest.

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