

## Research Article

# Quantitative Raman Spectroscopic Determination of the Composition, Pressure, and Density of CO<sub>2</sub>-CH<sub>4</sub> Gas Mixtures

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The Raman spectra for pure CO<sub>2</sub> and CH<sub>4</sub> gases and their ten gas mixtures were collected at pressures and temperatures ranging from 2 MPa to 40 MPa and room temperature (~24°C) to 300°C, respectively. A systematic analysis was carried out to establish a methodology for the quantitative determination of the composition, pressure, and density of CO<sub>2</sub>-CH<sub>4</sub> mixtures. The shift in the peak position of the  $\nu_1$  band for CH<sub>4</sub> was sufficiently large to enable the accurate determination of the pressure of pure CH<sub>4</sub> and CH<sub>4</sub>-dominated fluids (>50 mol% CH<sub>4</sub>). An equation representing the observed relationship of the peak position of the  $\nu_1$  band of CH<sub>4</sub>, density, and composition was developed to calculate the density of CO<sub>2</sub>-CH<sub>4</sub> mixtures. The Raman quantification factor  $F(\text{CH}_4)/F(\text{CO}_2)$  was demonstrated to be near a constant value of  $5.048 \pm 0.4$  and was used to determine the CH<sub>4</sub> to CO<sub>2</sub> molar ratio in an unknown CO<sub>2</sub>-CH<sub>4</sub>-bearing fluid with high internal pressure (>10 MPa) based on the Raman peak area ratio. The effect of temperature on the variation in Raman spectral parameters was also investigated at temperatures up to 300°C. The results showed that the effect of temperature must be considered when Raman spectral parameters are used to calculate the pressure, density, and composition of CO<sub>2</sub>-CH<sub>4</sub> gas mixtures. Raman spectroscopic analysis results obtained for six samples prepared in fused silica capillary capsules were validated by comparison with the results obtained from microthermometry measurements.

## 1. Introduction

Carbon dioxide and methane are two important gaseous components in many geological environments and play important roles during geological processes, such as metamorphic reactions, melting, and deformation [1]. CO<sub>2</sub>-rich fluid inclusions (FIs) are often found in metamorphic and mantle-derived rocks [2], and occur in close association with many ore deposits [3], such as gold [4], copper [5], and rare Earth element (REE) deposits [6], while CH<sub>4</sub>-rich FIs are commonly found in sedimentary basins [7, 8], low-grade metamorphic rocks [9], and mid-ocean ridge hydrothermal environments [10, 11]. However, water-free FIs that are approximated by the CO<sub>2</sub>-CH<sub>4</sub> system are common in some

metamorphic terranes [12–14]. The information acquired from the composition and density (pressure) of these FIs can be used to help reconstruct the physicochemical conditions under which they and their host minerals are formed [15].

Although several analytical techniques can be used to analyze FIs, only microthermometry (MT) and Raman spectroscopy can provide information on both the composition and pressure of individual inclusions [16, 17]. MT is a nondestructive method based on the observation of phase-transition temperatures. When the melting temperature  $T_m$  and homogenization temperature  $T_h$  are known, the  $VX$  (molar volume and composition) properties can be directly acquired from the  $VX$  diagram [18, 19]. However, the uncertainties in the composition and density obtained from the

VX diagram due to the error in  $T_m$  and  $T_h$  ( $\sim\pm 0.1^\circ\text{C}$ ) can reach up to  $\pm 5$  mol% and  $\pm 0.09\text{ g}\cdot\text{cm}^{-3}$ , respectively [19, 20]. In addition, the microthermometry technique is difficult to apply to small FIs with diameter of  $<5\ \mu\text{m}$ , low density, and more complex composition [20]. As an alternative to microthermometric measurements, Raman spectroscopy is an extreme useful tool to get quantitative analysis in individual inclusions, partly because it is fast, nondestructive, and high spatial resolution ( $\sim 1\ \mu\text{m}$ ).

Over the past decades, previous studies have been paid more attention to the quantitative analysis of FIs with Raman spectroscopy. Spectral parameters (peak position, peak width, and peak intensity ratio) could be used to determine the pressure, density, and composition of the fluid in FIs. For instance, the concentration of dissolved  $\text{CH}_4$  or  $\text{CO}_2$  in water solution has been determined based on their peak area (or height) ratio [21–25], and the pressure (or density) of  $\text{CH}_4$  can be determined by the peak position of the  $\nu_1$  band of  $\text{CH}_4$  [17, 26–33]. In addition, the  $\text{CO}_2$  density in FIs can be obtained by the measurements of the Fermi diad splitting of  $\text{CO}_2$  [3, 19, 20, 33–42]. The effect of composition on the Raman spectral parameters of  $\text{CO}_2$  and  $\text{CH}_4$  was reported earlier by Seitz et al. [17] with low spectral resolution ( $\sim 5\text{ cm}^{-1}$ ) [20]. Recently, calibration data for  $\text{CO}_2$ - $\text{CH}_4$  mixture and  $\text{N}_2$ - $\text{CO}_2$ - $\text{CH}_4$  mixture have been provided by Le et al. [20] and Sublett et al. [43], respectively. The quantitative relationship between fluid properties (compositions, density, or pressure) and Raman spectral parameters (peak area ratio, peak position, etc.) has to be calibrated carefully with standard samples, since an accurate calibration is fundamental for quantitative measurements.

For two Raman-active species, a and b, in a homogeneous phase, their molar ratios can be calculated from their peak areas [44] based on the following equations:

$$\frac{A_a}{A_b} = \left(\frac{C_a}{C_b}\right) \times \left(\frac{\sigma_a}{\sigma_b}\right) \times \left(\frac{\eta_a}{\eta_b}\right) = \left(\frac{C_a}{C_b}\right) \times \left(\frac{F_a}{F_b}\right), \quad (1)$$

where  $A$  is the peak area;  $C$  is the concentration;  $\sigma$  is the Raman scattering coefficient;  $\eta$  is the instrumental efficiency factor; and  $F$  is the Raman quantification factor [44].

Normally, the reported Raman quantification factors for gaseous are relative to  $\text{N}_2$  [45]. However, the reported values for  $F$  ( $\text{CH}_4$ ) and  $F$  ( $\text{CO}_2$ ) among different laboratories are quite different [19, 20, 38, 46]. Considering that two separate experiments are required to obtain  $F$  ( $\text{CH}_4$ ) and  $F$  ( $\text{CO}_2$ ), a more practical method is to determine the  $F$  ( $\text{CH}_4$ )/ $F$  ( $\text{CO}_2$ ) ratio; then, the molar ratio of  $\text{CH}_4$  to  $\text{CO}_2$  can be obtained directly according to the peak area ratio  $A$  ( $\text{CH}_4$ )/ $A$  ( $\text{CO}_2$ ). In addition, there are only a few Raman spectroscopy calibration data available in the literature for  $\text{CO}_2$ - $\text{CH}_4$  mixtures at high temperatures up to  $300^\circ\text{C}$ , except for the pure  $\text{CO}_2$  and  $\text{CH}_4$  system [30, 33, 37] and  $\text{N}_2$ - $\text{CO}_2$ - $\text{CH}_4$  system [43].

In this study, an in situ Raman spectroscopy system combined with a fused silica capillary-high-pressure optical cell (FSC-HPOC) is applied to collect Raman spectra for pure  $\text{CO}_2$  and  $\text{CH}_4$  gases and their  $\text{CO}_2$ - $\text{CH}_4$  binary mixtures at temperatures ranging from room temperature ( $\sim 24^\circ\text{C}$ ) to  $300^\circ\text{C}$  and pressures up to 40 MPa. The peak

position and peak area of the  $\nu_1$  band of  $\text{CH}_4$  and the  $\nu_1$  and  $2\nu_2$  bands (Fermi diad) of  $\text{CO}_2$  are systematically analyzed to establish a methodology for the quantitative determination of the composition, pressure, and density of  $\text{CO}_2$ - $\text{CH}_4$  FIs. The purpose of our work is to construct relationships between Raman spectral parameters and composition, pressure, and density with potential applications to the quantitative study of natural  $\text{CO}_2$ - $\text{CH}_4$  inclusions. In addition, the effects of temperature on Raman spectral parameters are also investigated. Six synthetic  $\text{CO}_2$ - $\text{CH}_4$  FIs, prepared in fused silica capillary capsules (FSCCs), were used to perform both laser Raman and MT studies to validate our established methodology.

## 2. Materials and Methods

**2.1. High-Pressure Optical Cell.** The instruments and procedures are similar to those used in the study of Lu et al. [29] (Figure 1). The experiments were carried out with a capillary high-pressure optical cell (HPOC [47]). Pure  $\text{CO}_2$  and  $\text{CH}_4$  gases (99.99%, Jiateng Air Production Inc.) and their ten gaseous mixtures with given compositions (10, 20, 30, 40, 50, 60, 70, 80, 85, and 90 mol%  $\text{CH}_4$ ) were investigated systematically. The HPOC was composed of a fused silica capillary with a circular section (Polymicro Technologies, LLC) with an inner (ID) and outer (OD) diameter of  $300\ \mu\text{m}$  and  $665\ \mu\text{m}$ , respectively, and a length of approximately 25 cm. The gas or gas mixture was loaded into the HPOC and sealed with a column of mercury ( $\sim 15\text{ mm}$  long) at one atmospheric pressure. The procedure used for the preparation of the gas samples was described previously in detail [48]. The sample temperature was controlled using a Linkam CAP500 heating-cooling stage. The pressure in the cell was maintained by a water medium in the pressure line and was adjusted using a pressure generator (HiP Model no. 37-6-30). The pressure was read using a Setra 204D digital pressure transducer with a Datum 2000™ manometer (69 MPa full scale, accurate to  $\pm 0.14\%$  of full scale).

**2.2. Fused Silica Capillary Capsule.** Fused silica capillary capsules (FSCCs;  $300\ \mu\text{m}$  OD,  $100\ \mu\text{m}$  ID, and approximately 1.0 cm long) were constructed to contain the  $\text{CO}_2$ - $\text{CH}_4$  mixture gas using the sample-loading system of Chou et al. [47], and they were used for verification of our quantitative Raman spectroscopic analyses of the  $\text{CO}_2$ - $\text{CH}_4$  system. First, the line was flushed twice with one gas ( $\text{CH}_4$  or  $\text{CO}_2$ ) after placing under vacuum. Then, gas at a pressure of slightly lower than 0.2 MPa was loaded into the line and silica tube, and the closed end of the tube was immersed in liquid nitrogen to form liquid or solid  $\text{CH}_4/\text{CO}_2$ . Second, the system was evacuated, while the closed end was still immersed in liquid nitrogen, and the silica tube was then pressurized with another gas ( $\text{CO}_2$  or  $\text{CH}_4$ ). Third, the system was evacuated again, and the open end of the tube was sealed with a hydrogen flame in vacuum, while the closed end remained immersed in liquid nitrogen. A higher loading concentration in  $\text{CO}_2$ - $\text{CH}_4$  mixture was obtained for the gas loaded first. Even though  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures

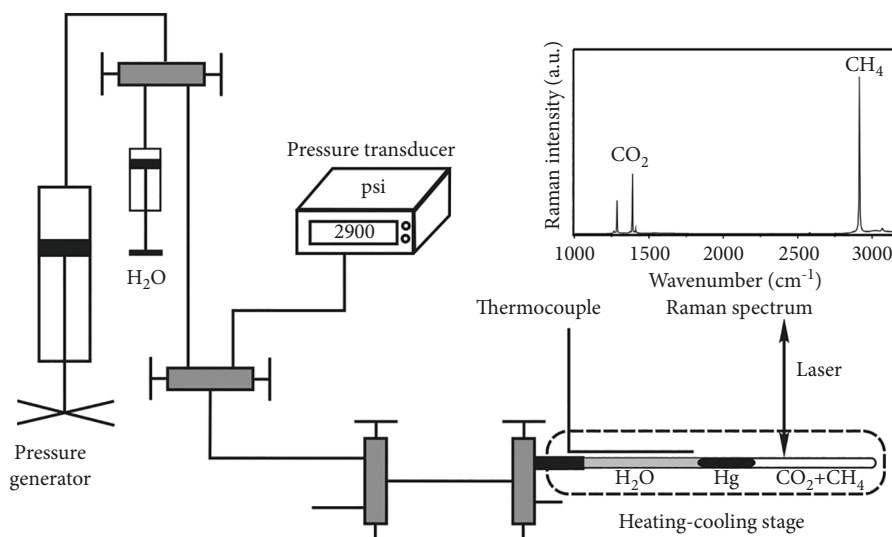


FIGURE 1: Schematic diagram of the experimental system (modified after Lu et al. [29]).

with known proportions could also be loaded directly into the FSC tubes, we found that their proportions changed slightly due to the different freezing temperature while freezing.

### 2.3. Microthermometry Analyses of FSCC Samples.

Microthermometric analyses of the FSCC samples were carried out using a Linkam CAP500 heating-cooling stage under the Olympus microscope in the laser Raman spectroscopic system. The temperature was calibrated with the triple points of H<sub>2</sub>O ( $\pm 0^\circ\text{C}$ ) and CO<sub>2</sub> ( $-56.6^\circ\text{C}$ ) and was accurate to  $\pm 0.1^\circ\text{C}$ . In the present study, measurements of the phase changes were carried out only below room temperature. The system was cooled rapidly at a rate of  $10^\circ\text{C}/\text{min}$  to  $-180^\circ\text{C}$ , and after holding for 10 minutes, it was then heated at a rate of  $10^\circ\text{C}/\text{min}$ . The heating rate was reduced to  $0.3^\circ\text{C}/\text{min}$  at a temperature of approximately  $5^\circ\text{C}$  below the phase transition point to precisely determine the temperature of the phase changes. The same rates were used for heating-cooling cycling. The CO<sub>2</sub>-CH<sub>4</sub> system shows three phases at low temperatures: solid CO<sub>2</sub>, liquid, and vapor. Two phase transition temperatures can be measured during heating from  $-180^\circ\text{C}$  to room temperature: the final melting temperature of solid CO<sub>2</sub> ( $T_m$ ) and the L-V homogenization temperature ( $T_h$ ) for the disappearance of the liquid phase or vapor bubble [49]. The measured values for  $T_m$  and  $T_h$  were then used to calculate the composition and density of the FSCC samples based on the VX diagram of Thiery et al. [18].

### 2.4. Collection and Calibration of Raman Spectra.

Raman spectra were acquired using a JY/Horiba LabRam HR Evolution Raman system equipped with a 532.06 nm (frequency-doubled Nd: YAG) laser, an SLWD 50 $\times$  Olympus objective with a numerical aperture of 0.35, and an 1800 groove/mm grating with a spectral resolution of approximately  $0.65\text{ cm}^{-1}$  or a 600 groove/mm grating with a spectral resolution of approximately  $2\text{ cm}^{-1}$ . Approximately 20 mW

of laser light was focused onto the center of the horizontal tube to generate Raman signals during the measurement.

Raman spectra for CO<sub>2</sub> and CH<sub>4</sub> are shown in Figure 2. CO<sub>2</sub> has five Raman bands, including the Fermi diad of CO<sub>2</sub> between 1200 and 1400  $\text{cm}^{-1}$  (the lower band at  $\sim 1285\text{ cm}^{-1}$  and the upper band at  $\sim 1388\text{ cm}^{-1}$ ), two hot bands flanking the Fermi diad at  $\sim 1265$  and  $1409\text{ cm}^{-1}$ , and a weak CO<sub>2</sub> band at  $1370.0\text{ cm}^{-1}$  due to the isotopic splitting of <sup>13</sup>CO<sub>2</sub> [37]. CH<sub>4</sub> has four Raman bands: the asymmetric bending overtone band ( $2\nu_4$  at  $\sim 2580\text{ cm}^{-1}$ ), the C-H symmetric stretching band ( $\nu_1$  at  $\sim 2917\text{ cm}^{-1}$ ), the asymmetric stretching band ( $\nu_3$  at  $\sim 3020\text{ cm}^{-1}$ ), and the asymmetric bending overtone band ( $2\nu_2$  at  $\sim 3070\text{ cm}^{-1}$ ) [50]; the other three bands have an intensity that is  $\sim 2\%$  of that of the  $\nu_1$  band. In this study, only the Fermi diad and hot bands of CO<sub>2</sub> and  $\nu_1$  band of CH<sub>4</sub> (approximately  $2917\text{ cm}^{-1}$ ) were investigated.

As shown in Figure 2, the 600 groove/mm grating, with a lower spectral resolution but wider spectral range, was used for collecting spectra in multiple windows ranging between  $1000\text{ cm}^{-1}$  and  $3200\text{ cm}^{-1}$  to cover the peaks of the Fermi diad of CO<sub>2</sub> and CH<sub>4</sub>. In addition, spectra employing the 1800 groove/mm grating with a high spectral resolution were also collected in two single windows to obtain precise peak positions. One window ranged between  $1170\text{ cm}^{-1}$  and  $1620\text{ cm}^{-1}$ , which covers the Fermi diad of CO<sub>2</sub> ( $1200\sim 1400\text{ cm}^{-1}$ ), and two well-established reference peaks of benzonitrile ( $1192.6$  and  $1598.9\text{ cm}^{-1}$ ) was used to calibrate the Fermi diad of CO<sub>2</sub>. The other window ranged between  $2718\text{ cm}^{-1}$  and  $3075\text{ cm}^{-1}$ , which covers the CH<sub>4</sub> peak at approximately  $2917\text{ cm}^{-1}$ ; neon lines at  $626.649\text{ nm}$  ( $2836.976\text{ cm}^{-1}$ ) and  $630.479\text{ nm}$  ( $2933.916\text{ cm}^{-1}$ ) were selected to calibrate the peak position of CH<sub>4</sub> [30].

The measured spectra with single windows for CO<sub>2</sub> and CH<sub>4</sub> were fitted with the program PeakFit v4.12 (AISN Software Inc.) using Gaussian-Lorentzian functions without smoothing or baseline correction before peak-fitting procedures. The real peak positions of CO<sub>2</sub> and CH<sub>4</sub> were derived from the following relations:

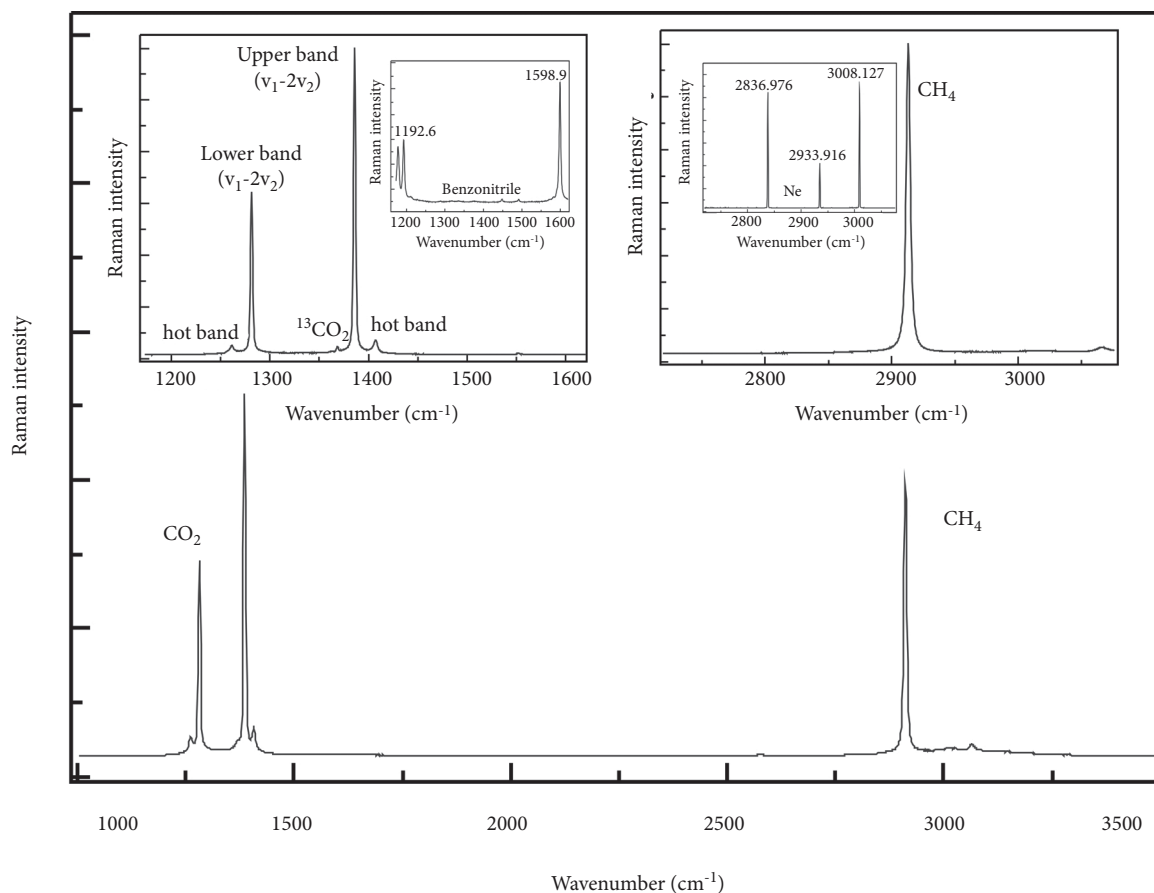


FIGURE 2: Raman spectrum of the CO<sub>2</sub>-CH<sub>4</sub> mixture showing the Raman lines for CO<sub>2</sub> and CH<sub>4</sub>. The top left insert shows the detailed spectrum of CO<sub>2</sub> containing Fermi diad bands, two hot bands, <sup>13</sup>CO<sub>2</sub>, and the spectrum for benzonitrile. The two marked peaks for benzonitrile (1192.6 and 1598.9 cm<sup>-1</sup>) were used to calibrate the CO<sub>2</sub> Fermi diad splitting. The right insert shows the C-H symmetric  $\nu_1$  band of CH<sub>4</sub> and Ne emission lines at 2836.976 cm<sup>-1</sup>, 2933.916 cm<sup>-1</sup>, and 3008.127 cm<sup>-1</sup>. The Ne emission lines at 2836.976 cm<sup>-1</sup> and 2933.916 cm<sup>-1</sup> bracket the  $\nu_1$  band of CH<sub>4</sub> and were used to calibrate the CH<sub>4</sub> $\nu_1$  peak position.

$$\frac{\nu_{\text{real}} - 1192.6}{\nu_m - \text{Benz}_{m1}} = \frac{1598.9 - 1192.6}{\text{Benz}_{m2} - \text{Benz}_{m1}}, \quad (2)$$

$$\frac{\nu_{\text{real}} - 2836.976}{\nu_m - \text{Ne}_{m1}} = \frac{2933.916 - 2836.976}{\text{Ne}_{m2} - \text{Ne}_{m1}}, \quad (3)$$

where  $\nu_{\text{real}}$  is the real position of the upper and lower bands of the Fermi diad of CO<sub>2</sub> in equation (2) and the C-H symmetric stretching band ( $\nu_1$ ) of CH<sub>4</sub> in equation (3);  $\nu_m$  is the measured peak position of CO<sub>2</sub> or CH<sub>4</sub>;  $\text{Benz}_{m1}$  and  $\text{Benz}_{m2}$  are the two reference peaks of benzonitrile; and  $\text{Ne}_{m1}$  and  $\text{Ne}_{m2}$  are the measured wavenumbers of the two reference neon lines.

The Raman peak area ratio (PAR) of the upper band of the Fermi diad of CO<sub>2</sub> to the C-H symmetric stretching band ( $\nu_1$ ) of CH<sub>4</sub> was determined by the software GRAMS/AI (Thermo Galactic). Three spectra were collected at each set  $P$ - $T$  condition.

### 3. Results and Discussion

**3.1. Determination of the Pressure of the CO<sub>2</sub>-CH<sub>4</sub> Gas Mixtures.** The calibrated peak positions for the CH<sub>4</sub> $\nu_1$  peak

and the lower and upper members of the CO<sub>2</sub> Fermi diad and the Fermi diad splittings for all samples at room temperature and various pressures collected in HPOC are listed in Table S1 and S2 (Supplemental material) and shown in Figure 3. Figure 3(a) shows the peak positions for the  $\nu_1$  peak of CH<sub>4</sub> in pure CH<sub>4</sub> and ten CO<sub>2</sub>-CH<sub>4</sub> mixtures as a function of total pressure; the addition of CO<sub>2</sub> can affect the peak position of CH<sub>4</sub>; however, the dominant factor affecting the peak position of CH<sub>4</sub> is the total fluid pressure, which is in agreement with the result of Seitz et al. [17]. A leap in  $\nu_1$  peak position for CH<sub>4</sub> in Figure 3(a) is evident for CO<sub>2</sub>-rich mixtures at low pressure. For the 10 mol% CH<sub>4</sub>-90 mol% CO<sub>2</sub> mixture, this abrupt shift in the peak position in response to pressure represents the liquid-vapor phase transition. The 30% CH<sub>4</sub>-70% CO<sub>2</sub> mixture, however, remains supercritical (but near-critical) at room temperature [17].

The lower and upper members of the Fermi diad peaks for CO<sub>2</sub> and Fermi diad splitting for CO<sub>2</sub> in CO<sub>2</sub>-CH<sub>4</sub> mixtures as a function of pressure are presented in Figures 3(b)-3(d), respectively. The peak shift with pressure from 2 MPa to 40 MPa is much greater for the lower Fermi diad member ( $\sim 5$  cm<sup>-1</sup>) than for the upper Fermi diad

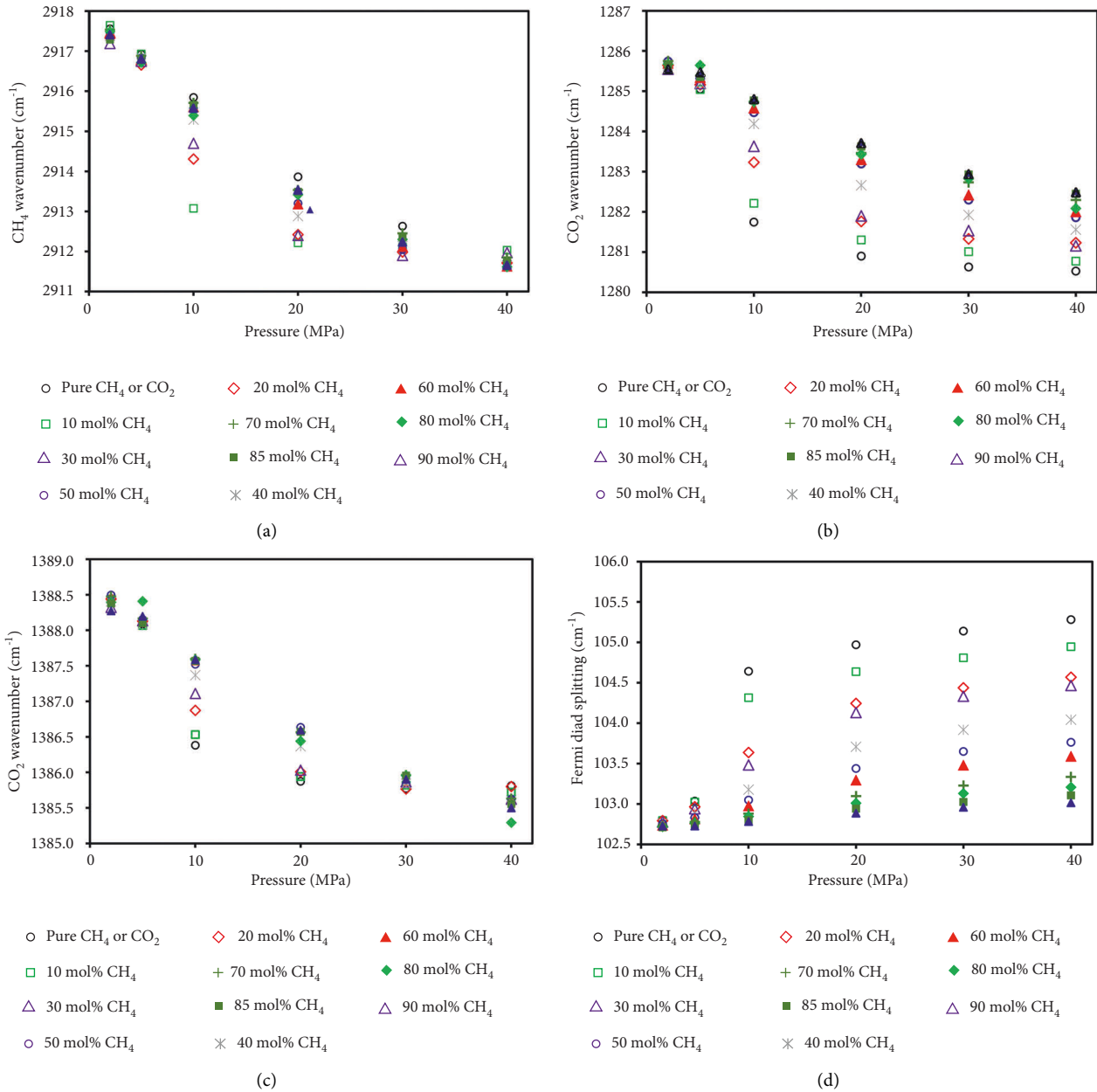


FIGURE 3: Raman peak positions for the  $\text{CH}_4 \nu_1$  peak (a) and the lower (b) and upper (c) members of the  $\text{CO}_2$  Fermi diad and the Fermi diad splitting (d) for different  $\text{CO}_2$ - $\text{CH}_4$  gaseous mixtures at room temperature as a function of pressure.

member ( $\sim 3 \text{ cm}^{-1}$ ). Thus, the peak position of the lower Fermi diad member is a more sensitive monitor of fluid pressure, which are consistent with previous studies [41]. However, the peak position of the lower Fermi diad member is also more sensitive to composition. Indeed, the Fermi diad splitting ( $\Delta$ , the separation between bands in the Fermi diad) of  $\text{CO}_2$  has been previously used to determine the pressure or density of  $\text{CO}_2$  in FIs [37, 40]. Our results show that at 2 MPa,  $\Delta$  is nearly identical ( $\sim 102.76 \text{ cm}^{-1}$ ) for any composition, while at higher pressures, the effect of the gas mixture composition becomes more obvious, which is similar to that observed for  $\text{CO}_2$  mixed with  $\text{N}_2$  and  $\text{CH}_4$  [19, 20]. Therefore, we can draw the same conclusion with

previous studies [17, 20] that the  $\Delta$  for  $\text{CO}_2$  is not only a function of pressure but also a function of the composition of the  $\text{CO}_2$ - $\text{CH}_4$  gas mixture (Figure 3(d)).

As mentioned above, it can be concluded that the  $\text{CH}_4 \nu_1$  peak could be the most suitable parameter for the pressure determination, although there is an abrupt shift in the  $\text{CH}_4$  peak positions due to the existence of the liquid-vapor transition in  $\text{CO}_2$ -rich mixtures. As shown in Figure 4, different laboratories have their own calibration curves for the  $\text{CH}_4$  peak position-pressure relationship. It can be ascribed to different measurement conditions, including excitation wavelength, the numerical aperture of the objective, the slit width, and the groove density of grating. However, all

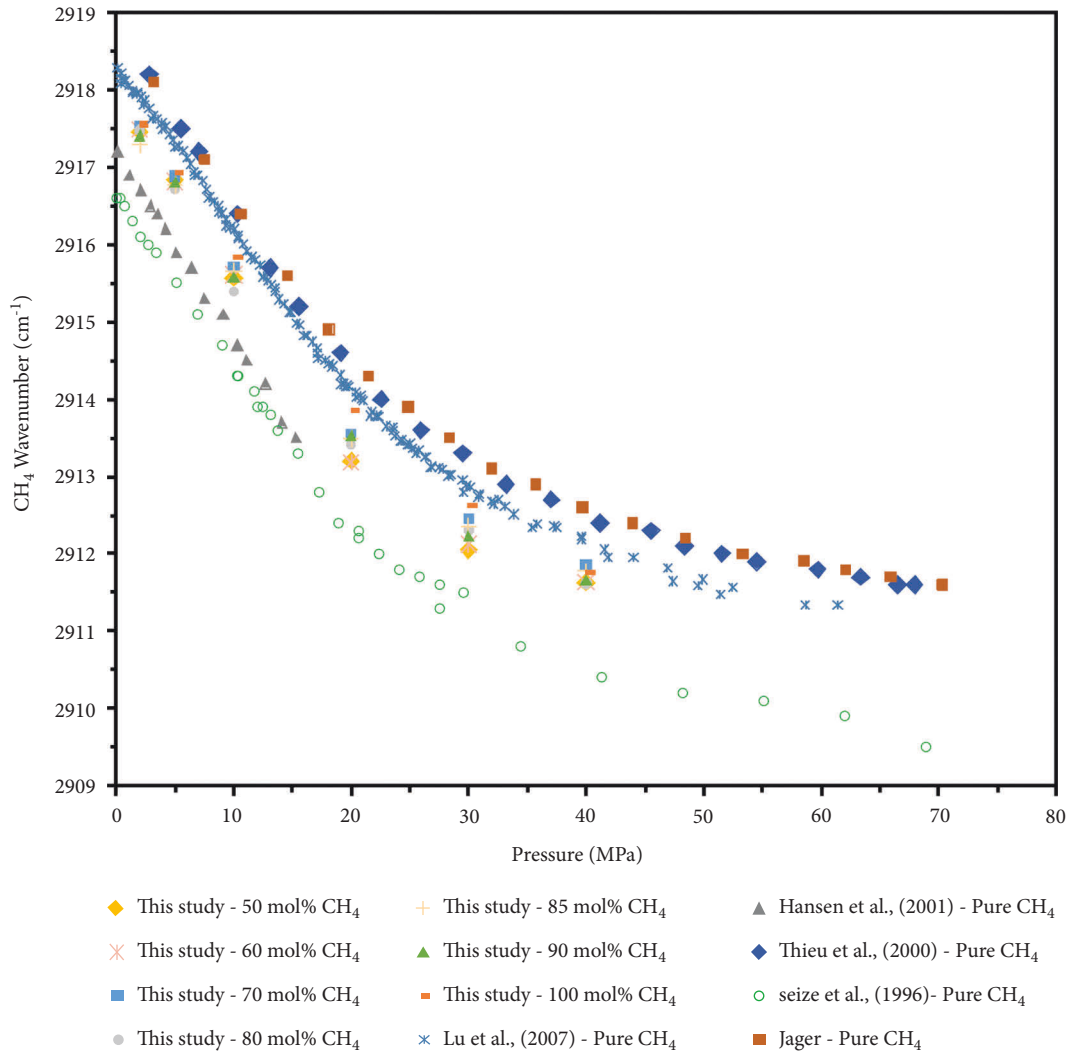


FIGURE 4: Raman peak position of C-H symmetric stretching band ( $\nu_1$ ) of  $\text{CH}_4$  from six available data sets as a function of pressure.

six data sets collected from the literature for pure  $\text{CH}_4$  and from this study for  $\text{CH}_4$ -rich mixture have similar relations between peak position and pressure, which are almost parallel to each other but with different intercepts at zero pressure, indicating that systematic discrepancies may exist among these data sets. Fortunately, these discrepancies can be eliminated by choosing  $D$  as a variable (Figure 5), where  $D = v_p - v_0$ , the difference between the peak positions at elevated pressure ( $v_p$ ) and near-zero pressure ( $v_0$ ) (in this study,  $v_0 = 2917.924 \text{ cm}^{-1}$  for  $<0.1 \text{ MPa}$ ). The same treatment was successfully demonstrated for pure  $\text{CH}_4$  many years ago by Lu et al. [29]. Our data ( $\text{CH}_4$ -rich mixtures with  $\text{CH}_4 \geq 50 \text{ mol}\%$  and pure  $\text{CH}_4$ ) at room temperature agree with previous studies [17, 27, 29, 51] for pure  $\text{CH}_4$  (Figure 5). After fitting the data obtained from the six sources, the relation between  $D$  (in  $\text{cm}^{-1}$ ) and pressure  $P$  (in MPa) can be described by the following equation:

$$D = 2.6953 \times 10^{-8} P^5 - 5.0314 \times 10^{-6} P^4 + 3.1297 \times 10^{-4} P^3 - 5.1191 \times 10^{-3} P^2 - 1.9752 \times 10^{-1} P, \quad (4)$$

with  $R^2$  (squared correlation coefficient) = 0.9972. The uncertainty ( $1\sigma$ ) in the calculated pressure was  $\pm 1.1 \text{ MPa}$ . This equation could be used to correlate  $P$  with  $D$  in other laboratories and calculate the pressure ( $\leq 70 \text{ MPa}$ ) of  $\text{CH}_4$ -rich mixtures ( $\text{CH}_4 \geq 50 \text{ mol}\%$ ) with the measured peak position of the  $\text{CH}_4 \nu_1$  band near room temperature, as long as  $v_0$  is known for the particular Raman spectrometer.

**3.2. Determination of the Density of  $\text{CO}_2$ - $\text{CH}_4$  Gas Mixtures.** The pressure of the  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures of known composition at room temperature ( $\sim 24^\circ\text{C}$ ) was converted to density using the GERG-2008 EoS (equation of state), which is a new wide-range EoS for natural gases and other mixtures of 21 natural gas components [52]. The variations in  $D$  and  $\Delta$  within  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures as a function of density and composition are plotted in Figure 6. The results show that  $D$  decreases with increasing  $\text{CH}_4$  concentration and density of the gas mixture, whereas  $\Delta$  increases with elevated  $\text{CO}_2$  concentration and density. The maximum value of the variation in  $\Delta$  is quite small ( $< \sim 3 \text{ cm}^{-1}$ ) compared to that for

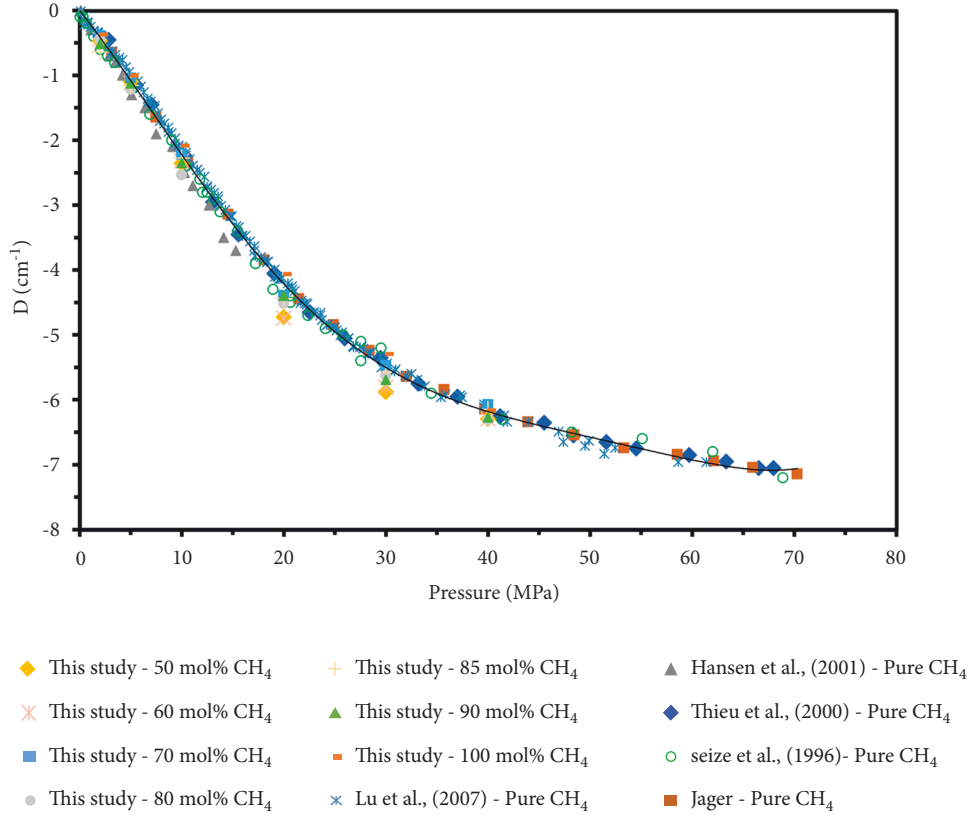


FIGURE 5: The P-D relationship determined in this study compared with other published values at near room temperature. Note that the original data of Jager's Colorado School of Mines are cited from Lu et al. [29].

$D$  ( $< \sim 6 \text{ cm}^{-1}$ ) over the same density range ( $0 \sim 1 \text{ g/cm}^3$ ). Thus, given an analytical uncertainty of  $\pm 0.01 \text{ cm}^{-1}$  for the measurement of the peak position,  $D$  will generally provide the most precise estimate of density. Therefore, the experimental data for  $D$  in  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures were fitted to the following equation by polynomial regression analysis (using the software 1stOpt):

$$\begin{aligned} \rho = & a_{00} + a_{10} \times C_{\text{CH}_4} + a_{01} \times D + a_{20} \times C_{\text{CH}_4}^2 \\ & + a_{11} \times C_{\text{CH}_4} \times D + a_{02} \times D^2 + a_{30} \times C_{\text{CH}_4}^3 \\ & + a_{21} \times C_{\text{CH}_4}^2 \times D + a_{12} \times C_{\text{CH}_4} \times D^2 + a_{03} \times D^3 \quad (5) \\ & + a_{31} \times C_{\text{CH}_4}^3 \times D + a_{22} \times C_{\text{CH}_4}^2 \times D^2 \\ & + a_{13} \times C_{\text{CH}_4} \times D^3 + a_{04} \times D^4. \end{aligned}$$

The fitted coefficients for equation (5) are listed in Table 1, and the root-mean-square error of equation (5) is  $0.012 \text{ g/cm}^3$  with an  $R^2$  of 0.997. This equation can be used to calculate the density of  $\text{CH}_4$ - $\text{CO}_2$  mixtures with the measured peak position of the  $\text{CH}_4 \nu_1$  band near room temperature.

**3.3. Determination of  $F(\text{CH}_4)/F(\text{CO}_2)$ .** The ratio of the peak areas of the  $\nu_1$  band of  $\text{CH}_4$  and the upper Fermi diad member of  $\text{CO}_2$  ( $A(\text{CH}_4)/A(\text{CO}_2)$ ) for  $\text{CO}_2$ - $\text{CH}_4$  mixtures increases with pressure below 10 MPa and then maintains the same levels at higher pressures (Table 2), indicating that

it is advantageous to use the ratios of the peak area of the  $\nu_1$  band of  $\text{CH}_4$  and the upper Fermi diad member of  $\text{CO}_2$  for quantitative determinations of the fluid composition in  $\text{CO}_2$ - $\text{CH}_4$  mixtures [17].

According to equation (1), if  $F(\text{CH}_4)/F(\text{CO}_2)$  has been determined by the known  $A(\text{CH}_4)/A(\text{CO}_2)$  and  $C(\text{CH}_4)/C(\text{CO}_2)$ , then it can be used to determine the molar ratio of  $\text{CH}_4$  to  $\text{CO}_2$  in an unknown  $\text{CO}_2$ - $\text{CH}_4$ -bearing fluid from the Raman peak area ratio.  $F(\text{CH}_4)/F(\text{CO}_2)$  in ten  $\text{CO}_2$ - $\text{CH}_4$  gaseous mixtures were calculated based on equation (1) and listed in Table 2. Within a given  $\text{CO}_2$ - $\text{CH}_4$  gaseous mixture, the  $F(\text{CH}_4)/F(\text{CO}_2)$  value rises at low pressures and remains constant (within analytical uncertainty) above 10 MPa (Table 2), which is in agreement with the previous work of Seitz et al. [17]. Therefore, the  $F(\text{CH}_4)/F(\text{CO}_2)$  values obtained at 20 MPa, 30 MPa, and 40 MPa were averaged and considered as the general  $F(\text{CH}_4)/F(\text{CO}_2)$  value of this mixture. It can be found that the general  $F(\text{CH}_4)/F(\text{CO}_2)$  value varies somewhat with different gaseous mixtures. Figure 7 shows a good linear relationship between the molar ratio ( $C(\text{CH}_4)/C(\text{CO}_2)$ ) and peak area ratio ( $A(\text{CH}_4)/A(\text{CO}_2)$ ) with an  $R^2$  of 0.9995. The slope shows a value of 5.048, which represents the  $F(\text{CH}_4)/F(\text{CO}_2)$  ratio based on equation (1). The uncertainty in the average values of  $F(\text{CH}_4)/F(\text{CO}_2)$  is  $\pm 0.4$ , which in turn can lead to an average error in the composition up to  $\pm 1.4 \text{ mol\% CH}_4$ .

Table 3 shows a comparison of our results with those obtained in six previous studies. The results show that  $F$

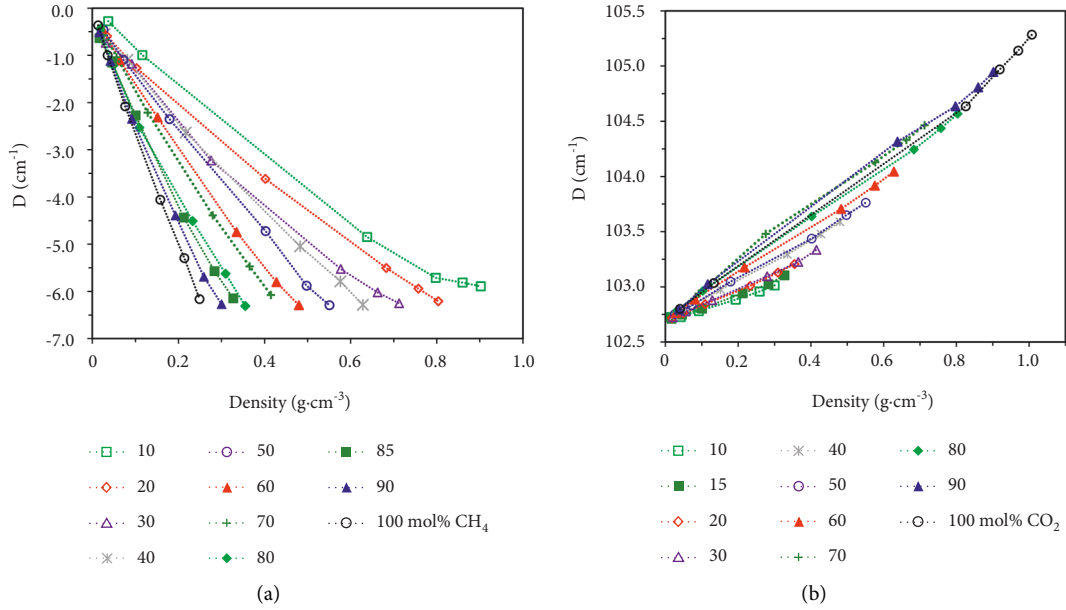


FIGURE 6: The variation in  $D$  (a) and  $\Delta$  (b) within  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures as a function of density and composition.  $D = \nu_p - \nu_0$  is the difference between the  $\nu_1$  peak positions of  $\text{CH}_4$  at elevated pressure ( $\nu_p$ ) and near-zero pressure ( $\nu_0$ );  $\Delta$  represents the separation between bands in the Fermi diad of  $\text{CO}_2$ .

TABLE 1: The fitting parameters for equation (5).

$a_{00}$	0.0458578972
$a_{10}$	0.0070465891
$a_{01}$	-0.0367435722
$a_{20}$	-0.3041825424
$a_{11}$	-0.0208569183
$a_{02}$	0.0264438359
$a_{30}$	0.2888501741
$a_{21}$	-0.0290891828
$a_{12}$	-0.0400853154
$a_{03}$	0.0014973436
$a_{31}$	0.0879585851
$a_{22}$	0.0270909951
$a_{13}$	0.0002772730
$a_{04}$	0.0000945387

$(\text{CH}_4)/F(\text{CO}_2)$  is quite different among different laboratories. Our value is in agreement with four previous results [17, 19, 20, 53]. Seitz et al. [17] stated that the variation in  $F(\text{CH}_4)/F(\text{CO}_2)$  is not significant within their analytical uncertainty, and the average value of  $F(\text{CH}_4)/F(\text{CO}_2)$  determined in five different  $\text{CO}_2$ - $\text{CH}_4$  mixtures at 20 MPa is  $5.16 \pm 0.18$ . Le et al. [19, 20] reported the  $F(\text{CH}_4)$  and  $F(\text{CO}_2)$  (relative to  $\text{N}_2$ ) to be  $7.73 \pm 0.16$  and  $1.40 \pm 0.03$ , respectively, yielding an  $F(\text{CH}_4)/F(\text{CO}_2)$  ratio of  $5.52 \pm 0.16$ . They also argued that the difference between different concentrations is not significant for the measured composition of the gas mixtures. Wopenka and Pasteris [44] obtained an average  $F(\text{CH}_4)/F(\text{CO}_2)$  of  $4.4 \pm 0.15$  in a variety of  $\text{CO}_2$ - $\text{CH}_4$ -bearing gas mixtures at low pressure ( $\leq 1.6$  MPa). Qiu et al. [32] determined a  $F(\text{CH}_4)$  of  $6.420 \pm 0.104$  and a  $F(\text{CO}_2)$  of  $1.690 \pm 0.042$  in a  $\text{CH}_4$ - $\text{CO}_2$ - $\text{CO}$ - $\text{N}_2$  mixture with a molar ratio of 1:1:1:2

from 10 MPa to 60 MPa, yielding an  $F(\text{CH}_4)/F(\text{CO}_2)$  ratio of  $3.80 \pm 0.10$ .

#### 4. Influence of Temperature on the Raman Spectral Parameters

To study the influence of temperature on the Raman spectra measured for  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures, HPOC was used to collect Raman spectra for five  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures (10 mol%  $\text{CH}_4$ , 30 mol%  $\text{CH}_4$ , 50 mol%  $\text{CH}_4$ , 70 mol%  $\text{CH}_4$ , and 90 mol%  $\text{CH}_4$ ) ranging from room temperature to 300°C and from 2 MPa to 40 MPa. The pressures and temperatures of the  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures of known composition were converted to density using the GERG-2008 EoS.

*4.1. Influence of Temperature on the Relationship between Peak Position and Pressure/Density for  $\text{CO}_2$ - $\text{CH}_4$  Gas Mixtures.* The peak positions of the  $\nu_1$  band of  $\text{CH}_4$  and Fermi diad splitting ( $\Delta$ ) of  $\text{CO}_2$  for pure  $\text{CH}_4$  or  $\text{CO}_2$  gas and their five  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures at various temperatures and pressures are listed in Table S3 (Supplemental material) and shown in Figures 8 and 9. As shown in Figure 8, as the temperature increases at constant pressure, the peak positions of the  $\nu_1$  band of  $\text{CH}_4$  and Fermi diad splitting ( $\Delta$ ) of  $\text{CO}_2$  for the five  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures shift to higher and lower wavenumbers, respectively. The results show that the effect of temperature on the peak positions of  $\text{CH}_4$  and  $\text{CO}_2$  in the mixed system has the same trend as demonstrated previously for the pure  $\text{CH}_4$  or  $\text{CO}_2$  system [30, 33, 37] and  $\text{N}_2$ - $\text{CO}_2$ - $\text{CH}_4$  system [43]. The pressure-induced change in the shift decreases at higher temperatures (Figure 8). Additionally, at room temperature, there is an abrupt shift in the measured position of  $\text{CH}_4$  and  $\text{CO}_2$  resulting from the



TABLE 2: The  $F(\text{CH}_4)/F(\text{CO}_2)$  in ten  $\text{CO}_2\text{-CH}_4$  gaseous mixtures at room temperature and pressures up to 40 MPa.

$P$ (MPa)	$X_{\text{CH}_4}$ (mol%)	${}^a A(\text{CH}_4)/A(\text{CO}_2)$	$F(\text{CH}_4)/F(\text{CO}_2)$
2		0.431	3.878
5		0.445	4.006
10		0.449	4.043
20	10	0.453	4.074
30		0.445	4.008
40		0.446	4.010
Avg.		0.448	4.031
2		0.994	3.978
5		1.037	4.147
10		1.078	4.314
20	20	1.055	4.219
30		1.059	4.235
40		1.063	4.250
Avg.		1.059	4.235
2		1.763	4.114
5		1.852	4.322
10		1.890	4.409
20	30	1.930	4.504
30		1.925	4.491
40		1.919	4.477
Avg.		1.925	4.491
2		2.873	4.309
5		2.976	4.463
10		3.021	4.531
20	40	3.021	4.531
30		3.027	4.540
40		3.044	4.565
Avg.		3.030	4.545
2		4.698	4.698
5		4.788	4.788
10		4.928	4.928
20	50	4.836	4.836
30		4.658	4.658
40		4.691	4.691
Avg.		4.728	4.728
2		6.643	4.429
5		6.919	4.613
10		7.072	4.714
20	60	7.165	4.777
30		7.176	4.784
40		7.198	4.799
Avg.		7.180	4.787
2		11.945	5.119
5		12.492	5.354
10		11.974	5.132
20	70	12.928	5.541
30		12.616	5.407
40		12.525	5.368
Avg.		12.690	5.430
2		18.170	4.542
5		18.541	4.635
10		18.808	4.702
20	80	19.427	4.857
30		19.652	4.913
40		19.516	4.879
Avg.		19.531	4.883

TABLE 2: Continued.

$P$ (MPa)	$X_{\text{CH}_4}$ (mol%)	${}^a A(\text{CH}_4)/A(\text{CO}_2)$	$F(\text{CH}_4)/F(\text{CO}_2)$
2		27.075	4.778
5		28.308	4.996
10		28.129	4.964
20	85	28.781	5.079
30		28.821	5.086
40		28.606	5.048
Avg.		28.736	5.071
2		39.812	4.424
5		41.345	4.594
10		42.173	4.686
20	90	46.208	5.134
30		45.496	5.055
40		44.950	4.994
Avg.		45.551	5.061

<sup>a</sup> $A$  is the Raman peak area.

phase transition, but this shift vanishes at higher temperatures (above 100°C). At the same pressure and temperature,  $\nu_{\text{CH}_4}$  and  $\Delta$  in the gas mixtures are shifted to lower wavenumbers compared with the values in the pure system (Figures 9(a) and 9(b)), which have also been reported in previous studies [17, 20]. However, in the  $\text{N}_2\text{-CO}_2\text{-CH}_4$  system,  $\nu_{\text{CH}_4}$  is located at higher wavenumbers relative to the pure system at the same  $P$ - $T$  conditions [43]. Figure 10 presents the variation in  $\nu_{\text{CH}_4}$  and  $\Delta$  as a function of density, composition, and temperature. The difference in density is notable at various temperatures. Taking 10 mol%  $\text{CH}_4$  at 40 MPa as an example, the density and Raman shift of  $\nu_{\text{CH}_4}$  at 24°C are 0.902 g/cm<sup>3</sup> and 2912.03 cm<sup>-1</sup>, respectively, and at 300°C, they are 0.338 g/cm<sup>3</sup> and 2914.0 cm<sup>-1</sup>, respectively. In conclusion, the effect of temperature must be considered when Raman shifts are used to calculate the pressure or density of  $\text{CO}_2\text{-CH}_4$  gas mixtures.

Effect of temperature on the variation of (a) the  $\nu_1$  band of  $\text{CH}_4$  and (b) Fermi diad splitting of  $\text{CO}_2$  ( $\Delta$ ) as a function of the density of  $\text{CO}_2\text{-CH}_4$  mixtures.

The relationship between the molar ratio and peak area ratio (average values above 10 MPa) is plotted at different temperatures in Figure 12 and listed in Table 4. This indicates that  $F(\text{CH}_4)/F(\text{CO}_2)$  increases with increasing temperature, and the peak area ratio obtained at high temperatures is not suitable for determining the composition at room temperature because it could lead to an overestimation of the molar ratio of  $\text{CH}_4/\text{CO}_2$ . Therefore, the temperature effect must be considered when  $\text{CH}_4$  to  $\text{CO}_2$  peak area ratios are used to calculate the fluid composition in  $\text{CO}_2\text{-CH}_4$  mixtures.

**4.2. Effect of Temperature on the Raman Peak Area for  $\text{CO}_2\text{-CH}_4$  Gas Mixtures.** The ratios of the peak area of the  $\nu_1$  band of  $\text{CH}_4$  and the upper Fermi diad of  $\text{CO}_2$  for the 50 mol%  $\text{CH}_4$ -50 mol%  $\text{CO}_2$  ( $C_{\text{CH}_4}/C_{\text{CO}_2} = 1:1$ ) mixture are plotted at different temperatures as a function of pressure in Figure 11.

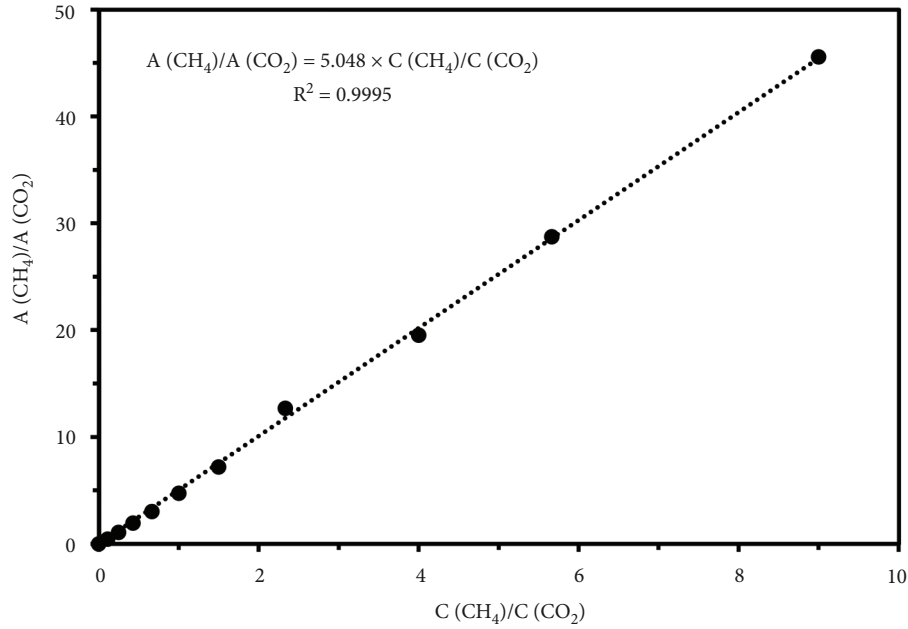


FIGURE 7: The linear relationship between the molar ratio ( $C(\text{CH}_4)/C(\text{CO}_2)$ ) and the Raman peak area ratio ( $A(\text{CH}_4)/A(\text{CO}_2)$ ) at room temperature.

TABLE 3: Comparison of  $F(\text{CH}_4)/F(\text{CO}_2)$  in this study with previous works.

	This study	Ref. [17]	Ref. [19, 20]	Ref. [32]	Ref. [44]	Ref. [53]	Ref. [54]
$\lambda$ (nm)	532.07	514.5	514.5	532	514.5	514.5	514
Pressure (MPa)	2–40	20	0.5–60	1–60	<1.6	0.1	5, 10
Temperature ( $^{\circ}\text{C}$ )	24	22	32	22	—	—	22
$C(\text{CH}_4)/C(\text{CO}_2)$	0.11~9	0.11~9	0.11~9	1	—	—	0.43~2.33
$F(\text{CH}_4)/F(\text{CO}_2)$	$5.05 \pm 0.4$	$5.16 \pm 0.18$	$5.52 \pm 0.16$	$3.8 \pm 0.1$	$4.40 \pm 0.15$	5.05	11.12

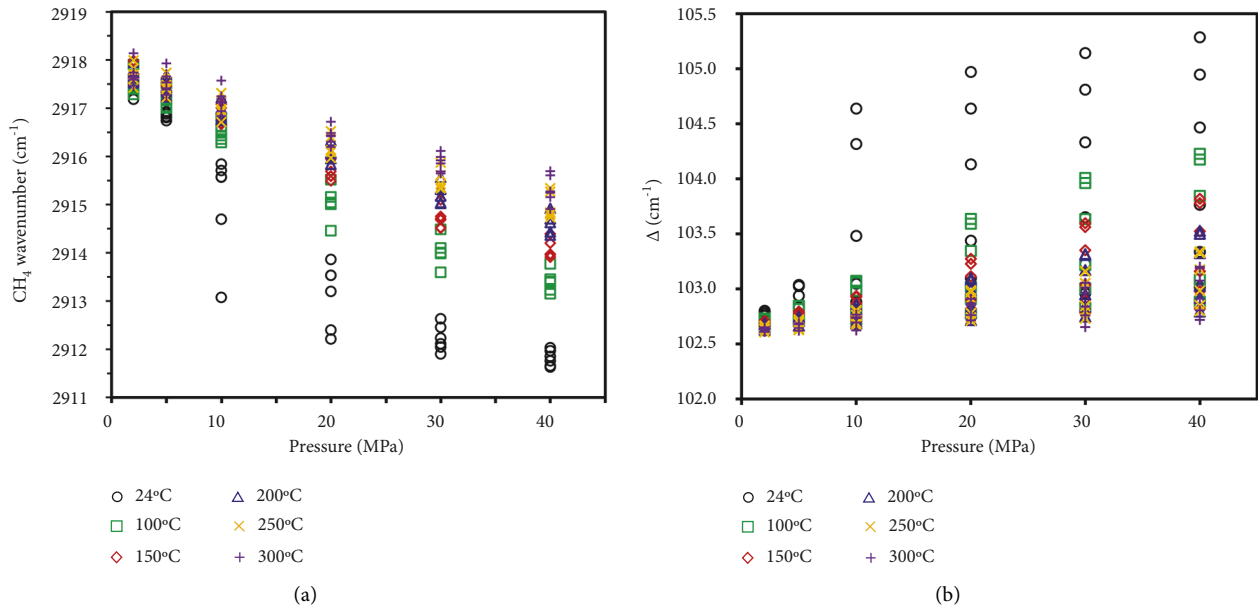


FIGURE 8: Effect of temperature on the Raman peak positions of (a) the  $\nu_1$  band of  $\text{CH}_4$  and (b) Fermi diad splitting of  $\text{CO}_2$  ( $\Delta$ ) for five  $\text{CO}_2$ - $\text{CH}_4$  gas mixtures as a function of pressure.

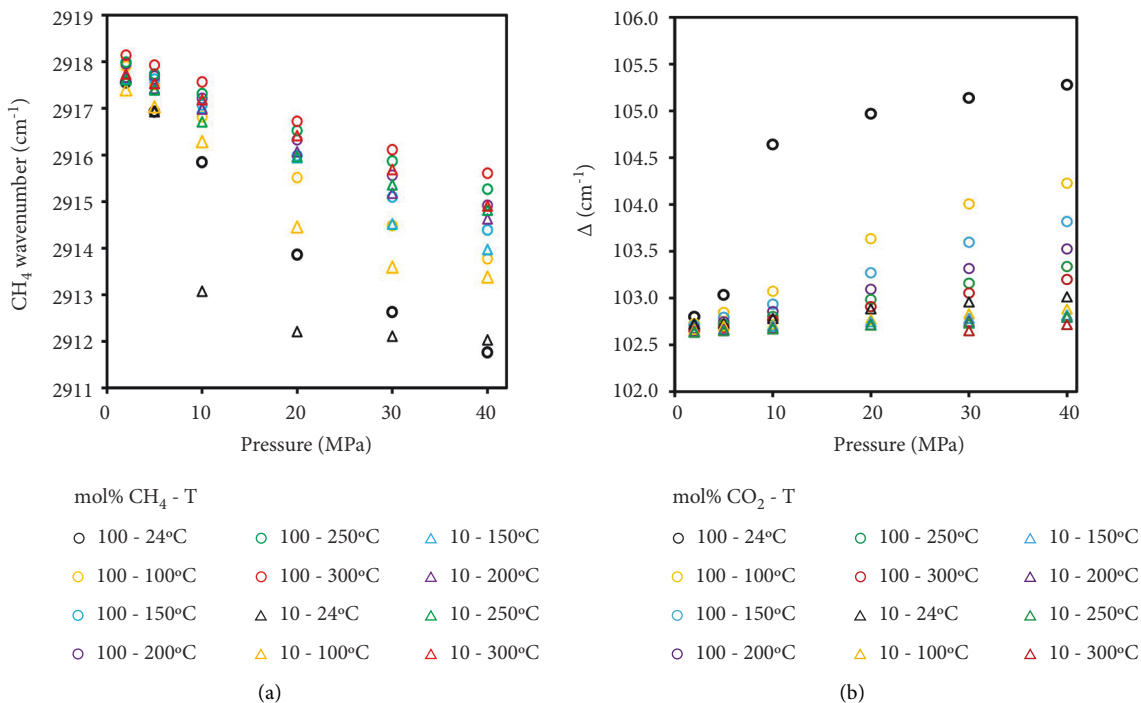


FIGURE 9: Effect of temperature on the Raman peak positions of (a) the  $\nu_1$  band of CH<sub>4</sub> for pure CH<sub>4</sub> and 10 mol% CH<sub>4</sub> and (b) Fermi diad splitting of CO<sub>2</sub> ( $\Delta$ ) for pure CO<sub>2</sub> and 10 mol% CH<sub>4</sub> as a function of pressure.

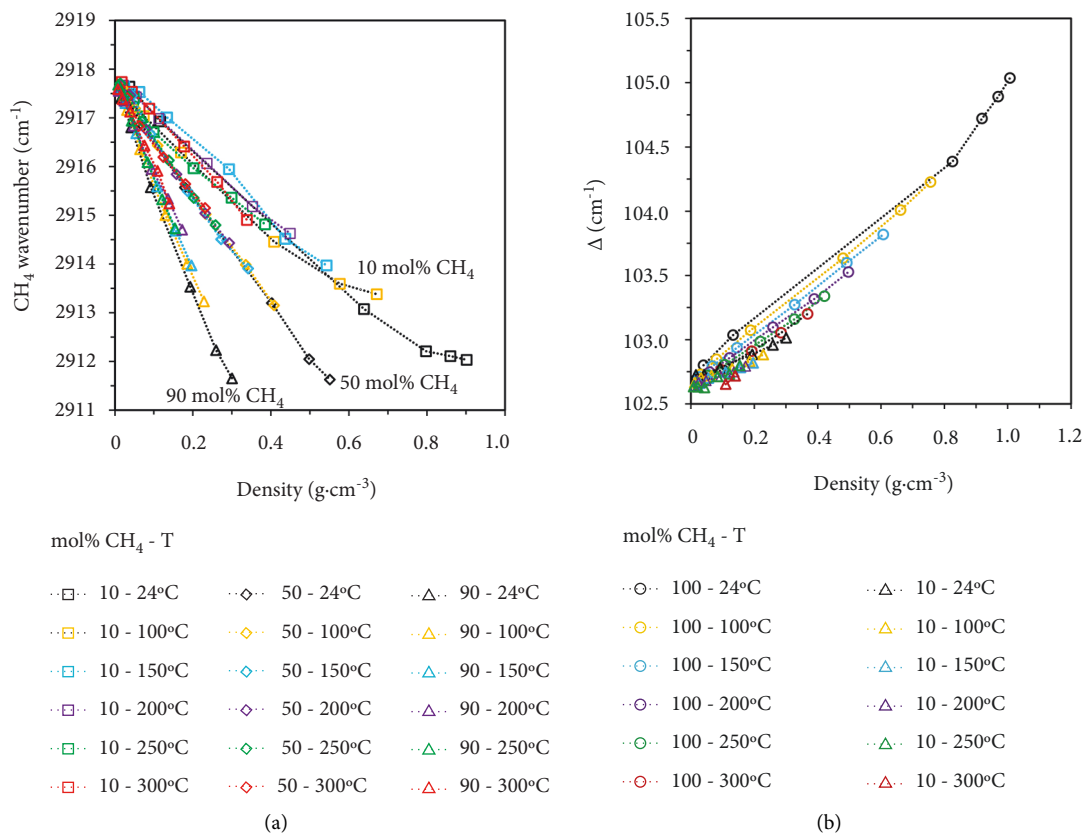


FIGURE 10: Effect of temperature on the variation of (a) the  $\nu_1$  band of CH<sub>4</sub> and (b) Fermi diad splitting of CO<sub>2</sub> ( $\Delta$ ) as a function of the density of CO<sub>2</sub>-CH<sub>4</sub> mixtures.

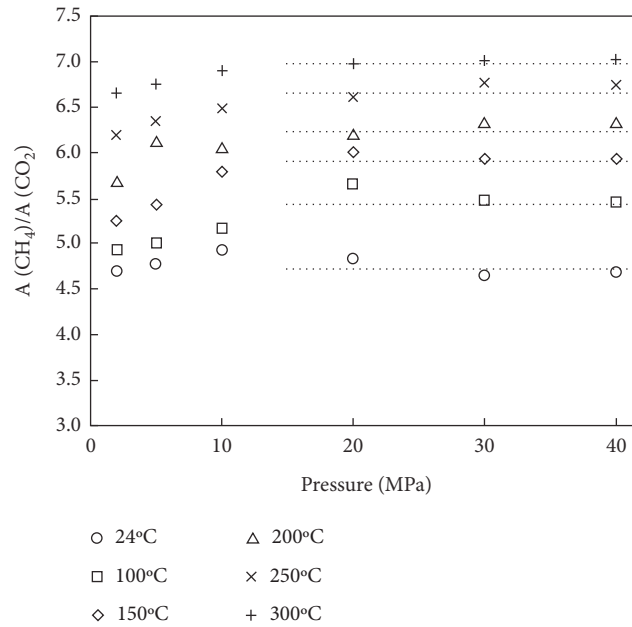


FIGURE 11: Relationship between the CH<sub>4</sub> to CO<sub>2</sub> peak area ratio and pressure at 50 mol% CH<sub>4</sub>-50 mol% CO<sub>2</sub> mixture and temperatures ranging from 24°C to 300°C.

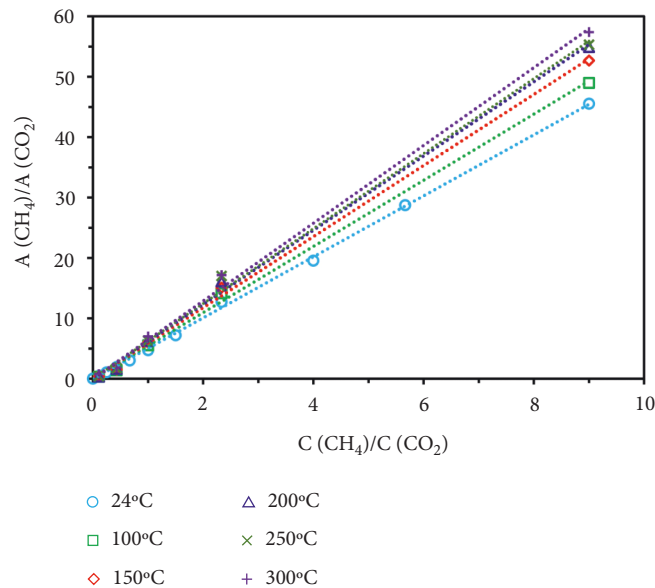


FIGURE 12: Raman peak area ratio between CH<sub>4</sub> and CO<sub>2</sub> as a function of their molar ratio at different temperatures.

TABLE 4: The fitting parameters for equation (1) at different temperatures.

Temperature (°C)	24	100	150	200	250	300
$F(\text{CH}_4)/F(\text{CO}_2)$	5.048	5.4770	5.8846	6.147	6.2119	6.4383
$R^2$	0.9995	0.9989	0.9990	0.9987	0.9977	0.9983

The results show that the  $A(\text{CH}_4)/A(\text{CO}_2)$  rises steeply with increasing pressure up to 10 MPa and then becomes approximately constant at higher pressures, ranging from 4.7 at 24°C to 7.0 at 300°C.

4.3. *Effect of Temperature on the Hot Band of CO<sub>2</sub> for CO<sub>2</sub>-CH<sub>4</sub> Gas Mixtures.* Earlier studies have confirmed that the intensity of CO<sub>2</sub> hot bands increases with increasing

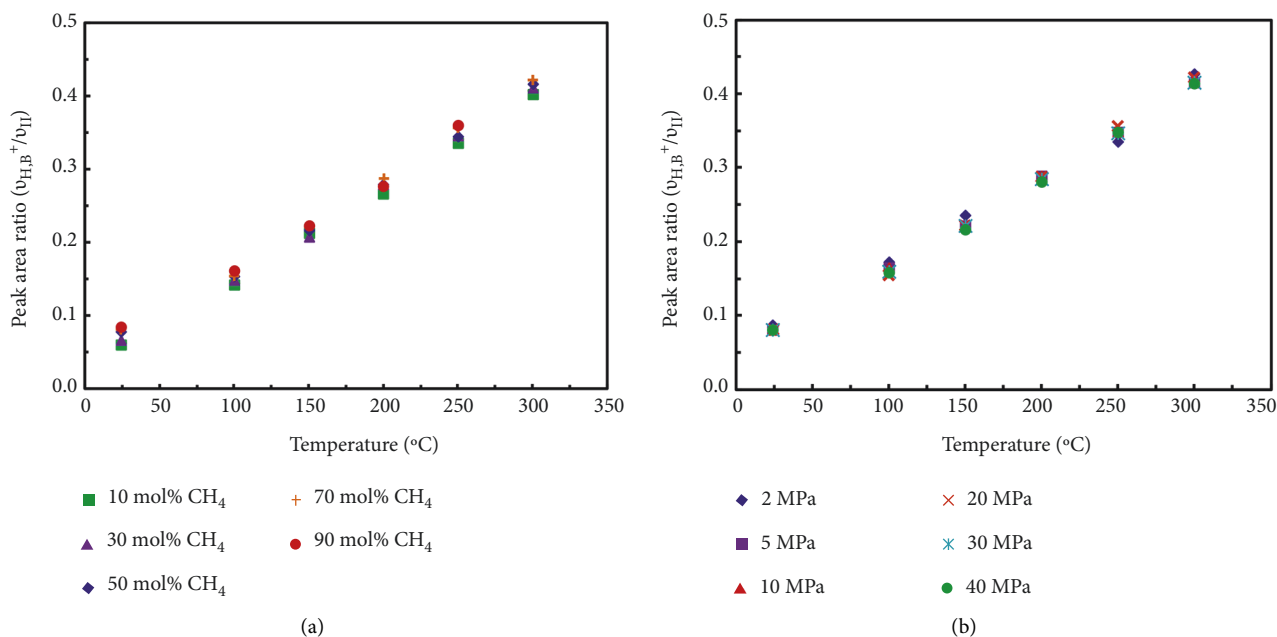


FIGURE 13: (a) Raman peak area ratio between the upper hot band and the upper Fermi domain versus temperature at 20 MPa for different CO<sub>2</sub>-CH<sub>4</sub> mixtures. (b) Raman peak area ratio between the upper hot band and the upper Fermi domain versus temperature at 70% CH<sub>4</sub>-30% CO<sub>2</sub> and different pressures.

temperature [37, 55, 56] and assessed the possibility of using the CO<sub>2</sub> hot band as a thermometer [57, 58]. In this study, we examined the relationship between the peak area ratio between the upper hot band ( $\nu_{H,B^+}$ ) and the upper Fermi diad member ( $\nu_{II}$ ) with temperature, pressure, and composition. Figure 13 shows that the peak area ratio between the upper hot band and the upper Fermi domain is a sensitive function of temperature, and almost independent of composition and pressure. This indicates that the CO<sub>2</sub> hot band can be used as a thermometer in the CO<sub>2</sub>-CH<sub>4</sub> gas mixtures.

## 5. Application to FSCC Samples

To verify the accuracy of the Raman quantitative analysis methods described above, six synthetic FIs in FSCC samples with unknown relative molar proportions of the components, total pressure, and density were prepared, and the results of microthermometric measurements for these samples were compared with laser Raman (LRM) spectroscopic data. Each FSCC sample was analyzed by MT to determine the melting temperature  $T_m$  and the homogenization temperature  $T_h$ , and they were used to calculate the composition and density of the fluid inclusion based on the VX diagram of Thiery et al. [18]. The total pressures were determined from MT using the GERG-2008 EoS. The selected FSCC samples were also analyzed three times with Raman spectroscopy. The average values of the CH<sub>4</sub> to CO<sub>2</sub> peak area ratio ( $A(\text{CH}_4)/A(\text{CO}_2)$ ) and peak position of CH<sub>4</sub> ( $\nu_{\text{CH}_4}$ ) were used to calculate the composition, pressure, and density through equations (1), (4), and (5). A comparison of the results obtained by LRM and MT is listed in Table 5 and plotted in Figure 14. The uncertainty for  $1\sigma$  is also provided,

and the procedure used for calculating the uncertainty propagation is the same as that reported by Le et al. [19].

For the LRM data, average errors in the values of mol% CH<sub>4</sub>, pressure, and density, based on the differences between these three sets of analyses on the same FSCC samples and the best-fitted equations (1), (4), and (5), are  $\pm 1.2$  mol%,  $\pm 0.9$  MPa, and  $\pm 0.011$  g·cm<sup>-3</sup>, respectively. For the MT data, the uncertainty in  $\pm 0.1^\circ\text{C}$  of the heating-cooling stage for  $T_m$  and  $T_h$  and an error in graphic reading arising from the VX diagram can result in average errors of  $\pm 1.0$  mol% in composition and  $\pm 0.008$  g/cm<sup>3</sup> in density, which can lead to an average error of  $\pm 1.1$  MPa in pressure, as determined from the GERG-2008 EoS.

Overall, for the same FSCC samples, there is a very good agreement between the mol% CH<sub>4</sub> derived from LRM and MT analyses (Table 5, Figure 14(a)). The largest difference in mol% CH<sub>4</sub> is 2.1 mol%. The densities derived from LRM analysis are slightly smaller than those obtained from MT analysis, and the difference is within 0.062 g/cm<sup>3</sup> (Figure 14(c)). However, there is a large deviation between the pressure derived from LRM analysis and MT analysis (Figure 14(b)). To investigate the reason for the large deviation, the pressure ( $P_{\text{call}}$ ) calculated from  $X_{\text{Raman}}$  and  $\rho_{\text{Raman}}$  and the density ( $\rho_{\text{call}}$ ) calculated from  $X_{\text{Raman}}$  and  $P_{\text{Raman}}$  using the GERG-2008 EoS are also obtained (Table 5, Figures 14(b) and 14(c)), and they are very close to the values for  $P_{\text{Raman}}$  and  $\rho_{\text{Raman}}$  obtained directly using equations (4) and (5), respectively. The results show that our calibration equations could provide reasonable estimates for the composition, pressure, and density of fluid inclusions. Therefore, it is reasonable to believe that this deviation is derived from the deviation of density and composition between LRM analysis and MT analysis. Incredibly, the slight

TABLE 5: Comparison of the results obtained from Raman spectroscopy and microthermometry.

No.	$A$ ( $\text{CH}_4$ )/ $A$ ( $\text{CO}_2$ )	$\nu_{\text{CH}_4}$ ( $\text{cm}^{-1}$ )	Raman results <sup>a</sup>			Microthermometry results <sup>b</sup>						Difference <sup>c</sup>			
			$X_{\text{Raman}}$ mol% $\text{CH}_4$	$P_{\text{Raman}}$ (MPa)	$\rho_{\text{Raman}}$ ( $\text{g}\cdot\text{cm}^{-3}$ )	$P_{\text{cal1}}$ (MPa)	$\rho_{\text{cal1}}$ ( $\text{g}\cdot\text{cm}^{-3}$ )	$T_m$ ( $^{\circ}\text{C}$ )	$T_h$ ( $^{\circ}\text{C}$ )	$X_{\text{Microth}}$ mol% $\text{CH}_4$	$\rho_{\text{Microth}}$ ( $\text{g}\cdot\text{cm}^{-3}$ )	$P_{\text{cal2}}$ (MPa)	$\Delta X$ mol%	$\Delta P$ (MPa)	$\Delta\rho$ ( $\text{g}\cdot\text{cm}^{-3}$ )
S.1	31.821 $\pm 0.765$	2911.342 $\pm 0.006$	86.6 $\pm 0.1$	50.2 $\pm 0.8$	0.345 $\pm 0.003$	48.8	0.349	-79.5	-89.1	87.5 $\pm 0.8$	0.364 $\pm 0.005$	59.6 $\pm 1.1$	-0.9	-9.4	-0.019
S.2	24.901 $\pm 0.053$	2911.644 $\pm 0.029$	83.1 $\pm 0.5$	42.2 $\pm 1.2$	0.342 $\pm 0.006$	41.3	0.345	-75.3	-83.1	82.5 $\pm 1.2$	0.398 $\pm 0.006$	61.1 $\pm 0.8$	0.6	-18.9	-0.056
S.3	12.214 $\pm 0.487$	2913.341 $\pm 0.023$	70.8 $\pm 1.5$	22.4 $\pm 0.8$	0.284 $\pm 0.013$	20.7	0.301	-67.4	-44.6	70.7 $\pm 0.8$	0.286 $\pm 0.001$	20.9 $\pm 0.1$	0.1	1.5	-0.002
S.4	4.907 $\pm 0.036$	2911.464 $\pm 0.008$	49.3 $\pm 1.3$	46.8 $\pm 0.9$	0.569 $\pm 0.012$	43.2	0.582	-65.6	-59.8	49.0 $\pm 1.3$	0.631 $\pm 0.010$	62.8 $\pm 0.3$	0.3	-16.0	-0.062
S.5	3.573 $\pm 0.024$	2911.469 $\pm 0.008$	41.4 $\pm 1.9$	46.7 $\pm 0.8$	0.636 $\pm 0.017$	44.8	0.643	-64.8	-64.2	43.5 $\pm 1.1$	0.677 $\pm 0.012$	64.3 $\pm 3.0$	-2.1	-17.6	-0.041
S.6	1.495 $\pm 0.0011$	2911.811 $\pm 0.004$	22.8 $\pm 2.1$	38.6 $\pm 0.7$	0.766 $\pm 0.017$	37.3	0.772	-61.7	-24.8	20.9 $\pm 1.0$	0.795 $\pm 0.011$	39.9 $\pm 1.4$	1.9	-1.3	-0.029

<sup>a</sup> $X_{\text{Raman}}$ ,  $P_{\text{Raman}}$ , and  $\rho_{\text{Raman}}$  are determined from equations (1), (4), and (5) based on  $A$  ( $\text{CH}_4$ )/ $A$  ( $\text{CO}_2$ ) and  $\nu_{\text{CH}_4}$ ;  $P_{\text{cal1}}$  is the pressure calculated from  $X_{\text{Raman}}$  and  $\rho_{\text{Raman}}$  using the GERG-2008 EoS;  $\rho_{\text{cal1}}$  is the density calculated from  $X_{\text{Raman}}$  and  $P_{\text{Raman}}$  using the GERG-2008 EoS. <sup>b</sup> $X_{\text{Microth}}$  and  $\rho_{\text{Microth}}$  are obtained from the VX diagram of Thiery et al. [18] based on  $T_m$  and  $T_h$ ;  $P_{\text{cal2}}$  is the pressure calculated from  $X_{\text{Microth}}$  and  $\rho_{\text{Microth}}$  using the GERG-2008 EoS. <sup>c</sup> $\Delta X = X_{\text{Raman}} - X_{\text{Microth}}$ ;  $\Delta P = P_{\text{Raman}} - P_{\text{cal2}}$ ;  $\Delta\rho = \rho_{\text{Raman}} - \rho_{\text{Microth}}$ . Note that all data are measured at room temperature ( $\sim 24^{\circ}\text{C}$ ), and the uncertainty is provided for  $1\sigma$ .

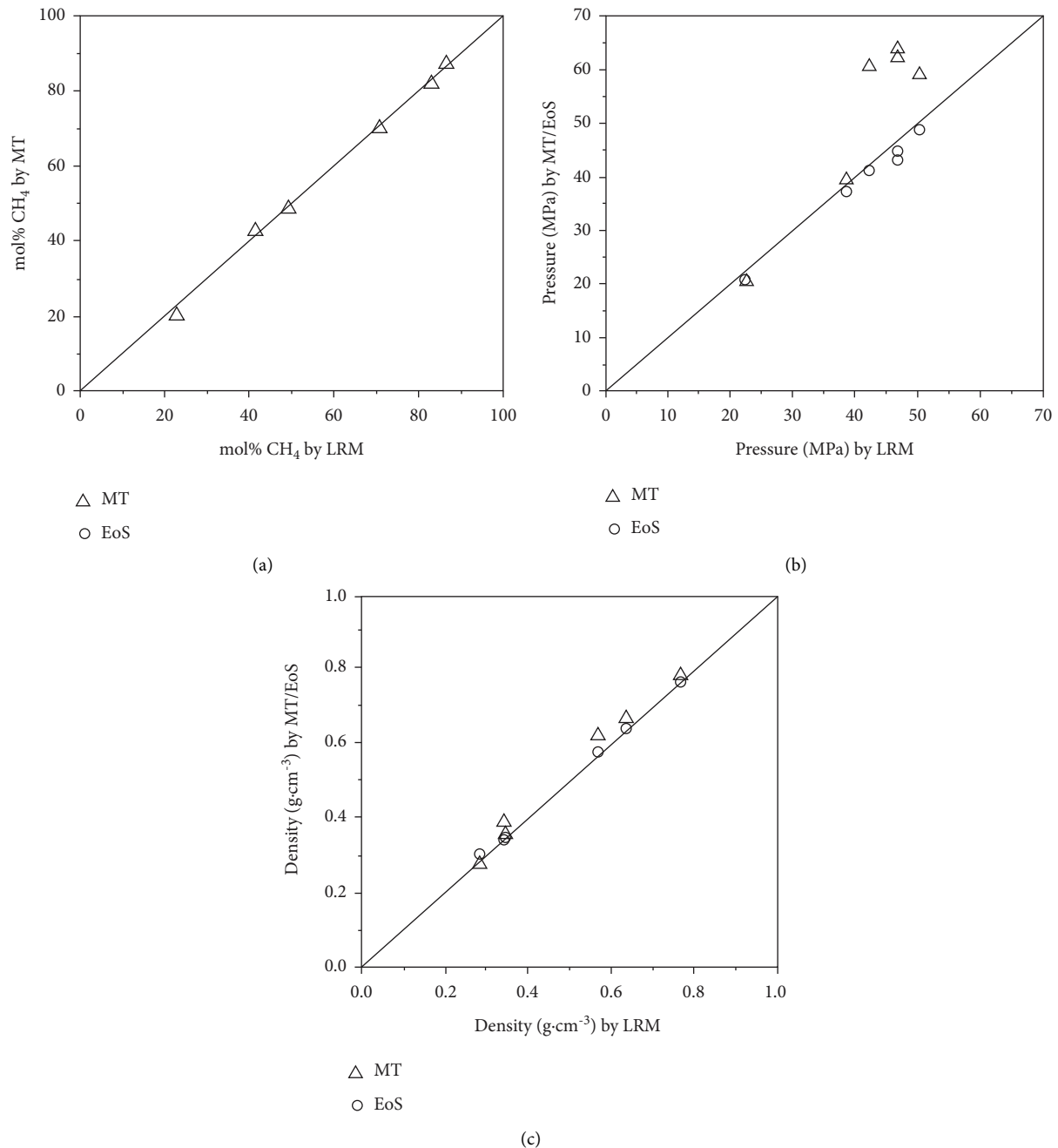


FIGURE 14: Comparison of (a) mol% CH<sub>4</sub>, (b) pressure, and (c) density of CO<sub>2</sub>-CH<sub>4</sub> mixtures in FSCCs based on laser Raman (LRM) and microthermometry (MT) data. Note that the pressure (b) and density (c) calculated by GERG-2008 EoS are based on  $X_{\text{Raman}}$  and  $\rho_{\text{Raman}}$ , and  $X_{\text{Raman}}$  and  $P_{\text{Raman}}$ , respectively.

deviation in density and composition can lead to such a large deviation in pressure. For example, a small difference between the composition of 0.6 mol% CH<sub>4</sub> and the density of 0.056 g/cm<sup>3</sup> can lead to a pressure deviation of up to 18.9 MPa (Table 5, Figure 14(b)). Although the uncertainty in the MT analysis is slightly better than that derived from the LRM analysis in this work, Raman spectroscopy is still a powerful tool for fluid inclusion investigations. In addition, for the analyses of natural fluid inclusions, MT may not be

more efficient than LRM analysis in some cases, such as small inclusions (<5 μm) and inclusions with poor optical properties (color, contrast, etc.) or the formation of clathrates [20]. It is worth noting that when analyzing natural fluid inclusions in anisotropic host minerals by using Raman spectroscopy, it must be placed at crystal extinction position to eliminate the effect of birefringence of the host mineral on the composition of the gas phase of fluid inclusions, as suggested by Caumon et al. [59].

## 6. Conclusion

This study investigated the Raman spectral parameters (peak position and peak area) for CO<sub>2</sub>-CH<sub>4</sub> mixtures at pressures up to 40 MPa and temperatures up to 300°C. The variation in each spectral parameter as a function of fluid pressure, density, temperature, and composition was discussed. The peak position of the  $\nu_1$  band for CH<sub>4</sub> can be used to determine the pressure of pure CH<sub>4</sub> and CH<sub>4</sub>-dominated fluids (>50 mol% CH<sub>4</sub>). The relationships for the peak position of the  $\nu_1$  band for CH<sub>4</sub>, density, and composition were established and used to calculate the densities for CO<sub>2</sub>-CH<sub>4</sub> mixtures. The Raman quantification factor  $F(\text{CH}_4)/F(\text{CO}_2)$  was found to remain almost constant ( $5.048 \pm 0.4$ ) with varying pressure/density and composition, and, therefore, was used for the determination of the CH<sub>4</sub> to CO<sub>2</sub> molar ratio for fluids with high internal pressures (>10 MPa) based on the Raman peak area ratio. The effects of temperature on the variations in Raman spectral parameters were first investigated at temperatures up to 300°C, and these effects should not be ignored when Raman spectral parameters are used to calculate the pressure, density, and composition of CO<sub>2</sub>-CH<sub>4</sub> gas mixtures. The results of our Raman spectroscopic measurements for six FSCC samples agree very well with those derived from MT analyses and, thus, verify our methodology for the quantitative determination of composition, pressure, and density of CO<sub>2</sub>-CH<sub>4</sub> mixtures.

## Data Availability

The data used to support the findings of this study are included within the supplementary materials.

## Conflicts of Interest

The authors declare that there are no conflicts of interest.

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## Supplementary Materials

Table S1. Calibrated Raman peak positions (cm<sup>-1</sup>) of the CH<sub>4</sub> $\nu_1$  peak for pure CH<sub>4</sub> gas and ten CO<sub>2</sub>-CH<sub>4</sub> mixtures at room temperature (~24°C) and various pressures. The uncertainties were derived from the standard deviation of the three consecutive measurements of spectra at each *P-T* condition. Table S2. Calibrated Raman peak positions (cm<sup>-1</sup>) of the lower and upper member of the CO<sub>2</sub> Fermi diad and the Fermi diad splits ( $\Delta$ ) of pure CO<sub>2</sub> gas and ten CO<sub>2</sub>-CH<sub>4</sub>

mixtures at room temperature (~24°C) and various pressures. The uncertainties were derived from the standard deviation of the three consecutive measurements of spectra at each *P-T* condition. Table S3. Calibrated Raman peak positions (cm<sup>-1</sup>) of CH<sub>4</sub> $\nu_1$  peak and Fermi diad splits ( $\Delta$ ) of CO<sub>2</sub> for pure CH<sub>4</sub> or CO<sub>2</sub> gas and their five CO<sub>2</sub>-CH<sub>4</sub> mixtures at various temperatures and pressures. The uncertainties were derived from the standard deviation of the three consecutive measurements of spectra at each *P-T* condition. (*Supplementary Materials*)

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