

# A MARKOV TIME RELATED TO A PRIORITY SYSTEM

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We consider a basic renewable duplex system characterized by cold standby and subjected to a priority rule. Apart from a general stochastic analysis presented in the previous literature, we introduce a Markov time called the recovery time of the system. In order to obtain the corresponding Laplace-Stieltjes transform, we employ a stochastic process endowed with transition measures satisfying generalized coupled differential equations. The solution is provided by the theory of sectionally holomorphic functions.

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## 1. Introduction

*Standby* provides a powerful tool to enhance the reliability, availability, quality, and safety of operational plants, for example, [4, 8, 16]. However, in practice, standby systems are often subjected to a *priority* rule. For instance, the *external* power supply station of a technical plant has usually overall priority in operation with regard to an *internal* (local) power generator kept in cold or warm standby, for example, [4]. The local generator is only deployed if the external power station is down.

Cold or warm standby systems subjected to a priority rule and attended by a repair facility have received considerable attention in the previous literature, for example, [2, 3, 5, 6, 9–15, 17–19, 21, 23, 24]. We consider a basic duplex system composed of a priority unit (the **p**-unit) and a nonpriority unit (the **n**-unit) kept in cold standby until the **p**-unit fails. The **p**-unit has overall (break-in) priority in operation with regard to the **n**-unit, that is, the **n**-unit is only deployed if the **p**-unit is down. In order to avoid undesirable delays in repairing failed units, we suppose that the entire system (henceforth called the **T**-system) is attended by *two* different repairmen. Each repairman has his own particular task. Repairman *N* is skilled in repairing the **n**-unit, whereas repairman *P* is an expert in repairing the **p**-unit. Both repairmen are jointly busy if, and only if, both units are down. Otherwise, at least one repairman is idle. Figure 1.1 displays a functional block diagram of the **T**-system operating in standby.

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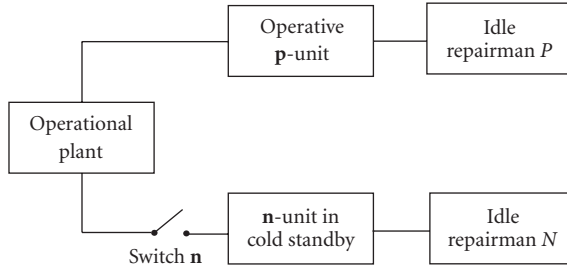


Figure 1.1. Functional block diagram of the T-system operating in standby.

Apart from a general stochastic analysis presented in the previous literature [19, 21, 23], we introduce a Markov time called the *recovery* time of the T-system. The recovery time is the total (random) time needed to restore the T-system from a prescribed risky state into the safe state (see the forthcoming formulation). In order to obtain the corresponding Laplace-Stieltjes transform, we employ a stochastic process endowed with transition measures satisfying *generalized* coupled partial differential equations. Our proposed *transient* equations are extending the *steady-state* equations presented by Vanderperre and Makhanov [23]. The explicit solution is provided by a refined application of the theory of sectionally holomorphic functions.

### 2. Formulation

Consider the basic T-system satisfying the following conditions. The **p**-unit has a constant failure rate  $\lambda > 0$  and a general repair time distribution  $R(\cdot)$ ,  $R(0) = 0$ . The corresponding failure-free time and repair time are denoted by  $f$  and  $r$ . The *operative n*-unit has a constant failure rate  $\lambda_s > 0$ , but a *zero* failure rate in standby (the so-called *cold* standby) and a general repair time distribution  $R_s(\cdot)$ ,  $R_s(0) = 0$ . The corresponding failure-free time and repair time are denoted by  $f_s$  and  $r_s$ . The random variables  $f$ ,  $f_s$ ,  $r$ ,  $r_s$  are statistically independent. Any repair is perfect [7]. The switch-over time from standby to the operative state is instantaneous. Characteristic functions are formulated in terms of a *complex* transform variable. For instance,

$$\mathbf{E}e^{i\omega r} = \int_0^{\infty} e^{i\omega x} dR(x), \quad \text{Im } \omega \geq 0. \quad (2.1)$$

Note that

$$\mathbf{E}e^{-i\omega r} = \int_{-\infty}^0 e^{i\omega x} d\{1 - R((-x) -)\}, \quad \text{Im } \omega \leq 0. \quad (2.2)$$

The corresponding Fourier-Stieltjes transforms are called *dual* transforms. Without loss of generality (cf. [21, page 361]) we may assume that both  $R$  and  $R_s$  have bounded density functions (in the Radon-Nikodym sense) defined on  $[0, \infty)$ .

In order to analyse the random behaviour of the **T**-system, we employ a stochastic process  $\{N_t, t \geq 0\}$  with arbitrary discrete state space  $\{A, B, C, D\} \subset [0, \infty)$  characterized by the following events:

- (i)  $\{N_t = A\}$ : “the **p**-unit is operative and the **n**-unit is in cold standby at time  $t$ ,”
- (ii)  $\{N_t = B\}$ : “the **n**-unit is operative and the **p**-unit is under repair at time  $t$ ,”
- (iii)  $\{N_t = C\}$ : “the **p**-unit is operative and the **n**-unit is under repair at time  $t$ ,”
- (iv)  $\{N_t = D\}$ : “both units are down at time  $t$ .”

State  $A$  is called the safe state. States  $B$  and  $C$  are called risky states and state  $D$  is called the system down state. The non-Markovian process  $\{N_t\}$  is defined on a filtered probability space  $\{\Omega, \mathbf{B}, \mathbf{P}, F\}$  where the *history*  $F := \{F_t, t \geq 0\}$  satisfies the Dellacherie conditions:

- (i)  $F_0$  contains the **P**-null sets of **B**;
- (ii) for all  $t \geq 0$ ,  $F_t = \bigcap_{u>t} F_u$ , that is, the family  $F$  is right continuous.

Consider the  $F$ -Markov time

$$\theta := \inf \{t > 0 : N_t = A \mid N_0 = B, Z_0 = 0\}, \quad (2.3)$$

where  $Z_t$  denotes the past repair time of the failed **p**-unit being under progressive repair at time  $t$ . Note that we take the instant of the first failure as time origin, that is,  $N_0 = B$ ,  $Z_0 = 0$ , **P**-a.s. Thus, from  $t = 0$  onwards,  $\theta$  is the total amount of time needed to restore the **T**-system from the risky state  $B$  into the safe state  $A$ .  $\theta$  is called the *recovery* time of the **T**-system. In addition, note that our priority rule implies that a transition from the safe state  $A$  into the risky state  $C$  is only possible via state  $D$ .

A (vector) Markov characterization of the process  $\{N_t, t \geq 0\}$  is piecewise and conditionally defined by

- (i)  $\{N_t\}$ , if  $N_t = A$  (i.e., if the event  $\{N_t = A\}$  occurs);
- (ii)  $\{(N_t, X_t)\}$ , if  $N_t = B$ , where  $X_t$  denotes the *remaining* repair time of the **p**-unit being under progressive repair at time  $t$ ;
- (iii)  $\{(N_t, Y_t)\}$ , if  $N_t = C$ , where  $Y_t$  denotes the *remaining* repair time of the **n**-unit being under progressive repair at time  $t$ ;
- (iv)  $\{(N_t, X_t, Y_t)\}$ , if  $N_t = D$ .

The state space of the underlying Markov process with *absorbing* state  $A$  is given by

$$\{A\} \cup \{(B, x); x \geq 0\} \cup \{(C, y); y \geq 0\} \cup \{(D, x, y); x \geq 0, y \geq 0\}. \quad (2.4)$$

Let

$$p_A(t) := \mathbf{P}\{N_t = A\}, \quad t \geq 0. \quad (2.5)$$

Finally, we introduce the measures:

$$\begin{aligned} p_B(t, x) dx &:= \mathbf{P}\{N_t = B, X_t \in dx\}, \\ p_C(t, y) dy &:= \mathbf{P}\{N_t = C, Y_t \in dy\}, \\ p_D(t, x, y) dx, dy &:= \mathbf{P}\{N_t = D, X_t \in dx, Y_t \in dy\}. \end{aligned} \quad (2.6)$$

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Note that, for instance,

$$p_D(t) = \int_0^\infty \int_0^\infty d_x d_y \mathbf{P}\{N_t = D, X_t \leq x, Y_t \leq y\} = \int_0^\infty \int_0^\infty p_D(t, x, y) dx dy. \quad (2.7)$$

*Notations 2.1.* The real line and the complex plane are denoted by  $\mathbb{R}$  and  $\mathbb{C}$ , with obvious superscript notations such as  $\mathbb{C}^+$ ,  $\mathbb{C}^-$ . For instance,  $\mathbb{C}^+ := \{\omega \in \mathbb{C} : \text{Im } \omega > 0\}$ .

The indicator (function) of an event  $\varepsilon \in \mathbf{B}$  is denoted by  $1_\varepsilon$ . The Heaviside unit-step function with the unit jump at  $t = t_0 > 0$  is denoted by  $U_{t_0}(t)$ ,  $t \geq 0$ . Finally, let  $[t]$  be the greatest integer function.

### 3. Differential equations

Applying Hokstad's supplementary variable technique in some time interval  $[t, t + \Delta]$ ,  $\Delta \downarrow 0$ , for example, Alfa and Srinivasa [1] and taking the absorbing state  $A$  into account, yields

$$\begin{aligned} p_A(t + \Delta) &= p_A(t) + p_B(t, 0)\Delta + p_C(t, 0)\Delta + o(\Delta), \\ p_B(t + \Delta, x - \Delta) &= p_B(t, x)(1 - \lambda_s \Delta) + p_D(t, x, 0)\Delta + o(\Delta), \\ p_C(t + \Delta, y - \Delta) &= p_C(t, y)(1 - \lambda \Delta) + p_D(t, 0, y)\Delta + o(\Delta), \\ p_D(t + \Delta, x - \Delta, y - \Delta) &= p_D(t, x, y) + \lambda_s p_B(t, x) \frac{d}{dy} R_s(y) \Delta + \lambda p_C(t, y) \frac{d}{dx} R(x) \Delta + o(\Delta). \end{aligned} \quad (3.1)$$

Invoking the definition of directional derivative entails that for  $t > 0$ ,  $x > 0$ ,  $y > 0$ ,

$$\begin{aligned} \frac{d}{dt} p_A(t) &= p_B(t, 0) + p_C(t, 0), \\ \left( \lambda_s + \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) p_B(t, x) &= p_D(t, x, 0), \\ \left( \lambda + \frac{\partial}{\partial t} - \frac{\partial}{\partial y} \right) p_C(t, y) &= p_D(t, 0, y), \\ \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} - \frac{\partial}{\partial y} \right) p_D(t, x, y) &= \lambda_s p_B(t, x) \frac{d}{dy} R_s(y) + \lambda p_C(t, y) \frac{d}{dx} R(x). \end{aligned} \quad (3.2)$$

Note that the initial condition  $N_0 = B$ ,  $Z_0 = 0$  implies that  $p_B(0, x) = d/dx R(x)$ ,  $x > 0$ . Moreover,  $\mathbf{P}\{\theta \leq t\} = p_A(t)$ .

Hence,

$$\mathbf{E}e^{-z\theta} = \int_0^\infty e^{-zt} dp_A(t), \quad z \geq 0. \quad (3.3)$$

#### 4. Solution procedure

It should be noted that our differential equations are well adapted to a Laplace-Fourier transformation. In fact, the  $p$ -functions are locally integrable with respect to  $t$  and bounded on  $[0, \infty)$ . Consequently, the derivatives with respect to  $t$  are also locally integrable.

Moreover, the integrability of the  $p$ -functions and the repair time densities with respect to  $x, y$  on  $[0, \infty)$  implies that the corresponding partial derivatives are also integrable on  $[0, \infty)$ . A Laplace-Fourier transform technique applied to the equations, taking the initial condition into account, reveals that for  $\text{Im } \omega \geq 0, \text{Im } \eta \geq 0, z > 0$ ,

$$\begin{aligned}
 \mathbf{E}e^{-z\theta} &= \int_0^\infty e^{-zt} p_B(t, 0) dt + \int_0^\infty e^{-zt} p_C(t, 0) dt, \\
 (\lambda_s + z + i\omega) \int_0^\infty e^{-zt} \mathbf{E}(e^{i\omega X_t} \mathbf{1}\{N_t = B\}) dt + \int_0^\infty e^{-zt} p_B(t, 0) dt \\
 &= \mathbf{E}e^{i\omega r} + \int_0^\infty \int_0^\infty e^{-zt} e^{i\omega x} p_D(t, x, 0) dx dt, \\
 (\lambda + z + i\eta) \int_0^\infty e^{-zt} \mathbf{E}(e^{i\eta Y_t} \mathbf{1}\{N_t = C\}) dt + \int_0^\infty e^{-zt} p_C(t, 0) dt \\
 &= \int_0^\infty \int_0^\infty e^{-zt} e^{i\eta y} p_D(t, 0, y) dy dt, \\
 (z + i\omega + i\eta) \int_0^\infty e^{-zt} \mathbf{E}(e^{i\omega X_t} e^{i\eta Y_t} \mathbf{1}\{N_t = D\}) dt \\
 &+ \int_0^\infty \int_0^\infty e^{-zt} e^{i\omega x} p_D(t, x, 0) dx dt + \int_0^\infty \int_0^\infty e^{-zt} e^{i\eta y} p_D(t, 0, y) dy dt \\
 &= \lambda \mathbf{E}e^{i\omega r} \int_0^\infty e^{-zt} \mathbf{E}(e^{i\eta Y_t} \mathbf{1}\{N_t = C\}) dt + \lambda_s \mathbf{E}e^{i\eta r_s} \int_0^\infty e^{-zt} \mathbf{E}(e^{i\omega X_t} \mathbf{1}\{N_t = B\}) dt.
 \end{aligned} \tag{4.1}$$

Adding (4.1) yields the functional equation:

$$\begin{aligned}
 (z + i\omega + i\eta) \int_0^\infty e^{-zt} \mathbf{E}(e^{i\omega X_t} e^{i\eta Y_t} \mathbf{1}\{N_t = D\}) dt \\
 + (z + \lambda(1 - \mathbf{E}e^{i\omega r}) + i\eta) \int_0^\infty e^{-zt} \mathbf{E}(e^{i\eta Y_t} \mathbf{1}\{N_t = C\}) dt \\
 + (z + \lambda_s(1 - \mathbf{E}e^{-i\eta r_s}) + i\omega) \int_0^\infty e^{-zt} \mathbf{E}(e^{i\omega X_t} \mathbf{1}\{N_t = B\}) dt \\
 = \mathbf{E}e^{i\omega r} - \mathbf{E}e^{-z\theta}, \quad \text{Im } \omega \geq 0, \text{Im } \eta \geq 0, z > 0.
 \end{aligned} \tag{4.2}$$

Inserting  $\omega = -\tau + iz, \eta = \tau, \tau \in \mathbb{R}$  into the functional equation entails that

$$\psi^+(\tau, z) - \psi^-(\tau, z) = \varphi(\tau, z), \tag{4.3}$$

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where

$$\begin{aligned}\psi^+(\tau, z) &:= \frac{1}{\gamma^+(\tau)} \int_0^\infty e^{-zt} \mathbf{E}(e^{i\tau Y_t} \mathbf{1}\{N_t = C\}) dt, \\ \psi^-(\tau, z) &:= \frac{\tau}{\tau - iz} \frac{1}{\gamma^-(\tau, z)} \int_0^\infty e^{-zt} \mathbf{E}(e^{-(i\tau+z)X_t} \mathbf{1}\{N_t = B\}) dt, \\ \varphi(\tau, z) &:= \frac{\mathbf{E}e^{-(i\tau+z)r} - \mathbf{E}e^{-z\theta}}{\gamma^+(\tau)\gamma^-(\tau, z)} \frac{1}{i\tau + z}, \\ \gamma^+(\tau) &:= 1 + \lambda_s \frac{\mathbf{E}e^{i\tau r_s} - 1}{i\tau}, \quad \gamma^+(0) := 1 + \lambda_s \mathbf{E}r_s, \\ \gamma^-(\tau, z) &:= 1 + \lambda \frac{1 - \mathbf{E}e^{-(i\tau+z)r}}{i\tau + z}, \quad z > 0.\end{aligned}\tag{4.4}$$

*Remarks 4.1.* Equation (4.3) constitutes a  $z$ -dependent *Hilbert* problem on the real line. Note that  $z$ -independent Hilbert problems, related to reliability engineering, have been solved by the theory of sectionally holomorphic functions. See [20, 21] for further details. A *similar* approach shows that the  $z$ -dependent function

$$\frac{1}{2\pi i} \int_\Gamma \varphi(\tau, z) \frac{d\tau}{\tau - \omega}, \quad \omega \in \mathbb{C}, z > 0,\tag{4.6}$$

is sectionally holomorphic in  $\mathbb{C}$ , provided that the singular Cauchy integral

$$\frac{1}{2\pi i} \int_\Gamma \varphi(\tau, z) \frac{d\tau}{\tau - u}, \quad u \in \mathbb{R},\tag{4.7}$$

is defined as a Cauchy principal value in a *double* sense, see [21, the Appendix]. Finally, note that our statement holds for general repair time distributions! (cf. [21, Remarks, page 361]).

### 5. The tail distribution

In order to obtain  $\mathbf{E}e^{-z\theta}$ , we first remark that  $\psi^+(\omega, z)$  is analytic in  $\mathbb{C}^+$ , boundedly continuous on  $\mathbb{C}^+ \cup \mathbb{R}$  and that

$$\lim_{\substack{|\omega| \rightarrow \infty \\ 0 \leq \arg \omega \leq \pi}} \psi^+(\omega, z) = 0.\tag{5.1}$$

Hence, by the Cauchy's theorem,

$$\frac{1}{2\pi i} \int_\Gamma \psi^+(\tau, z) \frac{d\tau}{\tau - \omega} = 0, \quad \omega \in \mathbb{C}^-.\tag{5.2}$$

On the other hand,  $\psi^-(\omega, z)$  is analytic in  $\mathbb{C}^-$ , boundedly continuous on  $\mathbb{C}^- \cup \mathbb{R}$  and

$$\lim_{\substack{|\omega| \rightarrow \infty \\ \pi \leq \arg \omega \leq 2\pi}} \psi^-(\omega, z) = 0.\tag{5.3}$$

Hence,

$$\psi^-(\omega, z) = -\frac{1}{2\pi i} \int_{\Gamma} \Psi^-(\tau, z) \frac{d\tau}{\tau - \omega}, \quad \omega \in \mathbb{C}^-. \quad (5.4)$$

By (4.3), (5.2), (5.4), we obtain

$$\psi^-(\omega, z) = \frac{1}{2\pi i} \int_{\Gamma} \varphi(\tau, z) \frac{d\tau}{\tau - \omega}, \quad \omega \in \mathbb{C}^-. \quad (5.5)$$

However, note that

$$\lim_{\substack{\omega \rightarrow 0 \\ \omega \in \mathbb{C}^-, z > 0}} \psi^-(\omega, z) = 0. \quad (5.6)$$

So that by (4.4)

$$\mathbf{E}e^{-z\theta} = \lim_{\substack{\omega \rightarrow 0 \\ \omega \in \mathbb{C}^-}} \frac{N^-(\omega, z)}{D^-(\omega, z)}, \quad (5.7)$$

where

$$N^-(\omega, z) := \frac{1}{2\pi i} \int_{\Gamma} \frac{\mathbf{E}e^{-(i\tau+z)r}}{(i\tau+z)\gamma^+(\tau)\gamma^-(\tau, z)} \frac{d\tau}{\tau - \omega}, \quad (5.8)$$

$$D^-(\omega, z) := \frac{1}{2\pi i} \int_{\Gamma} \frac{1}{(i\tau+z)\gamma^+(\tau)\gamma^-(\tau, z)} \frac{d\tau}{\tau - \omega}. \quad (5.9)$$

Consequently,  $\mathbf{E}e^{-z\theta}$  is completely determined.

*Example 5.1.* As an example, let  $R(t) = 1 - e^{-\rho t}$ ,  $\rho > 0$ , and  $R_s(t) = U_{t_0}(t)$ . Clearly,  $\mathbf{E}e^{-i\omega r} = \rho/(\rho + i\omega)$ ,  $\omega \neq i\rho$ , whereas  $\mathbf{E}e^{i\omega r_s} = e^{i\omega t_0}$ . Without loss of generality, we may take  $t_0$  as time unit. A straightforward application of the residue theorem entails that

$$\frac{1 - \mathbf{E}e^{-z\theta}}{z} = \frac{\alpha(z)}{\beta(z)}, \quad (5.10)$$

where

$$\begin{aligned} \alpha(z) &:= 1 + \lambda_s \frac{1 - e^{-z}}{z}, & \alpha(0) &:= 1 + \lambda_s, \\ \beta(z) &:= z + \rho + \lambda_s(1 - ae^{-z}), & a &:= \frac{\lambda + \rho e^{-(\lambda + \rho)}}{\lambda + \rho}. \end{aligned} \quad (5.11)$$

Hence,

$$\mathbf{E}\theta = \frac{1 + \lambda_s}{\rho + \lambda_s(1 - a)}. \quad (5.12)$$

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Note that the tail distribution  $\mathbf{P}\{\theta > t\}$  is *uniquely* determined by the Laplace transform:

$$\frac{1 - \mathbf{E}e^{-z\theta}}{z} = \int_0^\infty e^{-zt} \mathbf{P}\{\theta > t\} dt, \quad z > 0. \quad (5.13)$$

Moreover,  $\mathbf{P}\{\theta > t\}$  is Lebesgue absolutely continuous on  $(0, \infty)$ . Hence, by the inversion theorem

$$\mathbf{P}\{\theta > t\} = \frac{1}{2\pi i} \int_{C_\delta} e^{zt} \frac{\alpha(z)}{\beta(z)} dz, \quad t > 0, \quad (5.14)$$

where

$$\int_{C_\delta} \cdots dz := \lim_{T \rightarrow \infty} \int_{-iT+\delta}^{iT+\delta} \cdots dz, \quad \delta > 0. \quad (5.15)$$

A straightforward evaluation of the Cauchy integral, *similar* to the methodology introduced by Vanderperre [22], reveals that

$$\begin{aligned} \mathbf{P}\{\theta > t\} &= \sum_{k=0}^{\lfloor t \rfloor} a^k e^{-(\rho+\lambda_s)(t-k)} \frac{(\lambda_s(t-k))^k}{k!} \\ &+ \sum_{k=0}^{\lfloor t \rfloor} a^k \left( \frac{\lambda_s}{\rho+\lambda_s} \right)^{k+1} \left\{ 1 - e^{-(\rho+\lambda_s)(t-k)} \sum_{n=0}^k \frac{((\rho+\lambda_s)(t-k))^n}{n!} \right\} \\ &- \sum_{k=0}^{\lfloor t \rfloor - 1} a^k \left( \frac{\lambda_s}{\rho+\lambda_s} \right)^{k+1} \left\{ 1 - e^{-(\rho+\lambda_s)(t-k-1)} \sum_{n=0}^k \frac{((\rho+\lambda_s)(t-k-1))^n}{n!} \right\}. \end{aligned} \quad (5.16)$$

## 6. Conclusion

Our proposed priority system, subjected to general (bivariate) repair, can be analysed by elegant methods provided by the theory of sectionally holomorphic functions. However, the duplex system, subjected to general failure and repair time distributions, invokes an *open* (harsh) mathematical problem in the theory of statistical reliability engineering. The analysis of priority systems, subjected to arbitrary distributions, is far from complete.

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