

Research Article

Investigation into the Physical Mechanism and Influencing Factors of Sympathetic Inrush for Transformers in Series

Fang Peng^{1,2} and Houlei Gao^{1,2}

¹School of Electrical Engineering, Shandong University, Ji'nan, China ²Key Laboratory of Power System Intelligent Dispatch and Control of Ministry of Education, Shandong University, Ji'nan, China

Correspondence should be addressed to Houlei Gao; houleig@sdu.edu.cn

Received 20 July 2017; Revised 13 October 2017; Accepted 13 November 2017; Published 11 December 2017

Academic Editor: Guido Ala

Copyright © 2017 Fang Peng and Houlei Gao. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

Sympathetic inrush is a magnetizing inrush phenomenon generated in an in-service transformer, when a nearby transformer in series or in parallel is energized. This phenomenon lasts longer than the general inrush and can cause protection maloperation. The study focuses on sympathetic inrush in series transformers. Flux linkage expressions for the series transformers are derived using the Laplace transform. The residual flux, switching angle, and system impedance factors are included in the expressions. The physical mechanism of the sympathetic inrush for series transformers is explained based on the characteristics of the flux linkages. The interaction between the saturated transformers during sympathetic inrush is also investigated. Moreover, the influence of the switching angle, residual flux, system impedance, and load on the inrush currents is analyzed and validated.

1. Introduction

The magnetizing inrush phenomenon has been well known since the transformer was first used in energy conversion. This phenomenon may cause maloperation of the current differential protection of the transformer. Moreover, the large electromagnetic forces in the transformer windings can damage the insulation of the transformer [1], and the voltage decreases can lead to poor power quality [2, 3] during magnetizing inrush. Magnetizing inrush can be classified into energization inrush, sympathetic inrush, and recovery inrush [4]. Energization inrush occurs in the transformer being energized, as a result of the reapplication of the system voltage to the transformer. When one transformer is in service and the other (in series or in parallel) is energized, sympathetic inrush occurs in the in-service transformer. Energization inrush also occurs in the energized transformer and is called initial inrush hereinafter. Additionally, the recovery inrush is generated during the restoration of the voltage after clearance. The present study focuses on sympathetic inrush.

Several maloperation cases of current differential protection due to the sympathetic inrush have been reported [5]. Generally, the second harmonic component of the sympathetic inrush is high [6]. Thus, the second harmonic restraint algorithm is effective in preventing the protection of the in-service transformer from maloperation. However, the initial inrush and sympathetic inrush decay more slowly than the energization inrush in the single transformer due to the sympathetic interaction between the transformers [7]. Current transformer saturation caused by the long-lasting inrush current can decrease the second harmonic content, potentially causing the second harmonic restraint algorithm to fail [8–10]. Then, maloperation of the current differential protection of the in-service transformer will occur. Moreover, the sympathetic inrush may cause maloperation of other types of protection, such as the differential protection of the nearby generator and overcurrent protection, which leads to significant concern [8, 11].

The mechanism of sympathetic inrush for transformers in parallel has been investigated previously [12–14]. However, no study has been performed on transformers in series (called *series transformers* below). In this context, the present study aims to perform an analysis of sympathetic inrush for series transformers. Their flux linkage expressions are derived; then, the mechanism of sympathetic inrush and the interaction between the series transformers are explained using the characteristics of the flux linkages. Moreover, the impacts of several factors including the switching angle, residual flux, system impedance, and load are also investigated. A practical model is used in the case studies to demonstrate the theoretical analysis.

The rest of the paper is organized as follows. In Section 2, the equivalent model of series transformers is presented, and the flux linkage expressions of the series transformers are derived based on the Laplace transform. In Section 3, the physical mechanism of the sympathetic inrush for series transformers is explained. In Section 4, the characteristics of the flux linkages under various levels of magnetizing inductance and the effect of transformer saturation on the inrush currents are investigated. In Section 5, the influence of the switching angle, residual flux, system impedance, and load on the sympathetic inrush is analyzed and verified. Finally, conclusions are drawn in Section 6.

2. Equivalent Model of Series Transformers during Sympathetic Inrush

2.1. Series Transformers and Equivalent Model. Series transformers are generally used in a step-down distribution system connected to large and special power consumers, such as traction systems, steel mills, and electric boilers. The series transformers are shown in Figure 1. u_s is the power system equivalent source, and $Z_S = R_S + j\omega L_S$ is the equivalent impedance, where ω is the angular frequency. The circuit breaker (CB) is closed and T2 is energized. Initial inrush in T2 and sympathetic inrush in T1 are generated.

An equivalent T circuit is used to represent the transformer [15]. Thus, the equivalent circuit of series transformers is shown in Figure 2, where the symbols represent the following:

 R_{c1} and L_{m1} : the core loss resistance and magnetizing inductance of the iron core of T1,

 R_{c2} and L_{m2} : the core loss resistance and magnetizing inductance of the iron core of T2,

 R_{11} and R_{12} : the primary and secondary winding resistance of T1,

 R_{21} and R_{22} : the primary and secondary winding resistance of T2,

 $L_{\sigma 11}$ and $L_{\sigma 12}$: the primary and secondary winding leakage inductance of T1,

 $L_{\sigma 21}$ and $L_{\sigma 22}$: the primary and secondary winding leakage inductance of T2.

All parameters are normalized to the primary side of T1.

2.2. Flux Linkage Deviation. The equivalent circuit is simplified as in Figure 3 to investigate the transformer flux linkages. First, the current in the iron core consists of the core loss component and magnetizing component. The core loss resistance can be ignored since the core loss component is much smaller than the magnetizing component [13]. Second, the leakage inductance of the windings is far smaller than the magnetizing inductance of the iron core when the transformer is not saturated; therefore, L_{σ} can be neglected. Third, the secondary branch of T2 is inactive due to the lack of load. In addition, let $R_0 = R_S + R_{11}$, $R_1 = R_{12}$, $R_2 = R_{21}$,



FIGURE 1: Transformers connected in series.



FIGURE 2: Equivalent circuit of series transformers.



FIGURE 3: Simplified circuit of series transformers.

 $L_0 = L_s, L_1 = L_{m1}$, and $L_2 = L_{m2}$. Assume the source voltage varies sinusoidally with time; that is, $u_s = U_m \sin(\omega t + \varphi)$, where U_m and φ represent the amplitude and phase angle, respectively. The CB is closed at t = 0 s.

The nonlinear characteristics of the iron core can be modeled using a piecewise linear inductance with two different slopes [16, 17]. However, the magnetizing inductance is assumed to be linear to simplify the analysis [13]. The purpose of the study is to qualitatively understand the relationship between the flux linkages and inrush currents for series transformers. Equation (1) can be obtained by Kirchhoff's voltage law and Kirchhoff's current law:

$$U_{m}\sin(\omega t + \varphi) = R_{0}i_{0}(t) + L_{0}\frac{di_{0}(t)}{dt} + \frac{d\psi_{1}(t)}{dt}$$

$$\frac{d\psi_{1}(t)}{dt} = \frac{d\psi_{2}(t)}{dt} + (R_{1} + R_{2})i_{2}$$

$$i_{0}(t) = i_{1}(t) + i_{2}(t),$$
(1)

where $i_0(t)$, $i_1(t)$, and $i_2(t)$ represent the branch currents in L_0 , L_1 , and L_2 , respectively. $\psi_1(t) = L_1 i_1(t)$ and $\psi_2(t) = L_2 i_2(t)$ represent the flux linkages of T1 and T2, respectively.

The Laplace transform of (1) is used to calculate the flux linkages in the frequency domain. Let $R_3 = R_1 + R_2$. Then the flux linkage expressions in the time domain are obtained using the inverse Laplace transform, as presented in

$$\begin{split} \psi_1\left(t\right) &= -\frac{J}{2CFL_2} \left\{ e^{\tau_1 t} \left[1 \right. \right. \\ &\left. -\frac{A}{\sqrt{B^2/4 - AC}} \left(\frac{B}{2A} + \frac{AK}{J} \right) \right] + e^{\tau_2 t} \left[1 \right] \end{split}$$

$$+ \frac{A}{\sqrt{B^2/4 - AC}} \left(\frac{B}{2A} + \frac{AK}{J} \right) \right] \right\}$$

$$+ \frac{\sqrt{R_3^2 + (\omega L_2)^2}}{L_2 F} \left[-G \sin (\omega t + \varphi + \theta_1) + BD\omega \right]$$

$$\cdot \sin (\omega t + \varphi - \theta_1) \left[\psi_2 (t) = \frac{H}{2F} \cdot \left\{ e^{\tau_1 t} \left[1 - \frac{A}{\sqrt{B^2/4 - AC}} \left(\frac{B}{2A} - \frac{AI}{H} \right) \right] \right]$$

$$+ e^{\tau_2 t} \left[1 + \frac{1}{\sqrt{B^2/4 - AC}} \left(\frac{B}{2A} - \frac{AI}{H} \right) \right] + \frac{1}{F}$$

$$\cdot \omega \sqrt{(BD\omega)^2 + G^2} \sin (\omega t + \varphi + \theta_2), \qquad (2)$$

where

$$\begin{split} A &= L_0 L_1 + L_1 L_2 + L_0 L_2 \\ B &= R_0 L_1 + R_0 L_2 + R_3 L_0 + R_3 L_1 \\ C &= R_0 R_3 \\ D &= L_1 L_2 U_m \\ E &= -R_0 L_2 \left[\psi_1 \left(0_- \right) - \psi_2 \left(0_- \right) \right] \\ F &= A^2 \omega^4 - 2AC \omega^2 + B^2 \omega^2 + C^2 \\ G &= AD \omega^2 - CD \\ H &= \psi_2 \left(0_- \right) F + \omega G \\ I &= EF + C \left(G \sin \varphi - BD \omega \cos \varphi \right) \\ J &= -C \left[\left(GR_3 - BD \omega^2 L_2 \right) \sin \varphi \\ &+ \left(G \omega L_2 + BD \omega R_3 \right) \cos \varphi \right] + EFR_3 - CFL_2 \psi_2 \left(0_- \right) \\ K &= ACFR_3 \psi_2 \left(0_- \right) + EF \left(CL_2 - BR_3 \right) - \left(GL_2 \\ &+ BDR_3 \right) \sin \varphi + \left[AGR_3 + BD \left(BR_3 - CL_2 \right) \right] C \omega \\ \cdot \cos \varphi \\ \tau_1 &= \frac{-B + \sqrt{B^2 - 4AC}}{2A} \\ \theta_1 &= \arctan \frac{\omega L_2}{R_3} \\ \theta_2 &= \arctan \frac{-G}{BD \omega}. \end{split}$$

Here, τ_1 and τ_2 represent the decay coefficients of the DC components. $\psi_1(0_-)$ and $\psi_2(0_-)$ are the initial flux linkages of T1 and T2 at t = 0 s, respectively. $\psi_2(0_-)$ represents

the residual flux linkage of the energized transformer T2. Therefore, the effect of residual flux can be included in (2). Define $a_1, a_2, a_3, b_1, b_2, c_1$, and c_2 as in

$$\begin{aligned} a_{1} &= \frac{-G\sqrt{R_{3}^{2} + (\omega L_{2})^{2}}}{L_{2}F} \\ a_{2} &= \frac{BD\sqrt{R_{3}^{2} + (\omega L_{2})^{2}}}{L_{2}F} \\ a_{3} &= \frac{\omega\sqrt{(BD\omega)^{2} + G^{2}}}{F} \\ b_{1} &= -\frac{J}{2CFL_{2}} \left[1 - \frac{A}{\sqrt{B^{2}/4 - AC}} \left(\frac{B}{2A} + \frac{AK}{J} \right) \right] \quad (4) \\ b_{2} &= \frac{H}{2F} \left[1 - \frac{A}{\sqrt{B^{2}/4 - AC}} \left(\frac{B}{2A} - \frac{AI}{H} \right) \right] \\ c_{1} &= -\frac{J}{2CFL_{2}} \left[1 + \frac{A}{\sqrt{B^{2}/4 - AC}} \left(\frac{B}{2A} + \frac{AK}{J} \right) \right] \\ c_{2} &= \frac{H}{2F} \left[1 + \frac{A}{\sqrt{B^{2}/4 - AC}} \left(\frac{B}{2A} - \frac{AI}{H} \right) \right] . \end{aligned}$$

 a_1 , a_2 , and a_3 represent the amplitude of the AC components, while b_1 , b_2 , c_1 , and c_2 represent the amplitude of the DC components. Then, the transformer flux linkages are simplified as follows:

$$\psi_{1}(t) = a_{1} \sin (\omega t + \varphi + \theta_{1}) + a_{2} \sin (\omega t + \varphi - \theta_{1}) + b_{1} e^{\tau_{1} t} + c_{1} e^{\tau_{2} t}$$
(5)
$$\psi_{2}(t) = a_{3} \sin (\omega t + \varphi + \theta_{2}) + b_{2} e^{\tau_{1} t} + c_{2} e^{\tau_{2} t}.$$

Both $\psi_1(t)$ and $\psi_2(t)$ consist of the AC component and decaying DC component. The two AC components of $\psi_1(t)$ can be simplified as an AC component, which is explained in detail in Section 3. The DC components consist of the τ_1 component and τ_2 component for both transformers.

3. Transformer Flux Linkage Analysis

3.1. AC Component of the Flux Linkages. Before the CB is closed, the transformer T1 is not saturated. Since the impedance of the magnetizing branch is much larger than the system impedance in this situation, the former divides most of the source voltage. Thus, the flux linkage of T1 is approximately equal to the integral of the source voltage, as expressed in

$$\psi_1(t) = \int u_s(t) dt = \frac{U_m}{\omega} \cdot \sin\left(\omega t + \varphi - \frac{\pi}{2}\right).$$
 (6)

Only the AC component is included in the flux linkage of T1 before the CB is closed. $\psi_1(0_-) = -U_m/\omega$ can be obtained from (6).

After the CB is closed, the AC components of T1 and T2 can be obtained from (2) and (5). The magnetizing

inductance changes from several henries to several hundred henries for different types of transformers with different rated voltages, capacity, and so on. In comparison, the maximum winding resistance is several ohms, which means that $\omega L_2 \gg$ R_3 and $\theta_1 \approx \pi/2$. Thus, $\psi_{1AC}(t) = a_1 \sin(\omega t + \varphi + \theta_1) + a_2 \sin(\omega t + \varphi - \theta_1)$ can be written as

$$\psi_{1AC}(t) = \frac{D\sqrt{R_3^2 + (\omega L_2)^2} \left(A\omega^2 - C + BD\omega\right)}{L_2 \left[\left(A\omega^2 - C\right)^2 + \left(B\omega\right)^2\right]}$$
(7)
$$\cdot \sin\left(\omega t + \varphi - \frac{\pi}{2}\right) = a_4 \sin\left(\omega t + \varphi - \frac{\pi}{2}\right).$$

As for $\psi_{2AC} = a_3 \sin(\omega t + \varphi + \theta_2)$, $\theta_2 \approx -\pi/2$ can be obtained because $A\omega^2 - C \gg B$. Similarly, it can be verified that a_4 and a_2 are approximately equal to U_m/ω . Therefore, both ψ_{1AC} and ψ_{2AC} are approximately equal to the integral of the source voltage, that is, $U_m/\omega \cdot \sin(\omega t + \varphi - \pi/2)$. The variation of *R* and *L* will not change this feature.

Series transformers typically differ in their capacity and voltage level. After being normalized to the primary side of T1, a group of equivalent parameters is taken as an example:

$$R_{S} = 10 \Omega$$

$$L_{S} = 0.3 \text{ H}$$

$$R_{11} = 2 \Omega$$

$$L_{m1} = 67 \text{ H}$$

$$R_{12} = 2 \Omega$$

$$L_{m2} = 130 \text{ H}$$

$$R_{21} = 3 \Omega$$

$$\omega = 100\pi \text{ rad/s.}$$
(8)

The AC components of flux linkages for T1 and T2 can be obtained from (2) and (8), as shown in Figure 4. The base value is taken as the RMS value. The waveform is consistent with the analysis above.

3.2. Decaying DC Component. As shown in Figure 5, the decaying DC components of the flux linkages are also obtained from (2) and (8), where $\psi_{1\text{DC}} = b_1 e^{\tau_1 t} + c_1 e^{\tau_2 t}$ and $\psi_{2\text{DC}} = b_2 e^{\tau_1 t} + c_2 e^{\tau_2 t}$.

The DC components of the two flux linkages have opposite polarities. The amplitude of ψ_{2DC} continues decaying from the initial maximum ($|\psi_{2DC,max}| = 1.405$), while ψ_{1DC} increases first and then decreases ($|\psi_{1DC,max}| = 0.364$). The former has a peak value at the switching instant (t = 0 s), and the latter has a peak value after several seconds (t = 9.4 s). As shown in Figure 6, each flux linkage's DC component consists of two DC components with different decay coefficients (τ_1 component and τ_2 component, resp.).

Although the decay coefficients of ψ_{1DC} and ψ_{2DC} are identical, their polarities and amplitudes are different. As a result, ψ_{1DC} and ψ_{2DC} change in different ways. Two positive DC components ($b_2e^{\tau_1 t}$ and $c_2e^{\tau_2 t}$) exist in ψ_{2DC} . The superimposition of the two components is still positive and continually decays over time. In comparison, a positive



FIGURE 4: AC components of the flux linkages.







FIGURE 6: τ_1 and τ_2 components of the flux linkages.

TABLE 1: Parameters of T1 and T2.

	Rated capacity (MVA)	Voltage ratio (kV/kV)	Winding connection	Knee voltage (p.u.)	Magnetizing reactance (p.u.)
T1	31.5	220/6	Y/Y	1.12	0.14
T2	26	6.3/13.8	Δ/Y	1.15	0.3



FIGURE 7: Series transformers simulation model.

component $c_1 e^{\tau_2 t}$ and a negative component $b_1 e^{\tau_1 t}$ exist in ψ_{1DC} . Since the negative component initially decays more slowly than the positive component, the amplitude of the superimposition component ψ_{1DC} increases in the negative polarity. After several seconds, the negative component decays faster than the positive component; therefore, the amplitude of the superimposition component ψ_{1DC} decreases in the negative polarity. Furthermore, the initial value of ψ_{1DC} is approximately zero because the two components in opposite polarities have approximately equal amplitude at t = 0 s.

3.3. Sympathetic Inrush Case. A transformer typically operates near the knee point of the saturation curve. Only a single AC component is included in the transformer flux linkage when the transformer operates in the steady state. However, the superimposition of the DC component makes the transformer operate far above the knee point. Then, the transformer is driven into saturation and the magnetizing inrush is generated. If a positive DC component is added to the AC component in the positive half cycle, then positive inrush is generated. If a negative DC component is added to the AC component in the negative half cycle, negative inrush will be generated. During sympathetic inrush, the DC components of the two transformers have opposite polarities, while the AC components have the same polarity. Therefore, the initial inrush and sympathetic inrush are alternately generated in the positive and negative half cycle of the AC component. Further explanation of the process can be found in Section 4.1. The physical mechanism of the sympathetic inrush for series transformers is similar to that for parallel transformers [13].

A sympathetic inrush simulation model of series transformers is constructed as Figure 7 in Power Systems Computer Aided Design (PSCAD) according to a real maloperation case in China. The system parameters are $U_S = 235 \text{ kV}$ (RMS value of u_s), $\omega = 100\pi \text{ rad/s}$, and $Z_S = 30 + j283 \Omega$. The parameters of T1 and T2 are listed in Table 1, where the magnetizing reactance represents the reactance of the iron core that is fully saturated.

In the simulation, the CB is closed at t = 0.5 s. T2 is energized when T1 is in service. Taking Phase A as an example, the switching angle of Phase A is 0 ($\varphi_A = 0^\circ$). The magnetizing inrush currents corresponding to the inrush phenomenon can be obtained from the differential currents of T1 and T2. The differential currents I_{T1dA} and I_{T2dA}



FIGURE 8: Initial inrush and sympathetic inrush of series transformers.





are shown in Figure 8, which are normalized to the 6 kV side. Positive initial inrush is generated in T2 and negative sympathetic inrush is generated in T1. Only the Phase A currents are shown because the inrush currents of Phase A are the largest. The variation of the flux linkages is shown in Figure 9. ψ_1 consists of the AC component and negative DC component, while ψ_2 consists of the AC component and positive DC component.

3.4. Interaction between the Series Transformers. If there is a single in-service transformer in the power system, only an AC component will be included in the transformer flux linkage. ψ_{1DC} is generated due to the other energized transformer. Thus, it is the interaction between the series transformers that generates the sympathetic inrush.

The flux linkage of a single energized transformer $\psi_2(t)_{sg}$ can also be derived in a similar way to (2):

...

$$\psi_{2}(t)_{sg} = \frac{L_{2}U_{m}\sin(\omega t + \varphi - \theta_{3})}{\sqrt{(R_{4})^{2} + (\omega L_{3})^{2}}} + \left[\psi_{2}(0_{-}) - \frac{L_{2}U_{m}\sin(\varphi - \theta_{3})}{\sqrt{(R_{4})^{2} + (\omega L_{3})^{2}}}\right]e^{(-R_{4}/L_{3})t},$$
(9)

where $R_4 = R_0 + R_2$, $L_3 = L_0 + L_2$, and $\theta_3 = \arctan(\omega L_3/R_4)$.



FIGURE 10: DC component of the energized transformer.

An AC component ψ_{2ACsg} and a DC component ψ_{2DCsg} are included in the flux linkage. $\theta_3 \approx -\pi/2$ can be obtained, and ψ_{2ACsg} is from the sinusoidal source voltage. It can be seen from (2) and (9) that the in-service transformer changes the decay coefficient of the DC component and increases a DC component to the flux linkage of the energized transformer. ψ_{2DCsg} can be obtained from (8) and (9), as shown in Figure 10.

 ψ_{2DC} for series transformers is also shown in Figure 10 as a comparison. ψ_{2DC} decays more slowly than ψ_{2DCsg} . As a result, the initial inrush decays more slowly due to the inservice transformer.

4. Effect of Transformer Saturation on Inrush Currents

Although the linear magnetizing inductance assumption is made to simplify the transformer flux linkage analysis above, the magnetizing inductance of the saturated transformer is actually nonlinear. The magnetic characteristics of the transformer are typically represented using the piecewise lines. The magnetizing inductance will significantly decrease if the transformer is driven into saturation. In this section, the flux linkages of saturated transformers are analyzed to investigate the effect of transformer saturation.

4.1. Variation of Magnetizing Inductance during Sympathetic Inrush. The piecewise lines representing the magnetizing characteristics are illustrated in Figure 11(a). The magnetizing inductance below the knee point (nonsaturation zone) is L_{non} , and that above the knee point (saturation zone) is L_{sat} . Since the equivalent inductance around the knee point varies dramatically, L_{non} is much larger than L_{sat} . As shown in Figure 11(a), the addition of the DC components to the AC components during one cycle of the flux linkages can be used to explain the variation of the magnetizing inductance further. As analyzed in Section 3, the DC components of the two transformers have opposite polarities, and the AC components are approximately equal.

(1) Before t_1 , positive ψ_{2AC} and ψ_{2DC} are added together, but ψ_{2AC} is small, and thus $\psi_2 = \psi_{2AC} + \psi_{2DC}$ is below the knee point. Positive ψ_{1AC} and negative ψ_{1DC} also cause $\psi_1 = \psi_{1AC} + \psi_{1DC}$ to be below the knee point. Therefore, T1 and T2 are not saturated and the magnetizing inductance of each transformer is L_{non} . (2) As the instantaneous value of ψ_{AC} increases to the peak value during t_1 - t_2 , ψ_2 is above the knee point. T2 is driven into saturation while the positive ψ_{1AC} and negative ψ_{1DC} keep ψ_1 below the knee point. Therefore, the magnetizing inductance of T1 is L_{non} , while that of T2 is L_{sat} . (3) As the instantaneous value of ψ_{AC} decreases from t_2 , T2 enters the nonsaturation zone. As ψ_{AC} decreases further, T1 is driven into saturation in the negative half cycle in a similar way.

Therefore, T1 and T2 (T1/T2) go through four stages sequentially in each cycle during sympathetic inrush: nonsaturation/nonsaturation, nonsaturation/saturation, nonsaturation/nonsaturation, and saturation/nonsaturation. The corresponding four stages are for the periods $0-t_1$, t_1-t_2 , t_2-t_4 , and t_4 - t_5 . In the period t_5 - t_6 , T1 and T2 keep unsaturated to the next cycle. The magnetizing inductance values will change with different stages, as shown in Figure 11(b). When the sympathetic inrush is generated, the equivalent magnetizing inductance of the in-service transformer is alternately presented at L_{non} and L_{sat} . The same effect occurs for the magnetizing inductance of the energized transformer when the initial inrush is generated. In addition, L_{non} and L_{sat} of the two transformers are actually different. L_{non1} and L_{sat1} are the magnetizing inductance of T1 while L_{non2} and L_{sat2} are the magnetizing inductance of T2.

For the AC components, the decrease of the magnetizing inductance will not change the feature of ψ_{1AC} and ψ_{2AC} because the conditions $\omega L_2 \gg R_3$ and $A\omega^2 - C \gg B$ still hold true, even if the magnetizing inductance decreases. Therefore, ψ_{1AC} and ψ_{2AC} are both approximately equal to $U_m/\omega \cdot \sin(\omega t + \varphi - \pi/2)$. The saturation of the transformers has a slight impact on the AC components.

For the DC components, the magnetizing inductance is assumed to be linear in the DC component analysis in Section 3. This assumption can be explained further as follows. When the CB is closed, the DC component of the flux linkage is generated to prevent the flux linkage from changing discretely. At the switching instant, the magnetizing inductance is L_{non} , in the period $0-t_1$. Therefore, the variation feature of the DC component is determined predominantly by L_{non} at the switching instant. The main characteristic of the flux linkage will not change even if the transformer is driven into saturated saturation. The reason is that the circuit is no longer forced to change and there is no new DC component generated to prevent a discrete change in the flux linkage. The effect of transformer saturation is to change the amplitude and decay rate to some degree.

Further research is implemented to analyze the influence of the decrease of magnetizing inductance on the DC components.

4.2. DC Components of Flux Linkages When T2 Is Saturated. When T2 is saturated at the t_1 - t_2 stage, L_{m2} is reduced to L_{sat2} , while $L_{m1} = L_{non1}$. The DC components of the flux linkages are obtained by taking a smaller L_{m2} into (2), where $L_{sat2} = 60$ H or 20 H. Although the stage of $L_{m2} = L_{sat2}$ is only a short period in a cycle, the effect is analyzed qualitatively using L_{sat2} in the entire decay process. The DC components of the flux linkages are shown in Figure 12. ψ_{2DC} decays more quickly in the main decay process and ψ_{1DC} has an



FIGURE 11: (a) The flux linkages and piecewise linear magnetization curve. (b) The equivalent magnetizing inductance during sympathetic inrush.

TABLE 2: Decay coefficients varying with L_{m2} .

L_{m2} (H)	130	60	20
$ au_1$	-0.0241	-0.0349	-0.0455
$ au_2$	-0.2839	-0.4236	-0.9650

earlier and larger peak value with the decrease in L_{m2} . The reason is that the faster decaying component of ψ_{1DC} (τ_2 component) decays much faster and the amplitude of each component of ψ_{1DC} (τ_1 and τ_2 components) increases. The τ_1 and τ_2 components under different L_{m2} values are shown in Figure 13, and the decay coefficients are listed in Table 2. $|b_1|, |c_1|, \text{and } |c_2|$ increase, while $|b_2|$ decreases with decreasing L_{m2} .

As a result, the initial inrush and sympathetic inrush are influenced. The initial inrush will decay more quickly, and the sympathetic inrush will be strengthened.

4.3. DC Components of Flux Linkages When T1 Is Saturated. When T1 is saturated at the t_4 - t_5 stage, L_{m1} is reduced to L_{sat1} , while $L_{m2} = L_{non2}$. The DC components of the flux linkages are then obtained by taking a smaller L_{m1} into (2), where $L_{sat1} = 30$ H or 10 H. The DC components are shown in Figure 14. ψ_{2DC} decays more slowly, and ψ_{1DC} decreases with decreasing L_{m1} . When L_{m1} decreases to 10 H, ψ_{1DC} is too small to lead to the sympathetic inrush. τ_1 and τ_2 components of the flux linkages are shown in Figure 15. $|b_1|$, $|c_1|$, and $|c_2|$ decrease while $|b_2|$ increases with decreasing L_{m1} . As a result, the decay rate of each component increases. ψ_{2DC} decays more slowly with decreasing L_{m1} because the more slowly decaying component of ψ_{2DC} (τ_1 component) has a larger amplitude.

In regard to the inrush currents of the transformers, the sympathetic inrush is weakened and the initial inrush decays more slowly due to the saturation of T1.

In conclusion, the series transformers interact with each other due to the transformer saturation. The initial inrush



FIGURE 12: DC components of the flux linkages under different L_{m2} .

decays more quickly due to the saturation of the energized transformer, and the sympathetic inrush is weakened due to saturation of the in-service transformer. The saturation of the transformer resists further saturation, which can be thought of as a self-regulating mechanism. On the other hand, the initial inrush decays more slowly due to saturation of the in-service transformer, and the sympathetic inrush is strengthened due to saturation of the energized transformer. The saturation of one transformer strengthens the other transformer's saturation, which can be explained in terms of the energy balance of the two transformers.





FIGURE 13: τ_1 and τ_2 components of the flux linkages under different L_{m2} .



FIGURE 14: DC components of the flux linkages under different L_{m1} .

5. Influencing Factors of Sympathetic Inrush

External conditions for the series transformers that have been connected to the power system vary greatly during sympathetic inrush. The switching angle and residual flux of the energized transformer vary widely, and their values depend on the switching time of the CB. Moreover, the system impedance and the load vary greatly under different



FIGURE 15: τ_1 and τ_2 components of the flux linkages under different L_{m1} .

operating conditions of the power system. Therefore, their impacts on the sympathetic inrush are investigated in this section.

5.1. Switching Angle. For a single energized transformer without residual flux, the flux linkage has a maximum DC component when the CB is closed at $\varphi = 0^{\circ}$ because the prospective flux linkage is at the maximum, and the existing flux linkage cannot be changed suddenly. The flux linkage has a minimum DC component when the CB is closed at $\varphi = 90^{\circ}$.

Similarly, the largest initial and sympathetic inrush currents are generated at $\varphi = 0^{\circ}$ and the smallest are generated at $\varphi = 90^{\circ}$ without residual flux. By taking different φ in (2), the DC components of the flux linkages are obtained and are shown in Figure 16. $|b_1|$, $|b_2|$, $|c_1|$, and $|c_2|$ all decrease when φ varies from 0° to 90°. Therefore, the magnetizing inrush and sympathetic inrush are both weakened as φ increases and will cease at $\varphi = 90^{\circ}$.

A simulation model (shown in Figure 7) is constructed in PSCAD. The inrush currents can be represented by the differential currents, as shown in Figure 17, which are normalized to the 6 kV voltage side. A Y/Y energized transformer is used in this case instead of the Δ /Y transformer because the variation feature under different switching angles can be observed more clearly using a Y/Y transformer. It is observed that the variation of the inrush currents is the same as the variation of the DC components of the flux linkages.



FIGURE 16: Variation of the DC components of the flux linkages with the switching angle.

5.2. Residual Flux. The residual flux of the energized transformer is another important factor that influences the inrush phenomenon. The difference between the residual flux linkage and prospective flux linkage of the energized transformer determines the amplitude of the initial inrush and sympathetic inrush. If the difference is large, the initial inrush and sympathetic inrush are also large. As shown in Figure 18, the DC components of flux linkages vary with the residual flux ($\varphi = 0^{\circ}$). When the CB is closed at $\varphi = 0^{\circ}$, the prospective flux linkage is approximately equal to $-U_m/\omega$. The DC components are the largest when the residual flux is U_m/ω and are zero when the residual flux is $-U_m/\omega$. A simulation case is conducted in Real Time Digital Simulator (RTDS) to further demonstrate the analysis. The inrush currents are shown in Figure 19.

5.3. System Impedance. The system impedance $Z_S = |R_S + j\omega L_S|$ also has a significant effect on the inrush currents. The DC components of the flux linkages are obtained from (2) and are shown in Figure 20. ψ_{1DC} has larger and earlier peak values as Z_S increases. A larger Z_S is more beneficial for the generation of the sympathetic inrush. ψ_{2DC} is smaller in the main decay process as Z_S increases, particularly before t = 8 s. The system impedance has a greater effect on the sympathetic inrush than the initial inrush.

The inrush currents varying with the system impedance by simulation conducted in PSCAD are shown in Figure 21. The larger Z_S weakens the initial inrush and causes the sympathetic inrush to have a larger and earlier peak value, validating the theoretical analysis.

5.4. Load of the In-Service Transformer. The load of the inservice transformer is also an important factor, especially for series transformers. The load adds an impedance branch to the right side of $L_{1\sigma}$. The added branch restrains the influence



FIGURE 17: Variation of the initial inrush and sympathetic inrush with the switching angle.



FIGURE 18: Variation of the DC components of the flux linkages with residual flux.

of the initial inrush on the in-service transformer. Although the flux linkage expressions including the load parameter have not been obtained, a simulation can be conducted to analyze the effect of the load. As shown in Figure 22, the initial inrush and sympathetic inrush decrease with increasing load, and the sympathetic inrush may disappear under a heavy load.

In contrast, the maximum amplitude of the sympathetic inrush generated in the parallel transformers at $P_L = 30 \text{ MW}$



FIGURE 19: Variation of the initial inrush and sympathetic inrush with residual flux.



FIGURE 20: Variation of the DC components of the flux linkages with system impedance.

exhibits a 15% decrease relative to that at $P_L = 0$ MW. Here, the type of parallel transformers in the simulation is the same as T1. The results show that the influence of the load on the series transformers is more significant than that on the parallel transformers.

6. Conclusion

This paper presents a thorough study of the sympathetic inrush for series transformers. Flux linkage expressions of series transformers, including the switching angle, residual



FIGURE 21: Variation of the initial inrush and sympathetic inrush with system impedance.



FIGURE 22: Variation of the initial inrush and sympathetic inrush with the load.

flux, and system impedance factors, are derived using the Laplace transform. Flux linkages of series transformers consist of the AC and DC components when the sympathetic inrush occurs. The series transformers interact with each other during the sympathetic inrush. Saturation of each transformer resists further self-saturation and strengthens the other transformer's saturation.

The influencing factors of sympathetic inrush for series transformers, including the switching angle, residual flux, and system impedance, are analyzed and validated. The smaller switching angle $(0^{\circ}-90^{\circ})$, larger impedance, and lighter load are more beneficial for the generation of the

sympathetic inrush. If the difference between the residual flux in the energized transformer and the integral of the energizing voltage is large, the initial inrush and sympathetic inrush are also large.

Conflicts of Interest

The authors declare that there are no conflicts of interest regarding the publication of this article.

References

- M. Steurer and K. Fröhlich, "The impact of inrush currents on the mechanical stress of high voltage power transformer coils," *IEEE Transactions on Power Delivery*, vol. 17, no. 1, pp. 155–160, 2002.
- [2] M. Nagpal, T. G. Martinich, A. Moshref, K. Morison, and P. Kundur, "Assessing and limiting impact of transformer inrush current on power quality," *IEEE Transactions on Power Delivery*, vol. 21, no. 2, pp. 890–896, 2006.
- [3] J. Peng, H. Li, Z. Wang, F. Ghassemi, and P. Jarman, "Influence of sympathetic inrush on voltage dips caused by transformer energisation," *IET Generation, Transmission & Distribution*, vol. 7, no. 11, pp. 1173–1184, 2013.
- [4] V. Vaddeboina, G. Taylor, and C. Proudfoot, "Switching large transformers on weak transmission networks—a real time case study," in *Proceedings of the 47th International Universities Power Engineering Conference (UPEC '12)*, September 2012.
- [5] X. Qi, X. Yin, and Z. Zhang, "Sympathetic inrush current in a transformer and a method for its identification," *IEEJ Transactions on Electrical and Electronic Engineering*, vol. 11, no. 4, pp. 442–450, 2016.
- [6] H. S. Bronzeado, P. B. Brogan, and R. Yacamini, "Harmonic analysis of transient currents during sympathetic interaction," *IEEE Transactions on Power Systems*, vol. 11, no. 4, pp. 2051– 2056, 1996.
- [7] H. Bronzeado and R. Yacamini, "Phenomenon of sympathetic interaction between transformers caused by inrush transients," *IEE Proceedings - Science, Measurement and Technology*, vol. 142, no. 4, pp. 323–329, 1995.
- [8] X. Qi, X. Yin, Z. Zhang, D. Chen, Y. Wang, and F. Cai, "Study on the unusual misoperation of differential protection during transformer energization and its countermeasure," *IEEE Transactions on Power Delivery*, vol. 31, no. 5, pp. 1998–2007, 2016.
- [9] F. Wang, X. Qi, Y. Wang et al., "A sympathetic inrush detection method based on on-line magnet flux calculation," in *Proceedings of the 2014 International Conference on Power System Technology (POWERCON '14)*, pp. 700–704, October 2014.
- [10] T. Zheng, J. Gu, S. F. Huang, F. Guo, and V. Terzija, "A new algorithm to avoid maloperation of transformer differential protection in substations with an inner bridge connection," *IEEE Transactions on Power Delivery*, vol. 27, no. 3, pp. 1178– 1185, 2012.
- [11] U. Rudez and R. Mihalic, "A reconstruction of the WAMSdetected transformer sympathetic inrush phenomenon," *IEEE Transactions on Smart Grid.*
- [12] C. D. Hayward, "Prolonged inrush currents with parallel transformers affect differential relaying," *Electrical Engineering*, vol. 60, no. 12, pp. 1096–1101, 1941.

- [13] U. Rudez and R. Mihalic, "Sympathetic inrush current phenomenon with loaded transformers," *Electric Power Systems Research*, vol. 138, pp. 3–10, 2016.
- [14] P. Heretik, M. Kovac, A. Belan, V. Volcko, and M. Konicek, "Analytical analysis of transformer interaction caused by inrush current," in *Proceedings of the 15th International Scientific Conference on Electric Power Engineering (EPE '14)*, pp. 23–28, Brno-Bystrc, Czech Republic, May 2014.
- [15] A. Tokić, I. Uglešić, and G. Štumberger, "Simulations of transformer inrush current by using BDF-based numerical methods," *Mathematical Problems in Engineering*, vol. 2013, Article ID 215647, 10 pages, 2013.
- [16] J. A. Martinez and B. A. Mork, "Transformer modeling for lowand mid-frequency transients—a review," *IEEE Transactions on Power Delivery*, vol. 20, no. 2, pp. 1625–1632, 2005.
- [17] X. Zhang, "An efficient algorithm for transient calculation of the electric networks including magnetizing branches," *Mathematical Problems in Engineering*, vol. 2011, Article ID 479626, 11 pages, 2011.











Journal of Probability and Statistics







Submit your manuscripts at https://www.hindawi.com



International Journal of Differential Equations





Journal of Complex Analysis











Mathematical Problems

Journal of **Function Spaces**



Abstract and Applied Analysis



International Journal of Stochastic Analysis



Discrete Dynamics in Nature and Society

